Experimental (IR, Raman) and computational analysis of a series of PtBr₂ derivatives: vibrational coupling in the coordinated ethylene and Pt-Br modes

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Abstract

Compounds $(^n\text{Bu}_4\text{P})[\text{PtBr}_3(\text{C}_2\text{H}_4)]$ (**1**), *trans*- $[\text{PtBr}_2(\text{NH}_2\text{Ph})(\text{C}_2\text{H}_4)]$ (**2**), *cis*- $[\text{PtBr}_2(\text{NH}_2\text{Ph})(\text{C}_2\text{H}_4)]$ (**3**), $(\text{PBu}_4)_2[\text{PtBr}_4]$ (**4**), $(\text{PBu}_4)[\text{PtBr}_3(\text{NH}_2\text{Ph})]$ (**5**), and *cis*- $[\text{PtBr}_2(\text{NH}_2\text{Ph})_2]$ (**6**), as well as the trichlorido analogue of **1**, $(^n\text{Bu}_4\text{P})[\text{PtBr}_3(\text{C}_2\text{H}_4)]$ (**1**^{Cl}), have been investigated experimentally by both IR and Raman spectroscopy, and theoretically by geometry optimization and normal mode analysis by the DFT approach. An analysis of the normal modes of coordinated ethylene in compounds **1**, **1**^{Cl}, **2** and **3** followed by a potential energy distribution investigation shows extensive vibrational coupling between the $\nu(\text{C}=\text{C})$ and $\delta_s(\text{CH}_2)$ A_1 modes in two bands at around $1510\text{--}1520\text{ cm}^{-1}$ and $1230\text{--}1250\text{ cm}^{-1}$, the latter one having greater $\nu(\text{C}=\text{C})$ contribution. The $\rho_w(\text{CH}_2)$ A_1 mode, the contribution of which to the above two bands is negligible, is responsible for a lower frequency band at $995\text{--}1005\text{ cm}^{-1}$. A complete vibrational analysis, backed up by the DFT calculations, has also been carried out on the Pt-Br stretching vibrations of the tetrabromido complex **4**, the tribromido complexes **1** and **5**, and the dibromido complexes **2**, **3** and **6**, and on the Pt-Cl vibrations of the analogous complex **1**^{Cl}. The study illustrates the advantages of coupling high level computations to the vibrational analyses in order to make unambiguous band assignments in IR and Raman spectroscopy.

Keywords: Platinum; Chojnacki's anion; Zeise's anion; aniline complexes; olefin complexes; vibrational coupling

Introduction

Ethylene is a D_{2h} -symmetric planar molecule with 12 vibrational modes distributed as $3A_g(R) + A_u(\text{inactive}) + 2B_{1g}(R) + B_{1u}(\text{IR}) + B_{2g}(R) + 2B_{2u}(\text{IR}) + 2B_{3u}(\text{IR})$ (R = Raman active, IR = infrared active), the reference frame being chosen with the z axis perpendicular to the molecular plane.¹ These modes can be separated into 5 stretching modes for the $\text{C}=\text{C}$ and the 4 $\text{C}-\text{H}$ bonds ($2A_g + B_{1g} + B_{2u} + B_{3u}$), 4 *in-plane* bending modes involving the HCH and CCH angles ($A_g + B_{1g} + B_{2u} + B_{3u}$), and 3 *out-of-plane* bending modes (wagging; $A_u + B_{1u} + B_{2g}$). These modes, available in many textbooks, are also given in the Supporting Information for convenience. Two Raman-active² modes at 1623 and 1342 cm^{-1} are usually attributed³⁻¹⁰ to the A_g -type $\text{C}=\text{C}$ stretching vibration ($\nu_{\text{C}=\text{C}}$) and the *in-plane* CH_2 -scissoring deformation ($\delta_s(\text{CH}_2)$), respectively. However, vibrational coupling is expected between fundamental molecular motions that are bases of the same irreducible representation and close in frequency, thus the two above-mentioned A_g -type modes (labelled q_1 and q_2) include both $\nu_{\text{C}=\text{C}}$ and *in-plane* bending motions with different contribution as highlighted in Scheme 1. In the higher frequency band, the $\text{C}=\text{C}$ bond stretches while the HCH angles widen (and the CCH angles narrow), whereas the opposite combination (HCH angle narrowing and CCH angle widening together with $\text{C}=\text{C}$ bond stretching) is observed in the lower frequency band.¹¹ On the basis of an inelastic neutron scattering analysis, the contribution of $\nu_{\text{C}=\text{C}}$ in each of these band was quantified in terms of the potential energy distribution (PED): 69 and 46 in the high and low-frequency bands, respectively.¹² A high-level normal coordinate analysis (NCA) yields a PED of 69% $\nu_{\text{C}=\text{C}}$, 14% $\delta(\text{CCH})$ and 17% $\delta(\text{HCH})$ for the band at 1609 cm^{-1} and 33% $\nu_{\text{C}=\text{C}}$, 30% $\delta(\text{CCH})$ and 37% $\delta(\text{HCH})$ for the band at 1332 cm^{-1} .¹³ Similar values have been reported by another group.¹⁴ The A_g -type $\text{C}-\text{H}$ stretching combination, which mainly contributes to a higher frequency band (3022 cm^{-1})⁹, does not significantly contribute to these two bands.

An IR -active band observed at 949 cm^{-1} is attributed to the B_{1u} -type *out-of-plane* wagging mode ($\rho_w(\text{CH}_2)$). This normal mode (labelled as q_3) is the only B_{1u} -type mode in D_{2h} symmetry and is consequently not involved in vibrational coupling. Coordination of ethylene to a metal center, however, reduces the symmetry from D_{2h} to C_{2v} at the most (if the metal complex has two perpendicular mirror planes, as for instance in Zeise's salt). Thus, all three basic molecular motions (from now on simply abbreviated as ν , δ and ρ) become of A_1 -type and may mix with each other in all the above-described normal modes, the amount of mixing correlating in principle with the local symmetry perturbation, namely with the extent of deviation from planarity. In addition to an anticipated red shift for the modes with strong ν components (q_1 and

q_2), the symmetry reduction upon coordination opens the possibility of observing these modes by infrared spectroscopy ($A_g(R) \rightarrow A_1(R, IR)$).

Scheme 1. Three vibrational frequencies and normal mode assignments for ethylene.

Two models were initially proposed for coordinated ethylene and subsequently used in normal coordinate analyses: the first one based on the widely accepted¹⁵ Chatt-Duncanson approach^{16,17} whereby the ethylene H atoms move significantly away from the C_2H_4 plane; the second one, introduced by Bokii and Kukina, based on a perfectly planar arrangement for the coordinated ethylene.¹⁸ For a variety of Pt ethylene complexes (including Zeise's salt), the band observed at $1500\text{--}1520\text{ cm}^{-1}$ was assigned to ν , leading Chatt and Duncanson to conclude about the "conservation of the double bond character" upon olefin coordination.^{16,17} This assignment has been adopted in a few subsequent reports.¹⁹⁻²¹ Conversely, the same band was attributed by other authors²² to δ on the basis of comparisons with experimental and computational studies on (deutero) cyclopropane²³ and ethylene oxide.²⁴

Later, two NCA on Zeise's salt were based on the Bokii-Kukina's model.^{25,26} In the first one, the experimentally observed band at 1526 cm^{-1} (calc. value 1528 cm^{-1}) was assigned to the A_g -type q_1 combination,²⁵ whereas in the second one the observed band at 1518 cm^{-1} (calc. value 1644 cm^{-1}) was interpreted as a pure ν vibration.²⁶ In contrast with these reports, however, another NCA using the Urey-Bradley force field assigned ν , "probably coupled with δ ", to a band observed at 1243 cm^{-1} .²⁷ Numerous other vibrational studies of coordinated ethylene (to Pt or other metals) appeared later, occasionally without mention of vibrational coupling, with the assignment of ν to either the 1500 cm^{-1} or the 1200 cm^{-1} band.²⁸

In more recent work, the notion that ν and δ are coupled in two different normal modes q_1 and q_2 has become better appreciated. On the basis of new experimental measurements on Zeise's salt, Powell *et al.* concluded that ν contributes to both experimentally observed bands at ca. 1515 cm^{-1} and ca. 1240 cm^{-1} and suggested that the lower frequency one (q_2) includes a greater proportion of ν relative to δ .²⁹ This assumption was later confirmed by the already mentioned inelastic neutron scattering study,¹² which reported PEDs for ν and δ in Zeise's salt of 47 and 61 in the 1515 cm^{-1} band and 50 and 30 in the 1240 cm^{-1} band, and also by a theoretical

study.¹³ Note that the relative contribution of ν changes on going from free ethylene (major contribution in the higher-frequency band) to coordinated ethylene (major contribution in the lower frequency band).

However, the issue was not yet settled because until then the contribution from ρ had been neglected. A new PED analysis of Zeise's salt using a larger data set yielded decompositions into 60 (δ) + 20 (ν) + 20 (ρ) for the absorption at 1515 cm^{-1} and 60 (ν) + 25 (δ) (with no significant contribution from ρ) for that at 1240 cm^{-1} .³⁰ A later publication proposed that two sets of three bands (either at 1517, 1241 and 1023 cm^{-1} , or at 1517, 1410 and 1241 cm^{-1}) can be assigned to the strongly coupled ν , δ and ρ modes.³¹ However, it was concluded that "at this stage we are not able to decide which assignment is more realistic", although the second set was shown to give better agreement for complex $\text{Pt}(\text{C}_2\text{H}_4)_3$.

In more recent years additional publications have appeared in the literature where either one or two (but not three) bands were used to describe the C=C stretching vibration. For example, in a 1995 IR study of ethylene adsorption on $\text{Pt}/\text{Al}_2\text{O}_3$, *only one* band at 1205 cm^{-1} was chosen to represent ν (vibrationally coupled with δ).³² In another study limited to the 1515 and 1241 cm^{-1} bands,³³ it was pointed out that the strength of the ethylene-metal interaction determines which motion is dominant in each normal mode and that there is a crossover in relative importance. In a SERS (surface-enhanced Raman spectroscopy) study of ethylene on Pt electrodes, the observed bands at 1495 and 1210 cm^{-1} were *confidently* contributed to vibrationally coupled ν (+ δ) and δ (+ ν).³⁴ However, the authors state that "usual assignment of the higher and lower frequency bands to ν and δ , respectively, is largely formal and is sometimes reversed".³⁵ Finally, in two very recent contributions on Zeise's salt³⁶ and its tris- C_6F_5 analogue,³⁷ the authors reported three bands as having contribution from ν , but with no mention of the ρ mode. In the first of these contributions – a variable pressure study – the weak 1515 cm^{-1} mode collapsed upon application of pressure, while a band at 1418 cm^{-1} with a negative pressure dependence was assigned to the ν (+ δ) mode and a band at 1252 cm^{-1} with a positive pressure dependence was assigned to the δ (+ ν) mode.³⁶ In the second paper, no absorption in the IR spectrum of compound $[\text{Pt}(\text{C}_6\text{F}_5)_3(\text{C}_2\text{H}_4)]^-$ could be unambiguously assigned to ν , but B3LYP/LANL2DZ calculations predicted three bands associated with a combination of ν and δ at 1307 (symmetric), 1501 (asymmetric), and 1584 (symmetric).³⁷

Therefore, in spite of several publications on vibrational studies of coordinated ethylene, there are still many controversies with the incomplete application of previously established knowledge, plus the unresolved issue of the vibrational coupling between the ν and δ modes on one side and the ρ mode on the other side. Another problem affecting the correct assignment of the vibrational modes could be related to the nature of the metal atom, with isostructural

complexes possibly showing opposite trends for the different bands as was shown for example for Zeise's salt and its Pd-analogue.³⁸

In the course of a recent mechanistic investigation of the PtBr₂-catalyzed hydroamination of ethylene and higher olefins,³⁹⁻⁴² we have been drawn to synthesize and fully characterize a number of Pt^{II}-C₂H₄ compounds with different ligand sets, notably aniline, as well as other aniline complexes that do not contain ethylene.⁴³ A thorough investigation of these complexes by IR and Raman spectroscopy, coupled to high level DFT calculations, with particular focus on the vibrational modes that involve the ethylene ligand, allow us to shine more light on the vibrational coupling problem of coordinated ethylene. An analysis of other vibrations in the above compounds, notably the Pt-halogen stretching modes, will also be presented. This includes a previously unrecognized vibrational coupling and a band reassignment for the Pt-Br stretching vibrations.

Experimental Part.

IR spectra were recorded on a Perkin-Elmer Spectrum 100 FTIR spectrometer (neat/4000-600 cm⁻¹) and on a Magna 750 Nicolet FTIR spectrometer (nujol/polyethylene: 600-50 cm⁻¹) with 2 cm⁻¹ resolution. Raman spectra (solid) were recorded on a LabRAM HR 800 (HORIBA Jobin Yvon) spectrometer. ¹³C NMR investigations were carried out on a Bruker DPX300 spectrometer operating at 75.47 MHz. All complexes used in this study have been synthesized as described elsewhere.⁴³ Quantum chemical calculation were performed with the Gaussian03 suite of programs⁴⁴ using the B3LYP functional which includes the three-parameter gradient-corrected exchange functional of Becke⁴⁵ and the correlation functional of Lee, Yang, and Parr which includes both local and non-local terms.^{46,47} The basis set chosen was the standard 6-31+G*, which includes both polarization and diffuse functions that are necessary to allow angular and radial flexibility to the highly anionic systems, for all atoms of type H, C, N, and Br. The Pt atom was described by the LANL2TZ(f) basis, which is an uncontracted version of LANL2DZ and includes an f polarization function and an ECP.⁴⁸ Frequency calculations were carried out for all stationary points in order to verify their nature as local minima and to obtain the frequency, absorption intensity and composition of all vibrational normal modes, under the gas-phase and harmonic approximations. The Gaussian03 output files were analyzed using ChemCraft software.⁴⁹ The normal coordinate analysis and PED calculations were performed for all gas-phase optimized complexes with the program DISP⁵⁰ by using equilibrium geometries and Cartesian force constants from Gaussian output files. DISP was used for transformation of the equilibrium geometries and the Cartesian force constants into the redundant internal

coordinates and force constants with minimal off diagonal terms. PED was expressed as relative contribution of each internal coordinate to the potential energy of the given normal vibration.

Results and discussion.

Compounds $(^n\text{Bu}_4\text{P})[\text{PtBr}_3(\text{C}_2\text{H}_4)]$ (**1**), *trans*- $[\text{PtBr}_2(\text{NH}_2\text{Ph})(\text{C}_2\text{H}_4)]$ (**2**), *cis*- $[\text{PtBr}_2(\text{NH}_2\text{Ph})(\text{C}_2\text{H}_4)]$ (**3**), $(\text{PBu}_4)_2[\text{PtBr}_4]$ (**4**), $(\text{PBu}_4)[\text{PtBr}_3(\text{NH}_2\text{Ph})]$ (**5**), and *cis*- $[\text{PtBr}_2(\text{NH}_2\text{Ph})_2]$ (**6**) have been investigated experimentally by both IR and Raman spectroscopy, and theoretically by geometry optimization and normal mode analysis by the DFT approach. Compound $(^n\text{Bu}_4\text{P})[\text{PtCl}_3(\text{C}_2\text{H}_4)]$ (**1**^{Cl}) has also been investigated for comparative purposes. The theoretical study was also carried out on free C_2H_4 at the same theory level for comparison. The DFT results were subsequently used to perform potential energy distribution (PED) analyses. In the next section, we shall describe the vibrational features (with particular focus on the q_1 , q_2 and q_3 vibrations) of the $\text{Pt}(\text{C}_2\text{H}_4)$ moiety, which is present in compounds **1**^{Cl}, **1**, **2** and **3**. Subsequently, we shall analyze the Pt-halogen modes and finally other relevant vibrations. A full list of all computed vibrational frequencies and IR intensities is provided as Supporting Information.

(a) $\nu_{\text{C}=\text{C}}$ vibration of coordinated ethylene: two or three bands?

Experimental IR and Raman spectra in the $1600\text{--}980\text{ cm}^{-1}$ region for compounds **1**, **2** and **3** are presented in Figure 1, Figure 2 and Figure 3, respectively. Those of compound **1**^{Cl} are available in the Supporting Information. The frequencies of the bands that can be assigned to the three A_1 -type normal modes q_1 , q_2 and q_3 are reported in Table 1. As expected, these are located at *ca.* 1550 , 1250 and 1000 cm^{-1} . The IR spectrum of compound $(\text{Bu}_4\text{N})[\text{PtBr}_3(\text{C}_2\text{H}_4)]$ has been previously reported, with a weak band at 1510 cm^{-1} in CH_2Cl_2 solution being assigned to ν .³¹ Complex **2** was also previously investigated by IR spectroscopy.⁵¹ A band observed at 1252 cm^{-1} , in perfect agreement with our data, was attributed to the ν mode without mention of vibrational coupling.

Figure 1. Raman (a) and IR (b) spectra (neat) of compound **1**.

Figure 2. Raman (a) and IR (b) spectra (neat) of compound **2**.

Figure 3. Raman (a) and IR (b) spectra (neat) of compound **3**.

Analysis of the DFT results confirms the band assignment of Table 1. The three above mentioned normal modes of compound **1** are shown in Figure 4. Corresponding views of those of the other three complexes are available in the Supporting Information (Figure S1). The calculated frequencies are greater than those observed experimentally for each compound, as is usually the case for calculations done with density functional methods. However, use of an appropriate “scaling factor” should in principle afford a good agreement with the experimental data. Indeed, the computed frequency ratios (ν_1/ν_2 and ν_1/ν_3) are essentially identical with those obtained from the experimental data, further confirming the correctness of the band assignments. In addition, the frequencies measured for q_1 and q_2 of compound **1**^{Cl} are in good agreement with those described in several previous studies of Zeise’s salt (see Introduction). This also means that the nature of the cation does not affect the Pt(C₂H₄) vibration to a large extent. On the other hand, there was not universal agreement on the location of q_3 . As stated in the Introduction, two sets of bands (at 1517, 1241 and 1023 cm⁻¹ or at 1517, 1410 and 1241 cm⁻¹) were considered as possibilities for q_1 , q_2 and q_3 in Zeise’s salt, a preference being given to the latter set.³¹ Our computational results, however, clearly identify the q_3 mode as the lower frequency band (measured at 997 cm⁻¹ for the PBu₄⁺ salt).

Table 1. Observed and calculated^a frequencies and PED analysis for the q₁, q₂ and q₃ modes of the Pt(C₂H₄) moiety in C₂H₄ and in the Pt complexes **1**^{Cl}, **1**, **2** and **3**.

Vibr.	C ₂ H ₄ ^b			1 ^{Cl}			1			2			3		
	Freq.		%	Freq.		%	Freq.		%	Freq.		%	Freq.		%
q ₁	1623 R (1701.9)	ν δ _{HCH} δ _{CCH} ρ	62 23 15 0	1516vw IR 1516w R (1557.1)	ν δ _{HCH} δ _{CCH} ρ	21 36 39 5	1510w IR 1512w R (1556.6)	ν δ _{HCH} δ _{CCH} ρ	22 35 38 4	1519w IR 1520w R (1576.2)	ν δ _{HCH} δ _{CCH} ρ	23 36 36 5	1513w IR 1515w R (1568.7)	ν δ _{HCH} δ _{CCH} ρ	26 34 36 4
q ₂	1342 R (1388.3)	ν δ _{HCH} δ _{CCH} ρ	37 39 24 0	1230m (IR) 1230m (R) (1251.3)	ν δ _{HCH} δ _{CCH} ρ	76 10 9 2	1227w IR 1226m R (1252.3)	ν δ _{HCH} δ _{CCH} ρ	75 11 10 1	1252w IR 1254m R (1283.6)	ν δ _{HCH} δ _{CCH} ρ	74 11 9 2	1236w IR 1236m R (1265.6)	ν δ _{HCH} δ _{CCH} ρ	71 13 11 3
q ₃	949 IR (974.2)	ν δ _{HCH} δ _{CCH} ρ	0 0 0 100	997 (IR) 997w (R) (1019.0)	ν δ _{HCH} δ _{CCH} ρ	0 0 14 90	995m IR 995w R (1016.7)	ν δ _{HCH} δ _{CCH} ρ	0 0 14 90	1007m IR 1004m R (1050.3)	ν δ _{HCH} δ _{CCH} ρ	0 0 14 88	1005m IR 1002m R (1037.1)	ν δ _{HCH} δ _{CCH} ρ	0 0 12 91
ν ₁ /ν ₂	1.21 R (1.23)			1.23 IR 1.23 R (1.24)			1.23 IR 1.23 R (1.24)			1.21 IR 1.21 R (1.23)			1.22 IR 1.23 R (1.24)		
ν ₁ /ν ₃	1.71 R (1.75)			1.52 IR 1.52 R (1.53)			1.52 IR 1.52 R (1.53)			1.51 IR 1.51 R (1.50)			1.51 IR 1.51 R (1.51)		

^aIR = observed in the IR spectrum; R = observed in the Raman spectrum; the value in parentheses is that computed for the DFT optimized geometry.

^bThe Raman data for free C₂H₄ are taken from ref. ¹.

Figure 4. Two orthogonal views of the calculated normal modes q_1 , q_2 and q_3 for the anion of **1**.

The results of the PED analyses are also collected in Table 1. This analysis reveals that ν and δ are heavily mixed in all cases in the higher frequency modes q_1 and q_2 , and provide a small or negligible contribution to the lower frequency q_3 . Conversely, the ρ mode mostly contributes to q_3 and only to a very minor proportion to q_1 and q_2 . Thus, even though the deviation from planarity of the ethylene molecule upon coordination allows mixing between $(\nu + \delta)$ and ρ as discussed in the Introduction, this mixing remains in practice minimal. This is partly due to the small perturbation of the higher symmetry, and partly to the large frequency difference between the ν/δ and ρ modes, although the coordination process has the effect of raising the q_3 frequency by ca. 50 cm^{-1} and lowering the q_1 and q_2 frequencies by ca. 100 cm^{-1} , such as the $(\nu_1 - \nu_3)$ and $(\nu_2 - \nu_3)$ values are significantly reduced. As found in most previous studies, the major contribution of ν is found in q_2 , whereas δ contributes more heavily to q_1 . Thus, our analysis is fully consistent with previous results,^{12,13,29} but it additionally provides a clearer picture about the location and mixing of the ρ vibration.

An analysis of the relative frequencies of q_1 and q_2 for the various compounds may provide information on the Pt-C₂H₄ bonding. As noted above, the ν mode is more heavily contributing to q_2 , for which the relative frequency trend is **2** (1252 cm^{-1}) > **3** (1236 cm^{-1}) > **1^{Cl}** (1230 cm^{-1}) > **1** (1227 cm^{-1}) for a total frequency range of 25 cm^{-1} . This trend suggests a C=C bond weakening along the same series. The mode q_1 varies to a much smaller degree between the four

compounds (10 cm^{-1} frequency range, see Table 1), and the trend is not exactly the same ($\mathbf{2} > \mathbf{1}^{\text{Cl}} > \mathbf{3} > \mathbf{1}$), although the smaller range makes this assignment less reliable. Weakening of the C=C bond may result from either an increase of $\text{Pt} \rightarrow \text{C}_2\text{H}_4$ π back-bonding, or an increase of $\text{Pt} \leftarrow \text{C}_2\text{H}_4$ σ bonding, or both. It is not possible to make conclusions about the relative effect of each of these bonding components, although both are probably contributing because there is no clear correlation with the computed Mulliken charges on the ethylene C atoms. Note that the largest frequency change for q_2 is observed on going from $\mathbf{2}$ to $\mathbf{3}$, corresponding to a change of the ligand *trans* to ethylene from aniline to Br. The frequency changes are smaller when the same ligand variation occurs in the *cis* position or when Br is changed with Cl.

The absorption frequencies ν_1 and ν_2 are found to correlate with the $^{13}\text{C}\{^1\text{H}\}$ NMR resonance of the ethylene ligand in CD_2Cl_2 solution, see Figure 5. Compound $\mathbf{3}$ could not be used for this study since it is insoluble in CH_2Cl_2 . Although the validity of correlations based on only three points may be questioned, we note that analogous correlations have been highlighted between the ν_{CO} vibration of carbonyl compounds and the carbonyl ^{13}C NMR resonance.^{52,53} For both bands, the compound showing the lower frequency (smaller force constant) is associated to the most upfield-shifted resonance. This appears consistent with a greater importance of π back-bonding effects. Indeed, greater π back-bonding (as expected for the anionic compounds where the charge density on the metal is greater) would not only weaken the C=C bond (lower frequency for q_1 and q_2) but also increase the electron density on the olefin ligand, inducing an upfield shift of the ^{13}C NMR resonance.

Figure 5. Correlation between the normal mode frequencies of Band 1 (a) and Band 2 (b) and the olefin ligand ^{13}C NMR chemical shift for olefin complexes.

In conclusion, for all practical purposes a coordinated ethylene ligand has the ν mode distributed essentially in two normal modes, usually observed around 1500 and 1200 cm^{-1} . The contribution to the mode at ca. 1000 cm^{-1} may be neglected because too small. On the other hand, attributing the ν mode to only one band is incorrect, because mixing between ν and δ is always extensive. At any rate, the major contribution of ν is found in the lower-frequency band at ca. 1200 cm^{-1} and not in the higher frequency one at ca. 1500 cm^{-1} , at least for the family of Pt^{II} complexes investigated here.

The other 9 normal modes of the C_2H_4 are modified only in a minor way upon coordination to the Pt atom, see Supporting Information (Table S1). The band calculated at 1473 cm^{-1} (observed at 1406 cm^{-1}) for compound **1** and related to that previously considered as a possibility for ρ in Zeise's salt³¹ is clearly originating from an asymmetric CH_2 scissoring mode (B_1), derived from $\delta(B_{3u})$ of free ethylene. Note that this band is of B_1 type in C_{2v} symmetry, thus it cannot vibrationally couple with the A_1 -type q_1 , q_2 and q_3 modes.

(b) $\nu(\text{Pt-C})$ vibrations.

The symmetric and antisymmetric motions of the $\text{Pt-C}_2\text{H}_4$ moiety [$\nu(\text{PtC}_2)$] were assigned on the basis of the theoretical calculations and are in close agreement with those previously reported for other $\text{Pt}(\text{C}_2\text{H}_4)$ complexes.^{1,31} In all cases the measured frequencies were higher than the calculated ones, see Table 2. Both experimental and computed values are smallest for complex **2** and highest for **1^{Cl}**. As expected, a stronger $\text{Pt-C}_2\text{H}_4$ interaction should be associated to a greater weakening of the $\text{C}=\text{C}$ bond, thus yielding higher-frequency PtC_2 modes and lower frequencies for the q_2 mode where $\nu_{\text{C}=\text{C}}$ participates the most. Indeed, the overall trend in the two PtC_2 vibrational modes is opposite to that of ν_2 (**1^{Cl}** > **1** > **3** > **2**), cf. Table 2 and Table 1. As can be seen in Table 2, these trends also perfectly correlate with the optimized Pt-C and $\text{C}=\text{C}$ distances.

Table 2. Experimental (IR/Raman) and calculated frequencies (cm^{-1}) for the $\nu(\text{PtC}_2)$ vibrations and optimized Pt-C and $\text{C}=\text{C}$ bonding parameters ($\text{\AA}$).

	1^{Cl}	1	2	3
	413w IR	403m IR	382w IR	386w IR
$\nu_s(\text{Pt-C}_2)$ (A_1)	413m R (408)	407m R (393)	380w R (371)	393m R (377)
	507w IR	498w IR	475w IR	482m IR
$\nu_{as}(\text{Pt-C}_2)$ (B_1)	512m R (502)	504m R (487)	474w R (471)	482w R (475)

Pt-C	2.131	2.145	2.164	2.159
	2.131	2.145	2.164	2.163
C=C	1.410	1.408	1.398	1.404

^aIR = observed in the IR spectrum; R = observed in the Raman spectrum; the value in parentheses is that computed for the DFT optimized geometry.

(c) Pt-X vibrations

c1. Compound 4. The D_{4h} -symmetric square-planar $[\text{PtBr}_4]^{2-}$ complex is expected to give rise to 9 vibrational modes, $A_{1g}(\text{R}) + B_{1g}(\text{R}) + A_{2u}(\text{IR}) + B_{2g}(\text{R}) + B_{2u}(\text{inactive}) + 2E_u(\text{IR})$, which can be classified as 4 stretching modes ($A_{1g}+B_{1g}+E_u$), 3 in-plane bending modes ($B_{2g}+E_u$) and 2 out-of-plane (wagging) modes ($A_{2u}+B_{2u}$).⁵⁴ The far-IR and Raman frequencies of the $[\text{PtBr}_4]^{2-}$ complex in the PBU_4^+ salt (see Figure S2) are in good agreement with those previously reported for the $^n\text{Bu}_4\text{N}^+$ and K^+ salts^{55,56} (see Table 3). However, some assignments made previously need to be reconsidered.

Table 3. Calculated and observed frequencies and assignments for $[\text{PtBr}_4]^{2-}$

Calcd freq. cm^{-1}	symmetry	Assignment, PED	Observed frequency/ cm^{-1}				
			Compd 4 (IR and R)	NBU_4^+ salt ^a (IR)	K^+ salt ^b (R)	K^+ salt ^c (IR)	K^+ salt ^d (R)
192	$1E_u$ (IR)	94% ν + 6% δ	226vs	225vs (227)		231	
176	A_{1g} (R)	N	203s		208		214, 204 ^e (208)
159	B_{1g} (R)	N	190s		194		n.l. ^f
97	A_{2u} (IR)	P	110vwsh	104ms (105)		128	
96	$2E_u$ (IR)	6% ν + 94% δ	104vw	112sh (112)		129	
93	B_{2g} (R)	Δ	108m		106		189 (194)
43	B_{2u} (n.a.) ^g	P					

^a From ref. 56: powder data for Raman, Nujol mull data (CHCl_3 solution in parentheses) for IR. ^b From ref. 56: H_2O solution. ^c From ref. 57: single-crystal IR-reflectance (transverse optical mode, 295 K). ^d From ref. 55: powder data, with solution values (D_2O) in parentheses. ^e Observed as a nonsymmetrical doublet due to a Davydov splitting. ^f Not located. ^g Not active.

A previously unrecognized vibrational coupling in the two IR-active E_u modes is revealed by the computational results, as shown in Figure 6. The higher frequency band calculated at 192 cm^{-1} ($1E_u$) consists mostly of ν but also has a 6% component of δ , whereas the much weaker band calculated at 96 cm^{-1} ($1E_u$) has greater component from the bending mode and a 6% contribution from ν . The other 5 normal modes of the $[\text{PtBr}_4]^{2-}$ complex are shown in the

Supporting Information (Figure S3) and their calculated frequencies and symmetry labels are reported in Table 3. There are 3 higher frequency modes ($192\text{--}159\text{ cm}^{-1}$), of which only one is active in the IR, followed by a large gap with the next group of 3 modes in a narrow range ($97\text{--}93\text{ cm}^{-1}$), and a final inactive mode (B_{2u} at 43 cm^{-1}).

Figure 6. Calculated E_u normal modes for complex $[\text{PtBr}_4]^{2-}$ resulting from vibrational coupling between $\nu(\text{Pt-Br})$ and $\delta(\text{Br-Pt-Br})$.

The $1E_u$ mode is the most easily identified mode because it is the most intense band in the IR spectrum at 226 cm^{-1} , and corresponds to previously assigned bands for the K^+ and $n\text{Bu}_4\text{N}^+$ salts (see Table 3).^{56,57} The two other high-frequency modes (A_{1g} and B_{1g}) are visible in the Raman spectrum at 200 and 188 cm^{-1} . These assignments appear reasonable because the calculated and observed frequency ratios are quite comparable ($1E_u/A_{1g}$: calcd. 1.09; found 1.13. $1E_u/B_{1g}$: calcd. 1.21; found 1.20). In previous Raman studies, carried out only on the K^+ salt, the A_g mode was assigned also to a band in the same region of the spectrum at 208 cm^{-1} in aqueous solution,^{55,56} or to two bands at 214 and 204 cm^{-1} due to lattice effects (Davydov splitting).⁵⁵ A band around 190 cm^{-1} was also observed in the previous studies but wrongly assigned to the B_{2g} mode in one case.

For the next three bands, the assignment is more difficult given the narrow frequency range. Only one Raman active band is expected in this region and this was reported at 106 cm^{-1} for the K^+ salt in H_2O solution.⁵⁶ Compound **4** shows this band at 108 cm^{-1} . The other two bands in this group, A_{2u} and $2E_u$, are IR active. Our measured spectrum shows a band at 104 cm^{-1} with a shoulder at ca. 110 cm^{-1} , similar to the spectrum reported for the $n\text{Bu}_4\text{N}^+$ salt.⁵⁶ According to the calculation, the $2E_u$ band is more intense than the A_{2u} band (2.44 vs. 1.37), thus the assignment of the 104 cm^{-1} band to the $2E_u$ mode appears more consistent.

c2. The trihalo anions 1^{Cl} , **1** and **5**.

In these compounds, the three Pt-X ($\text{X} = \text{Cl}$ or Br) bonds afford three stretching modes, two in-plane bending modes, and two out-of-plane bending modes, the last two being mixed with the out-of-plane $\text{Pt-C}_2\text{H}_4$ bending motions. The experimental IR and Raman spectra of these

compounds in the low-frequency region are available in the Supporting Information (Figures S4 and S5, respectively). Only the stretching modes, shown in Figure 7 for compound **1**^{Cl} as an example, will be analyzed in detail (see Table 4). The in-plane bending modes, for which only partial and tentative assignment of the spectra could be made, are shown in the Supporting Information (Figure S6) together with the experimental IR and Raman spectra (Figures S4 and S5).

In several contributions reporting a vibrational analysis of T-shaped MX₃ moiety, the three stretching modes are described as one axial (ν_{axial} , A₁) and two equatorial (a symmetric A₁, $\nu_{\text{eq,s}}$, and an antisymmetric B₁, $\nu_{\text{eq,as}}$) modes. However, like the A₁-type $\nu_{\text{C=C}}$ and $\delta_{\text{s}}(\text{CH}_2)$ modes described in part *a*, vibrational coupling is to be expected and indeed the calculations yield two A₁ normal modes with contribution from all three Pt-X bonds. The higher frequency ν_1 mode is an in-phase motion of all three Pt-X bonds, whereas in ν_3 the two mutually *trans* Pt-X bonds move in-phase with respect to each other and out of phase with respect to the axial Pt-X bond (*trans* to the olefin). The B₂ stretching mode (ν_2) is an out-of-phase stretching of the mutually *trans* Pt-X bonds, with no contribution from the axial Pt-X bond. The contribution of axial and equatorial Pt-X vibrations to each of these modes is revealed by a PED analysis, the results of which are also reported in Table 4. It is interesting to compare the Cl/Br substitution (**1**^{Cl} vs. **1**) and also the C₂H₄/PhNH₂ substitution (**1** vs. **5**). The highest frequency stretching mode is ν_1 for **1**^{Cl} and ν_2 for **1**, according to the calculations. On going from **1** to **5** (C₂H₄/PhNH₂ substitution), however, ν_1 becomes again the higher frequency mode. The ν_3 mode is the lowest frequency stretching mode for all three complexes.

Figure 7. Normal modes of vibration for the PtX₃ moiety in compound **1**^{Cl}. Compounds **1** and **5** exhibit related modes.

Table 4. Experimental (IR/Raman) and calculated (gas phase) frequencies (cm⁻¹)^a for the $\nu(\text{Pt-X})$ vibrations in compounds **1**^{Cl}, **1** and **5**.

Vibr.	1 ^{Cl}			1			5		
	Freq.	Assign.	PED	Freq.	Assign.	PED	Freq.	Assign.	PED
$\nu_1(\text{PtX})$	339vw IR ^b 338s R	ν_{ax} $\nu_{\text{eq(s)}}$	58% 41%	221m IR 221m R	ν_{ax} $\nu_{\text{eq(s)}}$	86% 10%	237vs IR 238vw R	ν_{ax} $\nu_{\text{eq(s)}}$	89% 0%

(A ₁)	(308.9)			(208)			(224)		
v ₂ (PtX) (B ₂)	333m IR 332vw R ^b (300.2)	v _{ax} v _{eq(as)}	0% 99%	241vs IR 242vw R (214)	v _{ax} v _{eq(as)}	0% 86%	224vs IR 226m R (208)	v _{ax} v _{eq(as)}	0% 95%
v ₃ (PtX) (A ₁)	314w IR 313w R (286.8)	v _{ax} v _{eq(s)}	42% 58%	200m IR 201vs R (180)	v _{ax} v _{eq(s)}	11% 89%	212w IR 210s R (178)	v _{ax} v _{eq(s)}	3% 85%

^aIR = observed in the IR spectrum; R = observed in the Raman spectrum; the value in parentheses is that computed for the DFT optimized geometry. ^bObserved as a shoulder.

It is interesting to note that the axial and equatorial Pt-X bond stretches are heavily mixed in v₁ and v₃ for compound **1^{Cl}**, much less so for compound **1**, and not at all for compound **5**. The new assignments backed up by our DFT calculations are compared with those previously published by other authors in Table 5. Only one of the previously published assignments for the [PtBr₃(C₂H₄)]⁻ complex (*n*Bu₄N⁺ salt⁵⁶) is essentially in agreement with ours, except for the neglect of a small mixing of the symmetric equatorial PtBr₂ vibration in the middle frequency mode and a similarly small mixing of the unique PtBr vibration in the low frequency mode. The other two studies^{26,58} report only two of the three possible bands, with assignments in disagreement with ours. Concerning Zeise's salt, many different assignments have been made, all of them attributing the higher frequency band to either the symmetric or the asymmetric PtCl₂ vibration. On the other hand, there is high mixing between symmetric PtCl₂ and PtCl vibrations according to our calculations, and the strongest contribution to the higher frequency band is in fact provided by the latter.

Table 5. Comparison of Pt-X stretching band assignments.^a

Compound	High frequency	Middle frequency	Low frequency	Ref.
<i>Trichloride</i>				
K[PtCl ₃ (C ₂ H ₄)]·H ₂ O	339 PtCl ₂ as	331 PtCl ₂ s	310 PtCl	25
	338 PtCl ₂ s	327 PtCl ₂ as	309 PtCl	27
	335 PtCl ₂ s	328 PtCl	306 PtCl ₂ as	56
	336 PtCl ₂ s	330 PtCl	308 PtCl ₂ as	56b
	336 PtCl ₂ s	Not located	294 PtCl ₂ as	36c
	340 PtCl ₂ s	330 PtCl ₂ as	305 PtCl	26
K[PtCl ₃ (C ₂ H ₄)]	338 PtCl ₂ s	334 PtCl	312 PtCl ₂ as	56
<i>n</i> Bu ₄ N[PtCl ₃ (C ₂ H ₄)] ^d	338 PtCl ₂ s	332 PtCl	316 PtCl ₂ as	56
<i>n</i> Bu ₄ P[PtCl ₃ (C ₂ H ₄)]	338 PtCl (58%) + PtCl ₂ s (41%)	332 PtCl ₂ as	314 PtCl ₂ s (58%) + PtCl (42%)	This work
<i>Tribromide</i>				
K[PtBr ₃ (C ₂ H ₄)]·H ₂ O	246 PtBr ₂ as	227 PtBr ₂ s	Not located	58
C ₉ H ₇ NH[PtBr ₃ (C ₂ H ₄)] ^e	Not located	218 PtBr ₂ as	198 PtBr	26
<i>n</i> Bu ₄ N[PtBr ₃ (C ₂ H ₄)]	241 PtBr ₂ as	220 PtBr	201 PtBr ₂ s	56
<i>n</i> Bu ₄ P[PtBr ₃ (C ₂ H ₄)]	241 PtBr ₂ as	221 PtBr (86%) + PtBr ₂ s (10%)	201 PtBr ₂ s (89%) + PtBr (11%)	This work

^aMeasurements carried out in Nujol mull unless otherwise stated. ^bMeasured in H₂O-HCl solution. ^cMeasurement carried out under pressure. ^dMeasured in CDCl₃ solution. ^eC₉H₇NH = quinolinium.

c3. The neutral dihalo complexes 2, 3 and 6.

The vibrational properties of compound **2** have already been reported.^{51,59} However, the study was limited to IR and only one stretching vibration was reported at 248 cm⁻¹. We find (see Table 6 and supporting figures S4 and S5) an intense band in the IR at the same frequency (also visible as a weaker band in the Raman spectrum at 249 cm⁻¹, calculated at 223 cm⁻¹), which can be assigned to the asymmetric PtBr₂ vibration on the basis of the calculations. On the other hand, two bands at 199 and 189 cm⁻¹ are predicted by the calculations to be associated with the symmetric PtBr₂ vibration, because of vibrational coupling with the aniline and ethylene ligand deformations. These are seen as two intense bands at 202 and 212 cm⁻¹ in the Raman spectrum, whereas they are weaker in the IR spectrum (201 cm⁻¹ with a lower intensity shoulder on the higher frequency side). A relatively strong band observed in the IR spectrum at 195 cm⁻¹ (not visible in the Raman spectrum) can be assigned to a C₂H₄-Pt-NH₂Ph bending mode (calculated at 185 cm⁻¹). A band at 195 cm⁻¹ (but none at 201 cm⁻¹) was also reported in ref. 59 without assignment.

Table 6. Experimental (IR/Raman) and calculated (gas phase) frequencies (cm⁻¹) for $\nu(\text{Pt-X})$ and $\delta(\text{Pt-X})$ vibrations in compounds **2**, **3** and **6**.

	2	3	6
$\nu_{\text{as}}(\text{PtBr}_2)$	248vs (IR) 249vw (R) 233 (calc)	222m (IR) ^b 221m (R) ^b 213 (calc) ^b	222s (IR) 222s (R) 225 (calc)
$\nu_{\text{s}}(\text{PtBr}_2)$	201m (IR) 202m (R) 189, 199 (calc)	233vs (IR) ^c 232vs (R) ^c 237 (calc) ^c	228vs (IR) 228m,sh (R) 233 (calc)

^aObserved as a shoulder. ^btrans to ethylene. ^ctrans to aniline.

Complex **3** features two Pt-Br stretching vibrations that are seen as the most intense bands in the low-frequency region both in the IR and in the Raman spectrum (see Table 6). According to the calculations, the lower frequency band has the major component from the Pt-Br bond *trans* to ethylene (with the longer bond of 2.482 Å), with a small in-phase component from the other bond. The higher frequency band is an out-of-phase component with the major contribution from the Pt-Br bond *trans* to aniline (with the shorter bond of 2.449 Å). The small coupling between these two bonds is probably related to the electronic asymmetry of the molecule.

The vibrational properties of compound **6** have already been reported.⁶⁰ However, like in the case of compound **2**, the study was limited to IR and only one stretching vibration was

reported at 223 cm⁻¹. Like complex **3**, we find complex **6** to feature two high intensity Pt-Br stretching vibrations both in the IR and in the Raman spectrum, at frequencies very close to those of **3** (see Table 6). However, contrary to **3**, the electronic symmetry in **6** favors a strong coupling and thus symmetric and asymmetric combinations are observed, each having equal contribution from the two Pt-Br bonds. Note that for the *trans* arrangement in **2**, the in-phase mode ($\nu_s(\text{PtBr}_2)$) is seen at lower frequency relative to the out-of-phase mode ($\nu_{as}(\text{PtBr}_2)$), whereas the reverse is true for the *cis* arrangement of **6**.

(d) $\nu(\text{Pt-N})$ vibrations

In order to complete the analysis of Pt-ligand stretching vibrations, we also examined the Pt-N stretch for all complexes containing aniline. The calculated and experimental frequencies are summarized in Table 7. For all compounds, the Pt-N stretching vibration is strongly coupled with ring vibrations and thus substantially contributes to two different vibrational modes. These are shown in the Supporting Information for the specific case of compound **5** (Figure S7). For this reason, there is no clear correlation between either of the two frequencies and the calculated Pt-N bond length (also reported in Table 7). The previously reported vibrational analysis, limited to IR, of compound **2** attributes the Pt-N stretching vibration to a band observed⁵¹ at 434 cm⁻¹ whereas the analysis of compound **6** attributed to the Pt-N stretch bands at 367 and 290 cm⁻¹.⁶⁰ For each of compounds **2**, **3** and **5**, we observe a band in the 430-450 cm⁻¹ region and also two bands around 550 cm⁻¹, one of which is probably related to the combination mode with $\nu(\text{Pt-N})$ component. For compound **6**, two bands of similar intensity are observed at ca. 450 cm⁻¹, whereas the 550 cm⁻¹ region of the spectrum is similar to that of the other compounds.

Table 7. Experimental (IR/Raman) and calculated (gas phase) frequencies (cm⁻¹) for $\nu(\text{Pt-N})$ and $\nu(\text{C-N})$ vibrations in compounds **2**, **3**, **5** and **6**.

ν/cm^{-1} _{assnmt} or $d_{\text{AA}}/\text{\AA}$	2		3		5		6	
	calc	expr	calc	expr	calc	expr	calc	expr
$\nu(\text{Pt-N})$	555 421	430m (IR) 431m (R)	557 426	432m (IR) 431m (R)	557 415	442m (IR) 440m (R)	563 ν_{as} , 561 ν_s 442 ν_{as} , 438 ν_s	455s, 443s (IR) 445m (R)
d_{PtN}	2.137		2.147		2.116		2.122	

Conclusion

Our spectroscopic analysis of a few PtBr₂ derivatives, triggered by a theoretical study of these compounds in relation to their possible involvement in ethylene hydroamination catalysis, has led us to reconsider several spectra assignment made for these and related compounds in the

literature, including some relatively recent contributions. The main contributions of the present study have been the recognition of previously unsuspected vibrational couplings (notably for PtBr_n fragments with n = 2, 3 and 4) and an analysis of its limitations (notably for the q₁, q₂ and q₃ modes of coordinated C₂H₄). It has been shown that a satisfactory interpretation of the experimental spectra can be provided on the basis of high level DFT calculations. Given the ready availability of computational tools, it is now possible to reconsider previously dubious assignments of spectral features and put vibrational analysis on a more solid basis as a characterization tool for new compounds.

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