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# INVESTIGATION OF THE QUADRUPOLE DEFORMATION OF ${ }^{11}$ B BY MEANS OF $30 \mathrm{MeV} \cdot$ POLARIZED PROTON INELASTIC SCATTERING 

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#### Abstract

Résumé. - Les sections efficaces et pouvoirs d'analyse de la diffusion inélastique ${ }^{11} \mathbf{B}\left(\mathbf{p}, \mathrm{p}^{\prime}\right)$ à $E_{\mathrm{p}}=30,3 \mathrm{MeV}$ ont été analysés dans le formalisme des équations couplées. Ces calculs suggèrent la valeur positive de la déformation quadrupolaire $\beta_{2} \mathrm{du}{ }^{11} \mathrm{~B}$ (prolate) et donnent le résultat suivant $\beta_{2}=+0,52$.


#### Abstract

The cross-sections and resolving powers of the ${ }^{11} \mathrm{~B}\left(\mathrm{p}, \mathrm{p}^{\prime}\right)$ inelastic scattering at $E_{\mathrm{p}}=30.3 \mathrm{MeV}$ are analyzed in the coupled-channels formalism. These calculations suggest a positive value for the quadrupole deformation $\beta_{2}$ of ${ }^{11} \mathrm{~B}$ (prolate) and give the result $\beta_{2}=+0.52$.


In the understanding of 1 p shell nuclei, the investigation of their deformation plays an important role. For the ${ }^{11} \mathrm{~B}$ nucleus, Hartree-Fock calculations [1] do not give a prolate lower minimum compatible with the positive electric quadrupole moment obtained from experiments. For this nucleus a strong-coupling rotational model [2] has given a better result although a quantitative disagreement with the experimentally determined electric quadrupole moment still remains [3].

The above discrepancies have suggested that we need much more investigations about the quadruple deformation of the ${ }^{11} \mathrm{~B}$ nucleus by means of inelastic scattering. In particular a recent investigation of the quadrupole deformation of ${ }^{11} \mathrm{~B}$ by inelastic helion $\left({ }^{3} \mathrm{He}\right)$ scattering at $E_{3 \mathrm{He}}=74 \mathrm{MeV}$ [4] has shown, with analysis using the coupled-channels (CC) method, the possible existence of oblate-prolate effects of ${ }^{11} \mathrm{~B}$ in this reaction. It would therefore appear necessary to determine the quadrupole deformation of ${ }^{11} \mathrm{~B}$ by means of polarized proton inelastic scattering.

In view of the determination of the sign and the value for the quadrupole deformation $\beta_{2}$ of ${ }^{11} \mathrm{~B}$, we have analyzed, in the coupled-channels (CC) forma-
lism with the rotational model using the code ECIS 75 [5], the experimental data for the cross-sections and resolving powers in the ${ }^{11} B(p, p)$ and ${ }^{11} B\left(p, p^{\prime}\right)$ scattering to the lower two members of the $\mathrm{K}^{\pi}=\frac{3^{-}}{2}$ band of ${ }^{11} \mathrm{~B}$, i.e. the $\frac{3}{2}^{-}$ground state and the $\frac{5}{2}^{-}$secondexcited state $\left(E_{x}=4.46 \mathrm{MeV}\right)$ at $E_{\mathrm{p}}=30.3 \mathrm{MeV}[6]$. The optical parameters used as initial values for the optical model search procedure were taken from the analysis performed by Karban et al. [6] and are listed in table I. In the CC formalism, the nuclear radius is defined by

$$
R=R_{i}\left(1+\beta_{2} Y_{20}+\cdots\right)
$$

where the $\beta$ 's are the deformation parameters determined by the experiment, the $Y$ 's are spherical harmonics and $R_{i}$ corresponds to the various optical potential radii. The interaction potential arises from the deformation of the Coulomb potential, the complex central potential and the spin-orbit potential. The deformed spin-orbit potential was of the full Thomas form [7]. In the CC calculations, the states explicitly coupled are the lower two members of a

Table I
Optical model parameters used in the analysis of the ${ }^{11} \mathrm{~B}(\mathrm{p}, \mathrm{p})^{11} \mathrm{~B}$ scattering

| $V_{0}$ <br> $(\mathrm{MeV})$ | $r_{0}$ <br> $(\mathrm{fm})$ | $a_{0}$ <br> $(\mathrm{fm})$ | $W_{\mathbf{V}}$ <br> $(\mathrm{MeV})$ | $W_{\mathrm{D}}$ <br> $(\mathrm{MeV})$ | $r_{\mathrm{I}}$ <br> $(\mathrm{fm})$ | $a_{\mathrm{I}}$ <br> $(\mathrm{fm})$ | $V_{\mathrm{s0}}$ <br> $(\mathrm{MeV})$ | $r_{\mathrm{s} 0}$ <br> $(\mathrm{fm})$ | $a_{\mathrm{s} 0}$ <br> $(\mathrm{fm})$ | $r_{\mathrm{C}}$ <br> $(\mathrm{fm})$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| - | - | - | - | - | - | - | - | - | - | - |
| 45.18 | 1.09 | 0.59 | 0 | 3.38 | 1.30 | 1.01 | 7.78 | 0.98 | 0.57 | 1.09 |

$\mathrm{K}=\frac{3}{2}^{-}$rotational band in ${ }^{11} \mathrm{~B}$. The results are presented in figure 1 and the corresponding parameters listed in table II.

The two values $\beta_{2}=+0.43$ and $\beta_{2}=-0.50$ obtained from reference [4] by analyzing only the cross-sections of the ${ }^{11} \mathrm{~B}\left({ }^{3} \mathrm{He},{ }^{3} \mathrm{He}\right)$ inelastic scattering at $E_{3_{\mathrm{He}}}=74 \mathrm{MeV}$ with the CC method give equally low $\chi^{2}$ values. But it should be mentioned [4] that $\beta_{2}=+0.43$ agrees quite well with the experimental value of +0.0372 b [3] for the electric quadrupole moment. The results we have obtained by analyzing simultaneously the cross-sections and resolving powers of the ${ }^{11} \mathbf{B}\left(\mathbf{p}, \mathrm{p}^{\prime}\right)$ inelastic scattering at $E_{\mathrm{p}}=30.3 \mathrm{MeV}$ using the CC calculations suggest also a positive value for the quadrupole deformation $\beta_{2}$ of ${ }^{11} \mathrm{~B}$ (prolate) and give the result $\beta_{2}=+0.52$.

We are grateful to Dr. R. de Swiniarski for valuable discussions and his interest in this work.


Fig. 1. - Experimental results of the ${ }^{11} \mathrm{~B}\left(\mathrm{p}, \mathrm{p}^{\prime}\right)^{11} \mathrm{~B}^{*}$ scattering compared to the results of the coupled-channels calculations corresponding to the two parameter sets of table II.

Table II
Coupled-channel parameters used in the analysis of the ${ }^{11} \mathrm{~B}\left(\mathbf{p}, \mathrm{p}^{\prime}\right)^{11} \mathrm{~B}^{*}$ inelastic scattering

| $\beta_{2}$ | $V_{0}$ <br> $(\mathrm{MeV})$ | $r_{0}$ <br> $(\mathrm{fm})$ | $a_{0}$ <br> $(\mathrm{fm})$ | $W_{\mathbf{V}}$ <br> $(\mathrm{MeV})$ | $W_{\mathrm{D}}$ <br> $(\mathrm{MeV})$ | $r_{\mathrm{I}}$ <br> $(\mathrm{fm})$ | $a_{\mathbf{I}}$ <br> $(\mathrm{fm})$ | $V_{\mathrm{so}}$ <br> $(\mathrm{MeV})$ | $r_{\mathrm{s0}}$ <br> $(\mathrm{fm})$ | $a_{\mathrm{s} 0}$ <br> $(\mathrm{fm})$ | $r_{\mathrm{C}}$ <br> $(\mathrm{fm})$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| - | - | - | - | - | - | - | - | - | - | - | - |
| -0.52 | 46.65 | 1.09 | 0.59 | 0 | 3.22 | 1.30 | 1.01 | 8.38 | 0.98 | 0.57 | 1.09 |
| -0.60 | 46.98 | 1.09 | 0.59 | 0 | 3.34 | 1.30 | 1.01 | 8.34 | 0.98 | 0.57 | 1.09 |

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