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ROI-Wise Material Decomposition in Spectral Photon-Counting CT

Bingqing Xie, Pei Niu, Ting Su, Valérie Kaftandjian, Loic Boussel, Philippe Douek, Feng Yang, Philippe Duvauchelle and Yuemin Zhu

Abstract—Spectral photon-counting X-ray CT (sCT) opens up new possibilities for the quantitative measurement of materials in an object, compared to conventional energy-integrating CT or dual energy CT. However, achieving reliable and accurate material decomposition in sCT is extremely challenging, due to similarity between different basis materials, strong quantum noise and photon-counting detector limitations. We propose a novel material decomposition method that works in a region-wise manner. The method consists in optimizing basis materials based on spatio-energy segmentation of regions-of-interests (ROIs) in sCT images and performing a fine material decomposition involving optimized decomposition matrix and sparsity regularization. The effectiveness of the proposed method was validated on both digital and physical data. The results showed that the proposed ROI-wise material decomposition method presents clearly higher reliability and accuracy compared to common decomposition methods based on total variation (TV) or L1-norm (lasso) regularization.

Index Terms—X-ray CT, Material decomposition, Photon-counting detector.

I. INTRODUCTION

SPECTRAL photon-counting X-ray computed tomography (sCT) is a new kind of multi-energy X-ray CT that offers new possibilities for getting insights into material components in an object, thanks to the advances in photon-counting detector (PCD) [1]–[3]. Compared to conventional energy-integrating CT or dual energy CT, sCT can count the number of photons in multiple energy bins with one single exposure, *i.e.* utilizing spectral information. This advantage enables efficient material decomposition that aims to quantitatively separate different materials (at least three) present in a pixel.

Different approaches were developed to realize material decomposition: decompose projection data acquired at different energy bins into different material sinograms, each of which corresponds to a material (*i.e.* the so-called basis material), based on their energy-dependent characteristics (*i.e.* mass attenuation coefficients), and then reconstruct individually each spatial material image containing one single material (projection-domain

approach) [4]–[6], or firstly reconstruct the spatial image from each energy bin sinogram and then decompose the reconstructed spatial images corresponding to different energy bins into spatial material images (image-domain approach) [7]–[15], or directly reconstruct spatial material images from projection data (one-step approach) [16]–[18]. The advantage of both projection-domain and one-step approaches is that they directly decompose projection data (raw data) rather than after the procedure of reconstruction. The image-domain material decomposition has its strong point that it allows us to work directly on abundant morphological features.

However, whatever the methods used, their reliability and accuracy are always impacted by unavoidable similarity between linear attenuation coefficients of basis materials [19], [20]. Such similarity makes it difficult to separate the basis materials. This difficulty is exacerbated by strong quantum noise. Quantum noise also limits the ability of PCD to recognize photons between adjacent energy bins [4], which renders the basis material separation still more difficult. To illustrate this, we give in Fig. 1 two linear attenuation coefficient (μ) curves corresponding to two mixtures with different components or different mass densities. In this case, mixtures 1 and 2 are dilutions of gadolinium (mass density 2 mg/cc) and iodine (1 mg/cc), respectively. Theoretical values are plotted for five different energy bins (*i.e.* 30-50 keV, 51-61 keV, 62-71 keV, 72-82 keV and 83-130 keV). Even in the ideal case, mixtures 1 and 2 cannot be totally distinguished due to their low mass densities. Moreover, strong quantum noise adds to uncertainty. As a result, the measured value for a fixed mixture and energy bin can vary with pixel position. This is reflected in Fig. 1 by the vertical bands representing measurement uncertainty. Where the uncertainty intervals overlap, we can no longer distinguish between the two mixtures. The performance of image-domain material decomposition also depends on image reconstruction-quality, which is impacted by photon counting statistics. To improve the reliability and accuracy of material decomposition, it is important to make full use of more information beside the aforementioned spectral information [12], [21]–[24]. A straightforward way of realizing this is to exploit morphological information embedded in the reconstructed sCT images. Image-domain material decomposition is then an approach of choice. In this paper, we investigate a novel image-domain material

B.Q. Xie, P. Niu and Y.M. Zhu are with Univ Lyon, INSA Lyon, CNRS, Inserm, CREATIS UMR 5220, U1206, F-69621, Lyon, France.

T. Su, V. Kaftandjian and P. Duvauchelle are with Univ Lyon, INSA Lyon, Laboratoire Vibrations Acoustique, F-69621 Villeurbanne, France.

L. Boussel and P. Douek are with Univ Lyon, CREATIS ; CNRS; Inserm; INSA-Lyon; Radiology Department, Hospices Civils de Lyon, Lyon, France.

F. Yang is with Beijing Jiaotong Univ, School of Computer and Information Technology, Beijing, China

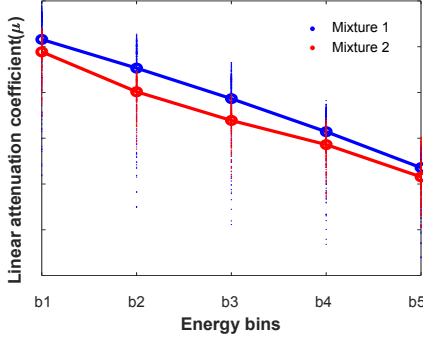


Fig. 1. Attenuation coefficient curves of two basis materials. The two curves represent dilutions of gadolinium (2 mg/cc) and iodine (1 mg/cc), respectively. The circle points represent theoretical values of the mixtures at five different energy bins. Uncertainty intervals due to measurement noise are represented by vertical tails at each point.

decomposition method by directly decreasing the impacts of similarity between basis materials with the help of multiple features extracted from the reconstructed multi-energy spatial sCT images. The idea is to exploit the abundant information and high correlations in sCT images suffering from serious reconstruction errors and artifacts. To do that, we perform basis material optimization by selecting basis materials according to their spatio-energy similarity in segmented region of interests (ROIs) of multi-energy sCT images, thus leading to so-called ROI-wise material decomposition. To our knowledge, this is the first work to improve the mathematical condition of material decomposition through optimizing basis materials by means of spatio-energy segmentation.

The rest of this paper is organized as follows. In Section II, we describe the proposed method of ROI-wise material decomposition. Section III presents experiments and results on both simulations and real data. Finally, Sections IV and V are respectively given discussion and conclusion.

II. MODEL AND METHOD

This section firstly presents typical models of image-domain material decomposition in sCT. Then, the proposed ROI-wise material decomposition method is described in detail.

A. Model of image-domain material decomposition

In the model of image-domain material decomposition, spatial images should first be reconstructed.

Spatial images are reconstructed at each separated polychromatic energy bin in sCT. The mean measured signal (number of photons penetrating materials) recorded by a PCD for the u -th ray within the i -th energy bin can be modeled by:

$$\bar{s}_i(u) = \int_{\mathbb{R}} n_0(E) d_i(E) e^{-\int_{L(u)} \mu(\vec{x}, E) dL} dE, \quad (1)$$

where $n_0(E)$ is the spectral X-ray photon fluence, $\mu(\vec{x}, E)$ the linear attenuation coefficient at position or pixel \vec{x}

for energy E , $L(u)$ the u -th ray, and $d_i(E)$ the detector response function or bin sensitivity function describing the ability of the detector to separate photons belonging to two adjacent energy bins. The measured signal is assumed to be corrupted by independent Poisson noise:

$$s_i(u) = \mathcal{P}(\lambda = \bar{s}_i(u)), \quad (2)$$

where $\mathcal{P}(\lambda)$ denotes the Poisson distribution of mean λ , and $s_i(u)$ the measured number of photons for the u -th ray in the i -th energy bin. The aim of reconstruction is to obtain the linear attenuation coefficients at each energy bin:

$$\mu(\vec{x}, i) = \underset{\mu}{\operatorname{argmin}} \mathcal{D}(\ln(\frac{s_i(u)}{\int_{\mathbb{R}} n_0(E) d_i(E) dE}), \mu(\vec{x}, i)) + \mathcal{R}, \quad (3)$$

where $\mu(\vec{x}, i)$ denotes the reconstructed linear attenuation coefficients at the i -th energy bin, which also represents the reconstructed spatial image at the i -th energy bin, \mathcal{D} the discrepancy function, and \mathcal{R} the regularization term. Reconstruction in sCT suffers from severe noise problem because of the limited photon flux, as an immediate consequence of preventing the pileup effect of PCD which describes the distortion of recorded energy spectrum by coincident pulses [25]. Therefore, we solve the above reconstruction model with simultaneous algebraic reconstruction technique and total variation (SART-TV) algorithm that is commonly used for low-dose CT image reconstruction [26], [27].

Once reconstructed, the multi-energy images are decomposed into the linear combination of mass attenuation coefficients weighted by the corresponding mass density, described by:

$$\mu(\vec{x}, i) = \sum_{\alpha=1}^M \dot{\mu}_{m\alpha}(i) \rho_{\alpha}(\vec{x}), \quad i = 1, \dots, B, \quad (4)$$

where $\dot{\mu}_{m\alpha}(i)$ designates the calculated effective mass attenuation coefficient of the α -th basis material at the i -th energy bin, M the total number of basis materials, B the total number of energy bins, and $\rho_{\alpha}(\vec{x})$ the mass density of the α -th basis material at pixel \vec{x} . For clarity, material decomposition can be formulated in matrix form as:

$$\mathbf{Y} = \mathbf{MX} + \mathbf{N}, \quad (5)$$

where \mathbf{N} denotes the noise and \mathbf{M} the decomposition matrix, each column of which represents the effective mass attenuation coefficients of one basis material for the B energies:

$$\mathbf{M} = \begin{bmatrix} \dot{\mu}_{m1}(1) & \dots & \dot{\mu}_{mM}(1) \\ \vdots & \ddots & \vdots \\ \dot{\mu}_{m1}(B) & \dots & \dot{\mu}_{mM}(B) \end{bmatrix}. \quad (6)$$

$\mathbf{Y} \in \mathcal{R}^{B \times N_p}$ and $\mathbf{X} \in \mathcal{R}^{M \times N_p}$ represent respectively the reconstructed multi-energy spatial images containing linear attenuation coefficients μ and the mass densities ρ of basis materials with N_p indicating the total number of

pixels or voxels. Theoretically, the decomposition matrix can be initialized by the effective mass attenuation coefficients, calculated using [8]:

$$\dot{\mu}_{m\alpha}(i) = \frac{\int_{E \in E_i} n_0(E) d_i(E) \mu_{m\alpha}(E) dE}{\int_{E \in E_i} n_0(E) d_i(E) dE}, \quad i = 1, \dots, B, \quad (7)$$

where $\mu_{m\alpha}(E)$ is the theoretical mass attenuation coefficient at energy E retrieved from NIST [28], and $\int_{E \in E_i} n_0(E) dE$ the total number of incident photons belonging to the i -th energy bin of width E_i . In other words, $\dot{\mu}_{m\alpha}(i)$ represents the averaged value of all the theoretical mass attenuation coefficients inside each single energy bin, which is an estimate of the true mass attenuation coefficient corresponding to that energy bin.

B. ROI-wise material decomposition

To make material decomposition more reliable and accurate, we propose to fully exploit their spatio-energy similarity in ROIs of multi-energy sCT images. To do that, we perform a basis material optimization through reducing the impact of material similarity by means of spatio-energy segmentation. The global diagram of the proposed ROI-wise material decomposition is illustrated in Fig. 2.

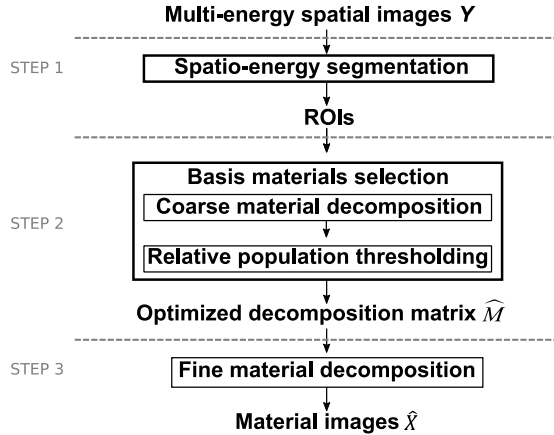


Fig. 2. Diagram of the ROI-wise material decomposition algorithm. Top to bottom: the spatio-energy segmentation based basis material optimization, and the fine material decomposition.

Basis material optimization is first processed based on the spatio-energy segmentation that separates multi-energy spatial images into different ROIs and more details will be discussed in Section II-C. Next, the basis materials in M are selected by performing a coarse material decomposition followed by a relative population thresholding and details will be presented in Section II-D. In the third step, the basis material optimization result, namely the optimized decomposition matrix \hat{M} , is finally utilized for the fine material decomposition.

C. Spatio-energy segmentation

We first obtain the ROIs of multi-energy images by spatio-energy segmentation, where each ROI represents

a homogeneous area containing similar materials. To achieve such spatio-energy segmentation, spectral and morphological features of multi-energy images are used.

Spectral features are extracted by regrouping the sCT images at all the energies as a single three-dimensional (3-D) image in which the energy is taken as a third dimension. Each pixel in the 3-D image has multi-energy values (energy-dependent μ). Pixels having different μ curves, *i.e.* different spectral features, belong to different ROIs.

Morphological features of sCT images are in fact energy-dependent because the characteristics of materials in sCT images are energy-dependent. As a result, structures imperceptible at certain energy bins may be distinguished more easily at another bin, depending on the physical characteristics of materials. Therefore, we take the morphological features from the energy bin having the most reliable segmentation by evaluating a pre-processing segmentation based on Gaussian mixture model (GMM) at each single energy bin. Then, the extracted structures are treated as the common constraint for images at all energy bins. The spectral and morphological features are then exploited jointly for segmentation based on clustering.

We use the kernel k-means method for the fusion of spectral and morphological features. The main advantage of kernel k-means is that it can make full use of kernel properties, which provides the ability to combine different features [29], [30]. More precisely, we utilize the kernel k-means method to automatically segment the pixels in the 3-D image:

$$\operatorname{argmin}_{m_k} \sum_{k=1}^K \sum_{y \in \pi_k} \|\Phi(y) - m_k\|_2^2, \quad \text{s.t. } m_k = \frac{\Phi(y)}{\|\pi_k\|_1}, \quad (8)$$

where Φ represents a non-linear transform function, π_k a partitioning of multi-energy pixel values $y \in \mathcal{R}^B$, K the total number of clusters (*i.e.* the total number of ROIs and is initialized by the total number of basis materials (M)), m_k the ideal cluster values, and $\|\cdot\|_2$ and $\|\cdot\|_1$ denote L-2 and L-1 norm, respectively. As an enhanced algorithm of normal k-means, kernel k-means can separate vectors in a high-dimensional feature space based on the non-linear transform. The non-linear transform can be calculated with a convenient kernel, and new kernel can be constructed by linearly adding two basic kernels, denoted by \mathcal{K}_1 and \mathcal{K}_2 respectively, *i.e.* $\theta_1 \mathcal{K}_1 + \theta_2 \mathcal{K}_2$, which is still a kernel. We utilize this property to combine spectral and spatial features which are separately selected by two kernels, named as $\mathcal{K}^{\text{spectrum}}$, $\mathcal{K}^{\text{space}}$. The new kernel is then obtained:

$$\mathcal{K}_{\sigma, \lambda} = (1 - \theta) \mathcal{K}_{\sigma}^{\text{spectrum}} + \theta \mathcal{K}_{\sigma}^{\text{space}}, \quad (9)$$

where $0 \leq \theta \leq 1$ and $1 - \theta$ designate the weights for spatial and spectral kernels, respectively. θ controls the relative impacts of spatial and spectral features on the final kernel.

More precisely, the kernel for spectral features is a Gaussian radial basis kernel:

$$(\mathcal{K}_\sigma^{\text{spectrum}})_{ij} = \kappa_\sigma(\mathbf{y}_i, \mathbf{y}_j) = \exp\left(-\frac{\|\mathbf{y}_i - \mathbf{y}_j\|_2^2}{2\sigma^2}\right), \quad (10)$$

where κ denotes the Gaussian radial basis kernel and σ the variance of the corresponding Gaussian distribution. The kernel for spatial features is also a Gaussian radial basis kernel:

$$(\mathcal{K}_\sigma^{\text{space}})_{ij} = \kappa_\sigma(\mathbf{y}_{si}, \mathbf{y}_{sj}) = \exp\left(-\frac{\|\mathbf{y}_{si} - \mathbf{y}_{sj}\|_2^2}{2\sigma^2}\right). \quad (11)$$

However, the pixel values $\mathbf{y}_{si} \in \mathcal{R}^B$ calculated in $\mathcal{K}_\sigma^{\text{space}}$ are those of the pre-processed 3-D image \mathbf{Y}_s containing the 'labels' of morphological information. The pre-processing aims to detect spatial features from the image having the most reliable morphological information. In practice, we firstly divide pixels in each spatial image by a classical classification method involving GMM, which is a probabilistic model that assumes that all the data points are generated from a mixture of a limited number of Gaussian distributions:

$$p_M(\mathbf{y}^b) = \sum_{i=1}^K \alpha_i p(\mathbf{y}^b | m_i^G, \sigma_i), \quad (12)$$

where $p(\mathbf{y}^b | m_i^G, \sigma_i)$ represents the i -th Gaussian distribution with means m_i^G and covariance σ_i , α_i the corresponding weights, K the total number of components, and \mathbf{y}^b the b -th bin value of pixel \mathbf{y} in the 3-D image. The models are known a priori; they are Gaussian. But, their parameters and mixture weights are unknown. The unknown parameters (mean m_i^G , covariance σ_i) in Gaussian models as well as the weights α were automatically estimated via the expectation maximization (EM) algorithm [31]. Meanwhile, the parameters are calculated independently at different bins. Then, the spatial image at each energy bin is segmented into K areas by classification. In our case the probability density function $p_M(\mathbf{y}^b)$ of GMM is used to estimate the reliability of segmentation results. The reliability of segmentation is evaluated by the optimal value of the loglikelihood cost function:

$$LL(\mathbf{Y}^b) = \ln \left(\prod_{i=1}^{N_p} p_M(\mathbf{y}_i^b) \right), \quad (13)$$

where N_p denotes the total number of pixels in the two-dimensional (2-D) image at each bin. The segmentation at certain energy bin with larger loglikelihood value is believed to be more reliable, and then the detected edges in the image at that energy bin are taken as the common edges for all the spatial images at different bins. Note that the images at different bins should be normalized (pixel value/maximum value at that bin) before comparing loglikelihood values. After that, the mean value of pixels inside each segmented area at each bin is calculated:

$$\bar{\mathbf{y}}_{\{\pi_k\}}^b = \frac{1}{N_k} \sum_{i=1}^{N_k} \mathbf{y}_i^b, \quad s.t. \quad \mathbf{y}_i^b \in \{\pi_k\}, \quad (14)$$

where $\mathbf{y}_i^b \in \{\pi_k\}$ represents the i -th pixel value in the k -th segmented area at the b -th bin, and N_k the total number of pixels inside the k -th area. Then, a new 3-D image \mathbf{Y}_s (the subscript 's' refers to 'spatio') is produced by assigning its pixel value the corresponding mean value $\bar{\mathbf{y}}_{\{\pi_k\}}^b$:

$$\mathbf{y}_{si}^b = \bar{\mathbf{y}}_{\{\pi_k\}}^b, \quad s.t. \quad \mathbf{y}_{si}^b \in \{\pi_k\}, \quad (15)$$

where $\mathbf{y}_{si}^b \in \{\pi_k\}$ represents the i -th pixel value of \mathbf{Y}_s in the k -th segmented area at the b -th bin. In other words, a 'label' is attributed to each pixel associated with its spatial features. Thus, each multi-dimensional pixel is now assigned to two different values: the original value (\mathbf{y}) containing spectral feature and the mean value (\mathbf{y}_s) containing morphological feature, of which the features can be extracted by different kernels and fused together by our final kernel given by (9).

D. Basis materials Selection

The second step of basis material optimization is to update the initialized decomposition matrix under the principle of keeping the minimum needed number of basis materials (with respect to ground-truth). To do this, we introduce a coarse material decomposition in each ROI, which exploits the sparse nature by L-1 norm (*lasso* [32]):

$$\argmin_{\mathbf{X}} \frac{1}{2} \|\mathbf{Y} - \mathbf{MX}\|_F^2 + \lambda \|\mathbf{X}\|_1, \quad (16)$$

where $\frac{1}{2} \|\mathbf{Y} - \mathbf{MX}\|_F^2$ is the data fidelity term calculated in terms of Frobenius Norm $\|\cdot\|_F$ and λ the Lagrange multiplier. The coarse material decomposition method is subject to obvious unreliability of detecting materials and poor accuracy [33]. It can nevertheless help indicate us which basis materials should be considered nonexistent, while the others are more plausible under certain selection criterion. Therefore, we propose a relative population thresholding (RPT) method to determine which basis material deserves to be selected. The RPT method is based on the population percentage defined as the ratio of the number of pixels containing a decomposed material N_x to the total number of pixels N_{ROI} in each ROI. Only the materials with percentage above a threshold will be preserved in each ROI. The physical meaning of RPT is partly involved with the aforementioned local property in both spectral and morphological domains. Local property limits the distribution of basis materials, which means that a corresponding minimum population percentage for all the ROIs should exist. In other words, the threshold of population percentage is associated with the aggregation degree of materials. After basis material optimization, we obtain a new $\hat{\mathbf{M}}$ from \mathbf{M} in the k -th ROI, as:

$$\begin{aligned} \hat{\mathbf{M}}_k &\leftarrow (\mathbf{M})_\alpha, \\ s.t. \quad \frac{N_{xak}}{N_{ROIk}} &\geq T, \end{aligned} \quad (17)$$

where $\hat{M}_k \in \mathcal{R}^{B \times M'}$ denotes the optimized decomposition matrix containing M' entries in the k -th ROI, $(M)_{\alpha}$ the α -th basis material in decomposition matrix M , \leftarrow the operator of assigning a material from M to \hat{M} , $N_{\alpha k}$ the number of pixels containing the α -th basis material in the k -th ROI, and T the population threshold.

E. Fine material decomposition

Fine material decomposition is the last step in the proposed ROI-wise material decomposition method, which involves two terms: data fidelity term based on the optimized decomposition matrix \hat{M} and sparsity regularization term.

Mathematically, we formulate the fine material decomposition as:

$$\begin{aligned} \argmin_{\mathbf{X}} \sum_k \frac{1}{2} \|\mathbf{Y}_k - \hat{\mathbf{M}}_k \mathbf{X}_k\|_F^2 + \lambda \|\mathbf{X}_k\|_1 \\ = \argmin_{\mathbf{X}} \frac{1}{2} \|\mathbf{Y} - \hat{\mathbf{M}} \mathbf{X}\|_F^2 + \lambda \|\mathbf{X}\|_1, \end{aligned} \quad (18)$$

where $\|\mathbf{Y} - \hat{\mathbf{M}} \mathbf{X}\|_F^2$ denotes the data fidelity term. The above fine material decomposition model is solved via the alternating direction method of multipliers (ADMM) iteration method [34], [35].

III. EXPERIMENTS AND RESULTS

The performance of the proposed ROI-wise material decomposition method was evaluated on both digital and physical phantom data.

A. Digital phantom data

1) *Digital phantom data generation*: The data of sCT was simulated using the software Virtual X-ray Imaging (VXI) [36] with a detector of 700 pixels (pixel size = 4×4 mm and pitch=0) by 1200 views over 360° . Projection data was obtained at tube voltage 100 kVp and current 0.1 mA without any filter. The X-ray energy bins were set as: 30-40 keV, 40-50 keV, 50-60 keV, 60-70 keV and 70-80 keV. The detector bin response function was considered ideal in the simulation. The reconstructed phantom (one slice) has 780×780 pixels and contains five basis materials: water, PMMA, gadolinium (Gd), iodine (I) and iron (Fe), as shown in Fig. 3a. The number on each disk indicates the concentration of materials (mg/cc). Note that the disk with '#' stands for mixture inserts that contain three basis materials (gadolinium, iodine and iron) with the same concentration in each column. The concentration of each material is the number on the disk in the same column. For example, the disk with '1#' denotes the mixture dilution of gadolinium, iodine and iron, each of them respectively having the concentration of 5, 5 and 15 mg/cc. The reconstructed image of the first bin is shown in Fig. 3b. The mass attenuation coefficients were retrieved from NIST [37].

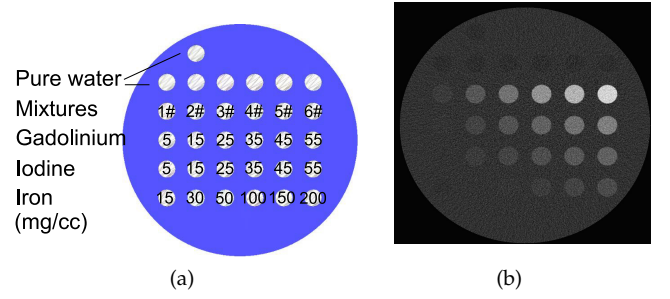


Fig. 3. (a) Digital phantom; (b) the reconstructed image of digital phantom at the first energy bin.

2) *Image quality metrics*: The performance of the proposed material decomposition was quantitatively evaluated using different metrics. The normalized Euclidean distance was utilized for the accuracy [6], considering its better evaluation for various concentrations compared to the common metric mean squared error. Given both the m -th decomposed basis materials \mathbf{x}_m and ground-truth \mathbf{x}_m^{gt} , the normalized Euclidean distance is:

$$error_m = \frac{\|\mathbf{x}_m - \mathbf{x}_m^{gt}\|_2}{\|\mathbf{x}_m^{gt}\|_2}. \quad (19)$$

The smaller the normalized Euclidean distance, the higher accuracy of the decomposition. To evaluate the reliability and sensitivity of material decomposition, we introduced two criteria: false positive (FP) and false negative (FN). The FP or FN rate is calculated as the ratio of the number of wrongly recognized pixels N_{FP} (i.e. for materials inexistent in ground-truth but in decomposition results) or unrecognized N_{FN} (i.e. for materials existent in ground-truth but not decomposed in results) to the total number of pixels in all the ROIs N_{ROIs} :

$$FP = \frac{N_{FP}}{N_{ROIs}}; \quad FN = \frac{N_{FN}}{N_{ROIs}}. \quad (20)$$

For sCT material decomposition, a smaller FP rate means smaller errors of confusing different materials, while a larger FN occurs when existing materials cannot be recognized. In other words, the smaller the FP and FN, the more reliable the material decomposition.

3) *Results*: Three other methods were compared to the proposed ROI-wise material decomposition method: (a) the common TV method as:

$$\argmin_{\mathbf{X}} \frac{1}{2} \|\mathbf{Y} - \mathbf{M} \mathbf{X}\|_F^2 + \lambda \|\mathbf{X}\|_{TV}; \quad (21)$$

(b) the Coarse method in (16) and (c) our formerly proposed DSR method [33]. Note that the 'Coarse' method is also an intermediate step of the ROI-wise method, which can be utilized to evaluate the impacts of the involved basis material optimization and energy averaging-based denoising. DSR is an image-domain material decomposition method which implicitly exploits both spectral and spatial information. For clarity, we denote the proposed ROI-wise material decomposition by 'ROI' method in the figures.

The decomposition results of five basis materials ($K=5$: water, PMMA, iron, iodine and gadolinium) are shown in Fig. 4. In contrast to the other three methods, the proposed ROI-wise material decomposition method shows better detection ability in terms of edge-preserving performance.

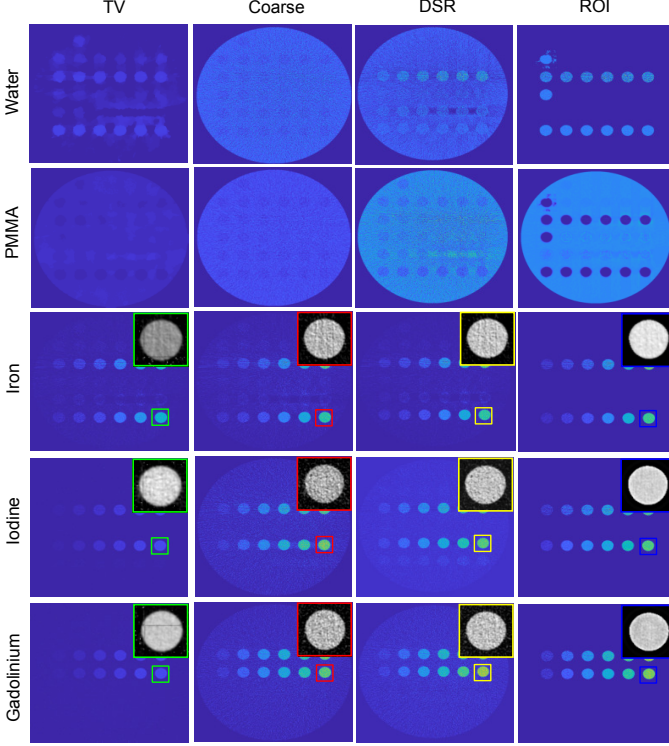


Fig. 4. The results of material decomposition using four methods on digital phantom. Left to right: TV, Coarse, DSR and ROI decompositions.

Visually, there are obviously more errors in water-like materials (e.g., water and PMMA, etc) than in iron, iodine and gadolinium for all the four methods. However, compared to TV, Coarse and DSR methods, the proposed ROI-wise method has better edge-preserving performance for water-like materials. In practical applications, the quantitative information of the last three materials are more useful, as a result of which we will pay more attention to the materials of interest: iron, iodine and gadolinium. As an illustration, we list $error_m$ of iron, iodine and gadolinium for the four methods in Table I. Clearly, ROI-wise method has the smallest $error_m$ for all the three materials compared to TV, Coarse or DSR method. We plot in Fig. 5 both measured and true concentrations of each vial along with linear regression lines for each set of points. The DSR and ROI-wise methods consistently underestimated concentrations (except at the lowest concentration of 5 mg/cc). In contrast, the other two methods underestimated the concentration for iron and gadolinium and overestimated the concentration for iodine. This consistency of both DSR and ROI-wise methods in terms of KID and HARE reliability. Table II. The proposed ROI-wise material decomposition method

TABLE I
THE NORMALIZED EUCLIDEAN DISTANCE $error_m$ OF DIFFERENT DECOMPOSITION METHODS ON DIGITAL PHANTOM.

Materials \ Methods	TV	Coarse	DSR	ROI
Iron	0.21	0.23	0.16	0.14
Iodine	0.28	0.37	0.18	0.15
Gadolinium	0.26	0.31	0.3	0.12

has a much smaller FP compared to the TV, Coarse and DSR methods. In terms of FN, the four methods exhibit similar performance for iodine and gadolinium, but the ROI-wise method leads to larger errors for iron.

TABLE II
THE FP AND FN RATES OF DIFFERENT DECOMPOSITION METHODS ON DIGITAL PHANTOM.

Materials \ Methods	TV	Coarse	DSR	ROI
Iron	13.6	18.4	0.2	0.02
Iodine	5.2	9.6	0.9	0.01
Gadolinium	4.9	6.1	7.1	0.01
Iron	9.1	14.8	16.1	19.5
Iodine	29.3	28.5	32.2	29.9
Gadolinium	28.0	26.3	26.8	30.3

B. Physical data

1) *Physical data acquisitions*: The physical phantom was acquired on a Phillips sCT prototype with a detector of 643 pixels (pixel size = 1×1 mm and pitch=0) [19], [38]. Projection data was obtained at tube voltage 120 kVp and current 220 mA with a Bowtie filter. The scan consists of 2400 projections; each projection has 643 parallel rays; each ray contains 5 energy bins: 30-50 keV, 51-61 keV, 62-71 keV, 72-82 keV and 83-130 keV. The incident photons $n_0(E)$ and detector bin response function $d_i(E)$ were provided by the manufacturer of the sCT prototype. The physical phantom is shown in Fig. 6a. The annotation is the same as for digital phantom. The substrate in the physical phantom is also PMMA. All solutions use water as the solvent. We reconstructed the spatial image at each energy bin using SART-TV, and the reconstructed image of the first bin is shown in Fig. 6b, where ring artifacts are obvious. The image reconstructed at each bin has a size of 380×380 . Due to the difficulty of separating water, PMMA, and low-concentration iron, we only consider three basis materials in physical experiments: water, gadolinium and iodine. Meanwhile, additional experiments on in vivo (rabbit) data have been performed. The reconstructed images of a rabbit 15 min after injection of gadolinium are shown in Fig. 7, in which the calcium (Ca) in the spine is identified as the basis material iodine, due to their similarity of mass attenuation coefficients in our energy thresholds setting. For the gadolinium, since the ground-truth is unknown for the in vivo data, the tubes of gadolinium dilutions nearby the rabbit are inserted as reference for

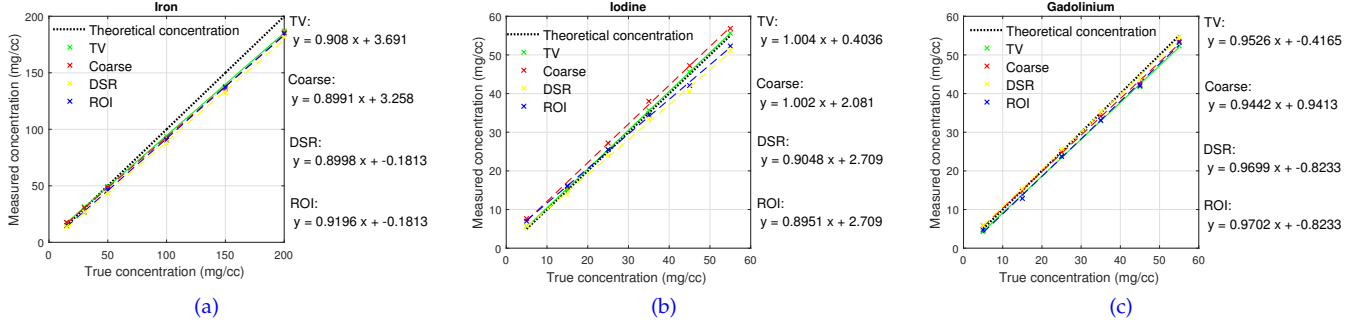


Fig. 5. The linear regressions of material decomposition using four methods for (a) iron, (b) iodine and (c) gadolinium.

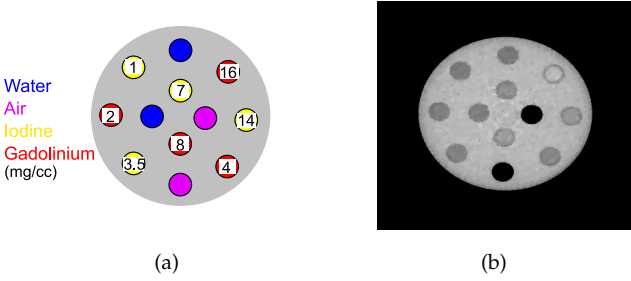


Fig. 6. (a) Physical phantom; (b) the reconstructed image of physical phantom at the first energy bin.

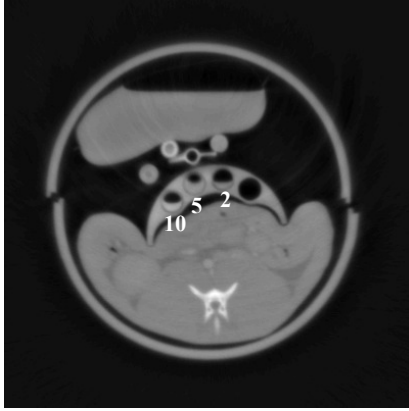


Fig. 7. The reconstructed image of an in vivo rabbit at the first energy bin 15 min after the injection of gadolinium. Labels (numbers) in ROIs refer to the nominal concentration present in each tube (mg/cc).

the evaluation. The concentrations of the gadolinium dilutions are labeled by the number.

2) *Results:* The decomposition results of three basis materials ($K=3$: water-like, iodine and gadolinium) on the physical phantom are shown in Fig. 8. The proposed ROI-wise material decomposition method has clearly better performance in detection and quantification compared to TV, Coarse and DSR methods. Firstly, our method gives much better morphological accuracy, even for water-like materials. The selected areas for water-like materials using all the four methods, shown in Fig. 9, illustrate that the edges of the selected areas are severely blurred by TV method, and mosaicked by Coarse and DSR methods. In contrast, the edges are substantially

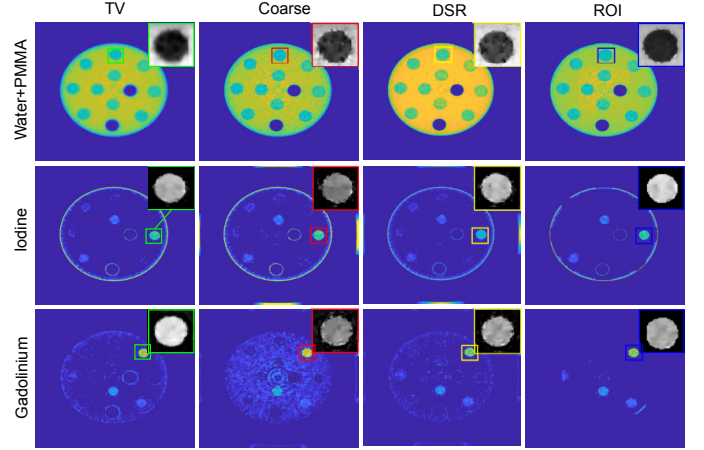


Fig. 8. The results of material decomposition on physical phantom using four methods. Left to right: TV, Coarse, DSR and ROI decompositions.

better preserved by ROI-wise method. As observed, none of the four methods has accurately decomposed iodine of concentration 1 mg/cc. For example, ROI-wise method was not able to recognize iodine inside the disk, while the other three methods were not able to separate the iodine from water or gadolinium of 2 mg/cc. Secondly, more quantitatively as shown in Table

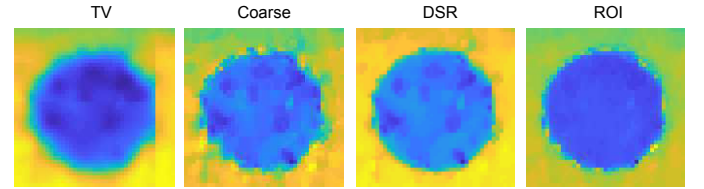


Fig. 9. The results of decomposed water-like materials in selected areas in Fig. 8.

III, ROI-wise method has clearly smaller normalized Euclidean distance for iodine and gadolinium. Specially, ROI-wise method has higher accuracy for gadolinium than for iodine. We also list the mean error in each disk (mean of each disk - ground-truth of each disk) in Table IV to further evaluate the decomposition performance of different concentrations. The mean errors of the ROI-wise method are always negative (i.e. underestimated) whatever the inserts, while the other three methods lead

to either positive (i.e. overestimated) or negative mean errors, depending on inserts.

TABLE III
THE NORMALIZED EUCLIDEAN DISTANCES $error_m$ OF DIFFERENT DECOMPOSITION METHODS ON PHYSICAL PHANTOM.

Materials	Methods	TV	Coarse	DSR	ROI
Water		0.25	0.24	0.23	0.24
Iodine		0.41	0.46	0.41	0.38
Gadolinium		0.26	0.98	0.26	0.23

TABLE IV
THE MEAN ERROR OF DECOMPOSITION IN EACH DISK ON PHYSICAL PHANTOM.

Materials	Methods	TV	Coarse	DSR	ROI
Iodine	1 (mg/cc)	0.4	0.1	0.3	-1
	3.5	-0.8	-1	-1.0	-0.6
	7	-3.1	-3.4	-3.6	-2.7
	14	-4.1	-4.0	-4.6	-4.5
Gadolinium	2	-0.7	1.7	-0.4	-0.3
	4	-0.9	1.4	-0.8	-0.3
	8	-2.2	1.3	-2.7	-2.2
	16	-2.4	0.8	-3.1	-3.0

Table V lists both FN and FP for different materials using the four methods among which the proposed ROI-wise method produces the smallest FP rate. The proposed ROI-wise method has 64% FP rate improvement for iodine and 96% for gadolinium compared to Coarse method. In terms of FN, ROI-wise method leads to obvious errors especially for iodine.

TABLE V
THE FP AND FN RATES OF DIFFERENT DECOMPOSITION METHODS ON PHYSICAL PHANTOM.

Materials	Methods	TV	Coarse	DSR	ROI
Iodine	FP (%)	16.7	21.2	9.6	7.7
Gadolinium	FP (%)	31.3	78.5	21	3.5
Iodine	FN (%)	3.9	2.8	5.3	25.7
Gadolinium	FN (%)	3.9	1.4	2.4	8.0

Concerning the influence of important kernel parameters including the weight of spatial kernel θ and the variance of the Gaussian distribution σ , the decomposition performance of iodine with different $[\theta \ \sigma^2]$ is listed in Table VI (noting that gadolinium has the same trend for $[\theta \ \sigma^2]$, no longer listed here). A too small or big θ can lead to larger normalized Euclidean distance $error_m$. An intermediate value of $\theta = 0.2$ was selected in the experiment. σ shows less impact on the decomposition, as a result of which we chose $\sigma^2 = 0.5$ according to its overall smaller $error_m$ for different θ .

Moreover, the influence of relative population threshold T is illustrated in Fig. 10. Both too small or big T leads to larger $error_m$. Nevertheless, the error varies relatively smoothly with T , especially for gadolinium. In our experiments, $T = 0.4$ corresponds to the smallest $error_m$.

TABLE VI
THE NORMALIZED EUCLIDEAN DISTANCES $error_m$ OF IODINE WITH DIFFERENT PARAMETERS ON PHYSICAL PHANTOM.

$\theta \backslash \sigma^2$	0.5	1	2	4
0	0.85	0.61	0.37	0.61
0.2	0.38	0.38	0.38	0.38
0.4	0.39	0.39	0.39	0.39
0.6	0.39	0.63	0.63	0.63
0.8	0.40	0.63	0.63	0.40
1	0.50	0.50	0.50	0.50

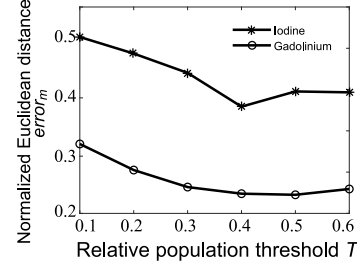


Fig. 10. Influence of relative population threshold on the normalized Euclidean distance for iodine and gadolinium.

The decomposition results of the in vivo data are shown in Fig. 11. Although all four methods can detect iodine in the spine, the ROI-wise method gives the highest accuracy considering it reduce obviously noise in the background. For gadolinium, Table VII lists the decomposed concentration of each tube. The ROI-wise method presents the highest accuracy (the smallest error) for all the measured concentrations.

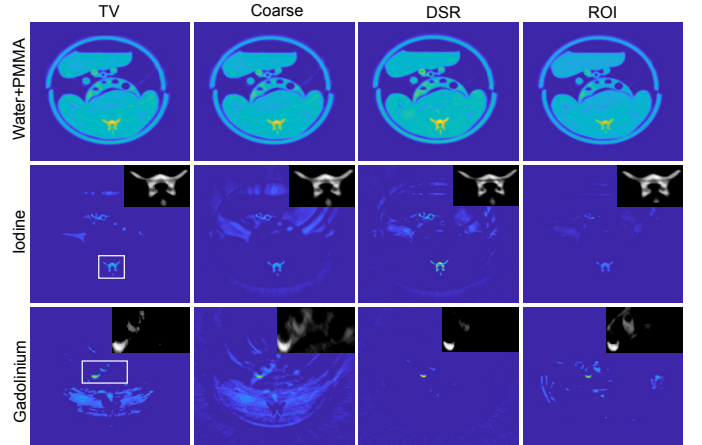


Fig. 11. The results of material decomposition based on in vivo data. Left to right: TV, Coarse, DSR and ROI decompositions.

TABLE VII
THE MEAN ERROR OF DECOMPOSED GADOLINIUM IN EACH TUBE ON IN VIVO DATA.

True Conc.	Methods	TV	Coarse	DSR	ROI
10 (mg/cc)		-3.3	-3.2	-4.4	-1.0
5		-3.8	-2.7	-4.0	-0.6
2		-2.0	-1.2	-2.0	-1

IV. DISCUSSION

We have proposed a ROI-wise material decomposition method. The proposed method enables materials to be more reliably and accurately decomposed. This is mainly due to the introduction of spatio-spectral segmentation that allows pertinent features encoded in multi-energy sCT images to be extracted for basis materials optimization.

1) *Generality of the method*: The proposed ROI-wise method is generic in the sense that it provides a framework for obtaining material decomposition from spatial reconstructed images. It provides promising perspectives for both medical applications and industrial nondestructive evaluation, considering its high accuracy and reliability. In addition to contrast agent imaging applications, our method can also be used for non-contrast imaging in medical applications. Indeed, in this work, we focused on five basis materials: water, PMMA, iron, iodine and gadolinium. Although none of the four methods can separate water and PMMA, the proposed ROI-wise method has nevertheless better edge-preserving performance for both water and PMMA, which implies that it can improve the performance of non-contrast imaging.

2) *Pre-classification*: A technique that involves pre-classification using multi-energy CT information has been shown in literature by Alessio et al [9]. However, there are two fundamental differences between that method and the proposed ROI-wise method. First, Alessio et al. use only energy information while our method exploits both spectral and morphological information. Moreover, the classification proposed by Alessio et al. demands additional measurement of mean difference in attenuation and weighted mean of linear attenuation coefficients for each class material, the precision of which varies with different concentrations. Second, four classes of materials are pre-determined in the classification of Alessio et al., whereas the classes of materials are automatically determined in our method.

The spatio-energy segmentation in our ROI-wise method can decrease the contamination between basis materials. Each ROI represents a homogeneous area containing similar materials which refer to materials having similar attenuation coefficients (related to both mass density and mass attenuation coefficients). Putting materials of similar attenuation coefficients (e.g., water and PMMA) in an ROI indeed cannot help separate them. However, removing other materials (which are different from the "similar materials") from ROI can improve the accuracy of decomposition. For example, separating gadolinium of 2 mg/cc and iodine of 1 mg/cc (the attenuation coefficients of which are shown in Fig. 1) into different ROIs can decrease the contamination between gadolinium and iodine. Meanwhile, putting similar materials in an ROI does not make it more difficult to separate similar materials, because we perform decomposition pixel by pixel. Indeed, as shown in Fig. 8, such strategy works; there is much less contamination

between iodine and gadolinium with the proposed method. More quantitatively, as observed in Table V, the proposed method has much smaller FP values than TV, Coarse and DSR methods that exhibit obvious contamination with much higher FP values, which is consistent with the results in Fig. 8.

3) *Reliability and sensitivity*: The results show that the ROI-wise material decomposition method favors more reliability while the sensitivity of detecting materials is somewhat sacrificed as the trade-off. This may partly explain why only materials with a concentration over certain limit could be accurately decomposed. This is the case for iodine with concentration 1 mg/cc compared to higher concentrations (Fig. 8). These results are consistent with the high improvement of FP (64% improvement for iodine and 96% for gadolinium compared to Coarse method on physical phantom) and worse performance of FN (Table V). Although Coarse, TV and DSR methods show better FN, they cannot separate materials of small concentrations or water-like materials. In other words, Coarse, TV or DSR method produced smaller FN at the cost of confusing water-like materials with materials of interest. For example, smaller FN of Coarse method compared to ROI method on physical phantom is ascribed to the fact that Coarse method was not able to detect gadolinium of 2 mg/cc from water or iodine of 1 mg/cc. The above detection limit of gadolinium and iodine for the proposed method is due to multiple factors, such as sCT image reconstruction quality and the performance of ROI segmentation. In clinical applications, although both FN and FP are preferred to be as small as possible, the focus depends on concrete situations. For example, false positive diagnoses of potential cancers create clinical uncertainty and often lead to multiple unnecessary biopsies or in certain cases surgical management of low-grade and low-volume disease [39]. Conversely, false negative diagnoses in terms of missing or delaying a diagnosis can also lead to real risks [40].

The trade-off between sensitivity and reliability of the proposed ROI-wise method is regularized by the threshold T in the relative population thresholding. As illustrated in Fig. 10, too small or too large T induced the increase of decomposition errors, because smaller T leads to poor reliability (but high sensitivity). In other words, when T is too small (e.g., 0.1 - 0.3), noise and reconstruction errors will have a strong impact on decomposition accuracy. On the opposite, too large T (e.g., 0.5 - 0.6) will degrade the decomposition ability for materials of small concentration. Fortunately, the results show a relatively large range (e.g., 0.3 - 0.5) for the choice of T around the optimal value, which implies that the proposed method is relatively little sensitive to T .

4) *Influence of image reconstruction quality*: Finally, it is worth noting that image reconstruction quality has dramatic influence on the performance of material decomposition in image domain. Because of excessively low dose allocated to detectors, sCT reconstruction at each

energy bin is a problem of low-dose CT reconstruction, which is also a challenging problem. A worse image reconstruction quality may deteriorate the performance of ROI segmentation, which in return may influence the accuracy of ROI-wise method, especially for images containing small structures (e.g., small blood vessels). We have chosen a common but efficient algorithm in the field of low-dose CT reconstruction (SART-TV) to reconstruct sCT images. Note that there are many other methods that jointly reconstruct sCT image to mitigate artifacts [12], [22], [41]–[44], which will in turn improve the performance of image-domain material decomposition. The results of reconstruction using SART-TV still show obvious artifacts and noise (Fig. 6b). Nevertheless, even in this situation, the materials were still correctly decomposed on both digital and physical phantoms, which demonstrates the reliability of the proposed ROI-wise method.

V. CONCLUSIONS

We have proposed a ROI-wise material decomposition method for sCT by optimizing basis materials. This is achieved through spatio-energy segmentation and exploiting both morphological and spectral information in the sCT images. The results on digital and physical phantoms showed that the ROI-wise material decomposition method presents clearly higher accuracy and reliability compared to common decomposition methods based on TV or *lasso* regularization, or another method using both spectral and spatial information (DSR). In the future work, the ability of detecting low-concentration materials will further be investigated to improve the sensitivity of the method while maintaining reliability.

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REFERENCES

- [1] C. Xu, M. Persson, H. Chen, S. Karlsson, M. Danielsson, C. Svensson, and H. Bornefalk, "Evaluation of a second-generation ultra-fast energy-resolved asic for photon-counting spectral ct," *IEEE Transactions on Nuclear Science*, vol. 60, no. 1, pp. 437–445, Feb 2013.
- [2] C. Xu, M. Danielsson, and H. Bornefalk, "Evaluation of energy loss and charge sharing in cadmium telluride detectors for photon-counting computed tomography," *IEEE Transactions on Nuclear Science*, vol. 58, no. 3, pp. 614–625, June 2011.
- [3] X. Liu, H. Bornefalk, H. Chen, M. Danielsson, S. Karlsson, M. Persson, C. Xu, and B. Huber, "A silicon-strip detector for photon-counting spectral ct: Energy resolution from 40 kev to 120 kev," *IEEE Transactions on Nuclear Science*, vol. 61, no. 3, pp. 1099–1105, June 2014.
- [4] J. P. Schlomka, E. Roessl, R. Dorscheid, S. Dill, G. Martens, T. Istel, C. Bäumer, C. Herrmann, R. Steadman, G. Zeitler, A. Livne, and R. Proksa, "Experimental feasibility of multi-energy photon-counting K-edge imaging in pre-clinical computed tomography," *Physics in Medicine and Biology*, vol. 53, no. 15, pp. 4031–4047, aug 2008.
- [5] K. C. Zimmerman and T. G. Schmidt, "Experimental comparison of empirical material decomposition methods for spectral CT," *Physics in medicine and biology*, vol. 60, no. 8, 2015.
- [6] N. Ducros, J. F. P.-J. Abascal, B. Sixou, S. Rit, and F. Peyrin, "Regularization of nonlinear decomposition of spectral x-ray projection images," *Medical Physics*, vol. 44, no. 9, pp. e174–e187, sep 2017.
- [7] K. Taguchi, M. Zhang, E. C. Frey, J. Xu, W. P. Segars, and B. M. W. Tsui, "Image-domain material decomposition using photon-counting CT," *Proceedings of SPIE*, vol. 6510, no. 1, pp. 651 008–651 008–12, 2007.
- [8] Q. Le Huy and S. Molloi, "Least squares parameter estimation methods for material decomposition with energy discriminating detectors," *Medical physics*, vol. 38, no. 1, pp. 245–55, jan 2011.
- [9] A. M. Alessio and L. R. MacDonald, "Quantitative material characterization from multi-energy photon counting CT," *Medical physics*, vol. 40, no. 3, p. 031108, mar 2013.
- [10] P. R. Mendonca, P. Lamb, and D. V. Sahani, "A flexible method for multi-material decomposition of dual-energy CT images," *IEEE Transactions on Medical Imaging*, vol. 33, no. 1, pp. 99–116, 2014.
- [11] Z. Li, S. Leng, L. Yu, Z. Yu, and C. H. McCollough, "Image-based Material Decomposition with a General Volume Constraint for Photon-Counting CT," *Proceedings of SPIE—the International Society for Optical Engineering*, vol. 9412, 2015.
- [12] Z. Li, S. Leng, L. Yu, A. Manduca, C. H. McCollough, A. Buades, B. Coll, and J. M. Morel, "An effective noise reduction method for multi-energy CT images that exploit spatio-spectral features," *Medical Physics*, vol. 44, no. 5, pp. 1610–1623, may 2017.
- [13] Z. Chen and L. Li, "Preliminary Research on Multi-Material Decomposition of Spectral CT Using Deep Learning," *The 14th International Meeting on Fully Three-Dimensional Image Reconstruction in Radiology and Nuclear Medicine*, 2017.
- [14] Y. Yao, L. Li, and Z. Chen, "Dynamic-dual-energy spectral CT for improving multi-material decomposition in image-domain," *Physics in Medicine and Biology*, vol. 64, no. 13, 2019.
- [15] Z. Chen and L. Li, "Robust multimaterial decomposition of spectral CT using convolutional neural networks," *Optical Engineering*, vol. 58, no. 01, p. 1, 2019.
- [16] R. Foygel Barber, E. Y. Sidky, T. Gilat Schmidt, and X. Pan, "An algorithm for constrained one-step inversion of spectral CT data," *Physics in Medicine and Biology*, vol. 61, no. 10, pp. 3784–3818, may 2016.
- [17] K. Mechlem, S. Ehn, T. Sellerer, E. Braig, D. Munzel, F. Pfeiffer, and P. B. Noel, "Joint statistical iterative material image reconstruction for spectral computed tomography using a semi-empirical forward model," *IEEE Transactions on Medical Imaging*, vol. 0062, no. c, pp. 1–13, 2017.
- [18] X. Hou, Y. Teng, Y. Kang, and S. Qi, "A separable quadratic surrogate total variation minimization algorithm for accelerating accurate CT reconstruction from few-views and limited-angle data," *Medical Physics*, vol. 45, no. 2, pp. 535–548, feb 2018.
- [19] S. Si-Mohamed, D. Bar-Ness, M. Sigovan, V. Tatard-Leitman, D. P. Cormode, P. C. Naha, P. Coulon, L. Rasclé, E. Roessl, M. Rokni, A. Altman, Y. Yagil, L. Bussell, and P. Douek, "Multicolour imaging with spectral photon-counting CT: a phantom study," *European Radiology Experimental*, vol. 2, no. 1, 2018.
- [20] S. Si-Mohamed, D. Bar-Ness, M. Sigovan, D. P. Cormode, P. Coulon, E. Coche, A. Vlassenbroek, G. Normand, L. Bussell, and P. Douek, "Review of an initial experience with an experimental spectral photon-counting computed tomography system," *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment*, vol. 873, pp. 27–35, nov 2017.
- [21] Y. Zhang, Y. Xi, Q. Yang, W. Cong, J. Zhou, and G. Wang, "Spectral CT Reconstruction With Image Sparsity and Spectral Mean," *IEEE Transactions on Computational Imaging*, vol. 2, no. 4, pp. 510–523, dec 2016.
- [22] Y. Zhang, X. Mou, G. Wang, and H. Yu, "Tensor-Based Dictionary Learning for Spectral CT Reconstruction," *IEEE Transactions on Medical Imaging*, vol. 36, no. 1, pp. 142–154, jan 2017.
- [23] M. J. Fadili, J.-L. Starck, J. Bobin, and Y. Moudden, "Image Decomposition and Separation Using Sparse Representations: An Overview," *Proceedings of the IEEE*, vol. 98, no. 6, pp. 983–994, jun 2010.
- [24] W. Xia, W. Wu, S. Niu, F. Liu, J. Zhou, H. Yu, G. Wang, and Y. Zhang, "Spectral CT Reconstruction ASSIST: Aided by Self-Similarity in Image-Spectral Tensors," *IEEE Transactions on Computational Imaging*, vol. 5, no. 3, pp. 420–436, sep 2019.

- [25] K. Taguchi, E. C. Frey, X. Wang, J. S. Iwanczyk, and W. C. Barber, "An analytical model of the effects of pulse pileup on the energy spectrum recorded by energy resolved photon counting x-ray detectors," *Medical Physics*, vol. 37, no. 8, pp. 3957–3969, jul 2010.
- [26] Z. Tian, X. Jia, K. Yuan, T. Pan, and S. B. Jiang, "Low-dose CT reconstruction via edge-preserving total variation regularization," *Physics in Medicine and Biology*, vol. 56, no. 18, pp. 5949–5967, sep 2011.
- [27] P. C. Hansen and J. S. Jørgensen, "AIR Tools II: algebraic iterative reconstruction methods, improved implementation," *Numerical Algorithms*, 2018.
- [28] J. H. Hubbell and S. M. Seltzer, "Tables of x-ray mass attenuation coefficients and mass energy-absorption coefficients 1 keV to 20 MeV for elements $Z = 1$ to 92 and 48 additional substances of dosimetric interest," National Inst. of Standards and Technology-PL, Gaithersburg, MD (United States), Tech. Rep., 1995.
- [29] I. S. Dhillon, Y. Guan, and B. Kulis, "Kernel k-means," in *Proceedings of the 2004 ACM SIGKDD international conference on Knowledge discovery and data mining - KDD '04*. New York, New York, USA: ACM Press, 2004, p. 551.
- [30] M. Fauvel, J. Chanussot, and J. Benediktsson, "A spatial-spectral kernel-based approach for the classification of remote-sensing images," *Pattern Recognition*, vol. 45, no. 1, pp. 381–392, jan 2012.
- [31] G. J. McLachlan and T. T. Krishnan, *The EM algorithm and extensions*. Wiley-Interscience, 2008.
- [32] R. Tibshirani, "Regression Shrinkage and Selection Via the Lasso," *Journal of the Royal Statistical Society: Series B (Methodological)*, vol. 58, no. 1, pp. 267–288, jan 1996.
- [33] B. Xie, T. Su, V. Kaftandjian, P. Niu, F. Yang, M. Robini, Y. Zhu, and P. Duvauchelle, "Material Decomposition in X-ray Spectral CT Using Multiple Constraints in Image Domain," *Journal of Nondestructive Evaluation*, vol. 38, no. 1, p. 16, mar 2019.
- [34] S. Boyd, N. Parikh, E. Chu, B. Peleato, and J. Eckstein, "Distributed Optimization and Statistical Learning via the Alternating Direction Method of Multipliers," *Foundations and Trends® in Machine Learning*, vol. 3, no. 1, pp. 1–122, 2010.
- [35] Y. Jiao, Q. Jin, X. Lu, and W. Wang, "Alternating Direction Method of Multipliers for Linear Inverse Problems," *SIAM Journal on Numerical Analysis*, vol. 54, no. 4, pp. 2114–2137, jan 2016.
- [36] P. Duvauchelle, N. Freud, V. Kaftandjian, and D. Babot, "Computer code to simulate X-ray imaging techniques," *Nuclear Instruments and Methods in Physics Research, Section B: Beam Interactions with Materials and Atoms*, vol. 170, no. 1, pp. 245–258, 2000.
- [37] E. Saloman and J. Hubbell, "X-ray attenuation coefficients (total cross sections): Comparison of the experimental data base with the recommended values of Henke and the theoretical values of Scofield for energies between 0.1–100 keV," jul 1986.
- [38] D. P. Cormode, S. Si-Mohamed, D. Bar-Ness, M. Sigovan, P. C. Naha, J. Balegamire, F. Lavenne, P. Coulon, E. Roessl, M. Bartels, M. Rokni, I. Blevins, L. Boussel, and P. Douek, "Multicolor spectral photon-counting computed tomography: in vivo dual contrast imaging with a high count rate scanner OPEN."
- [39] J. S. Quon, B. Moosavi, M. Khanna, T. A. Flood, C. S. Lim, and N. Schieda, "False positive and false negative diagnoses of prostate cancer at multi-parametric prostate MRI in active surveillance," pp. 449–463, aug 2015.
- [40] J. F. M. Verbeek and M. J. Roobol, "What is an acceptable false negative rate in the detection of prostate cancer?" *Translational Andrology and Urology*, vol. 7, no. 1, pp. 54–60, feb 2018.
- [41] H. Gao, H. Yu, S. Osher, and G. Wang, "Multi-energy CT based on a prior rank, intensity and sparsity model (PRISM)," *Inverse Problems*, vol. 27, no. 11, p. 115012, 2011.
- [42] Z. Yu, S. Leng, Z. Li, and C. H. McCollough, "Spectral prior image constrained compressed sensing (spectral PICCS) for photon-counting computed tomography," *Physics in Medicine and Biology*, vol. 61, no. 18, pp. 6707–6732, 2016.
- [43] J. Liu, H. Ding, S. Molloy, X. Zhang, and H. Gao, "TICMR: Total Image Constrained Material Reconstruction via Nonlocal Total Variation Regularization for Spectral CT," *IEEE Transactions on Medical Imaging*, vol. 35, no. 12, pp. 2578–2586, dec 2016.
- [44] S. Leng, L. Yu, J. Wang, J. G. Fletcher, C. A. Mistretta, and C. H. McCollough, "Noise reduction in spectral CT: Reducing dose and breaking the trade-off between image noise and energy bin selection," *Medical Physics*, vol. 38, no. 9, pp. 4946–4957, 2011.