New microwave power control technique by light shift detection in the DM-CPT clock

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Abstract - Atomic clocks based on coherent population trapping (CPT) and constructive polarization modulation are promising for high performance and compact devices. We propose a new method to act on the power repartition between the two laser frequency sidebands via the microwave power. We show an improvement of the frequency stability by a factor 10 after 1000 s integration time. Now we are mainly limited by the temperature variations of the polarization modulator.

Keywords: Compact atomic clock, CPT, Metrology, vapor-cell, frequency standards

1. INTRODUCTION

Compact atomic clocks are necessary where there is a need of a precise time and frequency reference for on-boards applications. They are widely used in satellites, inertial navigation, data transfer, telecommunication synchronization, etc.[?]

The Coherent Population Trapping (CPT) clock is a promising way to implement a high performance clock into on-board systems such as the next generation GNSS. Among the different vapor cell clocks, the Double Modulation CPT (DM-CPT)[?] offers an alternative compact optical technology without any microwave cavity. We modulate synchronously the phase and the polarization of a bichromatic laser beam in order to trap more atoms in the so called dark state, improving the atomic signal. Thanks to the linear architecture (reported in figure ??), the DM-CPT clock offers interesting perspectives towards a high performance integrated device, which could be miniature with a low power consumption.

In this proceeding, we present results obtained with a new method that acts on power repartition between

the two laser sidebands used for the CPT via the microwave power. We show it was one of the two main limitations of the clock frequency stability. We demonstrate a frequency stability better by a factor 10 at 1000 s integration time using this method and show the new limitation is mainly our polarization rotator.

Figure 1. DM-CPT compact atomic clock simplified scheme, laser power and frequency servo-loops not represented. DFB laser source emits at 895 nm. The microwave synthesis generates all the modulation signals and corrects the local oscillator frequency according to the detected signal after the Cs cell. The synthesis also controls the microwave power at 4.6 GHz transmitted to the EOPM that produces sidebands from the input laser beam. The EOAM modulates the laser polarization.

2. CONTEXT

Previous results has been demonstrated using the set up described in in the figure ???. It is a simple in line optical scheme. We modulate a Distributed Feed-Back laser source at 4,6 GHz around the central frequency in order to generate a bichromatic laser field where each side-band is separated by 9.2 GHz one from the other. Note this is done by feeding a 4.6 GHz to an EOPM where the microwave power influences the optical power repartition between the sidebands. The laser source is then power and frequency locked. We then modulate synchronously the phase and the polarisation of our laser beam before the clock cell, hence the name double modulation (DM). The polarisation modulation

keeps the atoms out of the extreme Zeeman sublevels and then let them participate more to the clock signal. The phase modulation allows the constructive superposition of the CPT dark states.[?] Therefore, the DM scheme allows us to have a good contrast of the atomic transition we use to lock the oscillator on.

Figure 2. Time sequence used. Frequency is modulated at $F_m = 250$ Hz at $\Delta_{F_M} = 130$ Hz depth. Pumping time and detection window are noted respectively t_p and t_w .

The polarisation and phase are modulated synchronously according to the time sequence depicted in figure ?? and said previously. At each phase/polarisation, after a time t_p where the atoms are pumped in the CPT state, we detect the signal during a time t_w .

Figure 3. Atomic signal (blue) and error signal (red), we obtained a contrast of 6.98 % for a full width at half maximum of 290 Hz represented in green.

The error signal used to lock the clock frequency is generated by square modulating the microwave frequency at a frequency of $F_m = 250$ Hz with a depth of $\Delta_{F_M} = 130$ Hz around the clock frequency. The obtained signal is also shown in the figure ??, it presents

a contrast of 6.98 % for a 290 Hz at full width at half maximum.

Figure 4. Frequency stability of the DM-CPT atomic clock. We measured a stability of 2.9×10^{-13} up to 100 s. At 1000 s integration time we are limited to 2.1×10^{-12} .

Figure 5. Contribution of the microwave power to the clock frequency stability measured at a level of 7.6×10^{-13} .

Frequency stability of 2.9×10^{-13} at 1 s up to 100 s and 2.1×10^{-12} at 1000 s has been measured with this set-up as reported in figure ?. Here, we will only describe the two main contributions since they are at least 10 time more important than the others. We also will limit to the stability at 1000 s since the frequency stability at that integration time was our main objective. First contribution we will address is the microwave power. It is responsible for the optical power repartition between the sidebands used for the CPT. Instabilities of the microwave power induced then inbalance between those sidebands and the carrier. Measuring a sensibility of clock frequency of -7.9 Hz.dB^{-1} , we can calculate the contribution reported in the figure ?? based on the

microwave power stability. We have then a contribution of 7.6×10^{-13} . Note that at this point, the microwave power is not regulated.

Figure 6. Clock frequency as a function of the temperature of the amplitude modulator modulator (in blue) and a model allowing to convert the temperature fluctuations of the EOAM in clock frequency fluctuations (in orange).

The other main contribution was the temperature of the polarization rotator (EOAM). In facts, we have observed the phenomenon depicted in the figure ???. We attribute this behaviour to the intern reflexions happening in the EOAM. The latter is composed of two triangular crystals placed in a staggered pattern a few millimeters apart. Despite our effort, we could not get rid of this effect, so we have decided to regulate its temperature on a maximum. In order to estimate its contribution to the clock frequency stability, we have modeled this behaviour with a polynomial set. Therefore we can now estimate its contribution based on a temperature measurement and we have determined a contribution of 6.9×10^{-13} .

3. MICROWAVE POWER CONTROL

This method has been inspired by a collaboration between Femto-ST, INRIM and SYRTE.^{?,?} Our idea is to use a particularity of the DM-CPT clock in order to control the laser power repartition between the sidebands via the microwave power. As shown in figure ??, the frequency shift induced by a total laser power variation depends also of the microwave. Therefore, by modulating the total optical power before the clock cell we can measure the light shift and then generate an error signal which can be used to correct the microwave power. In other words, we lock the microwave power on the atomic signal. This method differs from the one recently published[?] because it acts directly on the cause of the light shift.

Figure 7. DM-CPT clock frequency shift as function of the laser power and the microwave power.

The new time sequence used is given in figure ??. It is globally the same as the one in figure ??, except we modulate the laser power. At each of them, we modulate the modulate the phase and the polarisation as previously explained. We also apply a symmetry in the interrogation in order to get rid of a memory effect of the atoms described in.[?] The determination of the working parameters has been detailed in.[?] We chose those in a way that we can have the same error signal slope for the two laser power used with the same light induced shift. We also took care of choosing points where the stability at 1 s integration would not be too degraded.

Figure 8. Time sequence used to lock the local oscillator while correcting the microwave power. In pink and yellow the two laser powers we are working with. In grey we see the phase and polarisation modulations of the laser beam. The blue line represents the frequency modulation of the local oscillator and the green one the detection periods. Below are the summations of the signals, $S_{P\mu w}$ is the signal used for the microwave power control and S_0 is the signal used to control the oscillator.

4. RESULTS AND NEW LIMITATIONS

The frequency stability obtained with this method is reported in figure ??. We can see a slight degradation

at 1 s compared to previous results, from to 2.9×10^{-13} to 3.3×10^{-13} . Despite this, the stability at 1 s is still acceptable compared to other compact atomic clocks. At 1000 s integration time, we observe a stability of 2.1×10^{-13} , better by a factor 10 compared to the previous results, proving the efficiency of this new method.

Figure 9. Clock frequency stability. In black the free OCXO, in green the OCXO locked with the previous method and in red locked with the new one.

A more in-depth study has been conducted in order to estimate the new levels of the two previous main contributions. We first measured the microwave power stability as reported in figure ?? . We observe a stability better by a factor 8 at 1000 s. With a new sensibility of the clock frequency of -10.9 Hz.dB^{-1} we obtain a contribution of 1.6×10^{-13} .

Figure 10. Microwave power stability. We observe a gain of a factor 8 at 1000 s with the new method.

Secondly, we measured the stability of the polariza-

tion modulator (EOAM) temperature in order to calculate its contribution to the clock frequency stability. We did this by the mean of a polynomial set. This contribution is reported in the figure ?? . We measured a contribution of 1.8×10^{-13} , which is better by a factor 3. We explain this behaviour by the fact that the impact of the EOAM is now corrected by the mean of the microwave power.

Figure 11. Polarization modulator (EOAM) temperature contribution to the clock frequency stability measured at 1.8×10^{-13} after 1000 s integration time.

The polarization modulator being already temperature regulated, we will not be able to reduce its contribution as it is now. However, we plan to replace it with an other technology that should not present the same behaviour, reducing its contribution to the frequency stability.

5. CONCLUSION

We presented the results obtained with a new method controlling the optical power repartition in the sidebands generated by an EOPM by the mean of the microwave power. A stability better by a factor 10 has been measured at 1000 s integration time, proving the efficiency of this method. However, the two main limitations to the frequency stability remain the microwave power and the temperature of the polarization modulator. The first being now locked on the atomic signal, we should achieve a better microwave stability by reducing the contribution of the temperature of the polarization modulator. This can be achieved by changing the technology used to rotate the polarization. Thus, it should also allow a better clock frequency stability. However, studies about the dependencies between the two main

limitations should be conducted in order to confirm this assumption.

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