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Effect of intracortical bone properties on the phase velocity and cut-off frequency of low-frequency guided wave modes (20-85 kHz)

Daniel Pereira,^{1, a)} Guillaume Haiat,^{2, b)} Julio Fernandes,^{3, c)} and Pierre Belanger^{1, d)}

¹*Department of Mechanical Engineering, École de technologie supérieure,
1100 Rue Notre-Dame O, Montreal, Quebec, H3C1K3 Canada*

²*CNRS, Laboratoire Modélisation et Simulation Multiechelle,
UMR CNRS 8208, 61 avenue du Général de Gaulle, Cretéil Cedex, 94010,
France*

³*Centre de recherche l'Hôpital du Sacré-Coeur de Montréal, 5400 Boul Gouin O,
Montreal, Quebec, H4J1C5 Canada*

1 The assessment of intracortical bone properties is of interest since early-stage osteo-
2 porosis is associated with resorption in the endosteal region. However, understanding
3 the interaction between ultrasonic guided waves and the cortical bone structure re-
4 mains challenging. The purpose of this work is to investigate the effect of intracortical
5 bone properties on the ultrasonic response obtained at low-frequency (<100 kHz) us-
6 ing an axial transmission configuration. The semi-analytical finite element method
7 was used to simulate the propagation of guided waves in a waveguide with realistic
8 geometry and material properties. An array of 20 receivers was used to calculate the
9 phase velocity and cut-off frequency of the excited modes using the 2D Fourier trans-
10 form. The results show that the position of the emitter around the circumference of
11 the bone is an important parameter to control since it can lead to variations of up to
12 10 dB in the amplitude of the transmitted modes. The cut-off frequency of the high
13 order modes was, however, only slightly affected by the circumferential position of the
14 emitter, and was sensitive mainly to the axial shear modulus. The phase velocity and
15 cut-off frequency in the 20-85 kHz range are promising parameters for the assessment
16 of intracortical properties.

^{a)} daniel.pereira.1@ens.etsmtl.ca; Also at: Centre de recherche l'Hôpital du Sacré-Coeur de Montréal, 5400

Boul Gouin O, Montreal, Quebec, H4J1C5 Canada

^{b)} Also at: Department of Mechanical Engineering, École de technologie supérieure, 1100 Rue Notre-Dame

O, Montreal, Quebec, H3C1K3 Canada

^{c)} Also at: Department of Surgery, Université de Montréal, 2900 Boulevard Edouard-Montpetit, Montreal,

Quebec, H3T1J4 Canada

^{d)} Also at: Centre de recherche l'Hôpital du Sacré-Coeur de Montréal, 5400 Boul Gouin O, Montreal, Quebec,

H4J1C5 Canada

17 I. INTRODUCTION

18 Osteoporosis is associated with a reduction of the bone mass and microarchitectural bone
19 deterioration, leading to a reduction in bone quality ([Giangregorio *et al.*, 2006](#); [Papaioannou
20 *et al.*, 2004](#); [Ström *et al.*, 2011](#)) and to an increase in fracture risk ([Consensus development
21 conference: prophylaxis and treatment of osteoporosis, 1991](#)). Endosteal resorption of cor-
22 tical bone results in an increase in bone porosity at the inner part of the cortical shell
23 and in a reduction of the cortical bone thickness ([Ritzel *et al.*, 1997](#)). Dual-energy X-ray
24 absorptiometry (DEXA), which is currently the gold standard for osteoporosis diagnostics
25 ([Kanis, 1994](#)), nonetheless suffers from several limitations ([Gluer, 2008](#); [Haba *et al.*, 2016](#)).
26 Quantitative ultrasound (QUS) was developed to assess bone quality and has the advan-
27 tage of being non-radiative, non-invasive and relatively cheap ([Gluer, 1997](#)). It, therefore,
28 holds promise as a rapid screening method in a clinical setting. The mechanical nature of
29 ultrasonic waves allows the technique to be used to retrieve the biomechanical properties
30 of bone tissue ([Kaufman and Einhorn, 1993](#)). QUS techniques can be used to go beyond
31 a simple estimation of the bone mineral density ([Nicholson, 2008](#)), which can be achieved
32 with DEXA and which is not sufficient to assess fracture risk ([Office of the Surgeon General,
33 2004](#)).

34 Bone QUS was initially developed in the context of trabecular bone characterization using
35 transverse transmission devices ([Stein *et al.*, 2013](#)). However, the investigation of cortical
36 bone ([Rico, 1997](#)) has attracted significant interest since about 80% of the skeleton is made
37 of cortical bone, which supports most of the body load, and is involved in osteoporotic

38 fractures (Seeley *et al.*, 1991). QUS using axial transmission techniques have mostly been
39 used to study the cortical bone quality (Haiat *et al.*, 2011). Cortical bones, such as the
40 radius or the tibia, constitute suitable waveguides for the propagation of ultrasonic guided
41 waves, as has been shown in various studies, including for instance (Gluer, 2008; 1997; Haba
42 *et al.*, 2016). Ultrasonic guided waves have the advantage of being sensitive to both the
43 mechanical and geometrical properties of cortical bone (Foiret *et al.*, 2014; Muller *et al.*,
44 2005; Rozental *et al.*, 2013).

45 Most studies focusing on the assessment of cortical bone properties using axial transmis-
46 sion consider plate or cylinder waveguides (as an approximation of the actual bone geometry)
47 to simplify the interpretation of ultrasonic responses obtained numerically or experimentally
48 (Kilappa *et al.*, 2015; Le *et al.*, 2010; Minonzio *et al.*, 2015; Xu *et al.*, 2016). Recently, more
49 comprehensive numerical approaches were explored with the aim of improving the realism of
50 the effect of cortical bone features on the axial transmission propagation expected *in vivo*.
51 Bossy *et al.* (2004) (Bossy *et al.*, 2004) performed three-dimensional (3D) finite difference
52 simulations on geometries derived from a human radius tomography image to evaluate the
53 effect of bone curvature, anisotropy, and micro-porosity on the first arriving signal (FAS)
54 velocity. However, the actual anatomical variations of bone biomechanical properties were
55 not taken into account and more emphasis was brought to the analysis of numerical simula-
56 tions performed with idealized 3D objects (tubes and semi-infinite hemicylinder geometries).
57 Moilanen *et al.* (2007) (Moilanen *et al.*, 2007) developed a 2D numerical bone model to in-
58 vestigate the impact of realistic bone geometry specimens on the ultrasonic cortical thickness
59 evaluation using plate assumptions. Haiat *et al.* (2009) (Haiat *et al.*, 2009) (Haiat *et al.*,

60 2011) and Naili *et al.* (2010) (Naili *et al.*, 2010) assessed the effect of the heterogeneous
61 nature of the cortical bone on the axial transmission response at 1 MHz in the framework of
62 an anisotropic material by using 2D finite element (FE) simulation on a plate model. More
63 recently, Chen *et al.* (2014) (Chen and Su, 2014) proposed a quantitative compensation for
64 the effect of soft tissues based on an *in vitro* calibration, which facilitates the development
65 of high-precision measurements of guided wave modes. Moreau *et al.* (2015) (Moreau *et al.*,
66 2014) introduced a modified method to predict the dispersion curves of an isotropic plate
67 waveguide with a linearly varying thickness along the propagation direction. The method
68 allowed the detection of enhanced and more accurate wavenumbers in the context of the for-
69 mulation of inverse problems. Recently, Tran *et al.* (2018) (Tran *et al.*, 2018) performed a
70 numerical study to evaluate the effect of cortical thickness, stiffness coefficient, and thickness
71 of the overlying soft tissues on the responses of fundamental ultrasonic guided waves in the
72 frequency domain, using a homogeneous transversely isotropic tri-layered plate model. How-
73 ever, despite the improvements achieved so far concerning the modeling of cortical bone, a
74 more detailed numerical study remains needed to explain the influence of the cross-sectional
75 curvature and distribution of properties in the radial direction on the propagation of the
76 guided waves at low-frequency.

77 Low-frequency (typically below 200 kHz) axial transmission research (Egorov *et al.*, 2014;
78 Kilappa *et al.*, 2011; Sarvazyan *et al.*, 2009; Tran *et al.*, 2015) demonstrated that ultrasonic
79 guided waves are sensitive to changes in bone properties, such as the cortical thickness
80 and the porosity. Muller *et al.* (2005) (Muller *et al.*, 2005) introduced low-frequency axial
81 transmission as a promising method for assessing the cortical thickness. Tatarinov *et al.*

82 (2005) ([Tatarinov et al., 2005](#)) showed the potential of low-frequency (≈ 100 kHz) in assessing
83 changes of deep underlying spongy layers in bovine tibia. More recently, Killappa *et al.*
84 (2015) ([Killappa et al., 2015](#)) reported an improved performance of fundamental flexural
85 modes used to assess cortical bone thickness using phase-delayed excitation at very low-
86 frequency (50 kHz). Kassou *et al.* (2017) ([Kassou et al., 2017](#)) investigated the feasibility
87 of dry point-contact transducers to infer the thickness of the *in-vivo* subjects in the 50-150
88 kHz frequency range. Operating at low-frequency offers the advantage of a lower attenuation
89 as well as a reduced number of generated modes, which simplifies the signal analysis. In
90 addition, low-frequency guided wave modes tend to achieve greater penetration depths due
91 to the longer wavelengths as compared to high-frequency. As a consequence, the sensitivity
92 to variations in intracortical bone properties could be improved when using low-frequencies,
93 which is relevant to assess early stages of osteoporosis. Notwithstanding the remarkable
94 level of realism that has been introduced in simulations over the past few years, the physical
95 interaction between low-frequency guided waves and the cortical bone structure remains
96 unexplored because the cortical bone has i) a complex cross-sectional geometry and ii) a
97 heterogeneous distribution of material properties along the circumferential direction.

98 In a previous study by our group ([Pereira et al., 2017](#)), the propagation of guided waves
99 in an irregular, multi-layer and heterogeneous bone cross-section modeled with anisotropic
100 and viscoelastic material properties was investigated. The effect of the intracortical bone
101 properties was then evaluated using the first arriving signal (FAS) velocity obtained from
102 only five receivers. However, due to the interaction between several modes around the
103 FAS, the velocity associated with the FAS was shown to be a poor discriminator of the

104 intracortical bone properties. The goal of this paper is therefore to determine the effect of
105 intracortical bone properties on the phase velocity and cut-off frequency of low-frequency
106 guided wave modes using an array of receivers, as well as a time-spatial frequency analysis
107 technique instead of the FAS velocity. In order to do so, the excitation was generated by an
108 emitter while the acquisition was performed in the time domain by an array of 20 receivers
109 equally spaced along the bone surface. The bone ultrasonic responses were converted into the
110 frequency domain using the two-dimensional Fourier transform (2D-FFT) in order to obtain
111 the phase velocity and cut-off frequency of the propagating modes. Therefore, in this paper,
112 the individual contribution of each viscoelastic coefficient and of the density was evaluated
113 separately. The originality of this study when compared to our previous publication also
114 lies in the application of the excitation source at different positions on the bone surface.
115 This leads to a more comprehensive understanding of the physical interaction between the
116 cortical bone curvature and the excitability of low-frequency guided wave modes.

117 The paper is organized as follows: section II.A introduces the semi-analytical finite ele-
118 ment (SAFE) method used to simulate the signal of ultrasonic guided waves propagation in
119 an arbitrary cross-section waveguide; section II.B describes the emitter/receivers configura-
120 tion used to measure the ultrasonic responses on the bone surface in order to compute the
121 *frequency vs. wavenumber* diagram. This section also describes the procedure employed to
122 derive the phase velocity and cut-off frequency from the measured diagrams. In section II.C
123 and II.D, the cross-sectional geometry of the cortical bone model and the methodology used
124 to define the distribution of properties along the radial direction are presented respectively.
125 Section II.E focuses on the description of the experimental work performed on bone phan-

126 toms that was done in order to verify the robustness and accuracy of the SAFE simulations.
127 In section III.A, the effect of the excitation position and the physiopathological condition on
128 the phase velocity and cut-off frequency are presented. In addition, the individual contribu-
129 tion of each viscoelastic coefficient and density on the overall physiopathological conditions
130 are presented in the section III.B. Conclusion including potential clinical implications are
131 presented in section IV.

132 II. MATERIAL AND METHODS

133 A. Simulating time domain signals

134 The SAFE method has recently been used to simulate the modal properties of the ultra-
135 sonic guided wave in bone systems (Nguyen and Naili, 2014; Nguyen *et al.*, 2017; Thakare
136 *et al.*, 2017; Tran *et al.*, 2015). The computation of the time response from the modal
137 properties of waveguides in an infinite arbitrary cross-section waveguide was previously de-
138 scribed by Wilcox *et al.* (Wilcox *et al.*, 2001) and Loveday (Loveday, 2008). The resulting
139 time domain signals consist in a linear superposition of the propagated response of each
140 mode supported by the waveguide. A more detailed description of the implementation of
141 the SAFE method can be found in the literature (Fan *et al.*, 2008; Predoi *et al.*, 2007), while
142 the detailed procedure used to generate the time domain signal, can be found in a recently
143 published study (Pereira *et al.*, 2017).

144 The dispersive response of a given mode excited from a point source in the out-of-plane
145 direction (radial direction r in Fig. 1) can be calculated in the time domain using the

146 following expression (Wilcox *et al.*, 2001):

$$u(t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} F(\omega) E(\omega) H_0^{(1)}(k(\omega) z_0) e^{-i\omega t} d\omega \quad (1)$$

$$i = \sqrt{-1} \quad (2)$$

147 where $u(t)$ is the surface out-of-plane displacement for a given mode as a function of time t ,
 148 for an arbitrary propagation distance z_0 . The term $H_0^{(1)}$ is the zeroth-order Hankel function
 149 of the first kind and $F(\omega)$ is the frequency spectrum of a force input signal $u_{input}(t)$. The
 150 term $k(\omega)$ is the mode wavenumber, defined as a complex number by:

$$k(\omega) = k(\omega)_{real} + ik(\omega)_{imag} \quad (3)$$

151 where the real part denotes the propagating term and the imaginary part denotes the at-
 152 tenuation associated with the wavenumber. The term $E(\omega)$ is the frequency-dependent
 153 out-of-plane excitability of a given mode, and can be obtained with (Wilcox *et al.*, 2005):

$$E(\omega) = \frac{ik(\omega)\omega}{8} \left(\frac{u_{out}(\omega)^2}{P_z(\omega)} \right) \quad (4)$$

154 where $P_z(\omega)$ is the total power flow in the z -direction associated with the mode-shape as
 155 a function of the angular frequency, and $u_{out}(\omega)$ is the out-of-plane displacement on the
 156 surface of the mode-shape (at the position of excitation) as a function of the angular fre-
 157 quency. The dispersion curves of the waveguide (k , P_z and u_{out}) were obtained by solving the
 158 SAFE equations using the partial differential equation package in the Comsol Multiphysics
 159 and Matlab LiveLink environment. The complete time domain signal $u_{sum}(t)$ can then be

160 obtained with the linear superposition of n modes supported by the waveguide by:

$$u_{sum}(t) = \sum_{m=1}^n u^m(t) \quad (5)$$

161 where $u^m(t)$ is the surface out-of-plane displacement associated with the m -th mode and n is
 162 the total number of modes supported by the waveguide. The advantage of the SAFE method
 163 as compared to other simulation method such as 3D FE modeling is that with SAFE, only
 164 the cross-section is modeled and meshed, resulting in computationally efficient simulations,
 165 which may represent a significant gain of computation time. Furthermore, the dispersion
 166 curves are only calculated once using the SAFE method and then a number of time domain
 167 signals can be generated within minutes at minimal computational cost. The equivalent
 168 simulation using 3D FE would take hours, if not days, and use many times more memory
 169 ([Pereira *et al.*, 2017](#)).

170 The SAFE method can be considered as an accurate tool for simulating the ultrasonic
 171 guided wave response. The method has shown results in agreement with the 3D conventional
 172 FE simulations, as presented in our previous study ([Pereira *et al.*, 2017](#)). However, SAFE has
 173 the advantage of requiring reduced computing resources (e.g. 5.5 h and 2.5 Gb of memory
 174 against 4 h and 1.4 Gb of memory for the FE and SAFE methods receptively). Even though
 175 the advantages of SAFE in terms of resource requirements are modest, it is worth noticing
 176 that once the dispersion curves are calculated, the propagated responses can be computed
 177 for different excitation configurations without running the simulations again within seconds
 178 at minimal memory cost (<500 Mb).

179 **B. Axial transmission configuration**

180 The axial transmission configuration was modeled using the SAFE procedure described
 181 above. The waveguide consists of a 2D multilayer model composed of a viscoelastic het-
 182 erogeneous solid layer surrounded by two viscoelastic homogeneous fluids. The solid layer
 183 corresponds to the cortical bone while the two fluids correspond respectively to bone marrow
 184 and soft tissues. Because of the characteristics of the SAFE simulations, the multilayer 2D
 185 medium is assumed to have an infinite length along the axial direction, resulting in a 3D
 186 equivalent medium, as shown in Fig.1.

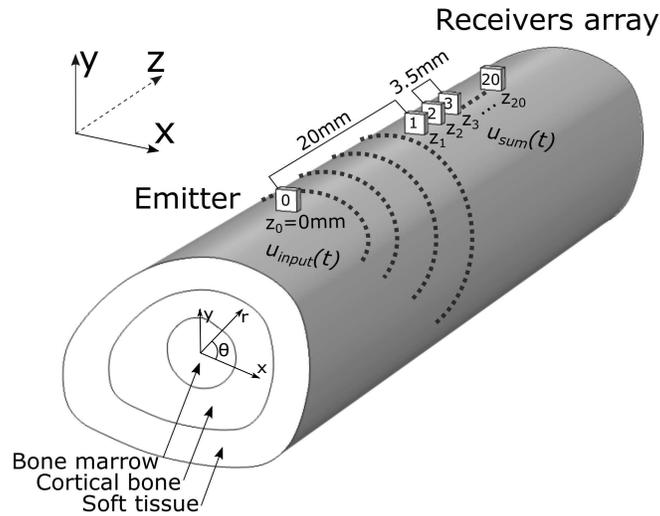


FIG. 1. Schematic of the equivalent 3D axial transmission configuration used to perform the SAFE simulations.

187 The acoustic excitation was performed using a point source located in a plane denoted
 188 $z_0 = 0$ mm in contact with the upper part of the periosteal region of the bone cross-section.
 189 A 5-cycle Hann windowed toneburst centered at different frequencies between 20-85 kHz was
 190 used as the input waveform to perform the excitation. Eight different excitation positions

191 around the circumference of the cortical bone were investigated in this study. The positions,
192 identified with the letters "A" to "H" in Fig. 2, were chosen based on the anatomical
193 accessibility of a probe in a left human forearm. As a limitation of our study, the simulation
194 of the excitability curves was made based on the displacement computed directly on the
195 cortical bone tissue instead of on the overlaying soft tissue (see Fig. 2). The reason for this
196 simplification is due to the higher mesh stability by considering the displacement in the inner
197 elements (cortical tissue) instead of in the outer elements (soft tissue). When considering
198 the displacement on the cortical bone tissue instead of the overlying soft tissue, a negligible
199 change in the excitability curves was observed. However, for modes traveling with most of
200 the energy in the cortical layer (which is the case for the majority of modes investigated in
201 this study), the effect of this assumption on the excitability curves is acceptable. As the
202 wavelengths of the modes of interest are long relative to the soft tissue thickness, a very
203 limited number of modes are traveling with significant energy in the soft tissue. For the
204 modes traveling in the cortical layer, the excitability curves have shown similar shape for
205 both cases so as the relative difference and the rank between modes remained the same. As
206 a consequence, no significant change on the relative variation obtained for the phase velocity
207 and cut-off frequencies in the sensitivity study was observed with or without the soft tissue."

208

210 In reception, the out-of-plane displacement of the propagating waves were simulated at
211 20 positions (#1, #2, #3...#20 in Fig. 1), denoting an array of 20 receivers. The acquisition
212 array was rotated in accordance with the excitation position around the circumference in
213 order to keep the emitter and the receivers aligned. The acquisition elements were separated

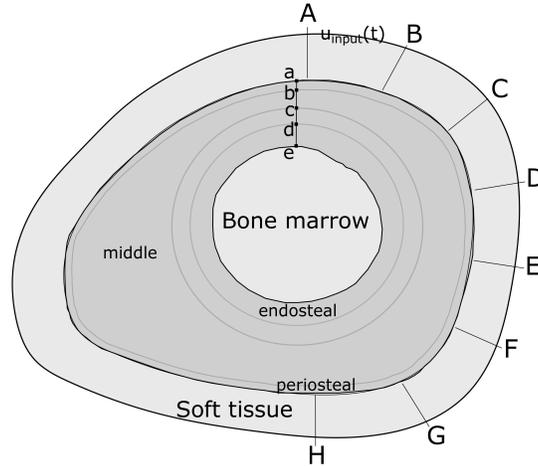


FIG. 2. Cross-section of the middle of the radius bone used for the SAFE analysis. Different regions were defined along the radial direction: periosteal (a-b), middle (b-c) and endosteal (d-e) regions. The letters "A" to "H" show the eight different positions used to perform the excitation on the bone surface.

214 from one another by a distance of 3.5 mm, while varied from $z = 20$ mm to $z = 90$ mm. The
 215 signals obtained were processed using the 2D-FFT (Alleyne and Cawley, 1990) in order to
 216 calculate the *frequency vs. phase-velocity* diagram. Figure 3a shows typical received time
 217 domain signals obtained for an excitation frequency centered at 45 kHz. Figure 3b shows
 218 the *frequency vs. phase-velocity* diagram obtained after performing the 2D-FFT.

219 For each diagram, the phase velocity of the low-order mode (named V_{ph}) was measured
 220 by taking the peak of energy at the central frequency of the excitation waveform, as shown
 221 in the black square in Fig. 3b. Similarly, the cut-off frequency of the higher-order mode
 222 (named $F_{cut-off}$) was measured by taking the peak of energy at a reference velocity equal
 223 to 4000 m/s (see black circle in Fig. 3b). Those two features (V_{ph} and $F_{cut-off}$) were used
 224 to evaluate the performance of each configuration in term of assessing the variations in the
 225 intracortical bone properties.

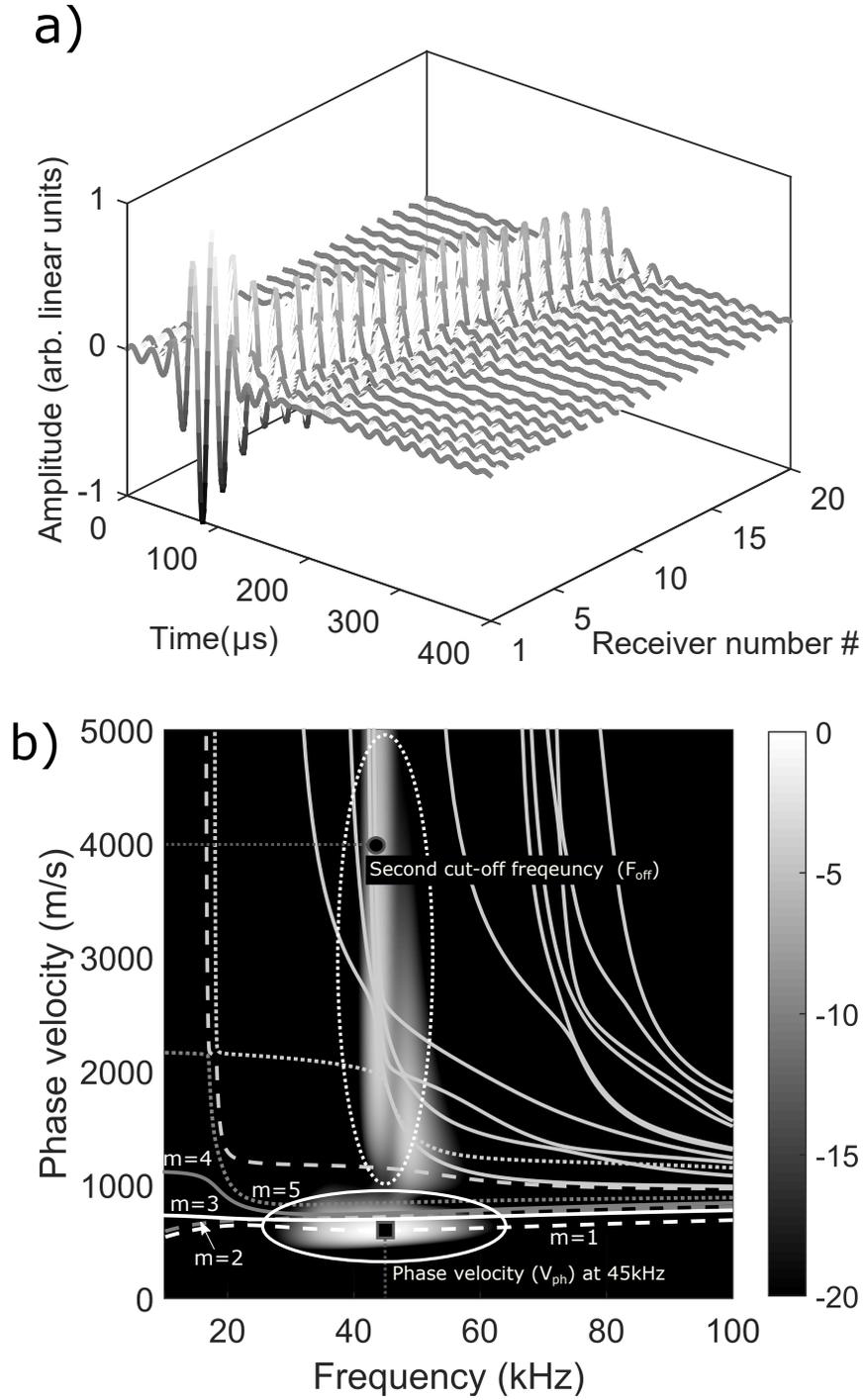


FIG. 3. a) Simulated time domain signals obtained at each receiving position. b) Typical phase velocity vs. frequency diagram showing the intensity of the excited modes (dB scale) and the measured phase velocity V_{ph} and cut-off frequency $F_{cut-off}$.

226 It should be noted that the 2D tri-layer model investigated in this study did not account
227 for the impacts of structural or geometrical irregularities along the axial direction since the
228 waveguide was assumed to be constant and infinite. Such an assumption may change the
229 signal of guided waves propagation in bone tissue. However, this limitation is balanced by
230 the fact that the bone properties are typically measured in the central region of the bone,
231 where the assumption of a constant profile along the axial direction holds because of the
232 long wavelength in the frequency range used in this study.

233 C. Modeling cortical bone

234 The geometry of the waveguide was extracted from a slice of a human radius geometry.
235 The radius geometry was provided by Sawbones ([Sawbone, 1997](#)), and was built based on
236 μ CT images of the human radius. In this work, the middle third region, which is composed
237 only of cortical bone, was used. Moreover, bone-marrow was added inside the cortical shell,
238 and a 3.5 mm layer of soft-tissue was added outer the cortical shell. The tri-layer cross-
239 section was imported into the COMSOL Multiphysics environment for execution of the
240 SAFE analysis. The mesh was built using triangular elements with a maximum size of 0.5
241 mm, subject to a constraint of at least 15 elements per wavelength. The bone marrow and
242 soft tissue were modeled as homogeneous viscoelastic fluids with the same properties used
243 by Naili *et al.* ([Naili *et al.*, 2010](#)) (bulk modulus=2.25e9 Pa, bulk viscosity=1.97 Pa.s and
244 density=1000 kg/m³), while the cortical bone was modeled as a viscoelastic transversally
245 isotropic material. Furthermore, cortical bone was divided into three regions along the radial
246 direction: periosteal (a-b), middle (b-c) and endosteal (d-e) regions (see Fig. [2a](#)). Each

TABLE I. Summary of the properties used to model cortical bone. The real part of the stiffness coefficients were adapted from the experimental results obtained by (Sansalone *et al.*, 2010), while the imaginary parts were adapted from the results presented by (Naili *et al.*, 2010), using equation 7.

Viscoelastic coefficient**	Periosteal (a-b)	Middle (b-c)	Endosteal	
			Healthy (d*-e*)	Osteoporotic (d-e)**
C_{11} (Pa)	13.17e9 + 94.95 <i>fi</i>	11.46e9 + 82.62 <i>fi</i>	4.78e9 + 34.46 <i>fi</i>	3.11e9 + 22.42 <i>fi</i>
C_{13} (Pa)	6.60e9 + 91.74 <i>fi</i>	6.03e9 + 83.81 <i>fi</i>	3.32e9 + 46.14 <i>fi</i>	2.65e9 + 36.83 <i>fi</i>
C_{33} (Pa)	16.04e9 + 109.2 <i>fi</i>	14.31e9 + 97.45 <i>fi</i>	6.19e9 + 42.15 <i>fi</i>	4.17e9 + 28.39 <i>fi</i>
C_{44} (Pa)	3.75e9 + 14.70 <i>fi</i>	3.15e9 + 12.34 <i>fi</i>	0.87e9 + 3.41 <i>fi</i>	0.30e9 + 1.17 <i>fi</i>
C_{66} (Pa)	3.62e9 + 14.19 <i>fi</i>	2.99e9 + 11.72 <i>fi</i>	0.78e9 + 3.05 <i>fi</i>	0.23e9 + 0.90 <i>fi</i>
Density (kg/m ³)	1850.00	1850.00	1299.04	1161.30

*The "Healthy" conditions were defined based on an arbitrary gain of 20% on the value of all stiffness coefficients and density. **Adapted from the experimental results obtained by Sansalone *et al.* (Sansalone *et al.*, 2010)

247 region was defined with different values for the transversally isotropic stiffness coefficients
 248 and density ($C_{11}, C_{13}, C_{33}, C_{44}, C_{66}$ and ρ), as summarized in Table I. The value of each
 249 region was chosen based on a simplification of experimental results reported by Sansalone *et*
 250 *al.* (Sansalone *et al.*, 2010). A more detailed description of the distribution of the material
 251 properties along the radial direction is provided in the following section.

252 In the viscoelastic model used in this study, the elastic stiffness constants were written
 253 as complex numbers:

$$C_c = C_c + \eta_c f i \quad (6)$$

254 where the subscripts $c=11, 13, 33, 44$ and 66 , and f corresponds to the frequency. The
 255 real part of the stiffness constants is assumed to be independent of the frequency, while
 256 the imaginary part results from the use of the linear theory of viscoelasticity, assuming
 257 frequency-dependent losses in the waveguide. Because relevant data was lacking, the viscos-
 258 ity coefficients $\eta_{11}, \eta_{13}, \eta_{33}, \eta_{44}$ and η_{66} shown in Table 1 were defined based on the viscosity
 259 coefficients found in the literature. The coefficients reported in (Naili *et al.*, 2010) were
 260 adjusted by the ratio between the elastic coefficients used in this study and those reported
 261 in (Naili *et al.*, 2010), as given by:

$$\eta_c = \frac{C_c}{C_c^*} \cdot \eta_c^* \quad (7)$$

262 where the subscripts $c=11, 13, 33, 44$ and 66 , η_c^* and C_c^* are respectively the viscosity and
 263 elastic coefficients reported by Naili *et al.* (Naili *et al.*, 2010) for the cortical bone.

264 **D. Heterogeneous distribution of bone properties**

265 The heterogeneous nature of the biomechanical properties of cortical bone tissue was
 266 adapted from the experimental results obtained by Sansalone *et al.* (Sansalone *et al.*, 2010)
 267 using 3-D synchrotron micro-computed tomography images. In order to simplify the spatial
 268 distribution of each material property, the original values (Fig. 4) were adapted by choosing
 269 three different constant values for the periosteal (between a-b), middle (between b-c) and

270 endosteal (between d-e) regions. Moreover, a linear variation was assumed in the region
271 between the middle and periosteal bone (between c-d). For the density, the same approach
272 was applied based on porosity distribution reported in the study. The conversion was made
273 considering a scale where 100% BV/TV (bone volume/ total volume) denotes a bone with
274 a 1850 kg/m^3 density and 0% of BV/TV denotes only bone marrow with a 1000 kg/m^3
275 density. The constant values, as well as the linear variation of each material property, were
276 chosen to minimize the gap between the original (Sansalone *et al.*, 2010) and simplified
277 values. Figure 4 shows the original and simplified (healthy and osteoporotic) distributions
278 for the coefficient C_{11} .

279 Two different physiopathological conditions, namely "Healthy" and "Osteoporotic", were
280 defined for the endosteal region (between d*-e* and d-e in Fig. 4, respectively). The
281 "Osteoporotic" condition was defined with the simplified distribution taken from the original
282 data (Sansalone *et al.*, 2010). This assumption was made because the data were obtained
283 from an osteoporotic 79-year old patient. The "Healthy" condition was then obtained by:
284 i) arbitrarily increasing the values of all stiffness coefficients and density in the endosteal
285 region by 20%, and ii) by reducing the length of the endosteal region by 20%.

286 As a limitation, the approach used to model the "Healthy" condition constitutes a simple
287 way of modeling a slightly less degraded condition, which does not necessarily represents the
288 condition associated with a healthy patient. However, since the properties are changing only
289 in the endosteal in our sensitivity study, the magnitude of change between the "Healthy"
290 and "Osteoporotic" condition can be considered very small, which may be similar to the
291 magnitude of degradation associated to the early-stage osteoporosis. Furthermore, the het-

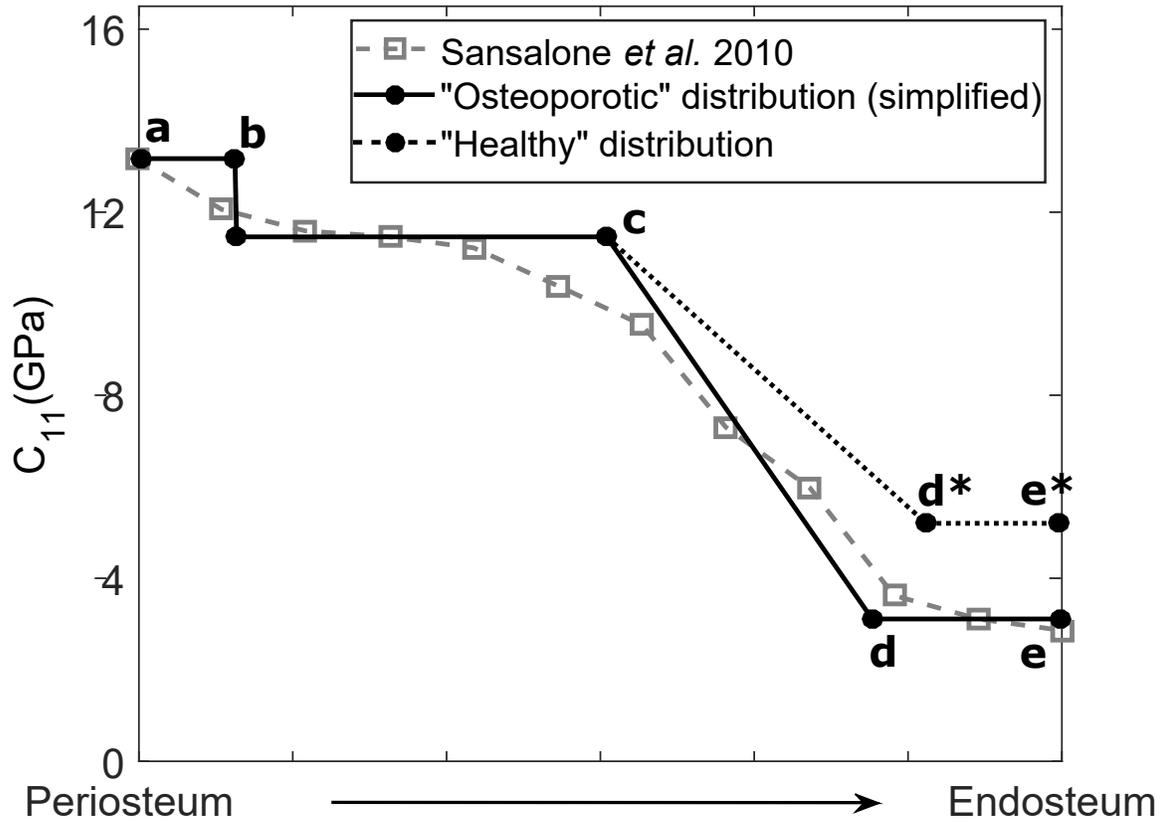


FIG. 4. Simplified distribution of the elastic coefficients and density varying with the distance from the periosteum for the healthy and osteoporotic condition.

292 erogeneous nature of cortical bone is based on homogenized local material properties, which
 293 are not easy to define due to the multi-scale nature of bone and its dependence on the
 294 microstructure at smaller scales. However, such discrepancies tend to be minimized by us-
 295 ing low-frequency excitation associated with long wavelengths of the excited modes. Table
 296 I shows a summary of the "Healthy" and "Osteoporotic" elastic coefficients and densities
 297 used to defined the aforementioned scenarios.

298 III. RESULTS AND DISCUSSION

299 1. *Effect of excitation position*

300 Figures 5a to h show the normalized excitability curves (in the 10-100 kHz range) for an
301 excitation performed at eight different positions on the bone surface (identified with letters
302 "A" to "H" in Fig. 2). The normalized excitability predicts the relative amplitude expected
303 to be observed by a receiver when an emitter at position "A" applies a load in the radial
304 direction (see Fig. 2). The curves show the modal excitability of five modes with similar
305 velocities ($m=1$, $m=2$, $m=3$, $m=4$, and $m=5$, marked with a solid line ellipse in Fig. 3).
306 A notable difference in the modal excitability can be observed between each configuration,
307 which is expected for non-symmetric waveguides such as the bone geometry modeled in
308 this study. For practical applications, a position that can mainly excite a single mode is
309 preferable in order to avoid interferences from other modes.

310 The modes $m=1$, $m=2$, $m=3$, $m=4$, and $m=5$ all present flexural-like mode shapes. For
311 instance, Fig. 6a and 6b show the out-of-plane component of the fundamental flexural
312 tube mode $F(1,1)$ and mode shape of $m=1$ respectively. Despite the similarities between
313 the displacement fields and the mode order, a direct association to the mode shape of
314 tubular waveguides is not possible due to the non-symmetric nature of the waveguide and
315 its arbitrary geometry

316 The modes $m=1$, $m=2$, and $m=3$ all show an excitability 10 dB higher than the other
317 modes, as shown in the highlighted regions in Fig. 3a, 3c, and 3h, respectively. Furthermore,
318 according to Figure 7, the modes $m=1$, $m=2$, and $m=3$ have a higher percentage of the total

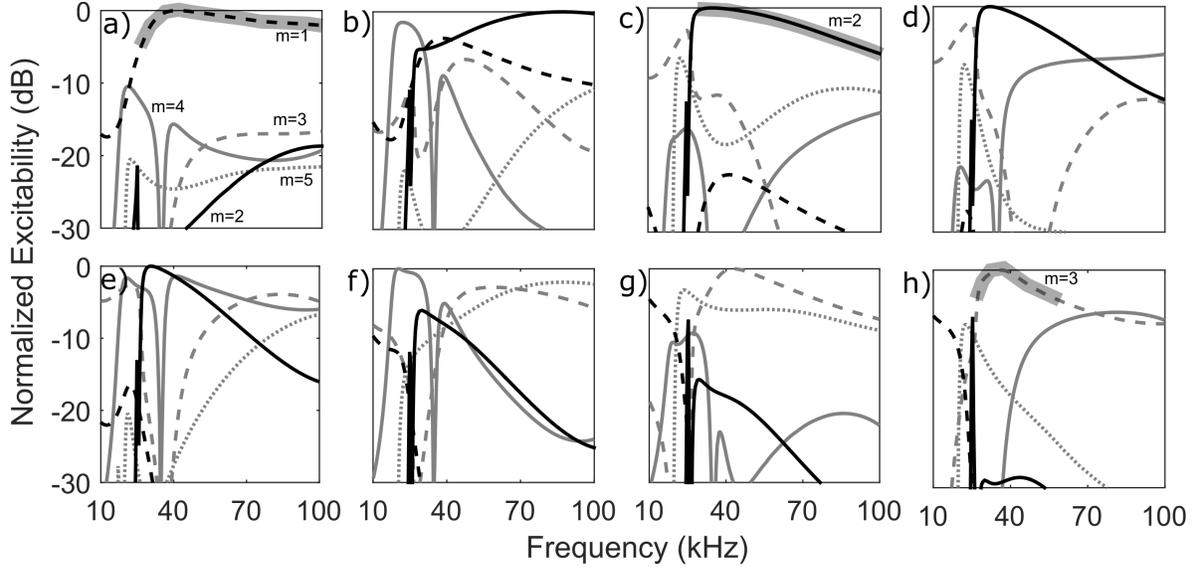


FIG. 5. Normalized modal excitability obtained for the excitation performed at eight different position on the bone surface from 10 kHz to 100 kHz. The letters a) to h) are associated to the excitation position "A" to "H" shown in Fig. 2, respectively.

319 power flow concentrated in the endosteal region as compared to the other modes, which
 320 indicates that they are likely to be sensitive to changes in this region. Thus, the excitation
 321 at position "A", "C" and "H" were chosen as the most suitable positions to excite a dominant
 322 mode; they were therefore further investigated in this paper. Since the attenuation of the
 323 modes is expected to be small at very low frequencies (<100 kHz), it was not taken into
 324 account to choose the most suitable excitation and detection positions

325 2. Effect of physiopathological conditions

326 Table II shows the percentage variation in the phase velocity (V_{ph}) and cut-off frequency
 327 ($F_{cut-off}$) for the excitation performed at position "A", "C" and "H". The variations was
 328 calculated based on the difference computed between the two physiopathological conditions

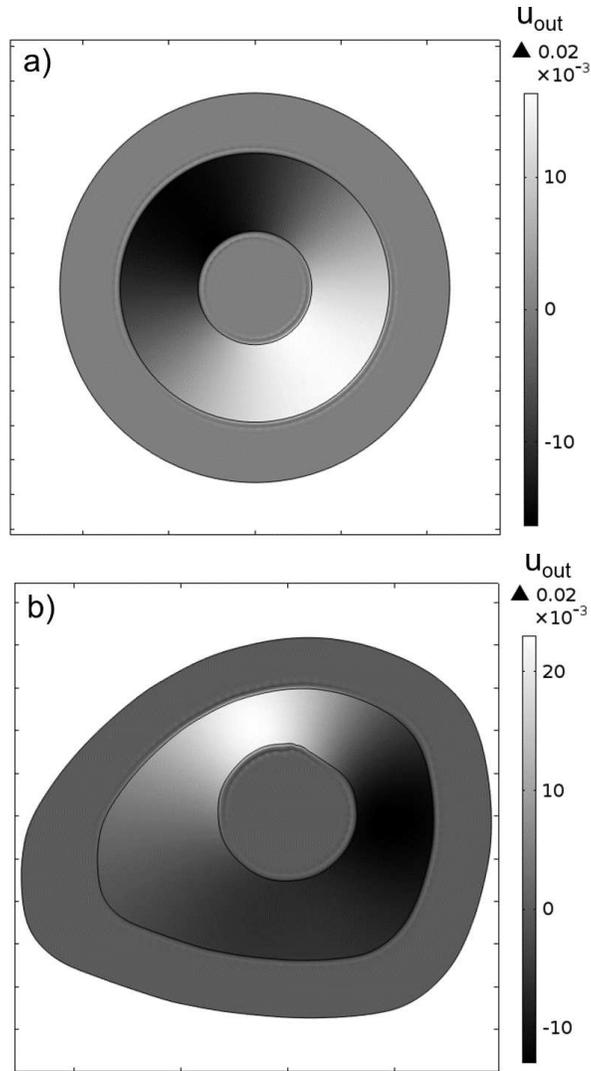


FIG. 6. Out-of-plane displacement field (radial direction r in Fig. 1) at 40 kHz for: a) fundamental flexural mode $F(1,1)$ on a tubular geometry waveguide and b) mode $m=1$ on a bone-like geometry waveguide.

329 ("Healthy" and "Osteoporotic"). The simulations were performed using input waveforms
 330 centered at different frequencies between 20-85 kHz. This choice was made in order to cover
 331 a frequency range that has the potential to excite separately one of the five modes with
 332 similar velocities on the excitability curves in Fig. 5.

TABLE II. Total variation of the phase velocity V_{ph} and cut-off frequency $F_{cut-off}$ for the excitations performed at position "A", "C" and "H".

Phase velocity (V_{ph}) variation (%)	Position A	Position C	Position H
35 kHz	5.59	4.91	4.21
45 kHz	7.98	6.01	4.23
55 kHz	7.36	6.23	4.03
65 kHz	6.72	5.44	2.18
75 kHz	6.21	4.99	2.50
85 kHz	5.99	4.48	2.07

Cut-off frequency ($F_{cut-off}$) variation (%)	Position A	Position C	Position H
First cut-off frequency (20 kHz)	17.37	16.31	16.31
Second cut-off frequency (40 kHz)	9.42	12.15	11.23
Third cut-off frequency (70 kHz)	5.45	6.99	5.38

*The variations were calculated based on the difference computed between the "Healthy" and the "Osteoporotic" condition.

333 The sensitivity of phase velocity (V_{ph}) was strongly affected by the excitation position.
 334 The highest sensitivity was achieved at position "A", for all frequencies. On the other hand,
 335 the cut-off frequency ($F_{cut-off}$) was only slightly affected by the position of excitation,
 336 showing similar sensitivities for all positions. However, the cut-off frequency was strongly

337 affected by the frequency of excitation, showing a significantly higher sensitivity at the first
 338 cut-off frequency (20 kHz) as compared to the second (40 kHz) and third (70 kHz). Similarly,
 339 an increase in the sensitivity of the phase velocity (V_{ph}) was observed at lower frequencies,
 340 reaching a maximum variation at 45kHz, for all positions.

341 Figure 7 shows the power flow in the endosteal region (normalized to the total power
 342 flow) for modes $m=1$, $m=2$, $m=3$, $m=4$, and $m=5$, varying from 20 kHz to 100 kHz. The
 343 results show that mode $m=1$ has a higher power flow in the endosteal region as compared to
 344 the other modes, achieving approximately 23% of the total power in this region at 45 kHz.
 345 This explains the higher sensitivity obtained for the excitation performed at position "A",
 346 in which the mode $m=1$ excitability is 15 dB higher than the other modes (see Fig. 5a)

347 *3. Effect of stiffness coefficients and density*

348 Table III shows the percentage variation computed for the V_{ph} and $F_{cut-off}$ for each
 349 viscoelastic coefficient and density separately. The condition varied from "Healthy" to the
 350 "Osteoporotic". The excitation was performed at position "A" using central frequencies
 351 ranging between 20-85 kHz.

352 The sensitivity of V_{ph} showed a positive variation associated with coefficients C_{11} , C_{33} ,
 353 C_{44} , C_{66} , while a negative variation was associated with the coefficient C_{13} and the density.
 354 Coefficient C_{44} showed the major contribution for most of the frequencies investigated (ex-
 355 cept at 35 kHz), accounting for approximately 50% of the variations in the degradation of all
 356 parameters taken together (identified as "Overall" in Table III). The reduction in sensitivity
 357 observed at 35 kHz was found to be associated with the sharp decrease in the power flow of

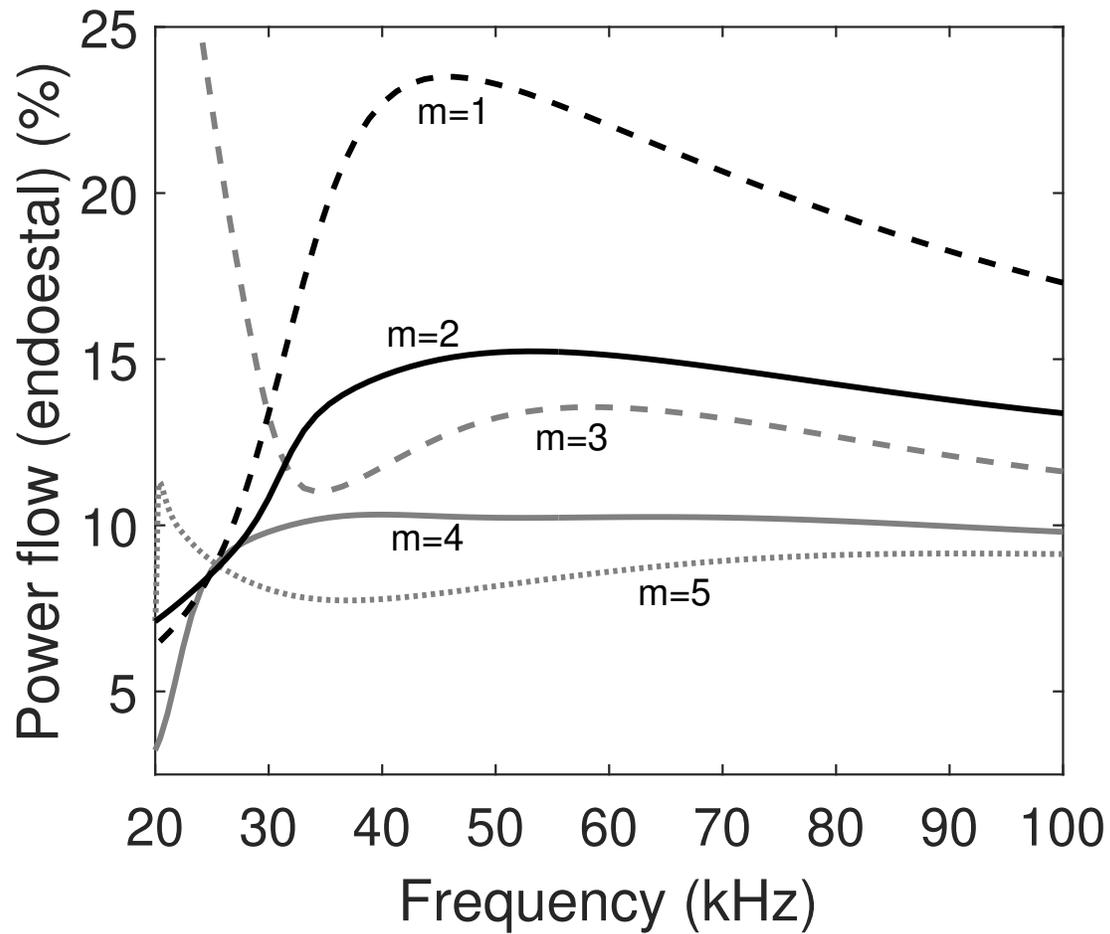


FIG. 7. Power flow in the endosteal region normalized by the total power flow for modes $m=1$, $m=2$, $m=3$, $m=4$, and $m=5$ from 20 kHz to 100 kHz.

358 mode $m=1$ seen in Fig. 7 at frequencies below 45kHz. In addition, it may be related to the
 359 reduction in excitability seen in mode $m=1$ at lower frequencies (see the excitability curves
 360 in Fig. 5a), which may increase the interference from other modes, and consequently, reduce
 361 the sensitivity.

TABLE III. Percentage variation computed for V_{ph} and $F_{cut-off}$ for each viscoelastic coefficient and density separately.

Phase velocity (V_{ph}) variation (%)	Position A						Overall
	C_{11}	C_{13}	C_{33}	C_{44}	C_{66}	Density	
35 kHz	0.32	-1.42	0.8	1.30	1.63	-0.79	5.59
45 kHz	0.37	-1.45	1.24	3.30	2.01	-0.611	7.98
55 kHz	0.31	-2.23	1.47	3.21	0.31	-0.72	7.36
65 kHz	0.36	-2.02	1.73	3.25	0.36	-0.62	6.72
75 kHz	0.16	-1.73	1.54	3.25	0.32	-0.71	6.21
85 kHz	0.08	-1.14	1.20	3.21	0.42	-0.67	5.99
Cut-off frequency ($F_{cut-off}$) variation (%)							
First cut-off frequency (20 kHz)	0	0	0	0	22.06	-2.34	17.37
Second cut-off frequency (40 kHz)	0	0	0.90	0.90	9.42	-0.90	9.42
Third cut-off frequency (70 kHz)	0.94	0.52	0.41	0.94	3.98	-1.05	5.45

362 For the sensitivity of $F_{cut-off}$, a positive variation associated with the coefficient C_{66} was
 363 observed, accounting for almost 100% of the "Overall" variation. Such a dependence could
 364 be used to assess the status of coefficient C_{66} separately.

365 **IV. CONCLUSION**

366 In this paper, the SAFE method was used to simulate the propagation of guided waves
367 at low frequency (20-85 kHz) in a bone waveguide in the context of an axial transmission
368 configuration. The SAFE method was shown to be an efficient tool for investigating the
369 effect of intracortical properties on the propagation of guided waves since a number of
370 configurations were simulated using reduced computing resources.

371 Considering the actual bone geometry, sixteen guided wave modes were found in the 10-
372 100 kHz frequency range (Figure 3b). However, only some of the modes have shown enough
373 excitability and consequently clinical interest. In addition, the excitability and sensitivity
374 of each mode were shown to vary according to the frequency and position of excitation on
375 the bone surface. The phase velocity showed a maximum sensitivity at 45 kHz, which would
376 appear to be associated with the peak of energy (power flow) concentrated in the endosteal
377 region at this frequency. The cut-off frequency was only slightly affected by the position of
378 excitation, but strongly by the frequency of excitation. A maximum sensitivity was achieved
379 for the first cut-off frequency at approximately 20 kHz.

380 The sensitivity of phase velocity was associated with the variation in the physiopatholog-
381 ical conditions of all the coefficients and density, but mainly with C_{44} . The sensitivity of the
382 cut-off frequency for its part was mainly associated with variation in the physiopathological
383 conditions of the stiffness coefficient C_{66} , allowing the axial shear modulus to be assessed
384 separately.

385 Thus, by selecting an adequate frequency and position of excitation, the low-frequency
386 axial transmission was shown to be a promising method for assessing intracortical bone
387 properties. The features identified in this study could be used as a metric to compare the
388 similarity between experimental and numerical data. A cost function based on these features
389 could then be implemented into an inversion scheme to retrieve reliable bone properties from
390 experimental data at low frequency. A parameterized bone-like geometry model, instead of
391 a plate or cylinder, could potentially reduce the errors commonly encountered with in-vivo
392 and ex-vivo experiments. For instance, the model could potentially compensate geometrical
393 variations between different patients that are unrelated to the bone quality, such as the
394 outer diameter and the external shape of the bone. The method could, therefore, be applied
395 to identify small changes associated with early-stage osteoporosis or gradual evolution of
396 the bone condition over time. For these extreme cases, since the sensitivity associated to
397 slightly different degraded conditions was found modest, a good basis for comparison (e.g.,
398 a baseline method or a large reference database) would have to be implemented in order to
399 guarantee the robustness of the method.

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