

Plasmonic photocatalysis applied to solar fuels

Steven Bardey, Audrey Bonduelle-Skrzypczak, Antoine Fecant, Zhenpeng Cui, Christophe Colbeau-Justin, Valérie Caps, Valérie Keller

▶ To cite this version:

Steven Bardey, Audrey Bonduelle-Skrzypczak, Antoine Fecant, Zhenpeng Cui, Christophe Colbeau-Justin, et al.. Plasmonic photocatalysis applied to solar fuels. Faraday Discussions, 2019, 214, pp.417-439. 10.1039/C8FD00144H . hal-02391217

HAL Id: hal-02391217 https://hal.science/hal-02391217

Submitted on 6 Dec 2019

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Plasmonic photocatalysis applied to solar fuels

Steven Bardey, ¹^{ab} Audrey Bonduelle-Skrzypczak, ^b Antoine Fecant, ^{bb} Zhenpeng Cui, ^c Christophe Colbeau-Justin, ^c Valerie Caps ^{b*a} and Valerie Keller^a

Received 2nd October 2018, Accepted 16th November 2018 DOI: 10.1039/c8fd00144h

The induction of chemical processes by plasmonic systems is a rapidly growing field with potentially many strategic applications. One of them is the transformation of solar energy into chemical fuel by the association of plasmonic metal nanoparticles (M NPs) and a semi-conductor (SC). When the localized surface plasmon resonance (LSPR) and the SC absorption do not match, one limitation of these systems is the efficiency of hot electron transfer from M NPs to SC through the Schottky barrier formed at the M NP/SC interfaces. Here we show that high surface area 1 wt% Au/TiO_{2-UV100}, prepared by adsorption of a NaBH₄-protected 3 nm gold sol, readily catalyzes the photoreduction of carbon dioxide with water into methane under both solar and visible-only irradiation with a CH₄ vs. H₂ selectivity of 63%. Tuning Au NP size and titania surface area, in particular via thermal treatments, highlights the key role of the metal dispersion and of the accessible Au-TiO₂ perimeter interface on the direct SC-based solar process. The impact of Au NP density in turn provides evidence for the dual role of gold as co-catalyst and recombination sites for charge carriers. It is shown that the plasmon-induced process contributes up to 20% of the solar activity. The plasmon-based contribution is enhanced by a large Au NP size and a high degree of crystallinity of the SC support. By minimizing surface hydroxylation while retaining a relatively high surface area of 120 m² g⁻¹, pre-calcining TiO_{2-UV100} at 450 $^{\circ}$ C leads to an optimum monometallic system in terms of activity and selectivity under both solar and visible irradiation. A state-of-the-art methane selectivity of 100% is achieved in the hot electron process.

Introduction

The direct conversion of solar energy into valuable chemical fuels is a great challenge for clean energy production. In this context, the photo-reduction of CO_2

^bIFP Energies Nouvelles (IFPEN), Rond-point de l'échangeur de Solaize, 69360 Solaize, France

^aInstitut de Chimie et des Procédés pour l'Energie, l'Environnement et la Santé (ICPEES), UMR7515, 25 rue Becquerel, 67087 Strasbourg, France. E-mail: caps@unistra.fr

^cLaboratoire de Chimie Physique, CNRS UMR 8000, Univ. Paris-Sud, Université Paris-Saclay, 91405 Orsay, France

Faraday Discussions

into methane and methanol has received increasing attention over the last few years.¹⁻³This way of producing synthetic fuels would indeed represent a step forward towards sustainability, due to the nature of the carbon source, a major byproduct of energy consumption, and its increasing availability. However, the conversion rate of photonic energy into chemical energy by semi-conductor-based photocatalysts is still too low, as current energy carriers suffer from a short lifetime (fast recombination) and use only a fraction of the solar spectrum. Hence, the development of photocatalytic materials with improved and broadened light absorption is needed in order to efficiently use sunlight as a primary energy source. Amongst the possible materials, multi-component systems containing metal nanoparticles (M NPs) interacting with a semi-conductor (SC) appear promising.⁴⁵ In addition to their beneftcial role as co-catalyst and electron traps,⁶⁷ MNPs with plasmonic properties⁸ may indeed improve the photocatalytic performances of the semi-conductor by plasmon-derived physical phenomena, including hot electron generation⁸ and utilization.⁹⁻¹³ Hot electron injection into the SC conduction band is however dependent on the Schottky barrier formed at $the metal-support interface^{14,15} and little is known about the nanostructuration of the support interface and the sup$ the optimum interaction. Energy transfer from the plasmonic NP to the SC remains low,¹⁶ although efficient hot electron transfer by plasmon-induced interfacial charge transfer has been reported for cadmium selenide nanorods with gold tips,¹⁷ highlighting the importance of atomic description of the interface.¹⁸ The interface in general has been found to play a key role in many catalytic processes.^{19,20} More structure-activity relationships are thus needed, which $implies \, a \, precise \, control \, of the \, M-SC \, interface \, at the nanoscale. \, We have recently$ designed an optimized heterojunction between titania and carbon nitride with improved activity for hydrogen production from water.²¹ Here we show the impact of the interface nanostructuration of gold-titania composites, in terms of accessibility of Au-TiO2 dual perimeter sites, Au NP size and density, and titania surface functionalization and crystallinity, on their compared performances in solar and visible photocatalytic gas phase reduction of CO₂ in an optimized swept ftxed-bed reactor.

Results and discussion

Under artiftcial solar conditions (see Experimental section), both methane and hydrogen are produced. No other products are detected, which is commonly observed in the gas phase process using water as a reductant.²² Fig. 1 shows the typical evolution of methane and hydrogen as a function of time. As can be seen, the production rates of both methane and hydrogen increase in the ftrst 60 to 90 minutes of reaction, at which point they reach a maximum. A decrease in methane and hydrogen production rates is then observed in the next 8.5–9h of reaction until a plateau is reached. The following discussion will be based on average production rates, using the average of all 45 production rates monitored during the 10 h reaction (one every 14 min). In our set-up, pure anatase TiO₂ UV100 readily produces 1 mmol h⁻¹ g⁻¹ of methane on average under solar irradiation. This low yet signiftcant activity obtained without metal deposition can be attributed to oxygen vacancies present on anatase crystallites²³⁻²⁶ and to the low yet signiftcant amount of UV irradiation present in the artiftcial solar spectrum (1%, see Experimental section), capable of efficiently activating anatase TiO₂.



Fig. 1 $H_2(-)$ and $CH_4(=)$ photocatalytic production rates per irradiated surface (dashed lines, right Yaxis) and per gram (continuous lines, left axis) *versus* solar irradiation time for Au_{3nm}/TiO_2 .

Oxygen vacancies are proposed to activate the particularly inert and stable CO_2 molecule by allowing CO_2 adsorption on the titania surface. Subsequent reduction of adsorbed CO_2 to CH_4 involves the *in situ* formation of the active H° reducing agent, ftrst by water oxidation on UV-generated holes, then by reduction of the resulting protons with UV-generated electrons, following the net reaction

$$CO_2 + 8H^+ + 8e^- \checkmark CH_4 + 2H_2O$$

Straight-forward addition of gold nanoparticles (1 wt%, 3 nm), using colloidal deposition on the high surface area semi-conductor, ²⁷ results in a major increase in the average solar production rate to 7.1 mmol h⁻¹ g⁻¹, *i.e.* 159 mL h⁻¹ g⁻¹. In addition, the selectivity of the process towards CH₄ vs. H₂ formation significantly increases from 45% to 60% (Fig. 2b). The starting gold-based composite thus exhibits close performance to a standard established in the literature in 2009 (160 mL h⁻¹ g⁻¹)²⁸ and a high selectivity to methane, as compared with other monometallic Au/TiO₂ photocatalysts.⁴ By acting as electron traps, Au NPs are known to enhance charge separation, *i.e.* decrease recombination occurrences, increase lifetimes of charge carriers and enhance *in situ* formation of reducing species. Au NPs are also well-known to act asco-catalysts.⁶ In addition, the presence of AuNPs on the titania surface is known to create oxygen vacancies at the perimeter interface, which could facilitate CO₂ activation.

Post-calcination effect

Calcining this Au_{3nm}/TiO_2 composite in air results in a significant decrease in methane and hydrogen solar production rates with increasing calcination temperature, until no product at all is detected when calcination is performed at 900 °C (Fig. 2a). Calcination actually causes a major loss of overall surface area, accompanied with a decrease in pore volume and modification of the mean pore diameters (Fig. 3).



Fig. 2 (a) H_2 and CH_4 production rates for TiO_2 , Au_{3nm}/TiO_2 and post-calcined Au/TiO_2 composites. The $[Au_{43nm}/TiO_2]$ -900 composite showing no activity is not shown on the graph. (b) CH_4 selectivity and production rates normalized by BET specific area for post-calcined samples, (c) H_2 and CH_4 production rates for pre-calcined Au/TiO_2 composites, and (d) CH_4 selectivity and production rates normalized by BET specific area for pre-calcined samples.

Only 43%, 6% and 1.3% of the initial surface area of the Au_{3nm}/TiO₂ composite (315 m² g⁻¹) is indeed retained a er calcination at 450 °C, 700 °C and 900 °C, respectively. Hence, by normalizing the photocatalytic productions with the specific surface area of the photocatalyst, average methane production rates are found to actually increase with increasing calcination temperature, from 23 to 43 and 85 nmol h⁻¹ m⁻², when calcining at 450 °C and 700 °C, respectively (Fig. 2b). This is attributed to the increasing degree of crystallinity of the semi-conductor particles, as evidenced by the narrowing of *e.g.* the (004) anatase relection at 2q 1/4 38° in the powder X-ray diffractograms of the composites (Fig. 4).

The Au–TiO₂ interface may also contribute to the increase in surface specific methane production rates, as shown by the change in the optical response of the composites upon calcination. The UV-vis spectra (Fig. 5a) indeed exhibit a major absorbance over the UV range, which is attributed to electronic transitions within the titania SC, as well as a smaller absorption in the visible range, typical of the localized surface plasmon resonance (LSPR) of gold nanoparticles. While the absorption related to bare titania ends at 400 nm, the remaining absorbance of the Au–TiO₂ composites in the 400–500 nm range is attributed to d–d interband transitions from the Au NPs. The most significant effect of calcination temperature on the optical response of the SPR position. Its maximum is indeed shi ed from 520 nm in Au_{3nm}/TiO₂, which is typical offree Au NPs, *i.e.* Au NPs in weak interaction with titania, to 545, 575 and up to 605 nm, when the composite is calcined at 450 °C, 700 °C and 900 °C, respectively. Although the increase in intensity is consistent with



Fig. 3 (a) N₂ physisorption isotherms of TiO₂, Au_{3nm}/TiO₂ and post-calcined samples with corresponding BET surface area values (BET surface area values for pre-calcined composites are in brackets). (b) Pore volume and (c) average pore diameters for TiO₂, Au_{3nm}/TiO₂, and pre- and post-calcined composites.

the increase in Au NP size from 3 nm to 43 nm as shown by TEM studies (Fig. 6 and Table 1), the extent of the red-shi cannot be accounted for solely by a size effect.

Such a shi implies a strengthening of the metal-support interaction, which was confirmed by HRTEM studies (Fig. 7). Thermal treatment of the composite in

Paper



Fig. 4 Diffractograms of post-calcined samples (450 $^\circ C$ is purple, 700 $^\circ C$ orange and 900 $^\circ C$ green).



Fig. 5 UV-visible spectra of (a) TiO₂, Au_{3nm}/TiO₂ and post-calcined samples and (b) TiO₂, Au_{3nm}/TiO₂ and pre-calcined samples.



Fig. 6 Representative TEM images and Au NP size distributions of (a) Au_{3nm}/TiO_2 , (b) $[Au_{5nm}/TiO_2]$ -450, (c) $[Au_{10nm}/TiO_2]$ -700, and (d) $[Au_{43nm}/TiO_2]$ -900.

air indeed induces, from 700 °C, an interaction between gold and titania akin to $the {\it so-called strong metal-support interaction (SMSI), except it is here occurring}$ upon oxidizing, and not reducing, heating. Inclusion of gold in the titania lattice is clearly observed in the composite treated at 700 °C, while some particles become fully encapsulated within a layer of titanium oxide upon further air treatment at 900 °C (Fig. 7). The former phenomenon has been described before for Pt/TiO_2 (ref. 29) and Au/TiO_2 (ref. 30) materials, while the latter (Au NP encapsulation) most likely contributes to the inactivity of [Au_{43nm}/TiO₂]-900, by lowering the metallic surface beyond the loss coming from the drastic growth of Au NPs (43 nm). It suggests that the metallic surface and the Au–TiO₂ perimeter interface sites play key roles in the photoreduction process, not from their interaction with light, but rather as direct adsorption sites for reactants or intermediates. We speculate that accessible Au-TiO₂ perimeter sites, by enhancing the number of oxygen vacancies and defects, increase the rate of CO₂ activation and that accessibility to the gold surface increases the rate of in situ formation of the reducing agent (H°).

TiO_2 pre-calcination effect

In order to evaluate the impact of the crystallinity of the SC (quasi-) independently from the metallic surface, titania was calcined at 450 °C, 700 °C and 900 °C before gold deposition. The NaBH₄-protected 3 nm gold sol was subsequently immobilized on pre-calcined SCs. Hence, variations in Au NP size with calcination temperature, as well as variations in Au–TiO₂ interactions, are minimized, as shown by the modest red shi in the LSPR position (530–560 nm, Fig. 5b). Average Au NP sizes are found to be 3.6 nm, 3.9 nm and 9.0 nm in Au_{3nm}/[TiO₂-450], Au_{4nm}/[TiO₂-700] and Au_{9nm}/[TiO₂-900], respectively (Experimental section, Table 1). Accordingly, the metallic dispersion, estimated on the basis of geometrical

Paper

Faraday Discussions Table 1 Summary of prepared samples

Sample name	Metal	TiO ₂ pre- calcination	Post- calcination	Metal loading ^a (wt%)	M NPaverage size ^b (nm)
A_{112nm}/TiO_2	A11	x	x	12T003	36T15
[A115nm/TiO2]-450	Au	X	450°C	1.3T0.03	5.1 T 0.9
$[Au_{10nm}/TiO_2]$ -700	Au	X	700 °C	1.3 T 0.03	10.0 T 1.6
$[Au_{43nm}/TiO_2]-900$	Au	Х	900 °C	1.3T0.03	42.3 T 6.0
$Au_{3nm}/[TiO_2-450]$	Au	450 °C	Х	1.05T0.03	3.6 T 1.3
$Au_{4nm}/[TiO_2-700]$	Au	700 °C	Х	0.49T0.01	3.9 T 1.0
Au _{5nm} /[TiO ₂ -700]-2	Au	700 °C	Х	0.74T0.01	4.7 T 1.7
Au _{9nm} /[TiO ₂ -900]	Au	900 °C	Х	0.91T0.02	9.0 T 4.2
Au _{15nm} /TiO ₂	Au	Х	Х	1.09T0.03	n.d.
Au _{50nm} /TiO ₂	Au	Х	Х	1.30T0.03	n.d.
Pt _{3nm} /TiO ₂	Pt	Х	Х	0.77T0.06	3.0 T 0.8
[Pt _{3nm} /TiO ₂]-700	Pt	Х	700 °C	0.86T0.06	n.d.
^{<i>a</i>} Determined by IC	P-AES.	^b Determined	l by TEM.		

calculations³¹ at 43% in the original sol and at 37% in the uncalcined Au_{3m}/TiO_2 composite, remains similar when the support is pre-calcined at 450 $^\circ C$ and 700 $^\circ C$ (37% and 35%, respectively). The Au–TiO₂ perimeter interface is thus expected to be unchanged in those samples. Pre-calcining TiO2 at 900 °C, however, induces a huge drop in surface area. Longer reaction times are thus needed to load the support with a similar amount of gold (see Experimental section), which favours aggregation and sintering of the deposited NPs (TEM, not shown). In this case, the metallic dispersion decreases to 16%, but remains much higher than in the postcalcined counterpart (ca. 4%) and the NPs are much more accessible. Structural characterization (XRD, not shown) shows that anatase crystallites grow with increasing temperature in the same manner as in the post-calcined Au/TiO₂ composites. The specific surface area decreases accordingly: only 40%, 5.4% and 0.95% of the initial surface area of the Au_{3nm}/TiO₂ composite (315 m² g⁻¹) is retained when TiO₂ is pre-calcined at 450 °C, 700 °C and 900 °C, respectively. Nevertheless, Au_{3nm}/[TiO₂-450] produces significantly more methane (and hydrogen) than the uncalcined, poorly crystallized Au_{3nm}/TiO₂ composite, setting a new standard in solar CH₄ production of 13 mmol $h^{-1}g^{-1}$, *i.e.* 291 mL $h^{-1}g^{-1}$ 1



Fig. 7 HRTEM images of (left) [Au_{43nm}/TiO₂]-900 and (right) [Au_{10nm}/TiO₂]-700.

Paper

(Fig. 2c). Crystallization of the SC is thus beneficial to the photoreduction of CO₂. In fact, surface speciftc (i.e. speciftc surface area-normalized) CH4 production rates increase with increasing pre-calcination temperature, *i.e.* crystallinity, of the titania support from 23 nmol $h^{-1} m^{-2}$ in the uncalcined, poorly crystallized Au_{3nm}/TiO_2 composite to 104 and 130 nmol $h^{-1}m^{-2}$, when pre-calcining TiO₂ at 450 °C and 700 °C, respectively (Fig. 2d). Additionally, the surface-normalized CH4 production of Au_{9nm}/[TiO₂-900] (33 nmol_{CH} $h^{-1}m^{-2}$) remains higher than that of the uncalcined composite. Crystallization of the titania SC is thus clearly favourable to the solar photocatalytic process. An increasing degree of crystallinity may indeed increase the lifetime of the charge carriers (Time Resolved Microwave Conductivity (TRMC) measurements) by avoiding defect-induced recombination.³² Yet, the surface-specific methane production rate observed over the most crystallized sample, Au_{9nm}/[TiO₂-900], is lower than those of the composites calcined at lower temperatures. It is attributed to the significantly decreased density of defects and oxygen vacancies in the highly crystallized structure (presence of the rutile phase, Fig. 4), which may induce less efficient adsorption of the carbon dioxide reactant per surface unit. In addition, the lower gold dispersion of Au_{9nm}/[TiO₂-900], as compared with Au_{3nm}/[TiO₂-450] and Au_{4nm}/[TiO₂-700], may also contribute to its lower BET-normalized activity.

Information about the impact of the actually accessible gold surface may be derived from comparison of the photocatalytic performances of the post-calcined and pre-calcined series. Indeed, the physico-chemical properties of the SCs are similar in both series (similar anatase crystallite size, nitrogen physisorption isotherms, surface areas, pore volumes and pore size distributions (Fig. 3)), while the accessible metallic surface is generally greater in the pre-calcined series. Clearly, the pre-calcined series is more active than the post-calcined composites under solar conditions, at each calcination temperature (Fig. 2). In particular, $Au_{3nm}/[TiO_2-450]$ exhibits more than twice the methane production rate of $[Au_{5nm}/TiO_2]$ -450, despite a moderate gain in accessible gold surface (0.39 vs. 0.35 wt%). Au_{9nm}/[TiO₂-900], despite its the low surface area and unlike its postcalcined counterpart, exhibits measurable photocatalytic activity (0.1 mmol_{CH_a} h^{-1} g^{-1}) which likely benefits from the much higher (0.15 vs. 0.05 wt%) and much more accessible Au surface content (no encapsulation). Finally, Au_{4nm}/[TiO₂-700] exhibits a slightly lower Au surface content (0.17 wt%) than its post-calcined counter-part (0.20 wt%), due to its peculiarly low Au loading (0.49 wt% instead of 1.1 T 0.2 wt% in all other Au/TiO₂ composites). It nevertheless exhibits an activity gain of about 30% (2.2 mmol $h^{-1}g^{-1}vs. 1.7$ mmol $h^{-1}g^{-1}$), likely due to the higher accessibility of gold (no inclusion in the TiO2 lattice). The accessible Au surface content could be marginally increased to 0.21 wt% by increasing adsorption time between the gold sol and TiO₂-700 $^{\circ}$ C (resulting in a more loaded Au_{5nm}/[TiO₂-700]-2 composite (0.74 wt%) with slightly increased NP size (4.7 nm)), with no improvement however in CH₄ production rate (1.8 mmol $g^{-1} h^{-1}$). Fig. 8 shows that both pre-calcined materials exhibit significantly higher Au NP densities (>400 NPs per mm²) than [Au_{10nm}/TiO₂]-700. This suggests that gold may play a dual role in the photocatalytic process and that the catalytic role of its available surface may be partially counterbalanced by its role as a recombination center when the density of M NPs on the titania surface is high.³³

Concerning selectivity to methane vs. H₂, the initial selectivity of the uncalcined Au_{3nm}/TiO₂ composite (60%) increases when applying calcination at 450 °C



Paper



() and post-calcined () samples.

to 63% for post-calcination and to 78% for pre-calcination. It may be attributed to the loss of terminating hydroxyl functional groups of the SC surface (dehydroxylation, TGA, not shown), which contribute preferentially to H₂ production.³⁴ The even higher selectivity of the pre-calcined composite is attributed to the more numerous population of Au–TiO₂ interface perimeter sites due to the presence of a higher density of Au NPs (197 *vs.* 74 NP per mm², Fig. 8), which results from the smaller Au NP size (3.6 nm), similar loading and slightly lower surface area (125 *vs.* 135 m² g⁻¹). The decrease in CH₄ selectivity upon further increasing the calcination temperature, *i.e.* the sites that favour CO₂ over water adsorption.²⁴

Au NP size effect

In order to probe the Au NP size effect, *i.e.* the effect of metallic dispersion, two commercially available PVP-protected gold colloids of 15 nm and 50 nm were loaded on the uncalcined TiO2 UV100 support (see Experimental section) in the same amount (1.1 T 0.2 wt%). In turn, the density of the Au NPs is also modified: it decreases significantly with increasing Au NP size. Au_{3nm}/TiO_2 and Au_{15nm}/TiO_2 lead to the same methane production rate (7.1 mmol $h^{-1}g^{-1}$, Fig. 9a), despite the much lower gold dispersion in Au_{15nm}/TiO₂ (10% instead of 37%). This suggests that methane production is not directly related to metal dispersion, which also implies that catalytic effects are not limiting. Such low impact of Au NP size (in the 3-12 nm range) on the photocatalytic efficiency has previously been observed in titania-based hydrogen production from water/ethanol mixtures.³⁵ Au_{15nm}/TiO₂ leads to the optimum internal quantum yield of this study, *i.e.* 0.012% with only 1% of UV photons emitted. Further increasing the Au NP diameter to 50 nm, *i.e.* lowering the metal dispersion to about 1%, causes a drop in the methane production rate to 1.3 mmol h⁻¹ g⁻¹. It is interesting that the activity of Au_{50nm}/ TiO₂ is close to that of the bare titania support. This is consistent with the fact that the density of Au NPs is particularly low in this composite (0.1 NP per mm²). By increasing the Au NP size on the same support (identical surface area), the



Fig. 9 (a) Comparison of visible (\Box) and solar (\blacksquare) CH₄ production rates as a function of Au NP size. (b) Comparison of CH₄ selectivity under visible (>420 nm, \triangleleft) and solar light (\blacklozenge) as a function of Au NP size.

decrease in Au dispersion is indeed accompanied by an even more significant decrease in Au NP density, from 90 NPs per mm² in Au_{3nm}/TiO₂ to 1 NP per mm² and $0.1 \text{ NP per } \text{mm}^2 \text{ in } Au_{15nm}/\text{TiO}_2$ and Au_{50nm}/TiO_2 , respectively. At such a low Au NP density, the positive impact of gold on the titania-based photocatalytic activity is modest. Nevertheless, the selectivity to methane continuously decreases with increasing Au NP size, from 60% in Au_{3nm}/TiO₂ to 15% in Au_{50nm}/TiO₂, suggesting that larger particles favour the reaction of H° with H° rather than with carbonaceous species. Larger Au NPs have previously been found active for hydrogen production from water, suggesting that H° + H° recombination may indeed occur on the surface of large Au NPs.^{36,37} More generally, larger Au NPs are reported active for radical recombination.^{38,39} On the other hand, the decreasing Au NP density decreases the number of Au-TiO2 perimeter interface sites, *i.e.* the number of CO₂ adsorption sites, which would decrease the probability for CO₂ activation and subsequent reduction to CH4. Such a decrease in CH4 vs. H2 selectivity with increasing Au NP size is thus consistent with a reaction pathway involving H° recombination over the gold surface and $H^{\circ} + CO_2*$ reaction at the Au-TiO₂ perimeter interface.

Faraday Discussions

Activity in the visible range

In order to probe the activity of the samples in the visible range, a solid ftl ter (GG420) and the samples in the visible range and the samples of the saFarbglass SCHOTT from ITOS) that cuts-off irradiation below 420 nm was applied. Under these conditions, the overall photon \Box ux (0.028 mol s⁻¹ m⁻²) is decreased by about 15%, to 0.024 mol s⁻¹ m⁻². Bare UV100 is no longer active for CO₂ photoreduction and hydrocarbon or hydrogen production (if any) falls below the detection limit. The gold-containing composites, however, all exhibit measurable activity, with average methane production rates ranging from 0.02 to 0.31 mmol h⁻¹ g^{-1} , *i.e.* 0.11 to 7.06 nmol $h^{-1} m^{-2}$ (Fig. 10). This occurs in the absence of measurable photon absorption from TiO₂ anatase above 400 nm, *i.e.* in the absence of direct light-mediated activation of the titania SC (Fig. 5). This suggests that the active charge carriers are not generated directly within the SC, but rather via the gold nanoparticles. The optical properties of Au NPs and more specifically their LSPR in the visible range are known to induce the formation of hot electrons upon excitation at the corresponding wavelengths.⁸⁻¹³ Such plasmon-derived hot electrons are reported to induce a variety of phenomena, including direct chemical reaction with adsorbates at the M NP surface^{10,15,40,41} and enrichment of the conduction band of the SC by transfer of these hot carriers from the Au NP to the SC,⁹⁻¹⁴ in addition to enhanced charge separation.^{10,42} Their lifetime, which is increased by the interaction of the metal nanoparticle with the SC via formation of a Schottky barrier, does indeed match the chemical reaction rates.

The plasmon-related activities observed here are typically one to two orders of magnitude lower than the direct, SC-based photocatalytic process, keeping in



Fig. 10 (a) H₂ and CH₄ production rates for Au_{3nm}/TiO₂ and post-calcined samples under visible light (>420 nm). The [Au_{43nm}/TiO₂]-900 composite showing no activity is not shown on the graph. (b) CH₄ selectivity and production rates normalized by BET specific surface area for post-calcined samples under visible light (>420 nm), (c) H₂ and CH₄ production rates for pre-calcined samples under visible light (>420 nm), and (d) CH₄ selectivity and production rates for pre-calcined samples under visible light (>420 nm), and (d) CH₄ selectivity and production rates normalized by BET specific area for pre-calcined samples under visible light (>420 nm).

mind that visible light irradiance with the UV-cut-off ftlter was 4600 W m⁻² compared to 5100 W m⁻² without the ftlter. Comparing the performance of all composites for visible CO₂ photoreduction, Au_{3nm}/[TiO₂-450] remains the most active of the series on a weight basis, producing $0.31 \text{ mmol}_{CH_2} h^{-1} g^{-1}$ under visible light (Fig. 10c), which is about ten times more than on Au_{3nm}/TiO₂ (Fig. 10a) followed by $[Au_{5nm}/TiO_2]$ -450 (0.25 mmol_{CH}, $h^{-1}g^{-1}$) and Au_{50nm}/TiO_2 (0.21 mmol_{CH}, h⁻¹g⁻¹) (Fig. 9a). The surface area remains an important factor, as methane production rates steadily decrease with decreasing surface area/ increasing Au NP density in the pre-calcined series to 0.12 and $0.02 \text{ mmol } h^{-1}$ g⁻¹ for Au_{4nm}/[TiO₂-700] and Au_{9nm}/[TiO₂-900], respectively. However, the most crystallized samples exhibit the highest surface-normalized activities of 6.67 and 7.06 nmol h^{-1} m⁻² for Au_{9nm}/[TiO₂-900] and Au_{4nm}/[TiO₂-700], respectively (Fig. 10d), highlighting the even higher importance of the crystallization of the SC in the plasmon-based process, as compared with the solar-induced, direct, SCbased process. Further impact of the degree of crystallization of titania can be seen by comparing "residual" activities (i.e. visible vs. solar activity ratios) in the pre-calcined series: residual activities indeed increase upon increasing calcination temperature (Fig. 11b), *i.e.* with increasing titania crystallinity (Fig. 4), with Au_{9nm}/[TiO₂-900] exhibiting the maximum residual methane production rate upon pure visible irradiation (20% of the solar activity).

The plasmon-based process also exhibits one major specificity: all calcined composites are apparently 100% selective towards methane, the amount of produced hydrogen (if any) being below the detection limit. On the other hand, uncalcined composites, except for Au_{50nm}/TiO₂, exhibit a much lower CH₄ selectivity under visible light than under solar irradiation (Fig. 9b). The beneftcial effect of the SC crystallization on the CH4 vs. H2 selectivity is thus exacerbated upon visible-only irradiation. In particular, the most crystallized composites (calcined at 700 °C and 900 °C) exhibit the same complete methane selectivity as the composites calcined at 450 °C, despite the lower density of defects which is unfavourable for CH4 adsorption. Hence the absence of unselective surface hydroxyl groups is not sufficient to account for the total selectivity to methane. Another phenomenon might contribute to the absence of hydrogen product. Molecular hydrogen is indeed known to readily dissociate over plasmon-activated Au NPs.^{15,40,43} Hence, even if it was formed in the present water-mediated CO₂ photoreduction process, it would not be stable. One may also speculate that it may simply not be formed as it is in the SC-based process and that H° + H° recombination simply does not occur in the gold-driven, plasmon-based process. In addition to the instability of the hydrogen molecule under plasmonic conditions, the absence of UV irradiation, which is known $to enhance formation of defects on titania, {}^{44,45} may also lower the possibility of$ CO₂ adsorbing on titania. Under visible-only irradiation, CO₂ adsorption and activation would thus be restricted to the gold surface, where it would more efficiently compete with H°, especially on dehydroxylated titania surfaces which likely lower the reverse spillover rate of H°.⁴⁶ Consumption of in situ formed hydrogen in the direct hydrogenation of CO₂ is less likely since gold is active for the reverse water-gas shi reaction leading to $CO(CO_2 + H_2 \checkmark CO + H_2O)$ and not for CO₂ methanation leading to methane (CO₂ + $4H_2 \checkmark CH_4 + 2H_2O$),⁴⁷ which could not account for the CH4 product formed here.

Paper



 $\frac{1}{2}$

Additionally, the pre-calcined materials (Fig. 10c and d) remain more active than their post-calcined counter-parts (Fig. 10a and b) under visible irradiation, showing that accessibility of the metal surface (and Au-TiO2 interface) remains key in the plasmon-based process. However, the visible CH₄ production rates (Fig. 9a), as well as the visible vs. solar activity ratio (Fig. 11a), increase with increasing Au NP size (i.e. with decreasing metallic surface). Au_{50nm}/TiO₂ retains 16% of activity upon pure visible irradiation, the second highest of the series, just a er the most crystallized composite and despite the poor crystallinity of the uncalcined titania support. Larger Au NPs are thus found to be more efficient for $plasmon-derived\ CO_2\ photoreduction, in agreement\ with\ previous\ reports.^{15,48,49}$ This suggests that the optical properties of gold prevail over the use of the gold surface in the plasmon-induced process. The intensity of the SPR indeed increases with increasing Au NP size, enhancing hot electron generation.^{49,50} In our composites, the plasmon-driven process thus appears limited by the initial charge generation, which may be due to the low metal loading used, as compared with those used in plasmonic studies.^{40,48,49}

Platinum vs. gold

Further evidence of gold involvement in the plasmon-based process is obtained by comparison with two Pt-based photocatalysts prepared under similar

Paper

conditions. Neither Pt_{3nm}/TiO₂ nor [Pt_{3nm}/TiO₂]-700 exhibit any detectable activity for CO₂ photoreduction under visible light-only conditions. The position of the Pt NP LSPR in the UV range (220 nm) indeed prevents any Ptmediated CO₂ photoreduction in the visible range. However, it may account for the much higher CH₄ production rate observed over Pt_{3nm}/TiO₂ under solar $irradiation of 91 \, mmol_{CH} \, h^{-1} g^{-1}$. The one order of magnitude higher activity of Pt vs. Au in the uncalcined composite may be explained by Plasmon Resonance Energy Transfer (PRET) i.e. the ability of the LSPR to enhance the electric fteld intensity in the SC, *i.e.* the ability of the LSPR to enhance light absorption by the SC,⁹ due to the overlap of the Pt LSPR and TiO₂ absorption bands. The higher workfunction of platinum vs. gold, which theoretically enhances charge separation, may also contribute to the higher solar activity of platinum. On the other hand, the activity of gold-based composites in the visible range is related to the inability of the SPR to transfer its energy to TiO₂ and to the hot carriers generated within the gold nanoparticle as a result. It is interesting that the PRET phenomenon in the SC-based photoreduction of CO₂ disappears with increasing crystallinity of titania and possible loss of accessible metallic surface via strong Pt-TiO2 interactions. In the composites post-calcined at 700 °C, surface-normalized CH4 production rates are indeed similar on Pt and Au (83 vs. 85 nmol h^{-1} m⁻²). The expected higher density of Pt vs. Au NPs in a low surface area [Pt_{3nm}/TiO₂]-700 (29 m² g⁻¹) may also inhibit the photocatalytic activity of this composite. Additionally, the initially higher selectivity of platinum towards $CH_4 \textit{vs}. H_2$ formation of 88% decreases below that of gold in the calcined composite (31% vs.46%).

TRMC measurements and charge carrier dynamics

The time resolved microwave conductivity method was previously used to determine the charge carrier lifetimes in TiO₂-based composites.^{21,33,51,52} The main data provided by TRMC are the maximum value of the signal (I_{max}), which indicates the number of excess charge carriers (mostly electrons) created by the pulse, including decay processes during the excitation by the laser (10 ns), and the decay (I(t)) due to the decrease of mobile electrons, either by recombination or trapping processes. Concerning the decay, *i.e.* the lifetime of charge carriers, a short and long range are usually analyzed. The short-range decay, arbitrarily fixed up to 40 ns a er the beginning of the pulse re \Box ects fast processes, recombination of charge carriers and/or trapping phenomena.^{51,52}

At 360 nm excitation (Fig. 12), one can observe that the [Au_{5nm}/TiO₂]-450 and Au_{3nm}/TiO₂ samples lead to the most intense signal, followed by the TiO₂ UV100 support without Au NPs. However, the ftrst material exhibits slightly higher charge carrier generation accompanied with a lower signal decay, characteristic of either slower charge recombination dynamics or a more pronounced electron trapping effect. Comparing those samples with the TiO₂ UV100 support, it must be underlined that Au NP deposition does not lead to an important signal intensity decrease, as very o en observed a er metal nanoparticle deposition and resulting from the strong electron trapping effect attributed to Au NPs. Nevertheless, this trapping effect induced by Au NP deposition may be argued for the two photocatalysts presenting the largest Au NPs, *i.e.* Au_{50nm}/TiO₂ and mainly [Au_{10nm}/TiO₂]-700. It is noted that absence of a signal does not mean that no



Fig. 12 TRMC signals at 360 nm on TiO₂ UV100 (pink), TiO₂ P25 (light green), Au_{3nm}/TiO_2 (dark blue), [Au_{5nm}/TiO_2]-450 (red), [Au_{10nm}/TiO_2]-700 (dark green), $Au_{3nm}/[TiO_2-450]$ (dark purple), $Au_{4nm}/[TiO_2-700]$ (light blue), Au_{15nm}/TiO_2 (orange) and Au_{50nm}/TiO_2 (light purple).

charges were created but it could be correlated with very rapid charge recombination a erlight excitation. It must also be noticed that $Au_{4nm}/[TiO_2-700]$ presents specific features, with the lowest signal decay amongst all the samples. This can be correlated either with low charge carrier recombination kinetics or with higher electron trapping effects. It might be supposed that the presence of the rutile phase due to the prior thermal treatment of TiO_2 UV100 at 700 °C, in intimate contact with the anatase phase, and the resulting formation of a nanoheterojunction between both phases lead to spatial charge separation, thus limiting their recombination.

At 450 nm excitation (in the blue region of visible light, Fig. 13), $Au_{4nm}/[TiO_2-700]$ also exhibits the lowest signal decay, which could be attributed to the same effect of limitation of charge recombination due to the presence of the TiO₂ rutile phase. In addition, this material yields the highest signal that may be explained by the highest generation of charges due to the tail of absorption of rutile at that wavelength of excitation, as compared with TiO₂P25 containing *ca.* 20% of rutile. Furthermore, the signal intensity is very close to that obtained under UV-A activation. This last explanation could also be considered for [$Au_{10nm}/[TiO_2]$ -700. However, in this case the decay in electron mobility is higher, perhaps due to more efficient electron trapping onto larger Au NPs. It must be mentioned that TiO₂ UV100 also shows a small signal that can be associated with charge generation that might result from a defective TiO₂ structure. TiO₂ UV100 is indeed not as well crystallized as TiO₂ P25, for example. Nevertheless, those charges rapidly recombine.

At LSPR maximum excitation (or at 500 nm excitation in the absence of LSPR), and as for activation at 450 nm, $Au_{4nm}/[TiO_2-700]$ reveals the highest signal associated with the highest charge carrier generation (Fig. 14). In this case, it must be underlined that the signal intensity is only slightly reduced compared to excitation at 450 nm. The other sample containing TiO₂ calcined at 700 °C, *i.e.*

Paper



Fig. 13 TRMC signals at 450 nm on TiO₂ UV100 (pink), TiO₂ P25 (light green), Au_{3nm}/TiO₂ (dark blue), [Au_{5nm}/TiO₂]-450 (red), [Au_{10nm}/TiO₂]-700 (dark green), Au_{3nm}/[TiO₂-450] (dark purple), Au_{4nm}/[TiO₂-700] (light blue), Au_{15nm}/TiO₂ (orange) and Au_{50nm}/TiO₂ (light purple).

 $[Au_{10nm}/TiO_2]$ -700, also seems to generate charge carriers, however, to a lesser extent and with a rapid decrease of the signal, probably due to the higher electron trapping effect on larger Au NPs. It must be noticed here that TiO₂ P25 does not show any signal, meaning that absorption due to the presence of the rutile phase is no more effective at 500 nm. From these observations, it may be assumed that LSPR effects may take place through the injection of hot electrons from the Au NPs to the conduction band of TiO₂ nanoparticles. Additionally, the enhanced surface-normalized CH₄ production rates observed over composites calcined at 700 °C may thus be attributed to an increased charge carrier lifetime,



Fig. 14 TRMC signals at 500 nm on TiO₂ UV100 (pink), TiO₂ P25 (light green), Au_{3nm}/TiO_2 (dark blue), [Au_{5nm}/TiO_2]-450 (red), [Au_{10nm}/TiO_2]-700 (dark green), $Au_{3nm}/[TiO_2-450]$ (dark purple), $Au_{4nm}/[TiO_2-700]$ (light blue), Au_{15nm}/TiO_2 (orange) and Au_{50nm}/TiO_2 (light purple).

Faraday Discussions

which may be related to the peculiar oblong shape of the titania crystallites present in these materials (Fig. 7).⁵³ The peculiar shape of the SC may result from restructuration related to the anatase / rutile transition and is consistent with the speciftcity of TiO_2 -700 °C towards Au NP deposition and interaction (Fig. 5b). Finally, the absence of a signal for most of the composites further supports the limitation of the plasmon-based process by charge generation. It is expected that the resulting modest photocatalytic activities may be improved by increasing the gold content.

Experimental

Synthesis of photocatalysts

All photocatalysts were prepared through a well-known colloidal deposition method^{27,54-56} but without any protecting agent, using commercially available 100% anatase titania (Hombikat UV100 from Sachtleben Chemie GmbH, 384 g m⁻²). For all photocatalysts, a gold loading of 1 wt% was targeted. First, an aqueous solution of sodium borohydride (0.1 M, 2.54 mL) was added to an aqueous solution of HAuCl₄ (0.25 M, 200 mL), with Au/NaBH₄ ¼ 1 : 5 mol mol⁻¹. The yellow gold solution immediately turned dark red upon introduction of NaBH₄, due to the formation of a gold sol, which is stabilized by the presence of NaBH₄-derived molecules in electrostatic interaction with the Au NP.⁵⁷ The support (1 g) was introduced directly a er the reduction step and the mixture was stirred for 30 min for total adsorption of the gold sol onto the support. A er ftltration and washing (H₂O/1 L), the pink powder was dried in air at 100 °C for 24 h. The resulting Au_{3nm}/TiO₂ product was calcined for 4 h in \Box owing air (100 mL min⁻¹) at different temperatures: 450 °C, 700 °C and 900 °C, yielding [Au_{5nm}/TiO₂]-700 and [Au_{43nm}/TiO₂]-900, respectively.

A "pre-calcined" sample series was also prepared using the same method, with the TiO₂ support previously calcined for 4 h in \Box owing air (100 mL min⁻¹) at the same temperatures (450 °C, 700 °C and 900 °C). For these samples, the contact time between the NaBH₄-stabilized gold sol and the support was extended in order to maximize adsorption of the gold sol on the lower surface area supports. The mixture was thus stirred for 2 h, 48 h and 48 h, when using TiO₂-450 °C, TiO₂-700 °C and TiO₂-900 °C, respectively, yielding Au_{3nm}/[TiO₂-450], Au_{4nm}/[TiO₂-700] and Au_{9nm}/[TiO₂-900], respectively. In the case of TiO₂-700 °C, deposition of the gold sol was subsequently tried for 96 h in order to increase the goldloading which appeared to be much lower (0.49 wt%) than in all the other composites (1.1 T 0.2 wt%). A new composite, yet still with a low gold loading of 0.74 wt%, was obtained as Au_{5nm}/[TiO₂-700]-2.

Colloidal solutions of gold nanoparticles of 15 nm and 50 nm were purchased (nanoComposix, Econix, Au nanospheres, PVP) and directly deposited onto the uncalcined UV100 TiO₂ support. The colloidal solution of gold (5 mg mL⁻¹, 1 mL) and the support (0.5 g) were added to 100 mL of methanol (Carlo ERBA reagents S.A.S, 99.9%) and stirred for 2 h. A er fltration and washing (MeOH/1 L), the powderwas dried at 100 °C for 24 h. Two different samples were obtained: Au_{15nm}/TiO₂ and Au_{50nm}/TiO₂.

Platinum-based materials were prepared by a similar colloidal deposition method, except reduction of the platinum salt was performed in the presence of the titania support.⁵⁸ Titania (1 g) was added to an aqueous solution of H₂PtCl₆

Paper

(0.256 M, 200 mL) and stirred for 20 min. Then, an aqueous solution of sodium borohydride (0.1 M, 2.54 mL, Pt/NaBH4 ¼ 1:5 mol mol⁻¹) was added and the mixture was stirred for 30 min. A er ftltration and washing (H₂O/1 L), the grey powder was dried in air at 100 °C for 24 h. Part of the resulting Pt_{3nm}/TiO₂ composite was then calcined in \Box owing air (100 mL min⁻¹) for 4 h at 700 °C, yielding [Pt_{3nm}/TiO₂]-700. All samples are summarised in Table 1.

Characterization

A Micromeritics Asap 2420 porosimeter was used to obtain Nitrogen adsorptiondesorption isotherms at 77 K. Materials were outgassed at 250 °C under primary vacuum for 4 h in order to desorb impurities from their surfaces and eliminate any water. Specific surface areas were calculated by the Brunauer, Emmett and Teller (BET) method in the relative pressure (p/p_0) range 0.05–0.3 and using t-plot methods.

X-ray diffraction (XRD) patterns were collected on a Bruker D8 Advance theta-theta diffractometer (Cu Ka radiation, $1 \frac{1}{4} 0.154$ nm), equipped with a LynxEye detector and operating at 40 kV and 40 mA. The datasets were acquired in step-scan mode over the $36-49^{\circ}$ 2q range, using a step interval of 0.024° and a counting time of 4 s per step. Powder samples were introduced in an alumina crucible, treated *in situ* at the desired temperature in \Box owing synthetic air using a heating chamber and cooled down to room temperature prior to data acquisition.

Transmission Electron Microscopy (TEM) was carried out on a JEOL 2100F microscope working at a 200 kV accelerating voltage and equipped with a probe corrector for spherical aberrations, giving a point-to-point resolution of 0.18 nm. The sample was dispersed in an ethanol solution for 2 min. A drop of the solution was subsequently deposited on a copper grid covered with a holey carbon membrane for observation. Statistical M NP size distributions were obtained on brightfield images using populations of more than 200 nanoparticles and ImageJ so ware.⁵⁹

UV-vis absorption spectra were recorded on a Perkin Elmer Lambda 950 Scan spectrophotometer equipped with a 100 mm diffuse redectance integrating sphere. UV-vis spectra were done on ftlms: powders were deposited on 2 cm diameter glass disks with the same surface concentration and using the same protocol as for the photocatalytic tests (25 g m⁻²).

Elemental analyses of the photocatalysts were performed using Inductively Coupled Plasma Atomic Emission Spectroscopy (ICP-AES). The limit of detection of the instrument is 0.1 mg L⁻¹ for Au. Analysis of the gold component of the Au/ TiO₂ composites allowed us to calculate deposition yields which are defined as the ratio between the actual deposited mass of Au (deduced from ICP-AES analysis) and the mass of Au introduced in the synthesis.

The charge carrier lifetimes in the TiO₂-based composites were determined by microwave absorption experiments using the Time Resolved Microwave Conductivity method (TRMC).^{51,52} The TRMC technique is based on measurement of the change of microwave power re \Box ected by a sample, D*P*(*t*), induced by its laser pulsed illumination. The relative difference D*P*(*t*)/*P* can be correlated, for small perturbations of conductivity, to the difference of the conductivity Ds(*t*) considering the following equation:

Faraday Discussions

$$\frac{DP\delta t^{\flat}}{P} \frac{1}{4} A Ds\delta t^{\flat} \frac{1}{4} A e^{-\sum_{i} Dn_{i}\delta t^{\flat} m_{i}}$$

where $Dn_i(t)$ is the number of excess charge-carriers *i* at time *t* and m_i is their mobility. The sensitivity factor *A* is independent of time, but depends on different factors such as the microwave frequency or the dielectric constant. Considering that the trapped species have a small mobility, which can be neglected, Dn_i is reduced to mobile electrons in the conduction band and holes in the valence band. And in the specific case of TiO₂, the TRMC signal can be attributed to electrons because their mobility is much larger than that of the holes.^{51,52}

The incident microwaves were generated by a Gunn diode of the K_a band at 30 GHz. The pulsed light source was an OPO laser (EKSPLA, NT342B) tunable from 225 to 2000 nm. It delivers 8 ns fwhm pulses with a frequency of 10 Hz. The light energy densities received by the sample were 0.8, 1.5 and 1.4–2.2 mJ cm⁻² at 360 nm, 450 nm and at the LSPR maximum (500 nm by default), respectively.

Photocatalytic tests

The photocatalytic set-up is composed of four main parts. (1) Generation of the reaction mixture was carried out by bubbling CO₂ (Air Product, 99.5%) through a thermo-controlled water saturator maintained at 35 °C (absolute humidity ¼ 4.3%) at $0.3 \,\mathrm{mLmin^{-1}}$. (2) The continuously \Box owing reaction mixture (CO₂+H₂O) \Box ow passes through a light-transparent photoreactor (6 mL) equipped with a Hg lamp top-irradiation system, simulating artificial solar light with light concentrator (150 W Ceramic-Metal-Halide Hg Lamp). The irradiance was 46 W m⁻² for wavelengths between 300 and 400 nm and 5100 W m⁻² for the whole lamp spectrum. The irradiated surface into the photoreactor is 18.8 cm². (3) Ventilators were fixed on the photoreactor to minimize temperature increase during photocatalytic tests. The maximum temperature of the reactor was about 40 °C during test. (4) Finally, products from the reactor were analysed online using a micro-GC analytical system (Agilent 3000A SRA instrument).

To perform photocatalytic tests, 50 mg of materials were homogeneously suspended in ethanol and slowly deposited on a 50 mm diameter glass disk until dry and further thermally treated at 100 °C for 2 h. The surface concentration of the photocatalyston the glass disc was 25 gm^{-2} . The glass disk was then inserted into the photoreactor and the set-up was purged for 5 min with a high CO₂ \square ow (>100 mL min⁻¹) to remove air and other gas impurities. Before starting the test, the set-up was run with a CO₂ gas \square ow (0.3 mL min⁻¹) without illumination to achieve dark absorption equilibrium. Finally, the lamp was switched on. The standard duration of a photocatalytic test was 10 h.

Production rates were calculated according to these equations:

$$r \quad \text{mol } h^{-1} \text{ g}^{-1} \overset{\text{X]} \times 10^{-6} \times \text{\delta flow rate} \triangleright \times 60}{V \times m_{\text{photocat}}}$$

$$r \cdot \mathrm{mol} \, \mathrm{h}^{-1} \, \mathrm{m}_{\mathrm{irradiated}}^{-2} \, \frac{\mathrm{X}}{\mathrm{X}} \, \frac{\mathrm{X} \times 10^{-6} \times \mathrm{\delta flow} \, \mathrm{rateP} \times 60}{V \times S_{\mathrm{irradiated}}}$$

where [X] is the concentration of CH₄ or H₂ in ppm, m_{photocat} is the mass of photocatalyst, $S_{\text{irradiated}}$ is the irradiated surface (0.001885 m²), the \Box ow rate is 0.0003 L min⁻¹ and V_{m} ¼ 24.79 L mol⁻¹ (STP conditions).

Selectivity was calculated using these equations:

$$CH_{4} \text{ selectivity } \frac{8 \times r_{CH_{4}}}{8 \times r_{CH_{4}} \not\models 2 \times r_{H_{2}}}$$

 H_2 selectivity $\frac{1}{4}$ 1 — [CH₄ selectivity]

The production rate per specific surface area of material (r_X /BET) was calculated by dividing the production rate per gram by the BET specific surface area of each compound.

Internal quantum yields h were calculated by this formula:

h
$$\frac{1}{4} \frac{8 \times r_{\mathrm{CH}_4} \not p \, 2 \times r_{\mathrm{H}_2}}{\mathbf{f}_{\mathrm{absorbed photons}}} \frac{1}{4} \frac{n_{\mathrm{e}}}{\mathbf{f}_{\mathrm{absorbed photons}}}$$

where $\mathbf{f}_{absorbed photons}$ is the absorbed photon \Box ux determined by UV-vis absorbance measurements.

Conclusions

By evaluating the solar and visible photocatalytic activities of a series of goldtitania nanocomposites for the reduction of carbon dioxide with water, we have shown the contribution of a plasmon-based process, in addition to the direct UVinduced SC-based process, in the overall solar transformation. Comparison with Pt composites and bare titania SCs suggests that the visible activity, induced by gold nanoparticles only, relies on the generation of hot carriers in the M NPs. Such a hot electron-mediated process may account for up to 20% of the overall solar activity in highly crystallized composites. It is more efficient over larger Au NPs (50 nm). One major specificity of the gold-mediated visible transformation is the 100% selectivity towards methane, *i.e.* the preferential reduction of CO_2 over H₂O. Dehydroxylation of titania is crucial to achieve such total selectivity.

Conflicts of interest

There are no con licts to declare.

Acknowledgements

Financial support from IFPEN is gratefully acknowledged. The authors also wish to thank Dris Ihiawakrim (IPCMS, CNRS UMR 7504) for performing TEM analyses. ZC and CCJ gratefully acknowledge the financial support of the Agence Nationale de la Recherche (ANR) through the UpPhotoCat project.

Paper

Notes and references

- 1 W.-N. Wang, J. Soulis, Y. J. Yang and P. Biswas, *Aerosol Air Qual. Res.*, 2014, 14, 533.
- 2 S. N. Habisreutinger, L. Schmidt-Mende and J. K. Stolarczyk, Angew. Chem., Int. Ed., 2013, 52, 7372.
- 3 J. K. Stolarczyk, S. Bhattacharyya, L. Polavarapu and J. Feldmann, *ACS Catal.*, 2018, 8, 3602.
- 4 S. Neaţu, J. A. Maciá-Agulló, P. Concepción and H. Garcia, J. Am. Chem. Soc., 2014, 136, 15969.
- 5 Q. Zhai, S. Xie, W. Fan, Q. Zhang, Y. Wang, W. Deng and Y. Wang, *Angew. Chem., Int. Ed.*, 2013, 52, 5776.
- 6 J. Ran, M. Jaroniec and S.-Z. Qiao, Adv. Mater., 2018, 30, 1704649.
- 7 F. Zaera, Top. Catal., 2018, 61, 336.
- 8 A. Manjavacas, J. G. Liu, V. Kulkarni and P. Nordlander, *ACS Nano*, 2014, 8, 7630.
- 9 S. C. Warren and E. Thimsen, Energy Environ. Sci., 2012, 5, 5133.
- 10 Y. Zhang, S. He, W. Guo, Y. Hu, J. Huang, J. R. Mulcahy and W. D. Wei, Chem. Rev., 2018, 118, 2927.
- 11 M. L. Brongersma, N. J. Halas and P. Nordlander, Nat. Nanotechnol., 2015, 10, 25.
- 12 G. V. Hartland, L. V. Besteiro, P. Johns and A. O. Govorov, ACS Energy Lett., 2017, 2, 1641.
- 13 P. Christopher and M. Moskovits, Annu. Rev. Phys. Chem., 2017, 68, 379.
- 14 Y. Shiraishi, N. Yasumoto, J. Imai, H. Sakamoto, S. Tanaka, S. Ichikawa, B. Ohtani and T. Hirai, *Nanoscale*, 2017, 9, 8349.
- 15 S. Mukherjee, L. Zhou, A. M. Goodman, N. Large, C. Ayala-Orozco, Y. Zhang, P. Nordlander and N. J. Halas, J. Am. Chem. Soc., 2014, 136, 64.
- 16 X.-C. Ma, Y. Dai, L. Yu and B.-B. Huang, Light: Sci. Appl., 2016, 5, e16017.
- 17 K. Wu, J. Chen, J. R. McBride and T. Lian, Science, 2015, 349, 632.
- 18 M. J. Kale and P. Christopher, Science, 2015, 349, 587.
- 19 V. Caps, S. Arrii, F. Morftn, G. Bergeret and J.-L. Rousset, Faraday Discuss., 2008, 138, 241.
- 20 I. Ro, J. Resasco and P. Christopher, ACS Catal., 2018, 8, 7368.
- 21 C. Marchal, T. Cottineau, C. Colbeau-Justin, V. Caps and V. Keller, *Adv. Energy Mater.*, 2018, 8, 1702142.
- 22 E. Karamian and S. Sharifnia, J. CO2 Util., 2016, 16, 194.
- 23 S. Bai, N. Zhang, C. Gao and Y. Xiong, Nano Energy, 2018, 53, 296.
- 24 Z. Xiong, Z. Lei, Y. Li, L. Dong, Y. Zhao and J. Zhang, J. Photochem. Photobiol., C, 2018, 36, 24.
- 25 Y. Ji and Y. Luo, J. Am. Chem. Soc., 2016, 138, 15896.
- 26 L. Liu, H. Zhao, J. M. Andino and Y. Li, ACS Catal., 2012, 2, 1817.
- 27 P. Lignier, M. Comotti, F. Schüth, J.-L. Rousset and V. Caps, *Catal. Today*, 2009, 141, 355.
- 28 O.K. Varghese, M. Paulose, T.J. LaTempa and C.A. Grimes, *NanoLett.*, 2009, 9, 731.
- 29 M.Zhang, Z.Jin, Z.Zhang and H. Dang, *Appl. Surf. Sci.*, 2005, 250, 29; Q.Li, K. Wang, S. Zhang, M. Zhang, J. Yang and Z. Jin, *J. Mol. Catal. A: Chem.*, 2006, 258, 83.

- 30 W. Gao, S. Sivaramakrishnan, J. Wen and J.-M. Zuo, Nano Lett., 2015, 15, 2548.
- 31 R. Van Hardeveld and F. Hartog, Surf. Sci., 1969, 15, 189.
- 32 M. Pelaez, N. T. Nolan, S. C. Pillai, M. K. Seery, P. Falaras, A. G. Kontos, P. S. M. Dunlop, J. W. J. Hamilton, J. A. Byrne, K. O'shea, M. H. Entezari and D. D. Dionysiou, *Appl. Catal.*, B, 2012, 125, 331.
- 33 K. M. Schindler and M. Kunst, J. Phys. Chem., 1990, 94, 8222.
- 34 P. Akhter, M. Hussain, G. Saracco and N. Russo, Fuel, 2015, 149, 55.
- 35 M. Murdoch, G. I. N. Waterhouse, M. A. Nadeem, J. B. Metson, M. A. Keane, R. F. Howe, J. Llorca and H. Idriss, *Nat. Chem.*, 2011, 3, 489.
- 36 C. Marchal, M. Behr, F. Vigneron, V. Caps and V. Keller, New J. Chem., 2016, 40, 4428.
- 37 C. Marchal, A. Piquet, M. Behr, T. Cottineau, V. Papae himiou, V. Keller and V. Caps, *J. Catal.*, 2017, 352, 22.
- 38 V. Mendez, V. Caps and S. Daniele, Chem. Commun., 2009, 3116.
- 39 A. Fkiri, A. Mezni, C. Robert, V. Caps and L. S. Smiri, *Colloids Surf.*, A, 2017, 530, 85.
- 40 S. Mukherjee, F. Libisch, N. Large, O. Neumann, L. V. Brown, J. Cheng, J. B. Lassiter, E. A. Carter, P. Nordlander and N. J. Halas, *Nano Lett.*, 2013, 13, 240.
- 41 U.Aslam, V.G. Rao, S. Chavez and S. Linic, Nat. Catal., 2018, 1, 656.
- 42 Y. Tian and T. Tatsuma, J. Am. Chem. Soc., 2005, 127, 7632.
- 43 G. Ganguly, M. Sultana and A. Paul, J. Phys. Chem. C, 2017, 121, 15326.
- 44 H. Zhang, S. Kawamura, M. Tamba, T. Kojima, M. Yoshiba and Y. Izumi, J. *Catal.*, 2017, 352, 452.
- 45 R. Wang, K. Hashimoto, A. Fujishima, M. Chikuni, E. Kojima, A. Kitamura, M. Shimohigoshi and T. Watanabe, *Nature*, 1997, 388, 431.
- 46 H.-Y. T. Chen, S. Tosoni and G. Pacchioni, ACS Catal., 2015, 5, 5486.
- 47 X. Zhang, X. Li, D. Zhang, N. Q. Su, W. Yang, H. O. Everitt and J. Liu, *Nat. Commun.*, 2017, 8, 14542.
- 48 J. S. DuChene, B. C. Sweeny, A. C. Johnston-Peck, D. Su, E. A. Stach and W. D. Wei, Angew. Chem., Int. Ed., 2014, 126, 8021.
- 49 K. Qian, B. C. Sweeny, A. C. Johnston-Peck, W. Niu, J. O. Graham, J. S. DuChene, J. Qiu, Y.-C. Wang, M. H. Engelhard, D. Su, E. A. Stach and W. D. Wei, J. Am. Chem. Soc., 2014, 136, 9842.
- 50 X.-T. Kong, Z. Wang and A. O. Govorov, Adv. Opt. Mater., 2017, 5, 1600594.
- 51 O.A. Tahiri Alaoui, C. Herissan, M. E. Le Quoc, M. Zekri, S. Sorgues, H. Remita and C. Colbeau-Justin, J. Photochem. Photobiol., A, 2012, 242, 34.
- 52 C. Colbeau-Justin, M. Kunst and D. Huguenin, J. Mater. Sci., 2003, 38, 24.
- 53 Y.Ohsaki, N. Masaki, T. Kitamura, Y. Wada, T. Okamoto, T. Sekino, K. Niihara and S. Yanagida, *Phys. Chem. Chem. Phys.*, 2005, 7, 4157.
- 54 S.Mishra, V.Mendez, E.Jeanneau, V. Caps and S. Daniele, *Eur. J. Inorg. Chem.*, 2013, 500.
- 55 V. Mendez, K. Guillois, S. Daniele, A. Tueland V. Caps, *Dalton Trans.*, 2010, 39, 8457.
- 56 S. Biella, F. Porta, L. Prati and M. Rossi, Catal. Lett., 2003, 90, 1.
- 57 C. Deraedt, L. Salmon, S. Gatard, R. Ciganda, R. Hernandez, J. Ruiz and D. Astruc, *Chem. Commun.*, 2014, 50, 14194.
- 58 K. Guillois, L. Burel, A. Tuel and V. Caps, Appl. Catal., A, 2012, 415–416, 1.
- 59 W. Rasband, http://rsb.info.nih.gov/ij/.