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Daniel Pereira, JASA

Effect of intracortical bone properties on the phase velocity and cut-off frequency of low-frequency guided wave modes (20-85 kHz)

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The assessment of intracortical bone properties is of interest since early-stage osteoporosis is associated with resorption in the endosteal region. However, understanding the interaction between ultrasonic guided waves and the cortical bone structure remains challenging. The purpose of this work is to investigate the effect of intracortical bone properties on the ultrasonic response obtained at low-frequency (<100 kHz) using an axial transmission configuration. The semi-analytical finite element method was used to simulate the propagation of guided waves in a waveguide with realistic geometry and material properties. An array of 20 receivers was used to calculate the phase velocity and cut-off frequency of the excited modes using the 2D Fourier transform. The results show that the position of the emitter around the circumference of the bone is an important parameter to control since it can lead to variations of up to 10 dB in the amplitude of the transmitted modes. The cut-off frequency of the high order modes was, however, only slightly affected by the circumferential position of the emitter, and was sensitive mainly to the axial shear modulus. The phase velocity and cut-off frequency in the 20-85 kHz range are promising parameters for the assessment of intracortical properties.

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I. INTRODUCTION

Osteoporosis is associated with a reduction of the bone mass and microarchitectural bone deterioration, leading to a reduction in bone quality ([Giangregorio *et al.*, 2006](#); [Papaioannou *et al.*, 2004](#); [Ström *et al.*, 2011](#)) and to an increase in fracture risk ([Consensus development conference: prophylaxis and treatment of osteoporosis, 1991](#)). Endosteal resorption of cortical bone results in an increase in bone porosity at the inner part of the cortical shell and in a reduction of the cortical bone thickness ([Ritzel *et al.*, 1997](#)). Dual-energy X-ray absorptiometry (DEXA), which is currently the gold standard for osteoporosis diagnostics ([Kanis, 1994](#)), nonetheless suffers from several limitations ([Gluer, 2008](#); [Haba *et al.*, 2016](#)). Quantitative ultrasound (QUS) was developed to assess bone quality and has the advantage of being non-radiative, non-invasive and relatively cheap ([Gluer, 1997](#)). It, therefore, holds promise as a rapid screening method in a clinical setting. The mechanical nature of ultrasonic waves allows the technique to be used to retrieve the biomechanical properties of bone tissue ([Kaufman and Einhorn, 1993](#)). QUS techniques can be used to go beyond a simple estimation of the bone mineral density ([Nicholson, 2008](#)), which can be achieved with DEXA and which is not sufficient to assess fracture risk ([Office of the Surgeon General, 2004](#)).

Bone QUS was initially developed in the context of trabecular bone characterization using transverse transmission devices ([Stein *et al.*, 2013](#)). However, the investigation of cortical bone ([Rico, 1997](#)) has attracted significant interest since about 80% of the skeleton is made of cortical bone, which supports most of the body load, and is involved in osteoporotic

fractures (Seeley *et al.*, 1991). QUS using axial transmission techniques have mostly been used to study the cortical bone quality (Haiat *et al.*, 2011). Cortical bones, such as the radius or the tibia, constitute suitable waveguides for the propagation of ultrasonic guided waves, as has been shown in various studies, including for instance (Gluer, 2008; 1997; Haba *et al.*, 2016). Ultrasonic guided waves have the advantage of being sensitive to both the mechanical and geometrical properties of cortical bone (Foiret *et al.*, 2014; Muller *et al.*, 2005; Rozental *et al.*, 2013).

Most studies focusing on the assessment of cortical bone properties using axial transmission consider plate or cylinder waveguides (as an approximation of the actual bone geometry) to simplify the interpretation of ultrasonic responses obtained numerically or experimentally (Kilappa *et al.*, 2015; Le *et al.*, 2010; Minonzio *et al.*, 2015; Xu *et al.*, 2016). Recently, more comprehensive numerical approaches were explored with the aim of improving the realism of the effect of cortical bone features on the axial transmission propagation expected *in vivo*. Bossy *et al.* (2004) (Bossy *et al.*, 2004) performed three-dimensional (3D) finite difference simulations on geometries derived from a human radius tomography image to evaluate the effect of bone curvature, anisotropy, and micro-porosity on the first arriving signal (FAS) velocity. However, the actual anatomical variations of bone biomechanical properties were not taken into account and more emphasis was brought to the analysis of numerical simulations performed with idealized 3D objects (tubes and semi-infinite hemicylinder geometries). Moilanen *et al.* (2007) (Moilanen *et al.*, 2007) developed a 2D numerical bone model to investigate the impact of realistic bone geometry specimens on the ultrasonic cortical thickness evaluation using plate assumptions. Haiat *et al.* (2009) (Haiat *et al.*, 2009) (Haiat *et al.*,

2011) and Naili *et al.* (2010) (Naili *et al.*, 2010) assessed the effect of the heterogeneous nature of the cortical bone on the axial transmission response at 1 MHz in the framework of an anisotropic material by using 2D finite element (FE) simulation on a plate model. More recently, Chen *et al.* (2014) (Chen and Su, 2014) proposed a quantitative compensation for the effect of soft tissues based on an *in vitro* calibration, which facilitates the development of high-precision measurements of guided wave modes. Moreau *et al.* (2015) (Moreau *et al.*, 2014) introduced a modified method to predict the dispersion curves of an isotropic plate waveguide with a linearly varying thickness along the propagation direction. The method allowed the detection of enhanced and more accurate wavenumbers in the context of the formulation of inverse problems. Recently, Tran *et al.* (2018) (Tran *et al.*, 2018) performed a numerical study to evaluate the effect of cortical thickness, stiffness coefficient, and thickness of the overlying soft tissues on the responses of fundamental ultrasonic guided waves in the frequency domain, using a homogeneous transversely isotropic tri-layered plate model. However, despite the improvements achieved so far concerning the modeling of cortical bone, a more detailed numerical study remains needed to explain the influence of the cross-sectional curvature and distribution of properties in the radial direction on the propagation of the guided waves at low-frequency.

Low-frequency (typically below 200 kHz) axial transmission research (Egorov *et al.*, 2014; Kilappa *et al.*, 2011; Sarvazyan *et al.*, 2009; Tran *et al.*, 2015) demonstrated that ultrasonic guided waves are sensitive to changes in bone properties, such as the cortical thickness and the porosity. Muller *et al.* (2005) (Muller *et al.*, 2005) introduced low-frequency axial transmission as a promising method for assessing the cortical thickness. Tatarinov *et al.*

(2005) ([Tatarinov et al., 2005](#)) showed the potential of low-frequency (≈ 100 kHz) in assessing changes of deep underlying spongy layers in bovine tibia. More recently, Killappa *et al.* (2015) ([Killappa et al., 2015](#)) reported an improved performance of fundamental flexural modes used to assess cortical bone thickness using phase-delayed excitation at very low-frequency (50 kHz). Kassou *et al.* (2017) ([Kassou et al., 2017](#)) investigated the feasibility of dry point-contact transducers to infer the thickness of the *in-vivo* subjects in the 50-150 kHz frequency range. Operating at low-frequency offers the advantage of a lower attenuation as well as a reduced number of generated modes, which simplifies the signal analysis. In addition, low-frequency guided wave modes tend to achieve greater penetration depths due to the longer wavelengths as compared to high-frequency. As a consequence, the sensitivity to variations in intracortical bone properties could be improved when using low-frequencies, which is relevant to assess early stages of osteoporosis. Notwithstanding the remarkable level of realism that has been introduced in simulations over the past few years, the physical interaction between low-frequency guided waves and the cortical bone structure remains unexplored because the cortical bone has i) a complex cross-sectional geometry and ii) a heterogeneous distribution of material properties along the circumferential direction.

In a previous study by our group ([Pereira et al., 2017](#)), the propagation of guided waves in an irregular, multi-layer and heterogeneous bone cross-section modeled with anisotropic and viscoelastic material properties was investigated. The effect of the intracortical bone properties was then evaluated using the first arriving signal (FAS) velocity obtained from only five receivers. However, due to the interaction between several modes around the FAS, the velocity associated with the FAS was shown to be a poor discriminator of the

intracortical bone properties. The goal of this paper is therefore to determine the effect of intracortical bone properties on the phase velocity and cut-off frequency of low-frequency guided wave modes using an array of receivers, as well as a time-spatial frequency analysis technique instead of the FAS velocity. In order to do so, the excitation was generated by an emitter while the acquisition was performed in the time domain by an array of 20 receivers equally spaced along the bone surface. The bone ultrasonic responses were converted into the frequency domain using the two-dimensional Fourier transform (2D-FFT) in order to obtain the phase velocity and cut-off frequency of the propagating modes. Therefore, in this paper, the individual contribution of each viscoelastic coefficient and of the density was evaluated separately. The originality of this study when compared to our previous publication also lies in the application of the excitation source at different positions on the bone surface. This leads to a more comprehensive understanding of the physical interaction between the cortical bone curvature and the excitability of low-frequency guided wave modes.

The paper is organized as follows: section II.A introduces the semi-analytical finite element (SAFE) method used to simulate the signal of ultrasonic guided waves propagation in an arbitrary cross-section waveguide; section II.B describes the emitter/receivers configuration used to measure the ultrasonic responses on the bone surface in order to compute the *frequency vs. wavenumber* diagram. This section also describes the procedure employed to derive the phase velocity and cut-off frequency from the measured diagrams. In section II.C and II.D, the cross-sectional geometry of the cortical bone model and the methodology used to define the distribution of properties along the radial direction are presented respectively. Section II.E focuses on the description of the experimental work performed on bone phan-

toms that was done in order to verify the robustness and accuracy of the SAFE simulations. In section III.A, the effect of the excitation position and the physiopathological condition on the phase velocity and cut-off frequency are presented. In addition, the individual contribution of each viscoelastic coefficient and density on the overall physiopathological conditions are presented in the section III.B. Conclusion including potential clinical implications are presented in section IV.

II. MATERIAL AND METHODS

A. Simulating time domain signals

The SAFE method has recently been used to simulate the modal properties of the ultrasonic guided wave in bone systems (Nguyen and Naili, 2014; Nguyen *et al.*, 2017; Thakare *et al.*, 2017; Tran *et al.*, 2015). The computation of the time response from the modal properties of waveguides in an infinite arbitrary cross-section waveguide was previously described by Wilcox *et al.* (Wilcox *et al.*, 2001) and Loveday (Loveday, 2008). The resulting time domain signals consist in a linear superposition of the propagated response of each mode supported by the waveguide. A more detailed description of the implementation of the SAFE method can be found in the literature (Fan *et al.*, 2008; Predoi *et al.*, 2007), while the detailed procedure used to generate the time domain signal, can be found in a recently published study (Pereira *et al.*, 2017).

The dispersive response of a given mode excited from a point source in the out-of-plane direction (radial direction r in Fig. 1) can be calculated in the time domain using the

146 following expression (Wilcox *et al.*, 2001):

$$u(t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} F(\omega) E(\omega) H_0^{(1)}(k(\omega) z_0) e^{-i\omega t} d\omega \quad (1)$$

$$i = \sqrt{-1} \quad (2)$$

147 where $u(t)$ is the surface out-of-plane displacement for a given mode as a function of time t ,
 148 for an arbitrary propagation distance z_0 . The term $H_0^{(1)}$ is the zeroth-order Hankel function
 149 of the first kind and $F(\omega)$ is the frequency spectrum of a force input signal $u_{input}(t)$. The
 150 term $k(\omega)$ is the mode wavenumber, defined as a complex number by:

$$k(\omega) = k(\omega)_{real} + ik(\omega)_{imag} \quad (3)$$

151 where the real part denotes the propagating term and the imaginary part denotes the at-
 152 tenuation associated with the wavenumber. The term $E(\omega)$ is the frequency-dependent
 153 out-of-plane excitability of a given mode, and can be obtained with (Wilcox *et al.*, 2005):

$$E(\omega) = \frac{ik(\omega)\omega}{8} \left(\frac{u_{out}(\omega)^2}{P_z(\omega)} \right) \quad (4)$$

154 where $P_z(\omega)$ is the total power flow in the z -direction associated with the mode-shape as
 155 a function of the angular frequency, and $u_{out}(\omega)$ is the out-of-plane displacement on the
 156 surface of the mode-shape (at the position of excitation) as a function of the angular fre-
 157 quency. The dispersion curves of the waveguide (k , P_z and u_{out}) were obtained by solving the
 158 SAFE equations using the partial differential equation package in the Comsol Multiphysics
 159 and Matlab LiveLink environment. The complete time domain signal $u_{sum}(t)$ can then be

obtained with the linear superposition of n modes supported by the waveguide by:

$$u_{sum}(t) = \sum_{m=1}^n u^m(t) \quad (5)$$

where $u^m(t)$ is the surface out-of-plane displacement associated with the m -th mode and n is the total number of modes supported by the waveguide. The advantage of the SAFE method as compared to other simulation method such as 3D FE modeling is that with SAFE, only the cross-section is modeled and meshed, resulting in computationally efficient simulations, which may represent a significant gain of computation time. Furthermore, the dispersion curves are only calculated once using the SAFE method and then a number of time domain signals can be generated within minutes at minimal computational cost. The equivalent simulation using 3D FE would take hours, if not days, and use many times more memory (Pereira *et al.*, 2017).

The SAFE method can be considered as an accurate tool for simulating the ultrasonic guided wave response. The method has shown results in agreement with the 3D conventional FE simulations, as presented in our previous study (Pereira *et al.*, 2017). However, SAFE has the advantage of requiring reduced computing resources (e.g. 5.5 h and 2.5 Gb of memory against 4 h and 1.4 Gb of memory for the FE and SAFE methods receptively). Even though the advantages of SAFE in terms of resource requirements are modest, it is worth noticing that once the dispersion curves are calculated, the propagated responses can be computed for different excitation configurations without running the simulations again within seconds at minimal memory cost (<500 Mb).

B. Axial transmission configuration

The axial transmission configuration was modeled using the SAFE procedure described above. The waveguide consists of a 2D multilayer model composed of a viscoelastic heterogeneous solid layer surrounded by two viscoelastic homogeneous fluids. The solid layer corresponds to the cortical bone while the two fluids correspond respectively to bone marrow and soft tissues. Because of the characteristics of the SAFE simulations, the multilayer 2D medium is assumed to have an infinite length along the axial direction, resulting in a 3D equivalent medium, as shown in Fig.1.

FIG. 1. Schematic of the equivalent 3D axial transmission configuration used to perform the SAFE simulations.

The acoustic excitation was performed using a point source located in a plane denoted $z_0 = 0$ mm in contact with the upper part of the periosteal region of the bone cross-section. A 5-cycle Hann windowed toneburst centered at different frequencies between 20-85 kHz was used as the input waveform to perform the excitation. Eight different excitation positions

around the circumference of the cortical bone were investigated in this study. The positions, identified with the letters "A" to "H" in Fig. 2, were chosen based on the anatomical accessibility of a probe in a left human forearm. As a limitation of our study, the simulation of the excitability curves was made based on the displacement computed directly on the cortical bone tissue instead of on the overlaying soft tissue (see Fig. 2). The reason for this simplification is due to the higher mesh stability by considering the displacement in the inner elements (cortical tissue) instead of in the outer elements (soft tissue). When considering the displacement on the cortical bone tissue instead of the overlying soft tissue, a negligible change in the excitability curves was observed. However, for modes traveling with most of the energy in the cortical layer (which is the case for the majority of modes investigated in this study), the effect of this assumption on the excitability curves is acceptable. As the wavelengths of the modes of interest are long relative to the soft tissue thickness, a very limited number of modes are traveling with significant energy in the soft tissue. For the modes traveling in the cortical layer, the excitability curves have shown similar shape for both cases so as the relative difference and the rank between modes remained the same. As a consequence, no significant change on the relative variation obtained for the phase velocity and cut-off frequencies in the sensitivity study was observed with or without the soft tissue."

In reception, the out-of-plane displacement of the propagating waves were simulated at 20 positions (#1, #2, #3...#20 in Fig. 1), denoting an array of 20 receivers. The acquisition array was rotated in accordance with the excitation position around the circumference in order to keep the emitter and the receivers aligned. The acquisition elements were separated

FIG. 2. Cross-section of the middle of the radius bone used for the SAFE analysis. Different regions were defined along the radial direction: periosteal (a-b), middle (b-c) and endosteal (d-e) regions. The letters "A" to "H" show the eight different positions used to perform the excitation on the bone surface.

from one another by a distance of 3.5 mm, while varied from $z = 20$ mm to $z = 90$ mm. The signals obtained were processed using the 2D-FFT (Alleyne and Cawley, 1990) in order to calculate the *frequency vs. phase-velocity* diagram. Figure 3a shows typical received time domain signals obtained for an excitation frequency centered at 45 kHz. Figure 3b shows the *frequency vs. phase-velocity* diagram obtained after performing the 2D-FFT.

For each diagram, the phase velocity of the low-order mode (named V_{ph}) was measured by taking the peak of energy at the central frequency of the excitation waveform, as shown in the black square in Fig. 3b. Similarly, the cut-off frequency of the higher-order mode (named $F_{cut-off}$) was measured by taking the peak of energy at a reference velocity equal to 4000 m/s (see black circle in Fig. 3b). Those two features (V_{ph} and $F_{cut-off}$) were used to evaluate the performance of each configuration in term of assessing the variations in the intracortical bone properties.

FIG. 3. a) Simulated time domain signals obtained at each receiving position. b) Typical phase velocity vs. frequency diagram showing the intensity of the excited modes (dB scale) and the measured phase velocity V_{ph} and cut-off frequency $F_{cut-off}$.

It should be noted that the 2D tri-layer model investigated in this study did not account for the impacts of structural or geometrical irregularities along the axial direction since the waveguide was assumed to be constant and infinite. Such an assumption may change the signal of guided waves propagation in bone tissue. However, this limitation is balanced by the fact that the bone properties are typically measured in the central region of the bone, where the assumption of a constant profile along the axial direction holds because of the long wavelength in the frequency range used in this study.

C. Modeling cortical bone

The geometry of the waveguide was extracted from a slice of a human radius geometry. The radius geometry was provided by Sawbones ([Sawbone, 1997](#)), and was built based on μ CT images of the human radius. In this work, the middle third region, which is composed only of cortical bone, was used. Moreover, bone-marrow was added inside the cortical shell, and a 3.5 mm layer of soft-tissue was added outer the cortical shell. The tri-layer cross-section was imported into the COMSOL Multiphysics environment for execution of the SAFE analysis. The mesh was built using triangular elements with a maximum size of 0.5 mm, subject to a constraint of at least 15 elements per wavelength. The bone marrow and soft tissue were modeled as homogeneous viscoelastic fluids with the same properties used by Naili *et al.* ([Naili et al., 2010](#)) (bulk modulus=2.25e9 Pa, bulk viscosity=1.97 Pa.s and density=1000 kg/m³), while the cortical bone was modeled as a viscoelastic transversally isotropic material. Furthermore, cortical bone was divided into three regions along the radial direction: periosteal (a-b), middle (b-c) and endosteal (d-e) regions (see Fig. [2a](#)). Each

TABLE I. Summary of the properties used to model cortical bone. The real part of the stiffness coefficients were adapted from the experimental results obtained by (Sansalone *et al.*, 2010), while the imaginary parts were adapted from the results presented by (Naili *et al.*, 2010), using equation 7.

Viscoelastic	Periosteal	Middle	Endosteal	
coefficient**	(a-b)	(b-c)	Healthy (d*-e*)	Osteoporotic (d-e)**
$C_{11}(\text{Pa})$	$13.17\text{e}9 + 94.95fi$	$11.46\text{e}9 + 82.62fi$	$4.78\text{e}9 + 34.46fi$	$3.11\text{e}9 + 22.42fi$
$C_{13}(\text{Pa})$	$6.60\text{e}9 + 91.74fi$	$6.03\text{e}9 + 83.81fi$	$3.32\text{e}9 + 46.14fi$	$2.65\text{e}9 + 36.83fi$
$C_{33}(\text{Pa})$	$16.04\text{e}9 + 109.2fi$	$14.31\text{e}9 + 97.45fi$	$6.19\text{e}9 + 42.15fi$	$4.17\text{e}9 + 28.39fi$
$C_{44}(\text{Pa})$	$3.75\text{e}9 + 14.70fi$	$3.15\text{e}9 + 12.34fi$	$0.87\text{e}9 + 3.41fi$	$0.30\text{e}9 + 1.17fi$
$C_{66}(\text{Pa})$	$3.62\text{e}9 + 14.19fi$	$2.99\text{e}9 + 11.72fi$	$0.78\text{e}9 + 3.05fi$	$0.23\text{e}9 + 0.90fi$
Density (kg/m^3)	1850.00	1850.00	1299.04	1161.30

*The "Healthy" conditions were defined based on an arbitrary gain of 20% on the value of all stiffness

coefficients and density. **Adapted from the experimental results obtained by Sansalone *et al.* (Sansalone *et al.*, 2010)

region was defined with different values for the transversally isotropic stiffness coefficients
 and density ($C_{11}, C_{13}, C_{33}, C_{44}, C_{66}$ and ρ), as summarized in Table I. The value of each
 region was chosen based on a simplification of experimental results reported by Sansalone *et al.* (Sansalone *et al.*, 2010). A more detailed description of the distribution of the material
 properties along the radial direction is provided in the following section.

In the viscoelastic model used in this study, the elastic stiffness constants were written as complex numbers:

$$C_c = C_c + \eta_c f i \quad (6)$$

where the subscripts $c=11, 13, 33, 44$ and 66 , and f corresponds to the frequency. The real part of the stiffness constants is assumed to be independent of the frequency, while the imaginary part results from the use of the linear theory of viscoelasticity, assuming frequency-dependent losses in the waveguide. Because relevant data was lacking, the viscosity coefficients $\eta_{11}, \eta_{13}, \eta_{33}, \eta_{44}$ and η_{66} shown in Table 1 were defined based on the viscosity coefficients found in the literature. The coefficients reported in (Naili *et al.*, 2010) were adjusted by the ratio between the elastic coefficients used in this study and those reported in (Naili *et al.*, 2010), as given by:

$$\eta_c = \frac{C_c}{C_c^*} \cdot \eta_c^* \quad (7)$$

where the subscripts $c=11, 13, 33, 44$ and 66 , η_c^* and C_c^* are respectively the viscosity and elastic coefficients reported by Naili *et al.* (Naili *et al.*, 2010) for the cortical bone.

D. Heterogeneous distribution of bone properties

The heterogeneous nature of the biomechanical properties of cortical bone tissue was adapted from the experimental results obtained by Sansalone *et al.* (Sansalone *et al.*, 2010) using 3-D synchrotron micro-computed tomography images. In order to simplify the spatial distribution of each material property, the original values (Fig. 4) were adapted by choosing three different constant values for the periosteal (between a-b), middle (between b-c) and

endosteal (between d-e) regions. Moreover, a linear variation was assumed in the region between the middle and periosteal bone (between c-d). For the density, the same approach was applied based on porosity distribution reported in the study. The conversion was made considering a scale where 100% BV/TV (bone volume/ total volume) denotes a bone with a 1850 kg/m³ density and 0% of BV/TV denotes only bone marrow with a 1000 kg/m³ density. The constant values, as well as the linear variation of each material property, were chosen to minimize the gap between the original (Sansalone *et al.*, 2010) and simplified values. Figure 4 shows the original and simplified (healthy and osteoporotic) distributions for the coefficient C_{11} .

Two different physiopathological conditions, namely "Healthy" and "Osteoporotic", were defined for the endosteal region (between d*-e* and d-e in Fig. 4, respectively). The "Osteoporotic" condition was defined with the simplified distribution taken from the original data (Sansalone *et al.*, 2010). This assumption was made because the data were obtained from an osteoporotic 79-year old patient. The "Healthy" condition was then obtained by: i) arbitrarily increasing the values of all stiffness coefficients and density in the endosteal region by 20%, and ii) by reducing the length of the endosteal region by 20%.

As a limitation, the approach used to model the "Healthy" condition constitutes a simple way of modeling a slightly less degraded condition, which does not necessarily represents the condition associated with a healthy patient. However, since the properties are changing only in the endosteal in our sensitivity study, the magnitude of change between the "Healthy" and "Osteoporotic" condition can be considered very small, which may be similar to the magnitude of degradation associated to the early-stage osteoporosis. Furthermore, the het-

FIG. 4. Simplified distribution of the elastic coefficients and density varying with the distance from the periosteum for the healthy and osteoporotic condition.

erogeneous nature of cortical bone is based on homogenized local material properties, which are not easy to define due to the multi-scale nature of bone and its dependence on the microstructure at smaller scales. However, such discrepancies tend to be minimized by using low-frequency excitation associated with long wavelengths of the excited modes. Table I shows a summary of the "Healthy" and "Osteoporotic" elastic coefficients and densities used to defined the aforementioned scenarios.

III. RESULTS AND DISCUSSION

1. *Effect of excitation position*

Figures 5a to h show the normalized excitability curves (in the 10-100 kHz range) for an excitation performed at eight different positions on the bone surface (identified with letters "A" to "H" in Fig. 2). The normalized excitability predicts the relative amplitude expected to be observed by a receiver when an emitter at position "A" applies a load in the radial direction (see Fig. 2). The curves show the modal excitability of five modes with similar velocities ($m=1$, $m=2$, $m=3$, $m=4$, and $m=5$, marked with a solid line ellipse in Fig. 3). A notable difference in the modal excitability can be observed between each configuration, which is expected for non-symmetric waveguides such as the bone geometry modeled in this study. For practical applications, a position that can mainly excite a single mode is preferable in order to avoid interferences from other modes.

The modes $m=1$, $m=2$, $m=3$, $m=4$, and $m=5$ all present flexural-like mode shapes. For instance, Fig. 6a and 6b show the out-of-plane component of the fundamental flexural tube mode $F(1,1)$ and mode shape of $m=1$ respectively. Despite the similarities between the displacement fields and the mode order, a direct association to the mode shape of tubular waveguides is not possible due to the non-symmetric nature of the waveguide and its arbitrary geometry.

The modes $m=1$, $m=2$, and $m=3$ all show an excitability 10 dB higher than the other modes, as shown in the highlighted regions in Fig. 3a, 3c, and 3h, respectively. Furthermore, according to Figure 7, the modes $m=1$, $m=2$, and $m=3$ have a higher percentage of the total

FIG. 5. Normalized modal excitability obtained for the excitation performed at eight different position on the bone surface from 10 kHz to 100 kHz. The letters a) to h) are associated to the excitation position "A" to "H" shown in Fig. 2, respectively.

power flow concentrated in the endosteal region as compared to the other modes, which indicates that they are likely to be sensitive to changes in this region. Thus, the excitation at position "A", "C" and "H" were chosen as the most suitable positions to excite a dominant mode; they were therefore further investigated in this paper. Since the attenuation of the modes is expected to be small at very low frequencies (<100 kHz), it was not taken into account to choose the most suitable excitation and detection positions

2. Effect of physiopathological conditions

Table II shows the percentage variation in the phase velocity (V_{ph}) and cut-off frequency ($F_{cut-off}$) for the excitation performed at position "A", "C" and "H". The variations were calculated based on the difference computed between the two physiopathological conditions

FIG. 6. Out-of-plane displacement field (radial direction r in Fig. 1) at 40 kHz for: a) fundamental flexural mode $F(1,1)$ on a tubular geometry waveguide and b) mode $m=1$ on a bone-like geometry waveguide.

329 ("Healthy" and "Osteoporotic"). The simulations were performed using input waveforms
 330 centered at different frequencies between 20-85 kHz. This choice was made in order to cover
 331 a frequency range that has the potential to excite separately one of the five modes with
 332 similar velocities on the excitability curves in Fig. 5.

TABLE II. Total variation of the phase velocity V_{ph} and cut-off frequency $F_{cut-off}$ for the excitations performed at position "A", "C" and "H".

Phase velocity (V_{ph}) variation (%)	Position A	Position C	Position H
35 kHz	5.59	4.91	4.21
45 kHz	7.98	6.01	4.23
55 kHz	7.36	6.23	4.03
65 kHz	6.72	5.44	2.18
75 kHz	6.21	4.99	2.50
85 kHz	5.99	4.48	2.07
Cut-off frequency ($F_{cut-off}$) variation (%)	Position A	Position C	Position H
First cut-off frequency (20 kHz)	17.37	16.31	16.31
Second cut-off frequency (40 kHz)	9.42	12.15	11.23
Third cut-off frequency (70 kHz)	5.45	6.99	5.38

*The variations were calculated based on the difference computed between the "Healthy" and the "Osteoporotic" condition.

The sensitivity of phase velocity (V_{ph}) was strongly affected by the excitation position. The highest sensitivity was achieved at position "A", for all frequencies. On the other hand, the cut-off frequency ($F_{cut-off}$) was only slightly affected by the position of excitation, showing similar sensitivities for all positions. However, the cut-off frequency was strongly

affected by the frequency of excitation, showing a significantly higher sensitivity at the first cut-off frequency (20 kHz) as compared to the second (40 kHz) and third (70 kHz). Similarly, an increase in the sensitivity of the phase velocity (V_{ph}) was observed at lower frequencies, reaching a maximum variation at 45kHz, for all positions.

Figure 7 shows the power flow in the endosteal region (normalized to the total power flow) for modes $m=1$, $m=2$, $m=3$, $m=4$, and $m=5$, varying from 20 kHz to 100 kHz. The results show that mode $m=1$ has a higher power flow in the endosteal region as compared to the other modes, achieving approximately 23% of the total power in this region at 45 kHz. This explains the higher sensitivity obtained for the excitation performed at position "A", in which the mode $m=1$ excitability is 15 dB higher than the other modes (see Fig. 5a)

3. *Effect of stiffness coefficients and density*

Table III shows the percentage variation computed for the V_{ph} and $F_{cut-off}$ for each viscoelastic coefficient and density separately. The condition varied from "Healthy" to the "Osteoporotic". The excitation was performed at position "A" using central frequencies ranging between 20-85 kHz.

The sensitivity of V_{ph} showed a positive variation associated with coefficients C_{11} , C_{33} , C_{44} , C_{66} , while a negative variation was associated with the coefficient C_{13} and the density. Coefficient C_{44} showed the major contribution for most of the frequencies investigated (except at 35 kHz), accounting for approximately 50% of the variations in the degradation of all parameters taken together (identified as "Overall" in Table III). The reduction in sensitivity observed at 35 kHz was found to be associated with the sharp decrease in the power flow of

FIG. 7. Power flow in the endosteal region normalized by the total power flow for modes $m=1$, $m=2$, $m=3$, $m=4$, and $m=5$ from 20 kHz to 100 kHz.

mode $m=1$ seen in Fig. 7 at frequencies below 45kHz. In addition, it may be related to the reduction in excitability seen in mode $m=1$ at lower frequencies (see the excitability curves in Fig. 5a), which may increase the interference from other modes, and consequently, reduce the sensitivity.

TABLE III. Percentage variation computed for V_{ph} and $F_{cut-off}$ for each viscoelastic coefficient and density separately.

Phase velocity (V_{ph}) variation (%)	Position A						Overall
	C_{11}	C_{13}	C_{33}	C_{44}	C_{66}	Density	
35 kHz	0.32	-1.42	0.8	1.30	1.63	-0.79	5.59
45 kHz	0.37	-1.45	1.24	3.30	2.01	-0.611	7.98
55 kHz	0.31	-2.23	1.47	3.21	0.31	-0.72	7.36
65 kHz	0.36	-2.02	1.73	3.25	0.36	-0.62	6.72
75 kHz	0.16	-1.73	1.54	3.25	0.32	-0.71	6.21
85 kHz	0.08	-1.14	1.20	3.21	0.42	-0.67	5.99
Cut-off frequency ($F_{cut-off}$) variation (%)							
First cut-off frequency (20 kHz)	0	0	0	0	22.06	-2.34	17.37
Second cut-off frequency (40 kHz)	0	0	0.90	0.90	9.42	-0.90	9.42
Third cut-off frequency (70 kHz)	0.94	0.52	0.41	0.94	3.98	-1.05	5.45

For the sensitivity of $F_{cut-off}$, a positive variation associated with the coefficient C_{66} was observed, accounting for almost 100% of the "Overall" variation. Such a dependence could be used to assess the status of coefficient C_{66} separately.

IV. CONCLUSION

In this paper, the SAFE method was used to simulate the propagation of guided waves at low frequency (20-85 kHz) in a bone waveguide in the context of an axial transmission configuration. The SAFE method was shown to be an efficient tool for investigating the effect of intracortical properties on the propagation of guided waves since a number of configurations were simulated using reduced computing resources.

Considering the actual bone geometry, sixteen guided wave modes were found in the 10-100 kHz frequency range (Figure 3b). However, only some of the modes have shown enough excitability and consequently clinical interest. In addition, the excitability and sensitivity of each mode were shown to vary according to the frequency and position of excitation on the bone surface. The phase velocity showed a maximum sensitivity at 45 kHz, which would appear to be associated with the peak of energy (power flow) concentrated in the endosteal region at this frequency. The cut-off frequency was only slightly affected by the position of excitation, but strongly by the frequency of excitation. A maximum sensitivity was achieved for the first cut-off frequency at approximately 20 kHz.

The sensitivity of phase velocity was associated with the variation in the physiopathological conditions of all the coefficients and density, but mainly with C_{44} . The sensitivity of the cut-off frequency for its part was mainly associated with variation in the physiopathological conditions of the stiffness coefficient C_{66} , allowing the axial shear modulus to be assessed separately.

Thus, by selecting an adequate frequency and position of excitation, the low-frequency axial transmission was shown to be a promising method for assessing intracortical bone properties. The features identified in this study could be used as a metric to compare the similarity between experimental and numerical data. A cost function based on these features could then be implemented into an inversion scheme to retrieve reliable bone properties from experimental data at low frequency. A parameterized bone-like geometry model, instead of a plate or cylinder, could potentially reduce the errors commonly encountered with in-vivo and ex-vivo experiments. For instance, the model could potentially compensate geometrical variations between different patients that are unrelated to the bone quality, such as the outer diameter and the external shape of the bone. The method could, therefore, be applied to identify small changes associated with early-stage osteoporosis or gradual evolution of the bone condition over time. For these extreme cases, since the sensitivity associated to slightly different degraded conditions was found modest, a good basis for comparison (e.g., a baseline method or a large reference database) would have to be implemented in order to guarantee the robustness of the method.

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