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Consuelo Guardiola, Ludovic de Marzi, Yolanda Prezado. Verification of a Monte Carlo dose calculation engine in proton minibeam radiotherapy in a passive scattering beamline for preclinical trials. *Brit.J.Radiol.*, 2020, 93 (1107), pp.20190578. 10.1259/bjr.20190578 . hal-02504664

HAL Id: hal-02504664

<https://hal.science/hal-02504664>

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Consuelo Guardiola, Yolanda Prezado, Ludovic de Marzi, Kathryn Held. Verification of a Monte carlo dose calculation engine in proton minibeam radiotherapy in a passive scattering beamline for preclinical trials. British Journal of Radiology, British Institute of Radiology, In press. hal-02475971

HAL Id: hal-02475971

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Received:
01 July 2019

Revised:
11 December 2019

Accepted:
16 December 2019

<https://doi.org/10.1259/bjr.20190578>

Cite this article as:

Guardiola C, Prezado Y, De Marzi L, Held K. Verification of a Monte Carlo dose calculation engine in proton minibeam radiotherapy in a passive scattering beamline for preclinical trials. *Br J Radiol* 2019; **92**: 20190578.

PROTON THERAPY SPECIAL FEATURE: FULL PAPER

Verification of a Monte Carlo dose calculation engine in proton minibeam radiotherapy in a passive scattering beamline for preclinical trials

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Objectives: Proton minibeam radiation therapy (pMBRT) is a novel therapeutic strategy that combines the benefits of proton therapy with the remarkable normal tissue preservation observed with the use of submillimetric spatially fractionated beams. This promising technique has been implemented at the Institut Curie-Proton therapy centre (ICPO) using a first prototype of a multi-slit collimator. The purpose of this work was to develop a Monte Carlo-based dose calculation engine to reliably guide preclinical studies at ICPO.

Methods: The whole “Y1”-passive beamline at the ICPO, including pMBRT implementation, was modelled using the Monte Carlo GATE v. 7.0 code. A clinically relevant proton energy (100 MeV) was used as starting point. Minibeam generation by means of the brass collimator used in the first experiments was modelled. A virtual source was modelled at the exit of the beamline nozzle and outcomes were compared with dosimetric measurements performed with EBT3 gafchromic films and a diamond detector in water. Dose distributions were

recorded in a water phantom and in rat CT images (7-week-old male Fischer rats).

Results: The dose calculation engine was benchmarked against experimental data and was then used to assess dose distributions in CT images of a rat, resulting from different irradiation configurations used in several experiments. It reduced computational time by an order of magnitude. This allows us to speed up simulations for *in vivo* trials, where we obtained peak-to-valley dose ratios of 1.20 ± 0.05 and 6.1 ± 0.2 for proton minibeam irradiations targeting the tumour and crossing the rat head. Tumour eradication was observed in the 67 and 22% of the animals treated respectively.

Conclusion: A Monte Carlo dose calculation engine for pMBRT implementation with mechanical collimation has been developed. This tool can be used to guide and interpret the results of *in vivo* trials.

Advances in knowledge: This is the first Monte Carlo dose engine for pMBRT that is being used to guide preclinical trials in a clinical proton therapy centre.

INTRODUCTION

Proton minibeam radiation therapy (pMBRT) is a novel therapeutic strategy¹ in which irradiation is performed by using arrays of narrow (≤ 1 mm) proton beams. pMBRT has already been implemented in our clinical centre (Institut Curie-Proton Therapy Center in Orsay, ICPO).^{2,3} The distinct dose delivery method used in pMBRT results in a net reduction in skin damage^{4,5} and neurotoxicity compared to conventional proton therapy.⁴ Moreover, superior tumour control compared to that obtained with standard proton therapy has been demonstrated in small animal experiments.⁶ Significant tumour control was also observed even with highly heterogeneous dose distributions,⁷ suggesting the participation of distinct biological mechanisms compared to standard RT. Comprehensive

biological experiments are needed to elucidate the underlying mechanisms in pMBRT. For this purpose, a suitable set of dosimetry tools in order to accurately guide and interpret preclinical trials at ICPO needs to be developed. The very small beam sizes used in pMBRT makes this a challenging and error-prone task, requiring thorough evaluation (by Monte Carlo simulations and experimental measurements).

In particular, the development of a validated Monte Carlo-based dose calculation engine for small animal experiments was necessary to obtain precise and detailed information about the dose distributions actually delivered in the rodent anatomy. As proof-of-concept, we assessed the dose distributions in water and CT images of a rat's head. To

Figure 1. Schematic view of the Y1-beamline with all beam modifiers.¹⁶

the best of our knowledge, this is the first complete dosimetric study, including a dose calculation engine, for preclinical trials of pMBRT.

METHODS

Monte Carlo simulations

Monte Carlo simulations based on the GATE v. 7.0 code⁸ were used to model the Y1-room beamline of ICPO, the multislit collimator, and to assess the dose distributions in the target. GATE v. 7.0 is based on the Geant4 toolkit (Geant4.9.6.p03).^{9–11} The physics lists and parameters recommended by the GATE collaboration for proton therapy applications were used.^{12–14} In particular, the Binary Cascade (BIC) model was used for the hadronic interactions adding the standard option 3 to include the low-energy electromagnetic processes (_EMY). Both options were included by using the QGSP_BIC_HP_EMY GATE builder. Range cuts of 1 and 0.025 mm were considered for all particles for the world and phantoms, respectively. Ionisation potentials for water and air were set to 75 and 85.7 eV, respectively.^{14,15} It is worth noting that these GATE recommendations include the use of 75 eV for ionisation potential of the water, which was the recommended value in earlier ICRU reports. However, the newer ICRU 73 recommends an I value of 78 ± 2 eV. For the sake of clarification, we run simulations with the virtual source (Virtual source modelling) with both 75 and 78 eV water ionisation potentials and no discernible differences in the percentage depth dose (PDD) were found in both cases for a 100 MeV proton energy. Consequently, we performed simulations following the GATE collaboration recommendation afterwards for this particular scenario.

Simulation of one passive scattering beamline at ICPO

In a first phase, the complete beamline was modelled and benchmarked against experimental data. In a second phase, for sake of computational time, a virtual source model was defined by means of an iterative process.

The ICPO facility is equipped with an isochronous cyclotron, a C230 from Ion Beam Applications company, that delivers 230 MeV protons at the accelerator exit. The treatment room chosen for our first studies was the Y1-room, which has a horizontal fixed beamline and uses a double scattering delivery system. Figure 1 shows a schematic view of the simulated Y1-room components. Further details can be found in Patriarca et al.¹⁷

Since the main targets of pMBRT are neurological lesions, *i.e.* brain tumours, and in order to consider the worst treatment scenario (centre of the human brain), we used 100 MeV Bragg peaks as starting-point, which corresponds to a proton range of 7.7 cm in water equivalent. As the first *in vivo* experiments of pMBRT focused on rodent brains with tumour diameters less than 5 mm, we used only simple Bragg peaks, which cover 5 and 3 mm with 95 and 98% of the prescription (Bragg peak) dose, respectively.

A virtual source was modelled to minimise computational time. For this purpose, beam characteristics at the nozzle exit, *i.e.* beam energy spread, spot size, and angular distribution, were stored in *Phase Space Files*, as a function of the beam energy required at

Figure 2. pMBRT brass multislit collimators comprising 3.5 cm long slits (left) and 2×2 cm² broad brass collimator (right). To achieve an irradiation area behind the collimator exit, an additional squared-shape brass collimator (1.6×1.6 cm²) was added to the beamline nozzle (right). Image from Cécile.¹⁶ pMBRT, proton minibeam radiation therapy.

Figure 3. Experimental PVDR values obtained with the PTW diamond detector and EBT3 gafchromic films when a 0.7 mm wide brass collimator with ctc of 3.5 mm was used to generate minibeam. PVDR, peak-to-valley dose ratio.

the isocentre. They were used along with the data coming from previous modelling of the same beamline at the Orsay proton therapy centre¹⁸ as starting point. In particular, a 100 MeV proton energy beam, 1.5 MeV energy spread and 15 mm spot size was used. The range of angular spread covered ranged from 0.1 to 5 mrad in our iterative process. Those included the typical range of values from 2 to 5 mrad for passive systems.^{19,20} An iterative comparison of simulated and measured dose distributions was performed for fine tuning of the parameters of the virtual source.

A gain of one order of magnitude in computational efficiency was achieved with this approach.

Multislit collimators

Multislit brass collimators described in the work by Peucelle et al^{2,16} were used for minibeam generation. These collimators are 5 cm thick (like the standard used for proton therapy at that

energy range), with widths of 0.4 and 0.7 mm and centre-to-centre (ctc) distances of 3.2 and 3.5 mm, respectively. This optimised design²¹ was manufactured at ICPO by electrical discharge machining (Figure 2).

Dose calculations

The phantom multislit collimator distance was always 7 cm. Dose distributions were recorded using voxel sizes of 2 mm × 100 μm × 1 mm in the vertical, lateral and beam directions, using the so-called GATE-DoseActor. PDD curve, lateral dose profiles, and peak-to-valley dose ratios (PVDRs) were evaluated. The PVDR assessment was performed by averaging over the tally bin size along the beam transverse axis, i.e. 100 μm, using the doses in the central peak and its adjacent valley.

The number of simulated primary particles was 10¹², resulting in an uncertainty below 1% in regions where more than 50% of the maximum dose was deposited. The total dose was the sum of all units used, calculated by a home-made C++ code. Dose uncertainty in a single voxel was computed according to the method described by Chetty et al.²²

The following phantoms were used: (a) water cylindrical phantoms with a diameter of 20 cm and a height of 20 cm, simulating a human head²³; and (b) voxelised CT images of a rat head (7-week-old male Fischer rat), corresponding to the age of the animals usually irradiated in our *in vivo* experiments. We used CT images obtained at Institut Curie from a representative Fisher rat in the same conditions as those of the irradiated set. DICOM images (CT raw format) must be converted to *MetalImage format* (.mhd). To save computational time, the rodent CT image was re-sampled with resolutions of 0.25, 2, and 1 mm in the transverse and longitudinal axes, respectively (with reference to the beamline). Image greyscales were converted into material definitions based on the correspondence between Hounsfield units, i.e. voxel values, and materials. The stoichiometric calibration procedure described in Lansonneur et al²⁴ was used to determine the tissue substitute calibration curves of the CT scanner using a set of materials with known elemental composition and physical density close to those of tissue samples. The output file

Figure 4. Left: PDD curve simulated and measured with EBT3 gafchromic films in solid water. Uncertainties of 8% were considered for measured data. Right: simulated PVDR values versus experimental data recorded with PTW diamond detector when a 0.7 mm wide brass collimator with a ctc of 3.5 mm was used to generate minibeam. PDD, percentage depth dose; PVDR, peak-to-valley dose ratio.

Figure 5. Lateral profile at 24 mm depth of MC simulated and experimental data recorded with PTW diamond detector. MC, Monte Carlo.

used to collect absorbed doses in the phantom was a binary file containing the absorbed dose in cGy with the same dimensions as the phantom.

Various geometries and energies (50, 70 and 100 MeV) were considered in the rat brain dose evaluations.

Experimental dosimetry

Experimental dose distributions were measured with EBT3 gafchromic films²⁵ and a PTW microDiamond 60019 detector.²⁶ Lateral and depth dose profiles, as well as PVDR values were assessed. Data were then used to validate the dose-calculation engine. Diamond detectors are nearly water equivalent for all beam energies and linear energy transfer-independent,^{27–29} which is important in proton therapy as linear energy transfer is expected to vary with beam depth. To take this into account

in the case of the EBT3 gafchromic films, we applied the corresponding measured quenching correction factors.²

Measurements were performed in a water tank (BluePhantom IBA) and in RW3 solid water phantoms. The microDiamond detector was placed inside a BluePhantom scanning system and oriented “in-edge”, *i.e.* perpendicular to the beam to obtain a resolution of 1 μm in the spatial fractionation direction of the minibeam. The signal was integrated over a depth of 2.2 mm, corresponding to its diameter, and the scanning step was 100 μm . The poor resolution in the longitudinal direction (2.2 mm) made the microdiamond detector unsuitable to measure depth dose curves in pMBRT. Uncertainties were derived from electrometer reproducibility (1%), scanning system accuracy (± 0.1 mm), and mean standard deviation in the average dose in the peak and valley regions (2%). It should be noted that, even after a low pre-irradiation dose, other authors have also reported a variation of the microDiamond signal stability by more than 1%, which could be due to non-reproducibility between devices at the time of the study.²⁹ A conservative overall uncertainty value of 3% was considered.

PDD curves were only characterised by EBT3 gafchromic films, as they provide the high spatial resolution required for the submillimetric minibeam. A flatbed scanner (Epson Perfection V750-M Pro Scanner) was used for film readout according to the methodology described in Devic et al³⁰ with a resolution of 1200 dpi. Film handling was performed according to the recommendations of Task Group 55 of the American Association of Physics in Medicine.³¹ Uncertainties in film dose measurements were evaluated according to the method described in Sorriaux et al.³² A conservative uncertainty value of 8% was considered. Scanner readings were calculated by using a home-made C++ code using the calibration curves obtained with the optical densities of the irradiated EBT3 gafchromic films.

Comparison of the measurements obtained with the two detectors confirmed the reliability of experimental dosimetry. Finally,

Figure 6. Left: Comparison of PVDR values measured with EBT gafchromic films in a solid water phantom and diamond detector with those simulated by using a fitted GPS virtual source. Right: PDD curves in the central peak measured with EBT gafchromic films and those simulated with the virtual source. GPS, *General Particle Source*; PDD, percentage depth dose; PVDR, peak-to-valley dose ratio.

Figure 7. Photograph of a 7-week-old male Fischer rat irradiated at the ICPO Y1-room beamline. The image shows the PMMA in front of the animal's head to target the tumour position when irradiated with 100 MeV protons. ICPO, Institut Curie-Proton therapy centre; PMMA, polymethylmethacrylate.

these measurements were used to benchmark our dose calculation engine using the dose calculation acceptability criteria compiled in Technical Report Series 430 (TRS 430).³³ In particular, the confidence limit (Δ) defined by Venseelar et al³⁴ was used. See Eq 1.

$$\Delta = |\text{average deviation}| + 1.5 \text{ SD} \quad [1]$$

SD: standard deviation.

RESULTS

In this section, we report the dosimetric data acquired, particularly validation of the dose calculation engine developed, with several examples of small animal treatment plans.

Experimental dosimetry

Figure 3 shows the peak-to-valley dose ratio measured with the two different detectors, which are in good agreement within

the uncertainty bars. As expected, multiple Coulomb scattering causes a continuous decrease of the PVDR with increasing depth. Note that the depth dose in water cannot be determined at depths of less than 20 mm due to the presence of the water tank entrance window.

Benchmarking of the dose calculation engine

Our Monte Carlo dose calculation engine was validated by comparing experimental and calculated dose distributions. Validation of the whole pMBRT modelling was performed by comparing the results obtained by simulating the Y1-room beamline plus one of the above brass collimators (0.7 mm wide with ctc of 3.5 mm) over a water phantom with those obtained with EBT3 gafchromic films and diamond detector with 100 MeV proton irradiations.

Figure 4 (left) shows the PDD curve within the solid water phantom in the central peak of the minibeam array. There is a globally satisfactory agreement for preclinical trials between simulated and measured doses in water (within the uncertainty bars). The discrepancy observed at depths of 55–65 mm could be due to slight inaccuracies in geometrical modelling of the beam-line, which have a greater effect at such small beam sizes. Another cause could be an imperfect multiple coulomb scattering model in the Monte Carlo code.

Similarly, Figure 4 (right) shows the simulated PVDR values versus the experimental PVDR values from the data collected with the PTW diamond detector.

Figure 5 shows a representative lateral profile for a water depth at 24 mm.

Δ (Eq. 1) complies with the acceptability criteria recommended in TRS 430.³³ Deviations less than 2% in the depth dose curves were obtained. A maximum Δ value of 5% was found in the high dose/small gradient regions of the profiles of the central peaks. The MC simulated processes are therefore in good agreement with the experimental data, *i.e.* the Geant4 MC code with the

Figure 8. PDD (left) and PVDR (right) in the rat brain (50 mm thickness of PMMA, no air gap). Note that the rat brain is about 14 mm long and the average diameter of the tumour along the beam direction 4 days after implantation (when animals are irradiated) is roughly 4 mm. Since the tumour implantation is approximately on the half of the second brain hemisphere in which the proton beam impinges, *i.e.* at a depth of approximately 10.5 mm of the rat brain, the tumour volume extends approximately from 8.5 to 12.5 mm. Note that the preferential direction of tumour growth is unpredictable. PDD, percentage depth dose; PMMA, polymethylmethacrylate; PVDR, peak-to-valley dose ratio.

Figure 9. Coronal 2D dose distributions on computed tomography images of 7-week-old Fischer 344 rats. From left to right, they correspond to broad beam (conventional) irradiation and proton minibeam irradiations targeting the tumour and crossing the rat head, respectively.^{37,38} Black and red lines indicate the approximate position of the centre of the tumour.

physics list recommended by GATE collaboration for proton therapy appears to accurately simulate the multiple coulomb scattering process for the particular conditions of small fields.

Independent *Phase Space Files* were scored at the nozzle exit (Figure 1) by means of the Phase Space GATE-actor⁸ that tallies information about particles entering the volume to which the actor is attached. They were scored as ROOT data analysis format.³⁵ On the basis of these *beam optics* parameters, the beam can be characterised at that depth (nozzle exit). Fitting was performed by converting raw data to ROOT two-dimensional histograms. The energy spread was characterised by a Gaussian standard deviation. These beam features were used to model a virtual source based on the predefined *General Particle Source*. As explained in "Simulation of one passive scattering beamline at ICPO".

Iterative comparison of the simulated and measured dose distributions was performed for fine tuning of the parameters of these sources. Figure 6 (left) shows the final best fitted *General Particle Source*-virtual source with the experimental data with a divergence of 4.5 mrad in agreement with previous values.¹⁸ The source-to-phantom distance was considered to 100 cm, as it was demonstrated in Cécile¹⁶ that the peak-to-entrance ratio becomes constant. A Gaussian spread was created by using a two-dimensional beam travelling along the local negative z axis.

Figure 10. Coronal 2D dose distributions on computed tomography images of 7-week-old Fischer 344 rats for ctc of 3 mm (left) and 2.8 mm (right). A smaller ctc was associated with greater homogeneity of the dose distribution over the tumour. Black lines indicate the approximate position of the centre of the tumour. 2D, two-dimensional.

Similarly, Figure 6 (right) shows the PDD curve in the central peak of the minibeam array. Both figures show good agreement between experimental and simulated results. A minor discrepancy at a depth of 0 cm was found, which could be due to minor inaccuracies in phantom alignment or phantom positioning or to the uncertainties in the energy spread used in the virtual source.

On the basis of these results and to speed-up simulations, forthcoming simulations were performed with that virtual source.

Note that the influence of the uncertainty of the ionisation potential on the proton range is out of the scope of this study. Nevertheless, although we used the GATE recommendation value for ionisation potential of the water (*i.e.* 75 eV), we corroborated that there is not significant differences in the PDD for a 100 MeV proton energy by using the current ICRU73 recommended one (*i.e.* 78 eV) with the virtual source described in "Virtual source modelling". Further studies regarding the influence of various sets of stopping power data with I values for water between 67 and 80 eV can be found in Luhr et al.³⁶ They found deviations of the order of 1% in the plateau region or range in water for that I value range.

Dosimetry for *in vivo* experiments: rat brain pMBRT irradiations

Two types of experiments were performed: animals were irradiated with the Bragg peak located at the tumour position and animals were irradiated in the plateau region. As the primary objective was to optimally reproduce clinically relevant conditions, a clinically relevant energy (*i.e.* 100 MeV) was used. Some attenuation material is required to place the Bragg peak in the tumour. Polymethylmethacrylate (PMMA) was used (Figure 7). The required thickness was calculated by Monte Carlo simulations taking into account irradiations of 7 week-old Fischer rats. We used a 0.4 mm wide brass collimator with a ctc of 3.2 mm, as it delivers a homogeneous dose distribution in water.²¹

Figure 8 shows the optimal value for PDD and PVDR inside the rat brain with no air gap between the rat head and PMMA slabs. Note that the rat brain is approximately 14 mm along on a midline section. Since the site of tumour implantation is approximately on the half of the second brain hemisphere in which the proton beam impinges, *i.e.* at a depth of approximately 10.5 mm

of the rat brain, the optimal PMMA thickness is between 50 and 48 mm. As the FWHM of the Bragg curve at that depth is approximately 6 mm, the tumour sizes considered (see materials and methods section) would be totally covered.

Figure 9 shows the dose profiles at the middle of the brain and the central part of the tumour for the two cases considered above. The PVDR ratios were 1.20 ± 0.05 and 6.1 ± 0.2 , respectively.^{37,38}

In the first case, with a ctc of 3.2 mm ctc, a homogeneous dose distribution inside the rat brain is not completely obtained due to the fact that the morphology and composition of the animal's head both differ from the water archetype. A homogeneous dose distribution in the tumour position can be achieved by reducing the centre-to-centre distance of the minibeam. For example, as shown in Figure 10, dose distributions become more homogeneous in the tumour position when a smaller ctc is considered, ranging from 3 to 2.8 mm.

These results illustrate the importance of reliable calculations to interpret results. For example, as those performed for understanding dose distributions inside the rat head in preclinical trials.^{37,38} In particular, the dose calculation engine for pMBRT implementation was used to guide the first two critical preclinical studies. We first compared animal survival when animals were irradiated with pMBRT and standard proton therapy in an *in vivo* test with RG2 glioma-bearing rats. In this study, we evaluated the efficacy of pMBRT for the treatment of high-grade gliomas and, although some spatial fractionation was still observed at the Bragg position, tumour eradication was observed in 67% of the animals treated by pMBRT.³⁷ We then reported the

first evaluation of tumour control efficacy in glioma-bearing rats with highly spatially modulated proton beams, *i.e.* large areas of the tumour received low doses (valley doses). In this case, we simulated the lateral irradiation of rat brain in the plateau region with a dose rate of 2 Gy/min at a depth of 1 cm. We showed that the heterogeneous dose distribution achieved significant tumour control in 22% of cases.³⁸

Simulations with 50 and 70 MeV protons, with ranges in water of approximately 2.3 and 4.1 cm, respectively, were also performed to compare these results with those obtained with usual preclinical energies. Figure 11 shows the dose distributions in the same intermediate axial plane when animals were irradiated with 50 and 70 MeV proton minibeam.

PVDRs in the tumour position were 10.4 ± 0.6 and 7.3 ± 0.4 for the 50 and 70 MeV proton irradiations, respectively. PVDRs were higher for low energies due to less homogeneous dose distribution, resulting in an increase of $\Delta = 4.3 \pm 0.8$ and 1.2 ± 0.6 over the PVDR compared to the PVDR with 100 MeV pMBRT.

CONCLUSIONS

This paper reports on the development and experimental validation of a dose calculation engine for pMBRT. The small beam sizes used in this study makes experimental and theoretical dosimetry particularly challenging. An overall good agreement between experimental and calculations was obtained, although further evaluations of the performance of the multiple coulomb scattering models implemented in the codes is needed and will be the subject of future works. This engine can be used to correlate the results on ongoing preclinical studies at the Orsay proton therapy centre with the dose distributions. In addition, it will be useful to design forthcoming experiments. To the best of our knowledge, this is the first dose calculation engine developed for pMBRT.

ACKNOWLEDGMENT

This project has received funding from the European Union's Horizon 2020 research and innovation programme under the Marie Skłodowska-Curie grant agreement No 74,5109. Calculation time was granted at Castilla y León Supercomputing Center (Calendula) from the national RES access programme (Spanish Supercomputing Network Area) and the MareNostrum Barcelona Supercomputing Center from the Partnership for Advanced Computing in Europe (PRACE Project Access Call 14th and 16th, proposals number 2016153507 and 2017174196, respectively).

Figure 11. Dose distribution on CT scans of the rat head irradiated with 50 MeV (left) and 70 MeV (right) proton minibeam. Spatial fractionation was observed in both cases.

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