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PHASE BOUNDARIES OF ANTIFERROMAGNETIC Cs$_2$[FeCl$_5$(H$_2$O)]

A. PADUAN FILHO (*), F. PALACIO (**) and R. L. CARLIN
Department of Chemistry, University of Illinois at Chicago Circle, Chicago, Illinois 60680, USA

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Résumé. — La susceptibilité à champ nul de Cs$_2$[FeCl$_5$(H$_2$O)] peut être exprimée par le développement en série à haute température du modèle de Heisenberg à trois dimensions pour $S = 5/2$, avec $J/k = -0.31$ K. Certaines portions des limites de phases magnétiques ont été mesurées, montrant que $H_{\text{sf}}(0) = 11.5$ kOe, $H_A/H_E = 1.2 	imes 10^{-2}$, et que le point bicritique se trouve à environ 6.3 K et 14.7 kOe. Puisque $T_c(0) = 6.57$ K, un ordre de grandeur plus bas que pour MnF$_2$ ou RbMnF$_3$, c'est un nouvel exemple important d'antiferromagnétisme de Heisenberg à trois dimensions.

Abstract. — The zero-field susceptibility of Cs$_2$[FeCl$_5$(H$_2$O)] may be fit by the 3-D Heisenberg, $S = 5/2$, high-temperature series expansion, with $J/k = -0.31$ K. Portions of the magnetic phase boundaries have been measured, with the result that $H_{\text{sf}}(0) = 11.5$ kOe, the bicritical point occurs at ~ 6.3 K and ~ 14.7 kOe, and $H_A/H_E = 1.2 	imes 10^{-2}$. Since $T_c(0) = 6.57$ K, which is an order of magnitude lower than that of either MnF$_2$ or RbMnF$_3$, this is an important new example of the 3-D Heisenberg antiferromagnet.

It was recently reported [1] that Cs$_2$[FeCl$_5$(H$_2$O)] is antiferromagnetic with a critical temperature at zero-field, $T_c(0)$, of 6.57 ± 0.05 K. Both zero-field susceptibility data, $\chi_0$, and specific heat results were presented for this $S = 5/2$, orthorhombic material. Some difficulty was experienced in fitting the susceptibility data to the various magnetic model calculations that are available [2], but it was clear that the compound has fairly small magnetic anisotropy. Data are presented here which allow a quantitative measure of this anisotropy, and which also show that this new material is a good example of the three-dimensional Heisenberg model [2, 3].

A careful analysis of the $a$-axis data [1] has led to the value $(T_{\text{max}} - T_c)/T_{\text{max}} = 0.088$, where $T_{\text{max}}$ is the temperature at which $\chi_0$ reaches its maximum value and $T_c(0)$ is that temperature at which $\partial \chi_0/\partial T$ has its maximal value, that is, the critical temperature.

The $a$-axis is the preferred axis of spin-alignment. This value of the above ratio is typical of that found with other three-dimensional Heisenberg antiferromagnets (3-D HA-F) such as RbMnF$_3$ (which is more nearly isotropic) and agrees well with the theoretical expectation of 0.08 [4]. More importantly, we illustrate in figure 1 the fit to the susceptibility data [1] of the high-temperature series expansion for the susceptibility of the s.c., $S = 5/2$ HA-F, as provided by de Jongh and Breed [4], the series being expanded in terms of $t = kT/J | S(S + 1)$. In order to observe the influence of the final term in the series for small values of the parameter $t$ and to find the minimum temperature down to which we might extend the calculation of the susceptibility, we have calculated it with all the terms available ($n = 7$) as well as with $n = 6$ and 5. Our results are essentially the same as those obtained in reference [4], hence we have extended our calculations down to $t = 2.4$ as shown in figure 1. The value of

$$J/k = 0.310 \pm 0.005 \text{ K}.$$
The points are the experimental data of O'Connor [1], and the solid curve is the calculated high temperature series expansion, as described in the text.

a value which may be compared (assuming \( z = 6 \)) with the values 0.29-0.33 K reported earlier [1], and obtained by molecular field and spin-wave (\( T = 0 \)) calculations. As may be seen from figure 1 there is good agreement between calculated and experimental points above \( T \approx 3.0 \). Below this point the agreement is within 2%.

No attempt was made to use extrapolation procedures for the calculated values of the susceptibility would not improve for \( t < 2.8 \), as they would be lower than the calculated ones. A possible reason for this small discrepancy is that the system is orthorhombic rather than an ideal simple cubic one.

Data were taken by sweeping the applied magnetic field while the temperature was kept constant or increasing the temperature at constant magnetic field. At constant temperature, the critical field was determined by a sharp peak in the susceptibility when the field crossed the phase boundary. In constant field, the phase transition was determined by the maximum value of \( \frac{d\chi}{dT} \) for \( H < H_b \) or by a discontinuity in the susceptibility for \( H > H_b \). In the figure, typical error bars reflect an uncertainty in the data estimated from the value of the critical field or the localization of the transition temperature.

The bicritical point is found at \( T_b = 6.3 \pm 0.2 \) K and \( H_b = 14.7 \pm 0.1 \) kOe. Taking the spin-flop field extrapolated to zero temperature as \( H_{sf}(0) = 11.5 \) kOe and \( \chi_{sp}(0) = 0.19 \) emu/mole [1], we estimate the anisotropy field \( H_A = 0.88 \) kOe and the exchange field \( H_E = 75.9 \) kOe. The ratio \( \alpha = H_A / H_E = 1.2 \times 10^{-2} \).

The AF-P boundary is quite vertical \( (T_a'/T_a(0) = 0.96) \) as has been observed in such other low-anisotropy antiferromagnets as CuCl\(_2\).2 H\(_2\)O [6] and MnF\(_2\) [7]. The main source of anisotropy in this material probably is the particular crystal structure, which is under detailed investigation [8]. Single-ion anisotropy in this type of compound appears to be small [1, 9].

We conclude that Cs\(_2\)[FeCl\(_4\)(H\(_2\)O)] is about as good an example of the 3-D HA-F as is the iso-electronic MnF\(_2\), and comparable in anisotropy to the \( S = 1/2 \) CuCl\(_2\).2 H\(_2\)O. What is interesting about this result is that \( T_c(MnF_2) \) is about an order of magnitude higher than that of the compound reported here, while \( T_c \) of RbMnF\(_3\), the least-anisotropic 3-D HA-F, is even higher, 83 K [4]. Thus, a variety of experiments can now be carried out on an \( S = 5/2 \) 3-D HA-F more conveniently than in the past; one such example is the behaviour of the magnetic specific heat, for the lattice contribution in

![Fig. 1.](image1)

![Fig. 2.](image2)
Cs₂[FeCl₃(H₂O)] at Tₑ is a small fraction of the total. Nuclear resonance and Mössbauer experiments also are particularly attractive. These and other experiments on Cs₂[FeCl₅(H₂O)] and its congeners are now in progress.

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