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# Photophysical Investigation of Iron(II) Complexes bearing Bidentate Annulated Isomeric Pyridine- NHC Ligands

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## ABSTRACT

The possibility of achieving luminescent and photophysically active metal-organic compounds relies on the stabilization of charge transfer states and kinetically and thermodynamically blocking non-radiative dissipative channels. In this contribution we explore the behavior of bidentate iron complexes bearing N-heterocyclic carbene ligands with extended conjugation systems by a multidisciplinary approach combining chemical synthesis, ultrafast time-resolved spectroscopy, and molecular modeling. Lifetimes of the metal-to-ligand charge transfer and metal-centered states reaching up to  $\approx 20$  picoseconds are evidenced, while complex decay mechanisms are pointed out, together with a possible influence of the facial and meridional isomerism. The structural degrees of freedom driving the non-radiative processes are highlighted and their rigidification is suggested as an effective way to further increase the lifetimes.

## INTRODUCTION

N-heterocyclic carbenes (NHCs) are highly popular ligands in catalysis due to their strong  $\sigma$ -donating properties and low dissociation rates.<sup>1-6</sup> Recently, these ligands have allowed to strongly improve the photophysical properties of iron(II) complexes making them appealing for optically active devices or as cheaper and non-toxic alternative to ruthenium-based sensitizers in solar energy applications.<sup>7-14</sup> However, controlling the photophysics of first-row transition metal complexes is challenging compared with heavier elements. For example, while ruthenium(II) polypyridyl complexes exhibit long-lived metal-to-ligand (MLCT) triplet states,<sup>15</sup> ultrafast relaxation to metal centred (MC) states and subsequent ground-state recovery are observed for iron

analogues.<sup>16–19</sup> However, the combination of the  $\pi$ -deficiency of pyridine and the  $\sigma$ -donation of NHC in pyridyl-NHC iron(II) complexes is effective in altering the energy order of the excited state manifold. Thus, important breakthroughs in increasing the MLCT lifetime have been achieved,<sup>8,20–25</sup> from tens of picoseconds for azinyl-NHC iron(II) complexes with tridentate C<sup>N</sup>C<sup>26–30</sup> and bidentate C<sup>N</sup> ligands<sup>31,32</sup> to the nanosecond domain for scorpionate-like C<sup>C</sup>C<sup>C</sup> ligands, this latter even showing luminescent properties.<sup>33</sup>

We have recently shown that for C<sup>N</sup>C ligands, framing the carbene moiety with extended  $\pi$ -conjugated systems notably stabilizes the MLCT excited state, and prevents spin crossover into the MC quintuplet <sup>5</sup>MC, by enhancing the  $\pi$ -back donation from the iron center to the ligand, hence leading to prolonged excited state lifetimes.<sup>28,30</sup> Bidentate C<sup>N</sup> ligands also brought an improvement by reducing the angular strain around the metal center upon coordination approaching a more ideal octahedral geometry and with a concomitant increase of the ligand field splitting.<sup>32</sup> Thus, these findings prompted us to investigate the design of  $\pi$ -conjugated C<sup>N</sup> ligands with the aim to exploit the combination of both geometric and electronic beneficial effects.

Nevertheless, bidentate complexes come in a mixture of *fac* and *mer* isomers, which may complicate the analysis of their spectroscopic and electrochemical properties. The effect of such isomerism has been studied in recent publications dedicated to the excited state dynamics of the parent **C0** complex (Figure 1) featuring less than 10% *fac* configuration,<sup>31,34</sup> and a 100% *fac* form of **C0** obtained by a tripodal interconnection of the three ligands.<sup>31,32</sup> Irrespective of the ligand arrangement, two excited states T<sub>1</sub> and T<sub>2</sub> are populated after intersystem crossing from the optically active <sup>1</sup>MLCT manifold. The computational analysis revealed that their character is still MLCT at the beginning of the excited-state relaxation pathway. However, the triplet manifold evolves elongating a Fe-N bond gradually and adiabatically increasing the MC character of the

states, reaching spin-crossover regions and farther singlet-triplet crossings that mediate the ground state recovery. In addition, femtosecond spectroscopy showed that the two T<sub>1</sub> and T<sub>2</sub> states differ in their recovery times, which are in the 2-3 and ≈15 ps range, respectively, regardless of the isomer considered. The isomerism solely modulates the relative population of these two excited-state deactivation channels, the *mer* form exhibiting a globally faster decay.

Herein we report the synthesis and the characterization of the mostly *mer* iron(II) complexes **C1**, **C2** and **C3**, with pyridyl-NHC bidentate ligands containing benzimidazole (BIm) and imidazo[1,5-*a*]pyridine-based (ImP) conjugated carbene moieties (Figure 1). The photophysical properties of the complexes have been investigated by ultrafast time-resolved spectroscopy and TD-DFT calculations. Unlike previous tridentate complexes studied by our group,<sup>28,30</sup> and despite significant modifications of the ground state absorption spectra, it turns out that the substitutions of imidazole by BIm or ImP do not have a significant effect on the excited state electronic structure and relaxation dynamics. The strategy based on altering the relative positions of the MC and MLCT energy levels by acting on the ligand field splitting at various levels has been instrumental in allowing the breakthrough of avoiding the spin-crossover transition and leading to the increase of the MLCT lifetimes from the sub-picosecond regime to tens of picoseconds.<sup>8,21,28–30,33</sup> However, it appears that, for complexes with 3 or 4 carbene bonds, this strategy is reaching its limit, as also shown by the present results. To further increase the lifetimes, a kinetic blocking of the geometrical degrees of freedom leading to the deactivation should be effective, as confirmed by the general better performance of the more rigid tridentate complexes. Indeed, iron compounds showing lifetimes at the nanosecond scale present multiple  $\pi$ -stacking interactions participating to the rigidification of the coordination sphere.<sup>22</sup> The results presented in the present article allow for a rationalization of the different factors leading to the excited states deactivation and can hence



three-electrode cell. The saturated calomel electrode (SCE) was separated from the test compartment using a bridge tube. The solutions of studied complexes (0.2 mM) were purged with argon before each measurement. The test solution was acetonitrile containing 0.1 M Bu<sub>4</sub>NPF<sub>6</sub> as supporting electrolyte. The working electrode was a vitreous carbon rod (1cm<sup>2</sup>) wire and the counter-electrode was a 1 cm<sup>2</sup> Pt disc. After the measurement, ferrocene was added as the internal reference for calibration. All potentials were quoted versus SCE. In these conditions the redox potential of the couple Fc<sup>+</sup>/Fc was found at 0.39V. In all the experiments the scan rate was 100mV/s.

### Synthesis of ligands

**L1.**<sup>35</sup> 2-Pyridylbenzimidazole (400 mg, 2.05 mmol) and iodomethane (2.20 mL, 20.51 mmol) were heated for 1h in 4 mL of acetonitrile at 100°C under microwave irradiation (60W). After evaporation of solvent, the resulting mixture was dissolved in water. Then, the metathesis was made by addition of saturated solution of KPF<sub>6</sub>. The solid was filtered, washed twice with Et<sub>2</sub>O, and dried under vacuum to give a yellow powder (469 mg, 66%). <sup>1</sup>H NMR (400 MHz, DMSO-d<sub>6</sub>) : δ = 10.40 (s, 1H), 8.77 (ddd, *J* = 0.80, 1.80, 4.80 Hz, 1H), 8.48 - 8.43 (m, 1H), 8.27 (dt, *J* = 8.10, 2.00 Hz, 1 H), 8.15 - 8.10 (m, 1 H), 8.01 (td, *J* = 8.300, 1.04 Hz, 1H), 7.82 - 7.75 (m, 2H), 7.71 (ddd, *J* = 7.48, 4.78, 1.04 Hz, 1H), 4.18 (s, 3H) ppm. <sup>13</sup>C NMR (100 MHz, ACN-d<sub>3</sub>) : δ = 149.5, 147.2, 142.5, 140.6, 132.2, 129.3, 127.3, 127.1, 125.0, 116.8, 115.6, 113.8, 33.7 ppm.

**L2.**<sup>36</sup> 2-aminopyridine (400 mg, 4.25 mmol) and paraformaldehyde (192 mg, 6.38 mmol) were heated for 30 min in 4 mL of acetonitrile at 80°C under microwave irradiation (60W). Then picolinaldehyde (0.41 mL, 4.25 mmol) and a solution of HCl 3M in acetonitrile (2 mL) were added

and heated for 1h at 100°C under microwave irradiation (60W). After evaporation of solvent, the resulting mixture was dissolved in water. Then, the metathesis was made by addition of saturated solution of KPF<sub>6</sub>. The solid was filtered, washed twice with Et<sub>2</sub>O, and dried under vacuum to give yellow powder (1.43 g, 98%). <sup>1</sup>H NMR (400 MHz, DMSO-d<sub>6</sub>) : δ = 10.57 (s, 1 H), 8.99 (s, 1 H), 8.74 (ddd, *J* = 4.80, 1.80, 0.80 Hz, 1 H), 8.60 (dd, *J* = 7.18, 0.88 Hz, 1 H), 8.29 (td, *J* = 7.50, 1.80 Hz, 1 H), 8.21 (d, *J* = 8.10 Hz, 1 H), 7.93 (d, *J* = 9.32 Hz, 1 H), 7.73 (ddd, *J* = 7.43, 4.78, 0.88 Hz, 1 H), 7.34 (ddd, *J* = 9.30, 2.30, 0.80 Hz, 1 H), 7.28 (td, *J* = 6.80, 1.30 Hz, 1 H) ppm.

**L3.**<sup>36</sup> 2-aminopyridine (400 mg, 4.25 mmol) and paraformaldehyde (192 mg, 6.38 mmol) were heated for 30 min in 4 mL of acetonitrile at 80°C under microwave irradiation (60W). Then 2-acetylpyridine (0.48 mL, 4.25 mmol) and a solution of HCl 3M in acetonitrile (2 mL) were added and heated for 1h at 100°C under microwave irradiation (60W). After evaporation of solvent, the resulting mixture was dissolved in water. Then, the metathesis was made by addition of saturated solution of KPF<sub>6</sub>. The solid was filtered, washed twice with Et<sub>2</sub>O, and dried under vacuum to give yellow powder (1 g, 67%). <sup>1</sup>H NMR (400 MHz, DMSO-d<sub>6</sub>) : δ = 10.16 (s, 1H), 8.90 - 8.74(m, 1H), 8.60 - 8.45 (m, 1H), 8.29 (dt, *J* = 7.81, 1.76 Hz, 1H), 8.03 (d, *J* = 8.56 Hz, 1H), 7.99–7.93 (m, 1H), 7.80 (ddd, *J* = 7.55, 4.78, 1.01 Hz, 1H), 7.32–7.19 (m, 2H), 2.71 (s, 3H) ppm.

### Synthesis of complexes

**C1.** Using Schlenk techniques, ligand **L1** (200 mg, 0.56 mmol) and FeCl<sub>2</sub> (24 mg, 0.19 mmol) were dissolved in 2 mL of DMF. Potassium *tert*-butoxide (70 mg, 0.62 mmol) was then added. The dark red solution was stirred at room temperature (rt) for 10 min. After evaporation of DMF, the residue was purified by column chromatography using water/acetone/KNO<sub>3</sub> (10/2/1) as eluent. The red fractions were evaporated, and a saturated solution of KPF<sub>6</sub> was added. The obtained solid

was filtered, washed twice with Et<sub>2</sub>O, and dried under vacuum to give a non-separable mixture of both *fac*- and *mer*-**C1** in 1:6 ratio (60 mg, 33%) as an orange solid. <sup>1</sup>H NMR (400 MHz, ACN-d<sub>3</sub>) : δ = 8.51 (d, *J* = 8.31 Hz, 1 H), 8.44 (d, *J* = 8.30 Hz, 1 H), 8.32 (d, *J* = 7.81 Hz, 1 H), 8.11–8.28 (m, 7 H), 8.08 (d, *J* = 4.78 Hz, 1 H), 7.91–8.00 (m, 2 H), 7.43 - 7.68 (m, 13 H), 7.21–7.29 (m, 2 H), 7.10 (d, *J* = 5.30 Hz, 1 H), 6.88–7.00 (m, 3 H), 3.41 (s, 3 H), 3.19 (s, 9 H), 2.94 (s, 3 H), 2.89 (s, 3 H) ppm. <sup>13</sup>C NMR (100 MHz, ACN-d<sub>3</sub>) : δ = 218.2, 218.2, 215.7, 215.7, 157.4, 156.9, 156.3, 155.9, 155.8, 155.0, 152.4, 150.8, 142.2, 141.5, 141.1, 140.0, 140.0, 139.7, 139.6, 139.2, 132.9, 132.5, 132.5, 132.4, 125.9, 125.9, 125.7, 125.7, 125.6, 125.5, 125.3, 124.9, 124.6, 124.0, 123.4, 122.8, 114.7, 114.4, 113.9, 133.5, 113.2, 113.1, 112.5, 122.4, 112.3, 111.9, 111.7, 111.2, 34.3, 34.0, 32.5, 32.4 ppm. HRMS calcd for C<sub>39</sub>H<sub>33</sub>FeN<sub>9</sub>P<sub>2</sub>F<sub>12</sub> *m/z* : 341.6099 [M – 2PF<sub>6</sub>]<sup>2+</sup>. Found 341.6104.

**C2.** The procedure was the same as for **C1** using **L2** (600 mg, 1.76 mmol), FeCl<sub>2</sub> (74 mg, 0.58 mmol) and potassium *tert*-butoxide (217 mg, 1.94 mmol) in 6 mL of DMF. The column chromatography was performed with water/acetone/KNO<sub>3</sub> (9/2/1) as eluent. A non-separable mixture of both *fac*- and *mer*-**C2** in 1:3 ratio (220 mg, 38%) was obtained as an orange solid. <sup>1</sup>H NMR (400 MHz, ACN-d<sub>3</sub>) : δ = 8.65 (s, 1 H), 8.55 (s, 1 H), 8.52 (s, 1 H), 8.46 (s, 1 H), 8.12–8.31 (m, 5 H), 7.94–8.04 (m, 4 H), 7.82 (dd, *J* = 5.80, 0.80 Hz, 1 H), 7.57 (td, *J* = 9.30, 1.30 Hz, 2 H), 7.47 (d, *J* = 9.30 Hz, 2 H), 7.29–7.37 (m, 3 H), 7.14 (d, 1 H), 6.99–7.10 (m, 3 H), 6.78–6.94 (m, 4 H), 6.32 (m, 4 H), 5.99 (dd, *J* = 7.30, 1.00 Hz, 1 H), 5.70 (dd, *J* = 7.50, 1.00 Hz, 1 H), 5.62 (dd, *J* = 7.30, 1.00 Hz, 1 H) ppm. <sup>13</sup>C NMR (100 MHz, ACN-d<sub>3</sub>) : δ = 200.2, 198.4, 195.7, 195.0, 194.1, 157.6, 156.6, 156.4, 155.9, 155.5, 154.6, 153.1, 151.6, 142.1, 141.3, 141.1, 140.1, 136.5, 136.4, 136.3, 136.2, 126.0, 125.9, 125.6, 125.3, 125.2, 125.0, 125.0, 124.9, 124.8, 124.8, 124.7, 124.1,

120.6, 120.5, 120.5, 120.5, 116.9, 116.7, 116.2, 115.9, 114.9, 114.4, 114.4, 113.9, 110.6, 109.6, 109.5 ppm. HRMS calcd for  $C_{36}H_{27}FeN_9P_2F_{12}$   $m/z$  : 320.5864  $[M - 2PF_6]^{2+}$ . Found 320.5856.

**C3.** The procedure was the same as for **C1** using **L3** (600 mg, 1.69 mmol),  $FeCl_2$  (70 mg, 0.56 mmol) and potassium *tert*-butoxide (209 mg, 1.86 mmol) in 6 mL of DMF. The column chromatography was performed with water/acetone/ $KNO_3$  (9/2/1) as eluent. A non-separable mixture of both *fac*- and *mer*-**C3** in 1:4 ratio (153 mg, 27%) as an orange solid was obtained.  $^1H$  NMR (400 MHz,  $ACN-d_3$ ) :  $\delta$  = 8.42 (d,  $J$  = 8.55 Hz, 1H), 8.31 (d,  $J$  = 8.55 Hz, 1H), 8.18–8.06 (m, 4H), 7.94–7.87 (m, 2H), 7.83 (dd,  $J$  = 5.54, 1.76 Hz, 1H), 7.46–7.39 (m, 3H), 7.35 (s, 1H), 7.33 (s, 1H), 7.29–7.24 (m, 2H), 7.15 (dd,  $J$  = 5.55, 1.46 Hz, 1H), 7.11 (dd,  $J$  = 6.05, 1.08 Hz, 1H), 7.02 (t,  $J$  = 6.54 Hz, 1H), 6.96 (t,  $J$  = 6.51 Hz, 1H), 6.77–6.62 (m, 4H), 6.28 (t,  $J$  = 6.24 Hz, 1H), 6.22 - 6.07 (m, 4H), 5.85 (d,  $J$  = 7.46 Hz, 1H), 5.55 (d,  $J$  = 7.46 Hz, 1H), 5.51 (d,  $J$  = 7.46 Hz, 1H), 3.02 (s, 3H), 2.92 (s, 6H), 2.86 (s, 3H) ppm.  $^{13}C$  NMR (100 MHz,  $ACN-d_3$ ) :  $\delta$  = 200.2, 198.5, 197.6, 195.8, 194.1, 157.9, 157.8, 156.8, 157.0, 156.7, 155.8, 153.3, 151.8, 141.5, 140.8, 140.5, 139.4, 133.2, 133.2, 133.0, 132.9, 125.6, 125.5, 125.0, 125.0, 124.7, 142.7, 124.4, 123.7, 122.9, 122.8, 122.8, 122.7, 122.4, 122.3, 121.4, 121.2, 119.8, 119.7, 119.7, 116.7, 116.6, 116.6, 116.1, 116.0, 116.0, 115.7, 115.4, 12.3, 12.1, 12.1, 12.0 ppm. HRMS calcd for  $C_{39}H_{33}FeN_9P_2F_{12}$   $m/z$  : 341.6102  $[M - 2PF_6]^{2+}$ . Found 341.6099.

### Computational details

The Franck-Condon geometries ( $S_0$  min) of the *fac* and *mer* isomers of **C1** and **C2** have been optimized using the DFT/B3LYP functional in combination with the 6-31+G(d,p), as implemented in the GAUSSIAN 09 (D.01 revision) software package.<sup>37</sup> This computational protocol has been able to describe the Franck-Condon geometry of related Fe(II) complexes bearing other NHC

ligands.<sup>23,28,30–32,34</sup> Note that only **C1** and **C2** complexes have been considered in the computational analysis since no significant differences, from a qualitative point of view, have been found in the photophysical properties of **C2** and **C3**, given the similarities in the chemical structure.

Vertical absorption energies at the Franck-Condon region have been computed through the TD-DFT method employing the Tamm-Dancoff approximation.<sup>38</sup> The exchange correlation functional HCTH407 (hereafter, TD-HCTH)<sup>39</sup> have been used given the good description of the excited states provided by this method in similar Fe-NHC systems as documented in previous works.<sup>8,28,31,27,40</sup> The theoretical spectra reported in ESI have been obtained by sampling the nuclear density of the Franck-Condon region by means of a semiclassical Wigner distribution of 20 structures.

Right after, the excitation energies have been convoluted with Gaussian functions with a full width at half-length of 0.3 eV. A total number of 30 singlet-singlet electronic transitions have been computed for each geometry.

Note that the states displayed in Figure 7 are labelled according to the multiplicity of the state and the adiabatic energetic ordering without considering the nature of them. The nature of the excitations progressively evolves along the potential energy surfaces<sup>31</sup> (see Figures S28-S30) and the nature of the excitations have been determined by computing the NTOs,<sup>41,42</sup> *via* post-processing of their density matrices obtained with the TD-HCTH method, using the NANCY\_EX code.<sup>42,43</sup> For the readers' convenience, we remind that NTOs may be considered as the best orbitals to describe a given electronic transition, and in contrast to the Kohn-Sham representation, often one or two couples of occupied/virtual orbitals are sufficient to fully describe the electronic density rearrangement in the excited state.

The S<sub>1</sub> min geometries have been computed using the TD-DFT/HCTH method, whereas the T<sub>1</sub> min and Q<sub>1</sub> min structures have been optimized using the unrestricted DFT/HCTH ansatz. No

symmetry or geometrical constraints have been imposed in the optimizations or single-point calculations.

### **Femtosecond transient absorption spectroscopy**

In the transient absorption setup, an amplified 5kHz Ti: sapphire laser generates 30 fs 0.5 mJ pulses to pump a commercial optical parametric amplifier (TOPAS: Light Conversion), from which 50-60 fs pulses are derived to excite the samples: 370 nm for **C1**, and 470 nm for **C2** and **C3**. A white light continuum is generated in a 2 mm thick CaF<sub>2</sub> crystal, mounted on an oscillating loudspeaker to reduce photo-damage. The white light is optimised for minimised fluctuations (rms < 0.4% at 220 Hz). The beam is split in two: the probe that is sent through the sample, and a reference beam is used for measuring and compensating the white-light intensity fluctuations. The polarization of the probe beam is set at magic angle (54.7°) with respect to the pump.

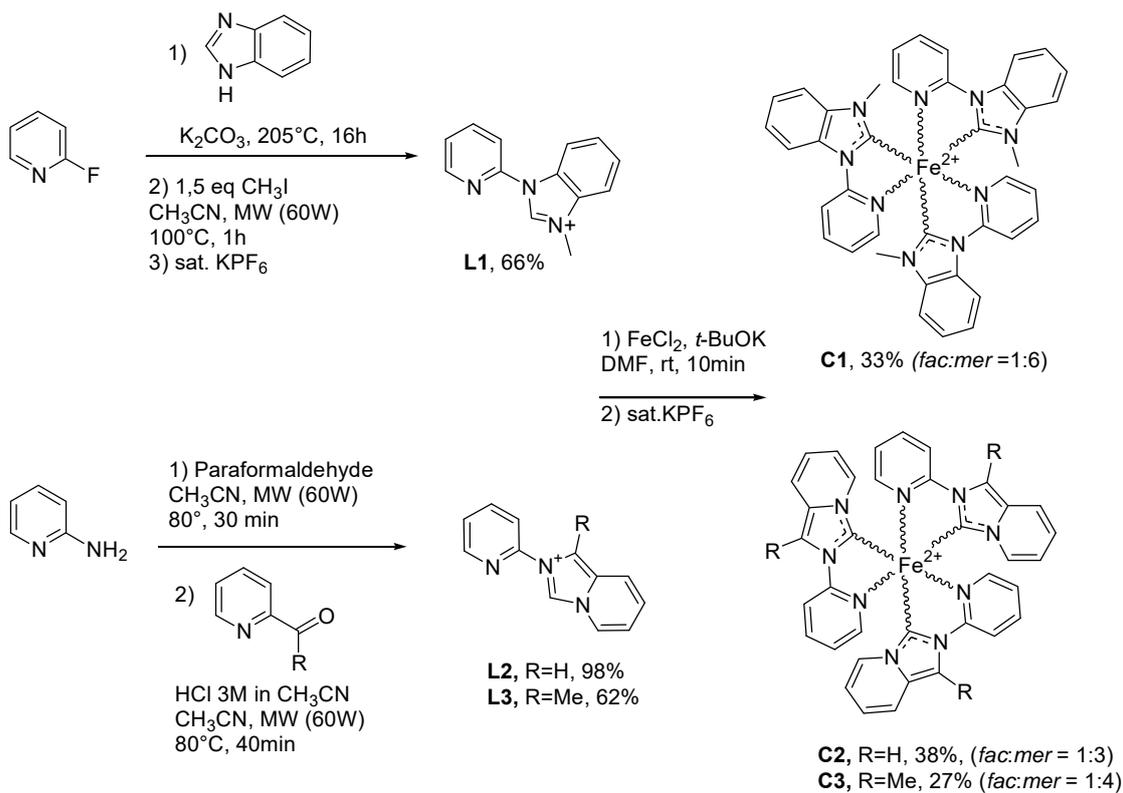
A 1 mm path length cuvette in fused silica contains the complexes dissolved in CH<sub>3</sub>CN. Time-resolved spectra  $I(\lambda)$  are acquired in the range  $\approx 300$  nm wide window with an adjustable central wavelength as a function of pump-probe delay, by a combination of 25 cm focal length spectrometer (resolution 2 nm) and a Peltier-cooled CCD with 220 Hz acquisition rate. A second chopper blocks the pump beam at 110 Hz and a home-made software computes and displays the differential spectra  $\Delta A = -\log_{10} \left( \frac{I_{\text{pump-on}}}{I_{\text{pump-off}}} \right)$ . A solvent-only sample is measured, and the data are processed to (i) remove the background at negative delay times, (ii) the solvent Raman signal and the coherent interactions of the pump and probe in the cell, (iii) to correct for the group velocity dispersion of the probe beam, characterized in the solvent-only data set. The temporal resolution is characterized

by the  $50\pm 5$  fs FWHM of the solvent Raman response. For more details see Ref 44. The multi-exponential fit procedure is described in the SI.

## RESULTS AND DISCUSSION

### Synthesis

The ligand precursors **L1**, **L2** and **L3** required for the preparation of the target complexes have been prepared according to scheme 1. **L1** was obtained in good yield by an amination-quaternarization sequence from 2-fluoropyridine. For the synthesis of **L2** and **L3** a construction of the imidazo[1,5-*a*]pyridine core was performed.<sup>45</sup> An efficient cyclizing condensation involving 2-aminopyridine and 2-formyl or 2-acetyl pyridine afforded respectively **L2** and **L3** in good yields. The three target complexes **C1–C3** were obtained in moderate yields by first reacting the ligand precursors with FeCl<sub>2</sub>, followed by addition of *t*-BuOK as the base for generation of the carbenes. The parent compound **C0** was also synthesised for the sake of comparison following our previously reported procedure.<sup>32</sup>

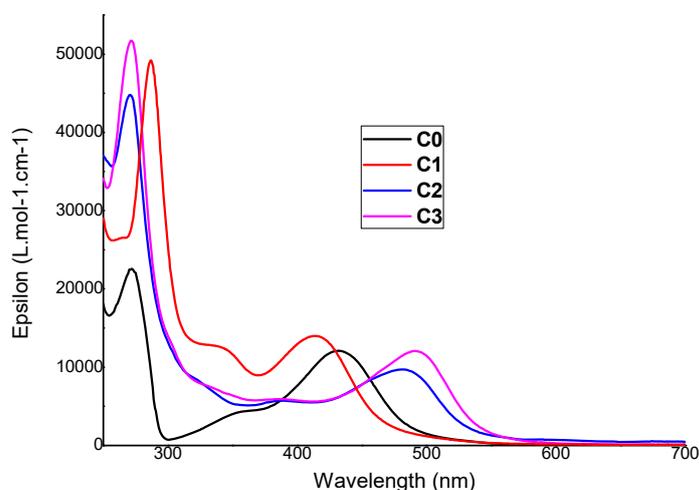


**Scheme 1.** Synthesis of ligands **L1**, **L2** and **L3** and the corresponding complexes **C1**, **C2** and **C3**.

Counter ion is  $\text{PF}_6^-$  in every case.

### Electronic and electrochemical properties

The electronic and optical properties of the **C0–C3** complexes were investigated by means of steady-state UV-vis spectroscopy (Figure 2) and cyclic voltammetry (Table 1).



**Figure 2.** UV-Vis spectra of **C0–C3** in acetonitrile. Only the 250–700 nm window is shown.

The UV/Vis spectra of all the complexes present three evident absorption bands. The bands in the UV region, below 300 nm, correspond to  $\pi-\pi^*$  transitions localized on the ligands, and are particularly intense. Two distinct broader and less intense bands are found at longer wavelengths, corresponding to MLCT transitions, with significant shifts of the absorption maximum ( $\lambda_{\max}$ ) as a function of the specific ligand. The weaker bands peaking in the 320–420 nm range are assigned to carbene-Fe MLCT transitions,<sup>27,28</sup> while the bands with  $\lambda_{\max}$  in the 430–600 nm area are due to the pyridine-Fe MLCT's. Comparison of the **C0** and **C1** absorption spectra shows that both MLCT bands are markedly blue-shifted when switching from the imidazolylidene (Im) to BIm due to the increased  $\pi$ -back donation induced by the annulation of the carbene moiety.<sup>28,32</sup> On the contrary, the inclusion of ImP significantly red-shifts both the carbene-Fe (30–35 nm) and the pyridine-Fe (50–60 nm) absorption as compared to **C0** and **C1**. Interestingly, the methyl substitution on ImP (**C3**) induces an additional slight red-shift with respect to the unsubstituted analogue (**C2**).

The TD-DFT electronic spectra for **C1** and **C2** are shown in Figure S27 and are in reasonably good agreement with the experimental measurements. Computational results also allow the disentanglement of the different contributions of the *fac* and *mer* isomers: for **C1**, absorption of *mer* isomer is slightly red-shifted with respect to the *fac* analogue, whereas **C2** exhibits an opposite trend. Indeed, the <sup>1</sup>MLCT band of *fac*-**C2** is red-shifted *ca.* 100 nm. This effect is masked in the experimental data by the *fac:mer* ratio of 1:3 and, as a matter of fact, absorption of *mer* isomer peaking at ~500 nm dominates the spectrum of the **C2** mixture. Natural transition orbitals (NTOs) confirm the <sup>1</sup>MLCT nature of the bright states responsible for the absorptions, as displayed in Figure S28 of the Supporting Information.

The cyclic voltammetry data (Table 1 and Figure S12–S14) globally confirm the trends underlined by the absorption spectra with a decrease of the electrochemical gap going from **C0** to **C2** and **C3** in agreement with the observed redshifts. The positive domain corresponds to the reversible electron transfer of one electron per complex leading to the oxidation of Fe(II) to Fe(III). The oxidation potential is only slightly affected by switching from Im to ImP, while the further inclusion of a methyl group (**C3**) induces a decrease of the oxidation potential consistently with the observed redshift. On the contrary, the BIm unit results in an increase of the oxidation potential of **C1** compared to **C0**, **C2**, and **C3**. Extending the conjugation upon annulation has been reported to increase the accepting character of BIm<sup>46</sup> and ImP<sup>47</sup> due to the delocalisation of the lone pairs of the nitrogen atoms, reducing the occupation of the  $p_{\pi}$  carbene orbital and allowing for an increased back-donation from the metal. However, these results highlight the remarkable differences in  $\pi$  back-donation between both isomeric carbenes as evidenced by the distinct stabilization of the lower occupied iron orbitals (commonly referred to  $t_{2g}$  for complexes with  $O_h$  symmetry), with BIm being the best  $\pi$ -acceptor.

In the negative domain, two separate irreversible mono-electronic transfers are present for **C0** while for **C1**, **C2** and **C3** the two waves coalesce into a bielectronic reversible process. It can be assumed that the LUMO levels are stabilized by about 0.25 eV by comparison with **C0**, suggesting that each electron is mainly injected into ImP or BIm moiety instead than pyridine as for **C0**. The increased conjugation of ImP and BIm would nicely justify the effective stabilisation of the LUMO levels in comparison to that associated to the pyridine moiety.

**Table 1.** Optical and electrochemical properties of the studied complexes

	$\lambda_{\text{abs-max}}$ (nm) ( $\epsilon(\text{M}^{-1}.\text{cm}^{-1})$ ) <sup>a</sup>	$E_{\text{ox}}(\text{Fe}^{\text{III}}/\text{Fe}^{\text{II}})$ (V/SCE) <sup>b</sup>	$E_{\text{red1}}$ (V/SCE)	$\Delta E$ (eV) <sup>c</sup>
<b>C0</b>	358(4300)	0.67 (rev)	-1.88 (irrev.)	2.55
	432(12100)		-1.93 (irrev)	2.60
<b>C1</b>	337 (12700)	0.87 (rev)	-1.63 (rev)	2.50
	414 (13980)			
<b>C2</b>	394 (5670)	0.70 (rev)	-1.68 (rev)	2.38
	481 (9700)			
<b>C3</b>	389 (5910)	0.61 (rev)	-1.66 (rev)	2.27
	490 (12100)			

<sup>a</sup> Measured in CH<sub>3</sub>CN at 25 °C. <sup>b</sup>First oxidation potential. Potentials are quoted vs SCE. Under these conditions,  $E_{1/2}(\text{Fc}^+/\text{Fc}) = 0.39\text{V/S.C.E.}$  Recorded in CH<sub>3</sub>CN using Bu<sub>4</sub>N<sup>+</sup>PF<sub>6</sub><sup>-</sup> (0.1M) as supporting electrolyte at 100 mV. s<sup>-1</sup>. <sup>c</sup> Electrochemical band gap ( $\Delta E = E_{\text{ox}} - E_{\text{red1}}$ )

### Time-resolved spectroscopy

We investigated **C1**, **C2** and **C3** with femtosecond transient absorption spectroscopy (TAS) in acetonitrile at room temperature. With the background of our first publication on complex **C0** and

a *fac*-only version of it,<sup>31,32</sup> these experiments were expected to reveal an extension of the excited state lifetime due to the fusion of a phenyl group to the imidazole part of the ligands, as observed for tridentates. Two main questions arise: How do these modified ligands change the excited state relaxation? Is the scenario of parallel population and decay of two excited states reported for **C0** and *fac*-only **C0** also valid for **C1–C3**?

Transient absorption measures the photo-induced changes of the optical density or absorbance  $\Delta A$  and can be expressed by:

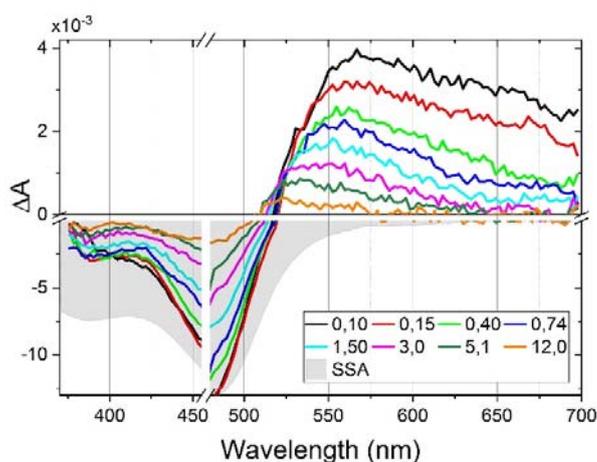
$$\Delta A(\lambda, t) = d \sum_i \epsilon_i(\lambda, t) \Delta c_i(t) \quad (1)$$

with  $d$  being the sample thickness. The sum runs over all molecular electronic states  $i$  whose population is changed by the laser excitation at  $t=0$ , and the subsequent photoreactions and their kinetics, as expressed by the molar concentration change  $\Delta c_i(t)$ . The latter is positive for excited state population, leading to positive excited state absorption (ESA), and negative for the ground state population,  $\Delta c_0(t=0) < 0$ , due to the partial depletion of the ground state by the laser excitation leading to ground state bleach (GSB) signals.

As an example, Figure 3 displays the transient absorption spectra of **C2** for time delays covering the first 12 ps after excitation with a 40 fs pulse at 470 nm. A 20 nm spectral range around this wavelength is dominated by scattered pump laser light, and therefore not displayed. A broad ESA band dominates the signal for  $\lambda > 520$  nm, while GSB is dominant for shorter wavelengths, leading to  $\Delta A < 0$ . The grey area is the inverted ground state absorption spectrum (SSA), normalised such as to fit the GSB amplitude at 485 nm. The neg.  $\Delta A$  portion is spectrally significantly sharper than the SSA, indicating, according to eq. (1), the contribution of ESA components, as is clearly visible for  $\lambda > 520$  nm, but hidden for shorter wavelengths where GSB dominates. Such an overlap of

excited state and ground state signals is a general limitation of TAS, leading to uncertainties as will be discussed below.

The time-dependent  $\Delta A$  spectra of **C3** are like the one of **C2** (Fig. S15). For **C1**, the ESA/GSB pattern is obviously blue-shifted according to the SSA (Fig. 2). The dynamic changes of the ESA spectra are more pronounced than for **C2** and **C3**, and resemble the evolution of **C0**.<sup>31</sup>

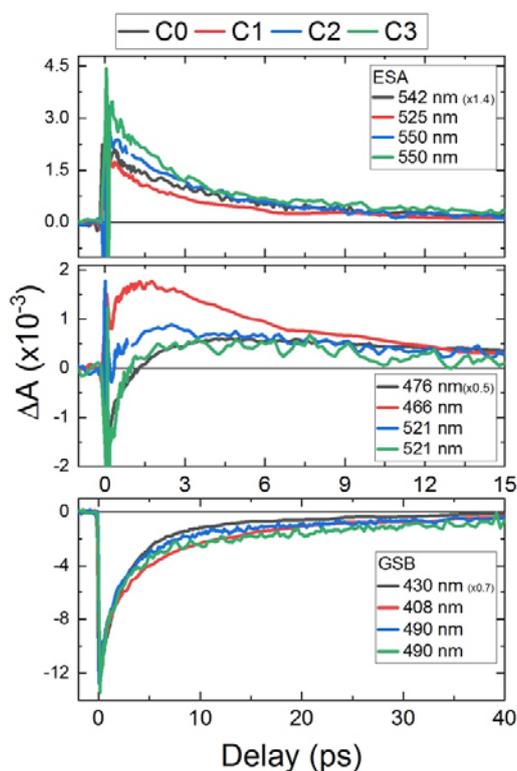


**Figure 3.** Transition absorption spectra of complex **C2** during the first 12 ps. Shaded area is the inverted steady-state absorption spectrum (SSA) scaled to match the GSB max. Note the break in the y-axis, and the different scales for the positive and negative signals. The 460–480 nm range is disregarded due to excessive pump light scattering. The progressive change in the spectral shape of the positive ESA is discussed in terms of the superposition and population of two excited states  $T_1$  and  $T_2$ . The signal decays dominantly on a  $\approx 3$  ps time scale.

A detailed inspection of the spectral features of **C2** reveals important information about the excited state evolution. We see that ESA and GSB decay on three different time scales. The ESA amplitude decays by half already in the first 0.4 ps, and the spectral width reduces since the decay is most prominent for  $\lambda > 600$  nm. The negative  $\Delta A$  part decreases by not more than 10% during this time. On a  $\approx 3$  ps time scale both ESA- and GSB-dominant parts decay by  $\approx 50\%$ , and the ESA

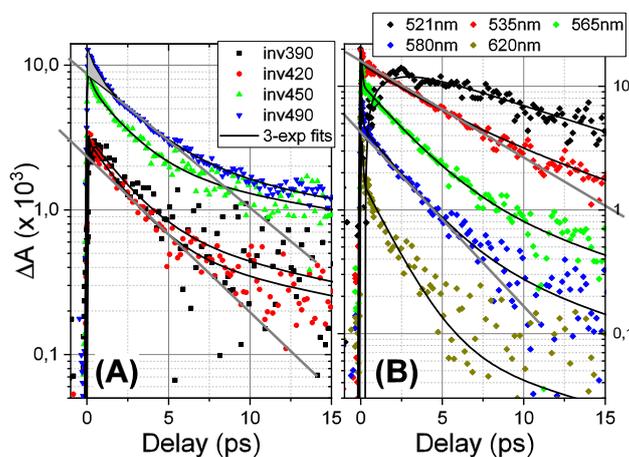
spectrum slightly blue shifts, as can also be seen by the ESA–GSB crossover point ( $\Delta A=0$ ) shifting by 5 nm to the blue. A third stage of  $\Delta A$  decay proceeds on a slower time scale of  $\approx 10$  ps, with the ESA further evolving in shape and energy. Qualitatively, the same behaviour is observed for **C1** and **C3** (Fig. S15).

Figure 4 depicts kinetic traces for all four complexes **C0–C3** at selected wavelengths. A quick glance shows that the BIm and ImP moieties have only minor effects on the excited state kinetics. In particular, the kinetic traces in the top and bottom panels highlighting dominant excited state decay (ESA) and ground state repopulation (GSB) almost overlay. The middle panel shows the GSB/ESA crossover region, highlighting the above described spectral blue shift. On a timescale of 2–3 ps, the positive ESA contribution rises turning the signal from dominant GSB ( $\Delta A < 0$ ) to dominant ESA ( $\Delta A > 0$ ).



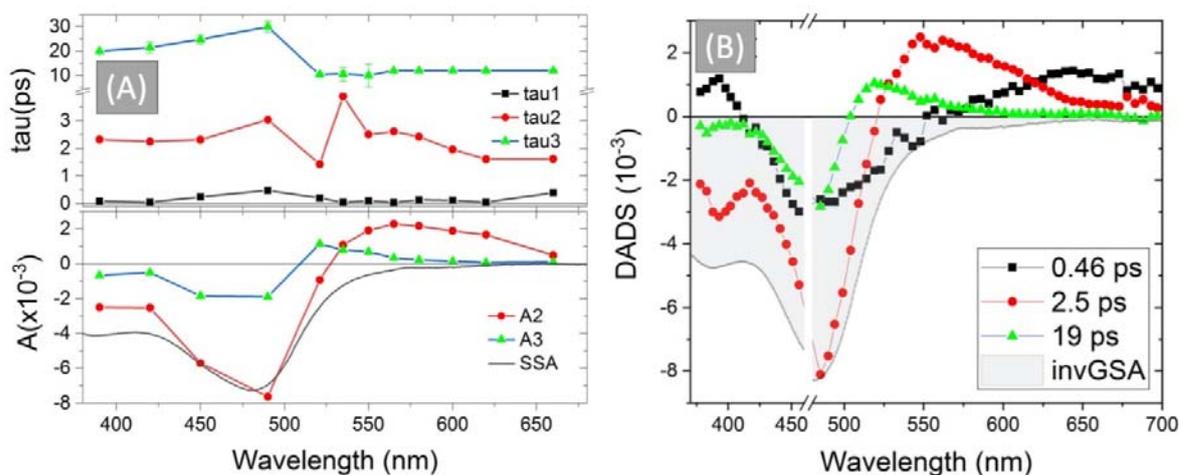
**Figure 4.** Comparison of kinetic traces for complexes **C0–C3** at selected wavelengths (**C0**-black; **C1**-red; **C2**-blue; **C3**-green). *Top:* early-delay ESA maximum, *middle:* GSB-ESA crossover region, *bottom:* long-wavelength GSB tail.

The kinetic transients of **C2** are reported in Figure 5 together with 3-exponential fits, convoluted by the instrument response function (50 fs FWHM, see SI for details). Plotted on a semi-log scale, the relative contributions of the three different time scales are particularly apparent for the GSB plotted with an inverse sign for the logarithmic intensity scale (Fig. 5A). A straight line is plotted as a guide-to-the eye, the slope of which corresponds to a  $\approx 3$  ps decay time, but a faster sub-picosecond decay and a slower  $>10$  ps decay time are clearly visible. For the dominant ESA wavelengths (Fig. 5B), the guide-to-the-eye indicates that the decay is wavelength dependent, with faster decays observed for longer wavelengths, consistent with the qualitative inspection of the time-resolved spectra (Fig. 3).



**Figure 5.** Kinetic transients for **C2** and 3-exponential fits in the GSB (A) and dominant ESA(B) regions, plotted on logarithmic intensity scale, highlighting the multi-exponential decay. (A): Selected traces of the GSB region, plotted with inverted sign. Lines correspond to a 3.0 ps decay. (B): Selected traces from the ESA long-wavelength region.

All kinetic traces can be well fitted by a sum of three exponentials with wavelength-dependent amplitudes and decay constants,  $\sum_{i=1}^3 A_i(\lambda)e^{-t/\tau_i(\lambda)}$ , convoluted with a Gaussian instrument response function (width 50 fs). The results of these fits are plotted in Figure 6A, and detailed plots of the fit curves and their residuals can be found in the SI, for **C1** and **C3** (Fig. S16 and S17). Since the spectral and kinetic evolution is very similar among the complexes, we discuss the fit results only for **C2**. The shortest decay time  $\tau_1$  is found in the 0.1–0.2 ps range and shows slightly larger values (0.4 ps) only for the range of maximum negative  $\Delta A$  (450–500 nm), where the signal-to-noise ratio is highest. The amplitudes of  $\tau_1$  are not displayed since their values are strongly wavelength-dependent and no clear interpretation can be given. For almost all wavelengths,  $\tau_2$  is the dominant decay component ( $A_2 > 3A_3$ ). Its value is slightly wavelength-dependent, with an average value of  $\tau_2 = 2.5 \pm 0.3$  ps, and dropping from 2.5 to 1.5 ps in the red. In the crossover region (520–540 nm), the fitted times are the results of compensating rising and decaying exponentials and thus fluctuate strongly, without any significance for the actual excited state lifetimes.



**Figure 6.** Results of 3-exponential fittings for **C2**. **A-Top:** Values of the 3 wavelength-dependent time constants. The symbol size indicates typical fit error bars.  $\tau_3$  is kept fixed at 12 ps for  $\lambda \geq 580$  nm. **A-Bottom:** Decay amplitudes  $A_2$  and  $A_3$ . **B:** Decay-associated difference spectra (DADS) obtained with a global fit approximation. The inverted and scaled ground state absorption spectra (SSA/invGSA) are indicated for comparison.

The spectral shape of  $A_2$  is in good agreement with the time-resolved spectra, displaying a dominant ESA for  $\lambda > 520$  nm. For  $\tau_3$ , the fits give values of 20–30 ps in the GSB region ( $\lambda < 500$  nm), and around 10 ps in the 520–550 nm ESA region, where the amplitude is large enough to determine this component reliably. The spectral distribution of  $A_3$  agrees with the long-delay  $\Delta A$  spectra, for  $t > 10$  ps (Fig. 3), with a maximum signal at 520–525 nm. Since the spectral shifts and the wavelength-dependence of the lifetimes are relatively small, we have performed a global analysis assuming wavelength-independent average lifetimes, using the shareware software OPTIMUS.<sup>48</sup> The validity of wavelength-independent average lifetimes is assessed in detail in the SI. It is found that the fit quality is very high except for the wavelengths of ESA/GSB crossover, displaying a dynamically shifting point of  $\Delta A = 0$ . The results are displayed in Fig. 6B and show that the shape and positions of the maxima of the decay-associated difference spectra (DADS) agree with the  $A_2$  and  $A_3$  distributions obtained by individual wavelength fits (Fig. 6A). For all the complexes **C1** to **C3**, the average lifetimes given in Table 2 are derived from such a global fit approach.

**Table 2.** Values of the three decay times and the  $A_2/A_3$  ratio at the wavelength of maximum GSB for the four complexes **C0–C3**. Values extracted from global fits for **C1–C3** (cf. SI).

Complex	$\tau_1$ (ps)	$\tau_2$ (ps)	$\tau_3$ (ps)	$A_2/A_3$
<b>C0</b> <sup>a</sup>	0.5±0.1	2.8±0.5	17±3	7.0

<b>C1</b>	0.6±0.1	3.1±0.5	15±3	1.3
<b>C2</b>	0.4±0.1	2.5±0.5	19±3	4.0
<b>C3</b>	0.3±0.1	1.8±0.5	24±3	2.3

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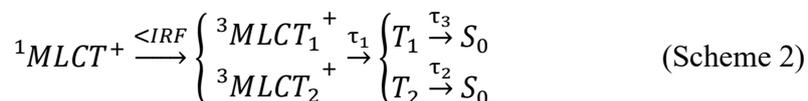
<sup>a</sup> Taken from ref. <sup>32</sup>.

The critical point is now to build a reaction scheme involving the three lifetimes. The sub-picosecond  $\tau_1$  is usually assigned to the transition from the optically excited <sup>1</sup>MLCT to <sup>3</sup>MLCT and vibrational relaxation of the latter, reflected in the ultrafast decay of ESA for the longest wavelengths. In **C0**, this phase of relaxation leads to a rise of the ESA signal at shorter wavelengths<sup>31</sup> and a similar behaviour is pronounced in **C1**, here in the 500–550 nm range (Fig. 6B, and fits in Figs. S24 and S25), indicating that a sequential scheme is appropriate.

Since the ratio of amplitudes  $A_2/A_3$  corresponds approximately to the *mer*-to-*fac* isomer ratio (see experimental section), one is tempted to attribute the faster  $\tau_2$  to *mer* and  $\tau_3$  to the *fac* isomer. However, as we have shown in our previous publication on the pure *fac*- and a dominant *mer*-forms of **C0**, both decay components are present no matter what isomer content.<sup>31</sup> Hence, interpretation of  $\tau_2$  and  $\tau_3$  and their related amplitude spectra requires a more careful inspection since, in the past, we have given different interpretations for seemingly similar spectral evolutions. Indeed, while in our recent work on tridentate complexes with pyrazine ligands we derived a scenario based on a single triplet excited state, with  $\tau_3$  being its lifetime and  $\tau_2$  representing slow picosecond structural relaxation,<sup>30</sup> the bidentate data for **C0** and its *fac*-only version revealed the population of two triplet states with  $\tau_2$  and  $\tau_3$  being their respective lifetimes.<sup>32</sup> Both scenarios differ, however, in two major points. For instance, for the former mono-state scenario: a) the amplitude of  $\tau_3$  is largely dominant, and b) the amplitude of  $\tau_2$  is close to zero or even positive in

the short wavelength part of the GSB (< 400 nm).<sup>30</sup> On the contrary, for the bidentate complexes under study, we observe  $A_2 > A_3$  and  $\tau_3$  showing up as an additional excited-to-ground state relaxation process since the amplitude of  $A_3$  is significant for all wavelengths below 450 nm. Note also that  $A_3$  follows the profile of the inverse SSA spectrum (Fig. 6).

In summary, as for **C0** and its *fac*-only version, the TAS data are consistent with the existence of two excited triplet states,  $T_1$  and  $T_2$ , populated via relaxation from  $^1MLCT$  and subsequent vibrational cooling during  $\tau_1$ , and decaying in parallel back to the ground state  $S_0$ , according to the following scheme:

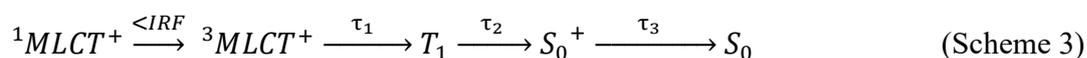


Here, the “+” index denotes vibrationally excited electronic states, inevitably present in the first hundred of femtoseconds after laser excitation. In our previous paper,<sup>31</sup> we assigned the  $A_2/A_3$  ratio at the GSB maximum to the ratio of populations of  $T_2 / T_1$ . However, due to the overlap of ESA and GSB and since the ESA absorption coefficients of  $T_1$  and  $T_2$  are unknown, an accurate determination of the excited state branching ratio is not possible.

The nature of the states  $T_1$  and  $T_2$ , namely MLCT vs. MC character, and the impact of isomerism, namely *fac* vs. *mer*, is rationalized in the following section via molecular modelling of the excited-state topology and relaxation pathways.

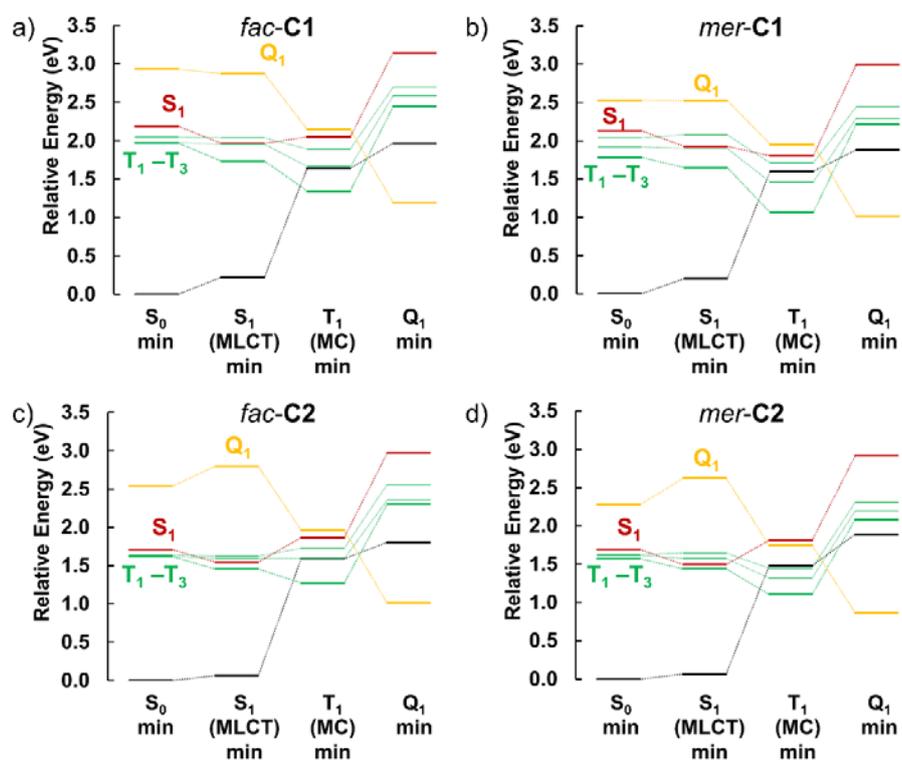
Since the global fitting procedure can be applied to a good approximation, we performed simulations of a purely sequential relaxation scenario. An analysis of the species-associated difference spectra (SADS) reveals, under this assumption and after subtraction of the GSB component, that the absorption spectrum decaying with  $\tau_3$  is actually very close to the GSA, displaying only a small red-shift of 8–10 nm with respect to the latter (Fig. S22). This finding leads

us to suggest a second alternative reaction scenario, with one excited state  $T_1$ , decaying within  $\tau_2 = 2\text{--}3$  ps, and the longer-lived species actually being assigned to an unrelaxed ground state,  $S_0^+$ . The  $\tau_3$  component would then be related to vibrational and dielectric relaxation of the molecules in the ground state. Such a scenario is commonly observed for ultrafast excited state reactions, where a large excited state excess energy ( $> 1\text{eV}$ ) is almost impulsively released into the ground state, leading to a non-equilibrated ground state population, which decays within several picoseconds.<sup>49,50</sup> This is the most plausible interpretation, but not a definitive proof the nature of  $S_0^+$ . Scheme 3 summarises this alternative purely sequential scheme.

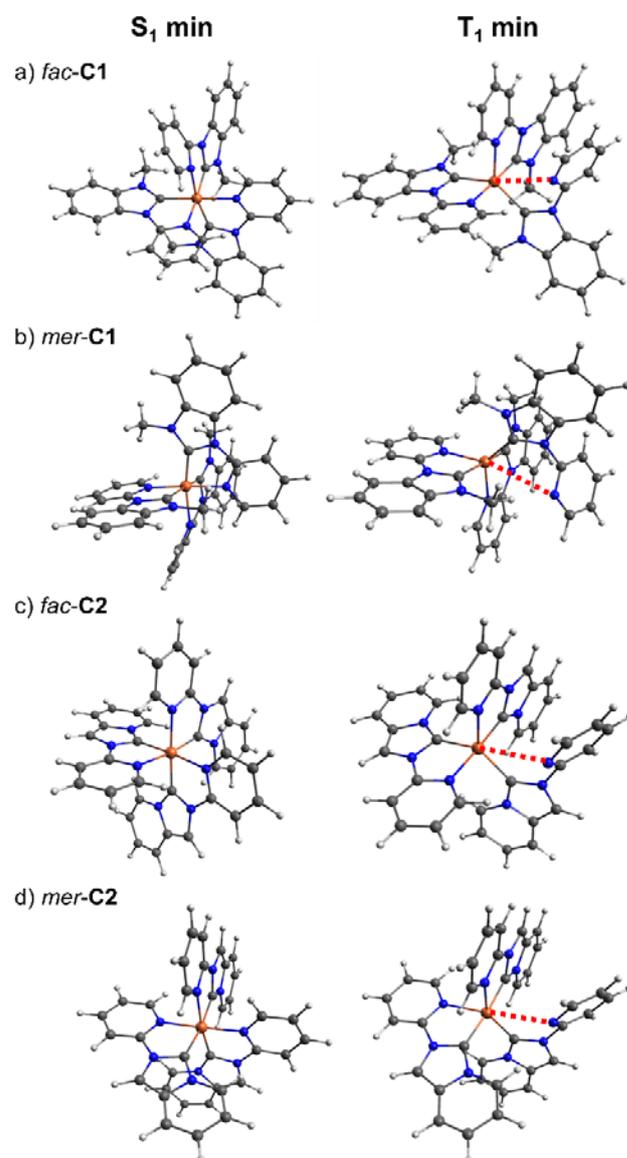


### Molecular modeling of the excited-state decay

Analysis of the excited state reveals that, as for **C0**,<sup>31,32,34</sup> the lowest-lying triplet states of **C1** and **C2** have  ${}^3MLCT$  character at regions close to the Franck-Condon as well as to  $S_1$  min area, whereas further relaxation of the triplet manifold is driven by the elongation of a single Fe-N bond up to distances larger than  $3.2 \text{ \AA}$  leading to spin-crossover regions with the ground state  $S_0$ , with dominant  ${}^3MC$  character for the triplet states. This is illustrated in Figures 7 & 8, and in the SI (Figs. S29–S31 and Table S1).



**Figure 7.** Relative energies of the most relevant electronic states at different equilibrium geometries for a) *fac-C1*, b) *mer-C1*, c) *fac-C2*, and d) *mer-C2*.



**Figure 8.**  $S_1$  min and  $T_1$  min optimized structures of the complexes. Red dashed lines highlight the Fe-N elongated bonds ( $> 3.2 \text{ \AA}$ ) in the triplet state.

These data reveal an adiabatic progression from  $^3\text{MLCT}$  to  $^3\text{MC}$  evidencing  $T_1/S_0$  and  $T_2/S_0$  singlet-triplet crossing regions close to the  $T_1$  min equilibrium geometries that mediate the non-

radiative decay to the ground state. The spin-crossover character of these T<sub>1</sub> min regions and the energy splitting between the T<sub>1</sub> and T<sub>2</sub> states (see Fig. 7) suggest an excited-state decay scenario very similar to the **C0** prototype, for which the triplet potential energy surfaces and the minimum energy paths were determined in much detail.<sup>31,32,34</sup> Therefore, we can safely identify two main deactivation channels, namely a fast T<sub>2</sub> pathway (few picoseconds) and a slower T<sub>1</sub> channel (tens of picoseconds), the longer trapping of the excited state in the T<sub>1</sub> channel is ascribed to the larger spin-crossover regions as discussed in detail elsewhere.<sup>32</sup>

The energy diagrams shown in Figure 7 also allow to compare the profiles for the different isomers of **C1** and **C2**. For *fac*-**C1** (Fig. 7a), the T<sub>1</sub> state at the corresponding T<sub>1</sub> equilibrium geometry lies at a relative energy of 1.34 eV, whereas in *mer*-**C1** (Fig. 7b) this energy is lower, 1.06 eV. The same trend is observed for **C2** (Fig. 7b and c) (1.28 eV for *fac*-**C2** vs. 1.11 eV for *mer*-**C2**). This finding suggests that, assuming very similar potential energy surfaces as those of **C0**,<sup>31,32</sup> *mer* isomers will decay preferentially via T<sub>2</sub>, since the triplet profiles are steeper and the corresponding <sup>3</sup>MC minima lie at lower energies. In addition, the T<sub>1</sub>-T<sub>2</sub> splitting is larger for the *mer* isomers (see Fig 7): this is especially relevant at the FC and S<sub>1</sub> (<sup>1</sup>MLCT) min regions since the T<sub>1</sub>/T<sub>2</sub> branching ratio is determined at these areas, as previously discussed for **C0**.<sup>32</sup>

Regarding the possible role of quintet states, the energy diagrams displayed in Figure 7 evidence a profile similar to that of **C0**, in which a significant participation of the Q<sub>1</sub> state was discarded on the basis of the energy barriers found to access the triplet-quintet crossing points along the potential energy surfaces. Thus, an analogous situation is expected for the current **C1** and **C2** complexes as also confirmed by TAS analysis discussed above.

## CONCLUSION

The present paper reports our recent results on the synthesis of novel complexes of Fe(II) with two types of pyridyl-NHC bidentate ligands. The prototypical imidazolylidene carbene moieties are substituted by benzimidazolylidene (BIm) and imidazo[1,5-*a*]pyridine-3-ylidene (ImP) units, providing dominantly *mer* isomers. The combined experimental and computational study focusses on the effect of these substitutions on the spectroscopic properties and on the excited state lifetimes (ESL) and relaxation pathways, since the aim of these substitutions is to extend the ESL in the context of designing photo-sensitizing and photo-functional complexes with earth-abundant metals.

While the substitutions have significant effects on the energies of the frontier orbitals, as measured by cyclic voltammetry, and on the absorption properties, the excited state lifetimes and photoresponse are only marginally affected. As for the previously published dominantly *fac* and *mer* isomers of  $[\text{Fe}(\text{PyIm})_3]^{2+}$ , the experimental transient absorption data are best described by a parallel population and decay of two excited states T<sub>2</sub> and T<sub>1</sub>, decaying with 2–3 and 15–18 ps lifetimes, respectively, according to the excited and ground state energetics found by TD-DFT calculations. The latter lifetime appears with smaller amplitude in the  $\Delta A$  signal, which may reflect a smaller population in T<sub>1</sub> as compared to T<sub>2</sub>. But it may also be due to the increased overlap and thus compensation of GSB and ESA in the case of T<sub>1</sub> (cf. Fig. S20).

The present study teaches us an important lesson, since the BIm substitution, which almost doubled the ESL for tridentate Fe(II) complexes with four carbenes,<sup>27,28,30</sup> turns out to be ineffective for the bidentate ligands. The initial design idea, to stabilize the MLCT orbitals via extended conjugation,

is validated by electrochemistry, consistent with shifts of the ground state absorption spectra, but it is not operational in extending the ESL. This conclusion holds whatever photoresponse scheme (2 or 3) is actually at work for these homoleptic bidentate complexes. One realizes that the molecular parameters such as enhanced ligand field splitting and stabilization of the MLCT energies, which were considered to be critical for the initial work fueling the development of NHC ligands for Fe(II),<sup>8,21,28-30,33</sup> are not relevant any more in the present complexes with 3 or 4 Fe-carbene bonds. So, which parameter rules the excited state relaxation, and the ESL, i.e. the time it takes to reach the <sup>3</sup>MC-S<sub>0</sub> spin crossover region?

It is clear from previous computational work and ours that, in the presence of a pyridine ligand, the elongation of the Fe-N bond is the key degree of freedom.<sup>31</sup> In other words, an increased stiffness is likely to create an excited state barrier leading to an increased ESL. This was realized already by the latest work of the McCusker group,<sup>51,52</sup> and it explains the difference between our bidentate and tridentate complexes. As in the former, the Fe-N bond is more flexible, there is less energetic penalty for reaching the <sup>3</sup>MC-S<sub>0</sub> crossing, and the effect of ligand substitutions does play a minor role for extending the ESL. Restricting the Fe-N motion via  $\pi$ - $\pi$  stacking with another pyridine, was also operational for increasing the ESL in a heteroleptic Fe(II) complex, designed by the Wärnmark group, with one bipyridine and two bis(triazole-phenyl) ligands,<sup>14,33</sup> and claimed to be important for a recently designed panchromatic Fe complex.<sup>12</sup> It is unclear how far one can get, in terms of ESL, with an increasingly hindered Fe-N elongation, but this can be a future direction of improvement.

From a purely computational and spectroscopic point of view, comparing homoleptic tri- and bidentate complexes with 4<sup>27,28,30</sup> and respectively 3<sup>31,32</sup> carbene bonds, it is remarkable to observe that the former are characterized by an excited state absorption with one dominant band and single

triplet state, while the latter display a parallel population of two excited triplet states. Since this seems to be a general observation for the numerous complexes with **tri- and bidentate ligands** investigated, further work is needed to understand the origin of this effect and how these transitions reflect the dynamically changing  $^3MC$  or  $^3MLCT$  character of the excited states. Note that the heteroleptic complex with tridentate ligands and three Fe-carbene bonds studied by Zimmer et al.,<sup>24</sup> displays yet another scheme, namely a sequential excited state relaxation:  $^3MLCT \rightarrow MC \rightarrow S_0$ . However, for a detailed and accurate study of the excited state relaxation scenario (parallel, sequential, with or without branching), transient absorption spectroscopy is not enough discriminating as we demonstrated with the validity of both Schemes 2 and 3. Hence, additional ultrafast spectroscopic and structural methods,<sup>14,53</sup> in combination with dynamic simulations, are expected to play an increasing role in the near future.

## ASSOCIATED CONTENT

Electronic Supplementary Information (ESI) available: [NMR, spectroscopic and electrochemical data for complexes, and additional computational and photophysical details]. See DOI: 10xxxxx

## Conflicts of interest

There are no conflicts to declare.

## Author Contributions

The manuscript was written through contributions of all authors. All authors have given approval to the final version of the manuscript.

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