Methods for the experimental study of ^{220}Rn homogeneity in calibration chambers

Krasimir Mitev, Philippe Cassette, D. Pressyanov, S. Georgiev, Chavdar Dutsov, Nathalie Michielsen, Benoît Sabot

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Manuscript Details

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Abstract

This work presents two experimental methods for the evaluation of ^{220}Rn homogeneity in calibration chambers. The first method is based on LSC of the ^{220}Rn decay products captured in silica aerogel. The second method is based on application of solid state nuclear track detectors facing the air of the calibration chambers. The performances of the two methods are evaluated by dedicated experiments. The repeatability of the method, estimated as relative standard deviation of the LSC measurements of ten silica aerogel samplers exposed under the same conditions is found to be 1.6%. Both methods are applied to study thoron homogeneity in a 50 L empty AlphaGuard emanation and calibration container with its fan turned on, and it was found that the ^{220}Rn distribution is homogeneous within 10%. Both methods are also applied to test the thoron homogeneity in the BACCARA chamber at IRSN during a thoron calibration exercise. The results show that, at the centre of the chamber where the inputs of the sampling systems of the instruments were put close to each other, the thoron inhomogeneity is less than 10%. However, regions of higher thoron concentrations are clearly identified near the walls and the upper part of the chamber, with ^{220}Rn concentrations being up to 60% higher compared to the concentration at the reference point. These results highlight the importance of the control and assessment of thoron homogeneity in thoron calibrations and in the cases when radon monitors are checked for thoron influence.

Keywords	Thoron (^{220}Rn), Thoron calibration, Thoron homogeneity, LSC, nuclear track detectors
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Corresponding Author	Krasimir Mitev
Corresponding Author's Institution	Sofia University "St. Kliment Ohridski"
Order of Authors	Krasimir Mitev, Philippe Cassette, Dobromir Pressyanov, Strahil Georgiev, Chavdar Dutsov, Nathalie Michielsen, Benoît Sabot
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Dear Editor,

Please find enclosed the revised version of our manuscript (Ref: ARI_2019_310) entitled: Methods for the experimental study of ^{220}Rn homogeneity in calibration chambers, by authors: K. Mitev, P. Cassette, D. Pressyanov, S. Georgiev, Ch. Dutsov, N. Michielsen, B. Sabot. This is the first revision of the manuscript.

I attach also a separate file, in which we have replied to all the comments and suggestions made by the reviewers. We would like to thank the reviewers for their comments and suggestions, which were very constructive and helped to improve the manuscript.

For all the correspondence regarding this manuscript, please do not hesitate to contact me as a corresponding author.

Sincerely yours,

Krasimir Mitev

/On behalf of all the co-authors/

Response to Comments & suggestions of Reviewer #1

A well written paper worthy of publication. I have made a few comments and suggestions for minor changes.

Thank you for this encouraging opinion.

Abstract

Repeatability of NTD ? - Our long-term experience with NTD of Kodak-Pathe LR-115/II indicates that the repeatability within a set of detectors exposed at identical conditions is within the Poisson track counting statistics.

Chapter 2. Evaluation of thoron homogeneity by LSC of silica aerogel

- The following logical rearrangement of test sessions is proposed – accepted, thank you.
- Ref to BIPM SI brochure. – corrected to 10 min, thank you.
- What do you mean? To enable LS cocktail to be fast absorbed in the aerogel for LSC measurement ? – That is exactly what was meant, Thank you.

“The silica aerogel was crushed and sieved (from ~0.5 mm to 3.0 mm) to get grains with suitable size – not too large so they can fit in the sampler without punctuating the closing filters at the same time not too small (e.g. dust) to enable LS cocktail to be fast absorbed in the aerogel for the LSC measurement after the exposure.”

- Results ? under which point of view ? which type among the three mentioned ?

Thank you for these questions, we have added some more comments in the text to elucidate how we judged the performance of the filters. In short, what we did is to compare the net LS counting rate, which is expected from the thoron activity in the sampler (available from Eq. 3) with the measured net LS counting rate. This way we found that the glass microfiber filters always give higher measured counting rates, which we interpret as inability to stop (filter) entirely the thoron progeny to penetrate in the sampler from the outside. The membrane filters (both types ADVANTEC and Millipore), however, showed good agreement of the net LS counting rates (when placed two of them on each side of the sampler), and that was our criterion to choose membrane filters.

- Usually the repeatability is given as standard deviation and not as maximum deviation. Furthermore, if you want to separate the variability components due to the thoron sampling and counting statistics, you should subtract (quadratically) the repeatability of the count per unit mass for a single sampler, that you can assume as a Poisson variable. So, I guess the real repeatability is lower than 2.3%.

Thank you for this suggestion, which we accepted and presented the repeatability as a relative standard deviation.

We agree with the referee that it is in principle possible to perform analysis of the different contributions to the overall variance. In this particular case we prefer to stay with the “overall” variance, since it is instructive what one can expect from the application of the method as a whole, including: careful preparation of the samplers, weighing, transfer of aerogel from the samplers to the cocktail, LSC measurement, etc.

Chapter 3. Evaluation of thoron homogeneity by SSNTDs

Please, add a reference. ; Please, add a reference - The corresponding references were added in the revised manuscript.

➤ What about the second one ??? - The second one is the calibration exercise at BACCARA.

Figure captions:

- Which filter type ? Millipore filters, that was added in the caption.
- What is the difference between the two vials? There is no difference between the vials (same vials, same aerogel, same LS cocktail), we simply made a snapshot of two of them.

Response to Comments & suggestions of Reviewer #2

Dear Authors,

You present results of a study of two different experimental methods to measure homogeneity of thoron concentrations in calibration chambers. I regard this as an important contribution to the quality assurance of thoron calibration and measurement. The quality of the first method that involves liquid scintillation counting seems to be high. However, I have a few comments regarding the presentation of the results, where I am not sure what is meant.

According to my understanding, the second method that makes use of solid-state nuclear track detectors seems to rely on several assumptions where I am not sure if they are justified. Please see my specific comments.

(Page numbers refer to the numbers of pages in the Word document, with the Introduction starting on page 3.)

page 3, "from thoron series": I suggest to write "decay series" or "decay chain" instead of "series" in order to clarify what is meant. - changed to thoron decay series, thank you.

page 3, Aerogel sampling in general: The aerogel acts as a sink to the thoron concentration (in contrast to the SSNTD and also in contrast to actively sampling measurement devices such as the AlphaGuard in pump mode if the outlet air goes back into the chamber). I guess that this effect is negligible. However, this should be assessed and mentioned. We think that each volume of the chamber, which is available for thoron, acts as a sink of the thoron concentration, as long as thoron decays happen everywhere. It is difficult to evaluate the degree to which a given volume in the chamber acts as a thoron sink, since this depends on the transport of the air through the volume, the air circulation in the chamber etc. However, given the high air circulation rate needed to ensure thoron homogeneity, we also think that the effect of the presence of the sampler on the thoron distribution in the chamber is negligible. In addition, the thoron activity distribution inside the sampler depends on the ratio of the thickness of the sampler to the diffusion length (Eqs 2 and 3). If $L/2L_D \ll 1$, then the thoron distribution in the sampler tends to uniform distribution. According to our experience, ratios $L/2L_D \ll 0.4$ are a good compromise, in this particular case of the samplers the ratio is $L/2L_D \approx 0.17$, which leads to a 1.4% deviation of the activity in the middle of the sampler compared to the activity on its surface. This is the maximal deviation and this result imply almost uniform thoron distribution inside the sampler, hence the sampler acts in the same way as the other air volumes in the chamber. This is also an argument for choosing thinner samplers with thickness 1 cm (see below).

page 5, "The best results were obtained": What does "best results" mean? What are the criteria to make a result a good one? And how was this tested? – we had the same question from Reviewer #1. Please see the reply to him, as well as the revised version of the manuscript in this part.

page 5, "Ultima Gold LLT LS cocktail": As far as I know there are LSC cocktails based on polar liquids ("water-based") and nonpolar liquids, where radon and thoron are retained to a higher degree in nonpolar cocktails. (For example, such nonpolar cocktails are used for measuring radon in water: The radon is transferred from the water to the cocktail because it is retained there much stronger.) What type of cocktail is the cocktail that you used? - We have tried both types of LS cocktails. Our non-polar cocktail was a home-made toluene-based (Toluene+PPO+Bis-MSB) cocktail, which is not miscible in water and which has excellent radon solubility. We have a long term experience with this cocktail for radon measurements, but in this case we observed something which looks like excitation of the cocktail (similar to chemi-luminescence) when it gets in contact with the aerogel (not shown in the manuscript). The Ultima Gold LLT cocktail, which is a commercially available polar cocktail and which contains surfactants, does not exhibit the drawback of the non-polar cocktail. It enters easily in the aerogel (the ultra sonic bath is necessary to remove the bubbles from the aerogel) and allows reliable LS counting. That is why we have chosen to use the polar cocktail UG LLT.

page 6, "The results are shown in Table 1 and indicate that the thinnest samples have highest net counting rate per unit mass.": Is this not obvious from Equations 2 and 3? The aerogel close to the surface gets the highest activity of thoron whereas the aerogel in the middle of the sampler gets the smallest activity. Therefore, thin packs of aerogel will have the highest specific activity (activity per mass). However, I think for the LSC measurement not the specific activity (or net counting rate per mass) but the total activity (or net counting rate) matters.

If we take a look at the equation which describes the thoron activity profile inside the sampler (Eq. 2), we see that $C(x)$ behaves like $\cosh(x)$, which states that the activity will be higher close to the surface and will tend to decrease in the middle of the sampler. Thus, the thinner the sample the higher the specific activity. We agree that what matters in the LSC is the total activity, thus larger volumes of aerogel would be beneficial as there is a tradeoff sensitivity vs. volume. However, we cannot make the volumes too large as we want to place the entire content of one sampler easily in one standard LS vial (which volume is 22 ml) and, for practical purposes, we do not want to fill the vial entirely with aerogel. That is why we believe that the optimal samplers are with 1 cm thickness.

Figure 3: It looks as if the aerogel gets translucent in the LSC cocktails. Is this true? I think that matters for the LSC measurement. If it was not translucent, decays of thoron and thoron progeny that stay in the aerogel, which is soaked with LSC cocktail, cannot be detected by the LSC counter. – Yes! The aerogel, when it is filled with LSC cocktail becomes translucent and the thoron progenies sucked in the aerogel are detected by the LSC counter.

Table 1: L is given in units of cm^3 . This is probably a typo and should read cm. – Yes, it was a typo and we corrected it to "L, cm". Thank you!

Table 2: I see that the uncertainty of each single measurement is larger than the deviation from the mean. I understand that the uncertainty of the mean is reduced by repeated or multiple measurements. However, in this case you compare single measurements with each other. Therefore, the uncertainty of the single measurements (e.g. 4.0% for AG 1) should be taken into account.

By the way, I would expect that measurement results with an uncertainty of that level should deviate from the mean more strongly, just because of the stochastic nature of uncertainty, even if the measurements are perfectly reproducible. The mean value of the relative uncertainties is 2.9%. I do not know if you applied a coverage factor (e.g. coverage factor $k=2$ in order to get an expanded uncertainty which covers approximately 95% of the results), but if the uncertainties are standard uncertainties, there should only be a not too big amount of the measurement

results within a range of the mean $\pm$ the uncertainty (about 68% in case the values are normally distributed). Maybe the uncertainties are too big?

Yes, we have been conservative in the estimation of the uncertainties of the net counting rate at the end of the exposure. This is due to the experimental method we use to deduce the net counting rates of the samples at the end of the exposure. The uncertainty of the net counting rate at the end of the exposure has several components which contribute to it, please see our reply to the next comment. We are confident that our experimental approach to determine the net LS counting rate at the end of the thoron exposure $n_0(t=0)$ is unbiased and robust and the estimate of the uncertainties is adequate. All the reported uncertainties in the tables are standard uncertainties, i.e. coverage factor was not applied ($k=1$).

page 7, "The net counting rate of each of the samples was followed for 60 h": Is there a reason why you performed the measurements for such a long time? How did you get the count rate at the beginning? Did you calculate it from the total counts within this measurement time or did you read the count rate in several short intervals and used only the first one?

Our procedure for these measurements is as follows: we place all the LS samples in the LS analyzer and we perform series of consecutive measurements of all the samples with fixed counting time for each measurement (5 minutes in each case). We repeat the measurements of all the samples in a row for 60h. Thus, we obtain a series of measurements, and we fit the data $\ln(\text{net counting rate})$ vs. t with a straight line. Hence, we determine the net counting rate at the beginning ($n_0(t=0)$) for each sample as well as the effective half-life of the nuclides in the sample. This procedure is necessary and advantageous for two reasons: 1.) Technically we cannot measure the net counting rate at the end of the exposure ($t=0$) for all the samples simultaneously; 2) We have to wait at least 4 hours after the end of the thoron exposure in order to attain secular equilibrium between Pb-212 and its short lived progeny. The uncertainty of the net counting rate at $t=0$ includes the uncertainty components due to counting statistics, background correction as well as the component due to the fitting.

page 8, "in this case ^{220}Rn and its progeny atoms": Is there a reason why you do not mention Tl-208? I see that it does not emit alpha radiation but Pb-212 does not either (at least not this isotope itself). Tl-208 was mentioned in the list, thank you.

page 8, "Parallel measurements of the activity concentrations of ^{222}Rn ": The amount of deposited progeny depends on the surface to volume ratio of the chamber and on the unattached fraction of the progeny. A 1 m^3 chamber has a smaller surface to volume ratio than the 200 L spherical chamber; therefore a smaller amount of progeny might be deposited on the walls. A small unattached fraction (because of aerosol particles in the air inside the chamber) reduces the deposition velocity; this also leads to a smaller amount of deposited progeny.

Although surface to volume ratio of the empty 1 m^3 cylindrical BACCARA chamber is somewhat smaller (within 50%) than that of the 200 L spherical chamber, we do not consider this to change significantly the conditions assumed (almost all of the ^{220}Rn progeny: ^{212}Pb , ^{212}Bi , ^{212}Po , ^{208}Tl is deposited on the internal surface of the chamber) for the following reasons: (1) In the experiment described radon/thoron monitors were put in the BACCARA chamber, which increased the surface at which progeny atoms may deposit and decreased the free internal volume, thus increasing the surface to volume ratio; (2) Experiments at much larger than 1 m^3 volumes (e. g. rooms, as in the cited reference of Harley et al. (2010)) demonstrate extreme disequilibrium between ^{220}Rn and its progeny in the air. The following sentence in the text: "Results of other authors (Harley et. al., 2010) also show extreme disequilibrium in air between ^{220}Rn and its progeny, even for much larger volumes (e.g. rooms) than the volumes of 200 L or 1 m^3 " was amended to make this more clear.

page 8, "They register alpha particles of energy within 1.5 – 4.0 MeV and incident angle $<55^\circ$ ": I wonder if this

assumption is justified. I know that the energy and angular range of LR-115 is limited but according to my understanding that only means that within this range LR-115 are best suited as a measurement device whereas outside this range the sensitivity is smaller but not zero. I do not have access to the references literature Pressyanov 2012 but this publication

Mheemeeed et al., Characterization of alpha-particle tracks in cellulose nitrate LR-115 detectors at various incident energies and angles. *Applied Radiation and Isotopes* 79 (2013) 48–55.

presents results that show that LR-115 is sensitive also to higher energies (in this publication presented up to 5 MeV), especially at large angles of incidence. I think the proposition that LR-115 is not sensitive to energies above 4 MeV should be justified more strongly.

The energy and angular window of the SSNTDs depends on the etching conditions and the mode of track counting. Essential for the application is the energy of alpha particles of ^{220}Rn progeny to be outside of this window, to avoid registration of alphas from the plate-out on the detector surface. This is valid as for the conditions we use, as well as for these reported by Mheemeeed et al. (2013). The corresponding paragraph in the text was revised, to make this more clear.

Highlights paper 78 Mitev et al.:

- Two new methods for checking Thoron homogeneity in chambers are suggested.
- The first method is based on trapping of Thoron progeny in aerogel and LSC.
- The second method is based on solid state nuclear track detectors.
- Both methods studied in laboratory experiments and found feasible.
- Both methods are applied in the BACCARA 2018 Thoron calibration exercise.

Methods for the experimental study of ^{220}Rn homogeneity in calibration chambers

K. Mitev^{a,*}, P. Cassette^b, D. Pressyanov^a, S. Georgiev^a, Ch. Dutsov^a, N. Michielsen^c,
B. Sabot^b

^a Sofia University “St. Kliment Ohridski”, Faculty of Physics, 1164 Sofia, Bulgaria.

^b CEA, LIST, Laboratoire National Henri Becquerel (LNE-LNHB), 91191 Gif-sur-Yvette Cedex, France.

^c Institut de Radioprotection et de Sûreté Nucléaire (IRSN), 92262 Fontenay aux Roses, France.

Abstract

This work presents two experimental methods for the evaluation of ^{220}Rn homogeneity in calibration chambers. The first method is based on LSC of the ^{220}Rn decay products captured in silica aerogel. The second method is based on application of solid state nuclear track detectors facing the air of the calibration chambers. The performances of the two methods are evaluated by dedicated experiments. The repeatability of the method, estimated as relative standard deviation of the LSC measurements of ten silica aerogel samplers exposed under the same conditions is found to be 1.6%. Both methods are applied to study thoron homogeneity in a 50 L empty AlphaGuard emanation and calibration container with its fan turned on, and it was found that the ^{220}Rn distribution is homogeneous within 10%. Both methods are also applied to test the thoron homogeneity in the BACCARA chamber at IRSN during a thoron calibration exercise.

* Corresponding author. Tel.: + 359 889 714 226.
E-mail address: kmitev@phys.uni-sofia.bg (K. Mitev).

The results show that, at the centre of the chamber where the inputs of the sampling systems of the instruments were put close to each other, the thoron inhomogeneity is less than 10%. However, regions of higher thoron concentrations are clearly identified near the walls and the upper part of the chamber, with ^{220}Rn concentrations being up to 60% higher compared to the concentration at the reference point. These results highlight the importance of the control and assessment of thoron homogeneity in thoron calibrations and in the cases when radon monitors are checked for thoron influence.

Keywords: Thoron (^{220}Rn), Thoron calibration, Thoron homogeneity, LSC, nuclear track detectors.

1. Introduction

Thoron (^{220}Rn) is an isotope of the noble gas radon with 55.8 s half-life. Studies in the last decade demonstrate that doses from thoron decay series cannot be considered negligible and there is a need for improvement of passive methods to measure thoron progeny and of the associated metrological assurance (McLaughlin, 2010; Hosoda et al., 2017). The quality assurance of thoron measuring instruments and the studies of the influence of ^{220}Rn on the radon measurement devices are areas of active research (Sabot et al., 2016; Röttger et al., 2009; Röttger et al., 2010; Röttger et al., 2014; Tokonami, 2005; He et al., 2017; Michielsen and Bondiguel, 2015). However, the short half-life of ^{220}Rn makes it difficult to ensure that it is homogeneously distributed in the chamber volume when thoron exposures are performed. Therefore, experimental methods able to probe thoron homogeneity are highly necessary.

The objective of this work is to present two newly proposed methods for evaluation of thoron homogeneity and their application in the thoron calibration exercise that has been carried out at IRSN in the framework of the MetroRADON Euramet EMPIR joint research project. The first method is based on a capture of thoron decay products in silica aerogel grains and subsequent liquid scintillation counting (LSC) of the silica aerogel. The second method is based on the measurement of the density of tracks formed by ^{220}Rn and ^{216}Po in Kodak Pathe LR-115/II solid state nuclear track detectors (SSNTDs). The two methods are applied successfully for the evaluation of ^{220}Rn homogeneity in small (50 L) calibration vessels as well as in the BACCARA chamber (1 m^3) at IRSN with seven ^{220}Rn measuring instruments inside.

2. Evaluation of thoron homogeneity by LSC of silica aerogel

This method makes use of a silica aerogel as thoron sampler and its subsequent mixing with a LS cocktail for LSC counting. The idea of the sampler is to allow thoron to enter freely from the environmental air into the cylindrical volume through the filters and to stop the thoron decay products on the filters. Thus, when ^{220}Rn decays inside the sampler, its decay products (^{216}Po , ^{212}Pb , ^{212}Bi , ^{212}Po and ^{208}Tl) will attach to the silica aerogel and their activity in the silica aerogel will be proportional to the ^{220}Rn activity that has entered in the cylinder. The latter is proportional to the ambient ^{220}Rn activity concentration in the air surrounding the sampler. The usage of silica aerogel provides a large / sufficient amount of free space in the sampler for thoron to diffuse in and at the same time allows the effective capture of its decay products.

The ^{220}Rn transport inside the sampler (assumed to have a cylindrical geometry) can be described by the diffusion equation with a term accounting for radioactive decay. As the thoron half-life is 55.8 s, for constant thoron concentration and exposures longer than 10 min, a steady state diffusion can be assumed:

$$D \frac{\partial^2 C}{\partial x^2} - \lambda C = \frac{\partial C}{\partial t}, \text{ with } \frac{\partial C}{\partial t} = 0 \quad (1)$$

where C is the ^{220}Rn activity concentration, D is the diffusion coefficient of ^{220}Rn in the material, λ is the ^{220}Rn decay constant, x is the space coordinate along the cylinder axis and t is the time variable. The solution of the above equation in plate parallel geometry along the axis of the sampler is:

$$C(x) = C_{out} \frac{\cosh\left(\frac{2x-L}{2L_D}\right)}{\cosh\left(\frac{L}{2L_D}\right)} \quad (2)$$

and the ^{220}Rn activity in the sampler (A_{Tn}) is given by

$$A_{Tn} = 2C_{out}V \frac{L_D}{L} \tanh\left(\frac{L}{2L_D}\right), \quad (3)$$

where C_{out} is the outside ^{220}Rn activity concentration (assuming that the filters are transparent to thoron), L and V are the height and the volume of the cylinder and $L_D=(D/\lambda)^{1/2}$ is the diffusion length of ^{220}Rn in the material (air or silica aerogel). Taking into account that the diffusion coefficient of ^{220}Rn in air is $D=10^{-5} \text{ m}^2/\text{s}$ (Ishimori et. al., 2013) which gives $L_D=3.0 \text{ cm}$, from Eq. 2 it follows that one can expect non-homogeneous thoron distribution inside the sampler. Thus, it is important, as shown hereafter, to choose carefully the thickness of the samplers to ensure best performance.

Specially designed thoron samplers were developed as shown in Fig 1. The sampler consists of a cylindrical body and two end caps, which serve to fix and support two air filters (bottom and top) at the ends of the cylinder (Fig. 1a). The inner and the outer diameters of the cylinder are 4.4 cm and 4.7 cm, respectively. The diameters were chosen to fit the diameter of the filters ($d=4.7 \text{ cm}$). The silica aerogel was crushed and sieved (from $\sim 0.5 \text{ mm}$ to 3.0 mm) to get grains with suitable size – small enough so they can fit in the sampler without punctuating the closing filters and, at the same time, large enough (e.g. larger than dust particles) to allow the aerogel to absorb the LS cocktail quickly in order to perform the LSC measurement as soon as possible after the exposure. The silica aerogel is placed in the cylinder (Fig. 1b) and the sampler is closed tight (Fig. 1c). All plastic parts of the thoron sampler are locally made with a 3D printer.

In order to evaluate the thoron homogeneity during the calibration exposures the thoron samplers can be positioned at any point of interest in the calibration chamber. After the end of the thoron exposure, the samplers are removed and the silica aerogel from each sampler is carefully transferred in a high performance LS glass vial. The mass of the transferred silica aerogel is determined by weighing the LS vials. The vials are then filled with 15 mL Ultima Gold LLT LS cocktail and placed for 10 min in an

296 ultrasonic bath in order to facilitate the full penetration of the scintillation cocktail in
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300 the silica aerogel and to remove air bubbles from it (Fig. 3).
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302 Different types of filters were tested - glass microfiber filter with an equivalent
303 pore size of 1.2 μm and thickness 260 μm (LLG-Labware 9045867), mixed cellulose
304 ester membrane filter with an equivalent pore size of 0.2 μm and thickness 130 μm
305 (ADVANTEC A20A047A) and mixed cellulose ester membrane filter with 0.3 μm
306 equivalent pore size and thickness 150 μm (Millipore PHWP04700). The performance
307 criterion for the choice of the filter is the agreement between the measured net LS
308 counting rate and the expected net LS counting rate from the activity, absorbed in the
309 aerogel, which is determined from Eq. 3.
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319 The results from the tests showed that net LS counting rate obtained with the glass
320 microfiber filters is always higher than the net LS counting rate expected from the
321 absorbed activity (Eq. 3). Therefore, we conclude that the used glass microfiber filters
322 do not completely stop the thoron progeny from penetrating in the samplers. The best
323 agreement between the measured and the expected net LS counting rate was obtained
324 with two membrane filters placed on each entry of the sampler. Both types of filters
325 ADVANTEC A20A047A and Millipore PHWP04700 show excellent performance.
326 Scanning electron microscope images of the used membrane filters are shown in Fig.
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337 2.

338 In order to choose the optimal thickness, three different samplers were produced
339 with $L=1.0$ cm, 1.5 cm and 2.0 cm. Six samplers (two of each thickness) were exposed
340 to thoron in a 50 L AlphaGuard calibration container with an AlphaGuard PQ2000 PRO
341 (Rn/Tn) reference monitor placed inside. The ^{220}Rn activity concentration during the
342 exposure was $C_{out}=612(61)$ kBq/m^3 and the exposure duration was 68.2 h. The samplers
343 were placed at the bottom of the vessel close to each other with a special focus on ^{220}Rn
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freely reaching each sampler and passing through the filters. After the exposure the silica aerogel from each sampler was transferred in a high performance glass vial and the vials were measured on a RackBeta 1219 LS counter (Wallac, Finland). The results are shown in Table 1 and indicate that the thinnest samples have the highest net counting rate per unit mass. Considering the results in Table 1 and noting that the net LS counting rate is due to several thoron progenies (^{212}Pb , ^{212}Bi , ^{212}Po and ^{208}Tl) it can be concluded that samplers with thickness $L=1$ cm which contain around 5 g of silica aerogel provide sufficiently good sensitivity of the method to be applied for the evaluation of ^{220}Rn homogeneity.

The most important characteristic of the proposed method from the point of view of practical applications is its repeatability. It is studied in this work with the experimental set-up shown schematically in Fig. 4 and on the photo in Fig. 5. The set-up consists of a powerful fan (gas-flow $2.5 \text{ m}^3/\text{min}$) mounted to a tube. Inside the tube there are 10 thoron samplers divided in two groups (AG1 to AG5 mounted closer to the fan and AG6 to AG10 just behind them, see Fig. 4a and Fig 5a). The samplers are positioned with their filter-sides parallel to the air-flow (see Fig. 4b and Fig 5b). The entire system is placed in a 50 L calibration container (Fig. 6) with one thoron sampler placed at the exit of the tube perpendicular to the gas-flow (AG11) and another sampler placed outside beside the tube (AG12), see Fig. 4a and Fig. 6. The thoron inlet is positioned right in front of the fan and the container is closed hermetically. The gas-flow of the fan ($150 \text{ m}^3/\text{h}$) is chosen large enough to guarantee that all the samplers in the tube (AG1-AG10) are exposed to the same thoron concentration and the volume refresh rate (50 times per min) is sufficient to assume that AG11 and AG12 are also exposed to the same concentration. The duration of the ^{220}Rn exposure was 66 h and the thoron activity concentration in the container was $474(47) \text{ kBq/m}^3$. After the exposure, the silica

aerogel from the samplers was transferred and measured on the RackBeta 1219 LS counter as described above. The net counting rate of each sample was followed for 60 h and the net counting rates at the moment of the end of exposure are evaluated. The results, presented in Table 2, show that the relative standard deviation of the net counting rates per unit mass of the samplers in the tube (AG1-AG10) is 1.6% and their variations are fully within the estimated uncertainties. Moreover, the sampler AG11, which is in front of the tube and is perpendicular to the air stream also agrees well with the mean value within its estimated uncertainty. The same is also true for the sampler AG12, which is behind the tube (see Fig. 4a and Fig. 6). These results suggest a repeatability of the method of the order of 1.6%, which is an excellent repeatability for the evaluation of thoron homogeneity in ^{220}Rn calibrations.

3. Evaluation of thoron homogeneity by SSNTDs

The other approach we investigated to evaluate the homogeneity of ^{220}Rn in chambers is based on the use of bare SSNTDs, placed at different points inside the chamber. These detectors register alpha particles that reach the detector surface with energy and incident angle within certain registration window specific for each type of SSNTDs. Normally, the air contains a mixture of isotopes (in this case ^{220}Rn and its progeny atoms ^{216}Po , ^{212}Pb , $^{212}\text{Bi}+^{212}\text{Po}/^{208}\text{Tl}$, ^{212}Po is always in equilibrium with ^{212}Bi). However, in an exposure chamber volume a substantial part of the progeny atoms is deposited on the internal walls (George et al., 1983). The deposited fraction is higher when a fan creates air turbulence inside the chamber (Cheng et al., 1990) (the described experiments were made in this mode). Parallel measurements (Pressyanov, 2002) of the activity concentrations of ^{222}Rn and its progeny in 200 L spherical volume with a fan operating inside showed that the deposited fractions are 94.0%, 99.7% and 99.9% for

^{218}Po (half-life 3.05 min), ^{214}Pb (half-life 26.8 min) and ^{214}Bi (half-life 19.9 min), respectively. When ^{220}Rn + progeny is created in the chamber, due to the longer half-life of ^{212}Pb (half-life 10.64 h) and ^{212}Bi (half-life 60.55 min) one can expect that practically all of the ^{212}Pb and $^{212}\text{Bi}+^{212}\text{Po}$ atoms are deposited on the walls and their air fraction is negligible. Results of other authors (Harley et. al., 2010) also show extreme disequilibrium in air between ^{220}Rn and its decay products ^{212}Pb and ^{212}Bi , even for much larger volumes (e.g. rooms) than the volumes of 200 L or 1 m³. Therefore, within the present approach, we assume that the isotopes in the air are ^{220}Rn and ^{216}Po and that ^{216}Po , due to its short half-life of 0.15 s, it has of the same volume distribution as ^{220}Rn .

In our experiments SSNTDs of Kodak-Pathe LR-115 type II were used. They register alpha particles within an energy and angular window of registration that depends on the etching conditions and the mode of counting. These conditions in our case were etching with 10% NaOH at 60 °C for 100 min, washing with water for 30 min and washing for 2 min in still 50% ethanol and visual counting by microscope - only of tracks that created holes through the 12 µm sensitive layer of this type of detectors were counted. As described elsewhere (Pressyanov, 2012), under these conditions the detectors register alpha particles of energy within 1.5 – 4.0 MeV and incident angle <55° to the normal. The air volumes from which the alpha particles of different isotopes can be detected are schematically shown in Fig.7. What is essential for this application is that the alpha particles from the ^{220}Rn progeny atoms deposited on the detector surface cannot be detected, as their energy is well above the upper energy threshold of the window for registration. From Fig.7 it follows that, to avoid contribution from the atoms deposited on the chamber internal surface, the SSNTD face should look to air being at a distance of at least 8 cm from any surface – i.e. outside the

range of ^{212}Po alphas. Under these conditions the tracks will be due only to the alpha particles from the sources in the air (^{220}Rn and ^{216}Po) and the registered tracks will originate from the isotopes that are in a small volume within a distance of 2.4 – 5.0 cm from the detector surface (“detection volumes”, see Fig. 7).

Respecting the above rule, Kodak-Pathe LR-115/II detectors were used in two experiments. The first was in 50 L cylindrical exposure chamber in which 33 pieces of detectors were placed. The map of the “detection volumes” is shown in Fig. 8. The exposure was made at the laboratory facility, described elsewhere (Pressyanov et. al., 2017). The exposure was made at an average ^{220}Rn concentration of 800(50) kBq m⁻³ for 220 min. During exposure the ^{220}Rn concentrations were followed by a reference instrument AlphaGUARD PQ2000 PRO (Rn/Tn). The analysis of the results showed that the distribution of ^{220}Rn in the air is homogeneous within 10%. The variations between the results can be explained by the detector reading uncertainty and sub-volumes of significantly higher or lower concentrations cannot be identified. Similar homogeneity (i.e. within 10%) of the empty 50 L AlphaGuard calibration volume with its own fan turned on is obtained by the LS counting of the silica aerogel (not shown).

4. Application of the methods during the thoron calibration exercise, performed at BACCARA chamber at IRSN

A thoron calibration exercise was carried out in May 2018 in the framework of the MetroRADON Euramet project using the IRSN reference radon chamber called BACCARA. The BACCARA chamber is a 1 m³ stainless steel chamber with a thoron reference instrument attached to it (Sabot et. al., 2016), which is used to create reference radon and thoron atmospheres for calibration of radon and thoron measuring devices. During the thoron calibration exercise, four AlphaGuards and three RAD7 instruments

were placed in the chamber and their thoron measurement performance was checked against the reference instrument (Fig. 9). The thoron calibration exercise and its results will be described in detail elsewhere. In order to test the thoron homogeneity in the BACCARA chamber during the calibration exercise 12 silica aerogel thoron samplers and 22 pieces of SSNTDs were placed at different positions in the chamber (Fig. 9). The SSNTDs were placed in all parts of BACCARA internal volume, but respecting the stated above rule for detector position. The thoron exposure duration was 48 h and the ^{220}Rn activity concentration in the chamber was 46 kBq/m^3 . In order to avoid the effect of possible ^{220}Rn inhomogeneity during the calibration, the inputs of the instruments sampling systems were put close to each other as much as possible to the extent of forming a sampling point in the chamber (Fig. 10). The silica aerogel samples and the SSNTDs were placed around the sampling and all other parts of the chamber, trying to cover the upper part of the chamber (Figures 11 and 12) as well as the space around the sampling points and around and between the detectors (Fig. 13).

After the end of the exposure, the silica aerogel samples were treated as described above. The activity in the samples is measured on a Wallac Guardian LS counter at LNHB and the results are shown in Table 3. For better visualization, the values obtained with the different samplers are shown relative to the reference point in Figures 11-13. The results in Table 3 and Fig. 11 show that the differences between the ^{220}Rn concentrations in positions “C”, near the center of the chamber, close to the instruments’ sampling points, do not exceed 10%. But, one can observe higher differences for positions “W”, far from the center of the chamber, for example the ^{220}Rn concentrations near the upper wall of the chamber (AT7) are up to 61% higher than those in the center (Fig. 11). Similarly, the thoron concentration near the walls of the chamber are higher than those at the reference point (see, for example AT4, AT5, AT11

and AT7 in Fig. 11 and 12). The ^{220}Rn concentration below the instruments (AT8) are also higher than those in the reference point (Fig. 13). Overall, it seems that the ^{220}Rn concentration near the sampling point of the instruments is homogeneous, but there are volumes in the chamber near the walls, the entry point of the thoron and in front of the fan where much higher ^{220}Rn concentrations are observed.

The SSNTD results reveal signs of inhomogeneity with up to about 60% deviation from the reference point (Fig. 14). It should be noted also that the silica aerogel sampling method and the SSNTD method give consistent results for the points where the two methods can be applied (see Fig. 12). For example, the detection volume of the highest deviation ratio 1.61(0.18) corresponds to SSNTD No 4 and it is close to the space where LSC of silica aerogel (AT5) also showed high deviation (ratio of 1.55). Similarly, SSNTD No 7 and AT11, which are close at the back of the chamber, give very close results (Fig. 12). The two methods clearly demonstrate that there can be high thoron inhomogeneity near the walls of the chamber.

5. Conclusions

Two experimental methods for evaluation of ^{220}Rn homogeneity in calibration chambers are presented. The first method is based on LSC of thoron short-lived decay products captured in silica aerogel and the second method is based on application of SSNTDs. The performance of the two methods is evaluated in several dedicated experiments where it is shown that the repeatability of the method based on LSC of silica aerogel is within 1.6%. Both methods are applied to study thoron homogeneity in a 50 L empty AlphaGuard emanation and calibration container with its fan turned on, and it was found that thoron distribution is homogeneous within 10%.

Both methods are applied successfully to test the thoron homogeneity in the BACCARA chamber during the thoron calibration exercise that was carried out at IRSN, France in the framework of the MetroRADON Euramet EMPIR project. The results show that, at the centre of the chamber, where the inputs of the instruments sampling systems were put close to each other, the thoron inhomogeneity is less than 10%. However, regions of higher thoron concentrations are clearly identified near the walls and the upper part of BACCARA, with ^{220}Rn concentrations being up to 60% higher compared to the concentration at the reference point. These results highlight the importance of the control and assessment of thoron homogeneity in thoron calibrations and in the case when radon monitors are checked for thoron influence. The assessment of ^{220}Rn homogeneity will be particularly important in the case of checking passive radon monitors for thoron influence. That is because, contrary to the case of active monitors, there are no inputs of the passive detectors, so they cannot be put together near a common sampling point in the calibration chamber. Thus, an experimental method for checking thoron homogeneity as those described in this work will be highly necessary.

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Figure captions:

Fig. 1. A photograph of a thoron sampler. a) – Empty sampler, with bottom air filter, ready to be filled with silica aerogel and the top end cap with top air filter mounted; b) the sampler filled with silica aerogel; c) closed sampler ready to be placed in thoron chamber.

Fig. 2. Scanning electron microscope images of the Millipore membrane filters used for the samplers. a) and b) –surface view, c) and d) – cross-sectional view.

Fig. 3. High performance LS glass vials (20 ml) filled with silica aerogel and Ultima Gold LLT LS cocktail. The photograph is taken after the vials had been placed for 10 min in a ultrasonic bath. When filled with LSC cocktail, the aerogel becomes translucent.

Fig. 4. Schematic view of the geometry used to test the repeatability of the silica aerogel method. a) Schematic view of the exposure vessel, the ^{220}Rn inlet, the fan and the tube with the thoron samplers; b) Illustration of the arrangement of the thoron samplers in the tube.

Fig. 5. Photographs of the experimental arrangement used to test the repeatability of the silica aerogel method: the fan, the tube and the thoron samplers can be seen in the picture.

Fig. 6. Photograph of the exposure set-up used to test the repeatability of the silica aerogel method. The fan and the tube with the samplers are placed in a 50 L vessel. Additional samplers are placed in front and near the tube.

Fig. 7. Air-volumes from which alpha particles can be detected by LR-115/II SSNTD. The volumes are shaped taking into account the fact that energy window for track registration get narrow when the incident angle is increasing.

Fig. 8. The picture shows the spots in the chamber volume from where activity can be detected by the placed in a grid Kodak Pathe LR-115/II SSNTD (“detection” volumes).

Fig. 9. A photograph of the BACCARA chamber showing the experimental set-up used in the thoron calibration exercise. The thoron inlet nozzle and the fan used to homogenize the air in the chamber are indicated as well as the positions of some of the thoron samplers and the SSNTDs.

Fig. 10. Photograph showing the sampling point, where the inlets of the thoron measuring instruments are located. There are thoron samplers (AT1, AT2 and AT6) and a SSNTD (No 16) around the sampling point. The percent in the boxes with the sampler number indicate the sampler’s readings relative deviation with respect to the reference (AT1).

Fig. 11. Position and relative deviation with respect to the reference position of the readings of the thoron samplers in the upper half of the BACCARA.

Fig. 12. Position and relative deviation with respect to the reference position of the readings of the thoron samplers in the upper half of the BACCARA and near the sampling point.

Fig. 13. Position and relative deviation with respect to the reference position of the readings of the thoron samplers located below and between the ^{220}Rn measuring instruments.

Fig. 14. The ratio of the SSNTD signal (net track density) to the signal of the detector that is at closest distance to the reference point (detector No 16).

Fig. 1

Fig. 2

Fig. 3

Fig. 4

Fig. 5

Fig. 6

Fig. 7

Fig. 8

Fig. 9

Fig. 10

Fig. 11

Fig. 12

Fig. 13.

Fig. 14

Table 1: Net counting rates per unit mass of thoron samplers with different heights exposed together to ^{220}Rn in air.

Sampler	L, cm	V, cm ³	Mass of the silica aerogel in the samplers, g	Net LS counting rate per unit mass of the silica aerogel, s ⁻¹ g ⁻¹
S1	1.0	15.21	5.421(5)	10.290(58)
S2	1.0	15.21	5.202(5)	10.415(60)
M1	1.5	22.81	8.649(5)	9.119(56)
M2	1.5	22.81	8.472(5)	9.054(58)
L1	2.0	30.41	11.932(5)	7.761(56)
L2	2.0	30.41	11.920(5)	7.714(55)

Table 2: Study of the repeatability of ^{220}Rn readings with thoron samplers. The numbers in the brackets indicate the overall estimated standard uncertainties.

Thoron sampler	Mass of the silica aerogel in the LS vial (g)	Net LS counting rate per unit mass of the silica aerogel at the end of the exposure ($\text{s}^{-1}\text{g}^{-1}$)	Difference from the mean (%)
AG 1	5.014(5)	6.44(26)	1.3%
AG 2	5.064(5)	6.24(22)	-1.9%
AG 3	5.054(5)	6.40(14)	0.6%
AG 4	4.914(5)	6.31(16)	-0.8%
AG 5	5.055(5)	6.43(16)	1.0%
AG 6	5.050(5)	6.51(18)	2.3%
AG 7	5.040(5)	6.21(21)	-2.3%
AG 8	5.042(5)	6.26(12)	-1.6%
AG 9	5.043(5)	6.37(19)	0.1%
AG 10	5.021(5)	6.44(20)	1.3%
AG 11	5.062(5)	6.48(25)	1.8%
AG 12	4.884(5)	6.22(14)	-2.2%
mean AG1-AG10 $\bar{x}$	5.030	6.36	
std. dev. s	0.059	0.10	
$u_{\text{bb}} = \text{rel. std. dev. } s/\bar{x} \text{ (%)}$	1.2 %	1.6 %	

Table 3: Homogeneity study with thoron samplers during the thoron calibration exercise performed in the BACCARRA chamber. The numbers in the brackets indicate the overall estimated standard uncertainties.

Sample	Decay corrected net counting rate, cpm	Mass of the silica aerogel in the LS vial (g)	Specific net LS counting rate at the end of the exposure($s^{-1}g^{-1}$)	Center (C) / Walls (W)	Normalized to AT1
AT1	108.2(17)	5.346(5)	20.24(32)	C	1.00
AT2	111.9(18)	5.546(5)	20.17(33)	C	1.00
AT3	139.1(20)	5.444(5)	25.56(37)	W	1.26
AT4	122.6(20)	5.471(5)	22.41(37)	C	1.11
AT5	169.6(24)	5.402(5)	31.39(44)	W	1.55
AT6	116.2(22)	5.361(5)	21.68(41)	C	1.07
AT7	179.2(26)	5.503(5)	32.57(48)	W	1.61
AT8	119.6(24)	5.550(5)	21.54(44)	W	1.06
AT9	119.5(26)	5.454(5)	21.91(47)	C	1.08
AT10	127.7(27)	5.550(5)	23.00(50)	C	1.14
AT11	125.8(29)	5.598(5)	22.47(51)	W	1.11

Declaration of interests

☐ The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

☐ The authors declare the following financial interests/personal relationships which may be considered as potential competing interests:

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Conceptualization: K. Mitev, P. Cassette, D. Pressyanov;

Methodology: K. Mitev, S. Georgiev, D. Pressyanov;

Validation: Ch. Dutsov, N. Michielsen, B. Sabot, K. Mitev, D. Pressyanov;

Formal analysis: S. Georgiev, K. Mitev, D. Pressyanov;

Writing: - Original Draft: K. Mitev, D. Pressyanov;

Writing: - Review & Editing: all authors

Funding acquisition: - K. Mitev and D. Pressyanov.