



Recovering underlying graph for networks of 1D waveguides by reflectometry and transferometry

Geoffrey Beck, Maxime Bonnaud, Jaume Benoit

► To cite this version:

Geoffrey Beck, Maxime Bonnaud, Jaume Benoit. Recovering underlying graph for networks of 1D waveguides by reflectometry and transferometry. WAVES 2019 - 14th International Conference on Mathematical and Numerical Aspects of Wave Propagation, Aug 2019, Vienna, Austria. hal-02414861

HAL Id: hal-02414861

<https://hal.science/hal-02414861>

Submitted on 16 Dec 2019

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Recovering underlying graph for networks of 1D waveguides by reflectometry and transferometry

Geoffrey Beck^{1,*}, Maxime Bonnaud², Jaume Benoit²

¹POEMS (CNRS-INRIA-ENSTA Paristech) Palaiseau, France

²CEA, LIST, 91191 Gif-sur-Yvette CEDEX, France

*Email: geoffrey.beck.poems@gmail.com

Abstract

We present a method for blind recovery of network made out of a tree of 1D homogeneous waveguides with the same physical characteristics using reflectogram and transferogram(s).

Keywords: Inverse problem, Topology recovering, Quantum graph, Reflectometry, Wire analysis

1 Introduction

We consider an unknown quantum graph \mathcal{G} (see [2]) equipped with a wave operator along its branches and some transmission conditions on its nodes connecting together the quantities evaluated on the branches. Our graph is a rooted tree graph, where all branches are oriented from a root-point Inp to end-points Out_k ($k = 1 \dots K$). We will consider a maximum of two consecutive nodes between Inp and any Out_k and at least a node in \mathcal{G} .

We will now explain how a wave V propagates along the graph \mathcal{G} :

- On each branch of \mathcal{G} the wave satisfy an homogeneous wave's equation

$$\partial_{tt}^2 V - c^2 \partial_{xx}^2 V = 0,$$

where t denotes the time and x the abscissa along the considered branch. The celerity c of the waves is supposed to be a known constant.

- Following Kirchhoff's rules, V is continuous on \mathcal{G} and at each node J

$$\partial_x V|_{\mathbf{e}_{j_0}}(J) = \sum_{j_k=j_1}^{j_K} \partial_x V|_{\mathbf{e}_{j_k}}(J),$$

where \mathbf{e}_{j_k} are the branches connected to J , with \mathbf{e}_{j_0} the branch closest to Inp .

- On Inp , we have an impedance boundary condition

$$\partial_t V(\text{Inp}, t) - c \frac{Z_u}{Z_c} \partial_x V(\text{Inp}, t) = (\partial_t u)(t)$$

where the constant Z_u and $u \in H_{\text{loc}}^1(\mathbb{R})$ are known. The unknown characteristic impedance Z_c is supposed to be constant.

- At each Out_k we have an impedance condition

$$\partial_t V + c \frac{Z_k}{Z_c} \partial_x V = 0,$$

where Z_k is an unknown constant.

2 Graph recovery problem

Reflectometry and transferometry methods can be applied to any practical electrical or acoustic network. The reflectogram is the following Steklov operator :

$$u(t) \mapsto \mathfrak{R}(t) := V(\text{Inp}, t),$$

whereas transferograms are operators for $k = 1 \dots K$:

$$u(t) \mapsto \mathfrak{T}_k(t) := V(\text{Out}_k, t).$$

We suppose that we can control u . With a known celerity c and the input load Z_u , the reflectogram and optionally some transferograms, we want to recover \mathcal{G} that is to say to determine

- the number of nodes and end-points, and their ordering (topology),
- the lenght ℓ_j of all branches,
- the end-points load Z_k ,
- the characteristic impedance Z_c .

3 Injectivity

There can exists several quantum graph with the same reflectogram, so we will make two hypothesis. Firstly, no scatterer (node or endpoints) have the same distance from Inp , to ensure they can be dissociated. Secondly, no Z_k is equal to Z_c to ensure waves are reflected on the end-point. We will suppose that $Z_k > Z_c$ is always satisfied.

4 Scattering

We can choose the excitation signal u such it is a peak function (sole local maximum). It propagates at celerity c along a branch until it meets a scatterer. \mathfrak{R} is the sum of attenuated u -shaped peaks, i-e of the form $\sum_p S_p u(t - t_p)$ where each (t_p, S_p) - called echo - corresponds to the duration and the amplitude attenuation of a propagation of u in \mathcal{G} through transmissions T and reflections Γ looping on Inp. T_k is similarly generated, with propagations from Inp to Out _{k} . An echo amplitude S_p gives the nature of its contained scatterers (order for a node, load for an end-point), its abscissa the path length between the observation point (Inp for \mathfrak{R} and Out _{k} for \mathfrak{T}_k) and the scatterer.

The algorithm presented in [1] identifies echoes in a complex reflectogram and associates them with unknown scatterers in \mathcal{G} , giving their nature and location. It runs iteratively, dispelling ambiguities from peaks overlapping and accumulated reflections.

But this method requires knowledge of Z_c and supposes that $Z_c = Z_u$ (no reflections at Inp). It can be enhanced by the use of transferograms.

5 Algorithm

5.1 Recovering Z_c

We simply recover Z_c from the reflectogram at origin where we see an echo (called mismatch echo) of amplitude $T_u = (1 - \Gamma_u)$ with $\Gamma_u = (Z_u - Z_c)(Z_u + Z_c)$. Of course if $Z_c = Z_u$ then the mismatch peak is null.

5.2 Recovering the first node

The first echo observe in \mathfrak{R} after the mismatch echo have for abscissa $2\ell_0/c$ and for amplitude $T_u(1 + \Gamma_u)\Gamma_0$ with $\Gamma_0 = (2/m_0 - 1)$ where m_0 is the order of the first node J_0 . We thus recover ℓ_0 and m_0 .

5.3 Using the transferograms

If the amplitude of the first echo (t_1, S_1) of \mathfrak{T}_k is above $4T_u(\Gamma_0 + 1)/3$, then Out _{k} is directly connected to J₀. Thus we have $S_1 = T_u(2/m_0)(1 + \Gamma_k)$ with $\Gamma_k = (Z_k - Z_c)(Z_k + Z_c)$, so we recover Z_k . The length of the J₀ to Out _{k} branch is $(ct_1 - \ell_0)$. If S_1 is under $T_u(\Gamma_0 + 1)$, a node exists between J₀ and Out _{k} . This recovered topology

can remove branch location ambiguities for an unknown scatterer with the reflectogram.

5.4 Using the reflectogram

We use the algorithm developed in [1] to continue the analysis of \mathfrak{R} . We changed the procedure to use informations from the transferograms, and adapt to $Z_u \neq Z_c$. Indeed, we need to apply a $T_u(1 - \Gamma_u)$ factor to all \mathfrak{R} echoes amplitude and to consider reflexions on Inp when discriminating between echoes.

This method achieves an error-free topology reconstruction if some technical hypothesis on u are fullfilled. ℓ_j are retrieved with an accuracy decreasing when farther from Inp (relative error under 5% from 350 simulations), as are Z_c (under 0.1%) and Z_k (under 10%). Better determination is possible by optimizing the all lengths $\tilde{\ell}$ and all loads $\tilde{\mathbf{Z}}$ vectors such that they minimize the functional

$$J(\boldsymbol{\ell}, \mathbf{Z}) = \int_0^{8\ell_{max}/c} \frac{|\mathfrak{R}(t) - \mathfrak{R}_{\boldsymbol{\ell}, \mathbf{Z}}(t)|^2}{3\ell_{max}/c + t^2} dt$$

where ℓ_{max} is the maximum ℓ_j from previous steps and $\mathfrak{R}_{\boldsymbol{\ell}, \mathbf{Z}}$ the simulated reflectogram using the recovered topology with $\boldsymbol{\ell}$ and \mathbf{Z} . We look for the minimum of J with Newton's algorithm, initializing $\tilde{\boldsymbol{\ell}}$ and $\tilde{\mathbf{Z}}$ with the previously recovered values.

6 Applications

This algorithm can be used for recovering unknown electrical networks, with one reflectometry device and optionnal transferometry transcievers on end-points. Removing the condition on Z_u makes the algorithm more resilient to real-life implementation limitations.

References

- [1] G. Beck, *Reconstruction of an unknown electrical network from their reflectogram by an iterative algorithm based on local identification of peaks and inverse scattering theory*. International Instrumentation and Measurement Technology Conference IEEE I2MTC, 2018 (to appear)
- [2] P. Kuchment, *Quantum graphs I. Some basic structures* Waves in Random Media 14, S107-S128 (2004)