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Improved ZT in ball-milled and spark plasma sintered \( \text{Cu}_{15}\text{As}_{30}\text{Te}_{55} \) glass-ceramics

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Abstract

The influence of ball-milling followed by Spark Plasma Sintering (SPS) on the thermoelectric properties of the \( \text{Cu}_{15}\text{As}_{30}\text{Te}_{55} \) glass composition has been investigated by means of X-ray diffraction, scanning electron microscopy and differential scanning calorimetry. The microstructure, composition and thermal stability of the ball-milled samples before and after SPS have been correlated with the electrical and thermal transport properties. Upon ball-milling, a glass-to-glass transformation is evidenced for short ball-milling times, followed by crystallization of the metastable \( \beta\)-\( \text{As}_2\text{Te}_3 \) phase within the glassy matrix. This glass-to-glass transformation favors the occurrence of the \( \beta\)-\( \text{As}_2\text{Te}_3 \) phase during the SPS process. A maximum figure of merit \( ZT \) of 0.29 at 400K is obtained in the sample exhibiting the largest \( \beta\)-\( \text{As}_2\text{Te}_3 \) crystalline fraction. This \( ZT \) value is twice as high as the value obtained for the glass-ceramic sintered from non-ball-milled powder.

Keywords: Thermoelectric materials, ball milling, chalcogenide glass-ceramics
1. Introduction

Thermoelectric devices can be used as refrigerators (Peltier effect) or as electric power generators (Seebeck effect) since they can either pump heat or convert it into electricity. Their conversion efficiency is related to the dimensionless figure of merit $ZT$ of the constituting materials defined by $ZT = \alpha^2 T/\rho \lambda$ where $\alpha$ is the Seebeck coefficient or thermopower, $\rho$ is the electrical resistivity, $\lambda$ is the total thermal conductivity and $T$ is the absolute temperature. This technology is currently considered for supplying low electrical power (~ 1 µW) to sensors by harvesting heat from ambient air or for medium power generation (~ 100 W) by recycling the heat wasted for instance by the exhaust pipe of cars.

Owing to their inherent disordered nature, chalcogenide glasses harbor extremely low thermal conductivity (~0.2 W.m$^{-1}$.K$^{-1}$) which is, however, at the expense of good electrical property with very large electrical resistivity (typically ~ 100 Ω.m) $^{1,2}$. Upon metal substitution, their electrical resistivity can be nevertheless reduced by several orders of magnitude (~ 1 mΩ.m). Since they display large Seebeck coefficients (~ 500 µV K$^{-1}$) $^2$, they have been considered as potential thermoelectric materials for room temperature applications $^{3-5}$. Modifying the composition of these glasses is nonetheless insufficient to ensure a large thermoelectric figure of merit competing with state-of-the-art materials such as crystalline Bi$_2$Te$_3$, which displays $ZT = 1$ at 300K. This ascertainment led to the partial crystallization approach of the glassy phase, which in the Cu-As-Te system, yields glass – ceramics materials with a complex multiscale nano/microstructure displaying $ZT$ values of up to 0.08 at 300K $^6$. In such glass-ceramic materials prepared from the peculiar Cu$_{15}$As$_{30}$Te$_{55}$ glass composition, the promoted crystalline phase is the metastable $\beta$-As$_2$Te$_3$ phase which by itself, is an interesting thermoelectric material with maximum $ZT$ of 0.3 at 300 K upon proper optimization $^7$. However, our recent study on glass-crystal composites (Si$_{10}$As$_{15}$Te$_{75}$ (glass) and Bi$_{10.4}$Sb$_{1.6}$Te$_3$ (crystal)) $^8$, combining experiments and generalized effective medium
theory-based (GEMT) calculations of the thermoelectric properties \textsuperscript{9-13}, ruled out the possibility for a composite to reach $ZT$ values larger than any of its component. Nonetheless, the GEMT was in most cases verified in composites which were constituted by mixtures of micrometric phases and very seldom checked on mixtures of nanometric phases, for which interface effects cannot be neglected. Hence, a reduction to the nanoscale of the typical size of the phases in the Cu-As-Te – $\text{As}_2\text{Te}_3$ glass - ceramic material was undertaken in the present work in order to determine whether improved thermoelectric properties could be achieved. To this end, high energy ball-milling was carried out in order to decrease the size of the Cu$_{15}$As$_{30}$Te$_{55}$ glass particles, prior to the crystallization process obtained by Spark Plasma Sintering \textsuperscript{6}. Our results show that the influence of the particle size reduction to the nanoscale on the thermoelectric properties can be still well understood within the EMTs. The maximum $ZT$ value of the composite is nonetheless strongly enhanced due to a fraction of crystalline phase larger than that achieved in non-ball-milled samples.

2. Experimental procedures

2.1. Synthesis and ball milling of the Cu$_{15}$As$_{30}$Te$_{55}$ glass powders

The Cu$_{15}$As$_{30}$Te$_{55}$ glass composition was synthesized by the conventional melt-quenching technique. A series of stoichiometric composition using highly-pure As (Ingots, Goodfellow, 99.99\%), Te (Pellets, 5N+, 99.999\%) and Cu (Beads, Sigma-Aldrich, 99.999\%) were melted at 1123 K in a silica tube sealed under secondary vacuum and water-quenched \textsuperscript{6}. The resulting ingots (~2 g) were hand-ground in an agate mortar, then analyzed by powder X-ray diffraction (PXRD) on a Bruker D8 diffractometer (Cu K\textalpha radiation) equipped with a Vantec 1D detector and by differential scanning calorimetry (DSC, TA Q100) in order to check that they are identical before being combined in a single batch. This Cu$_{15}$As$_{30}$Te$_{55}$ glass batch was dispatched in several 4g samples and milled using a Fritsch P7 premium planetary
machine equipped with 80 mL stainless steel jars at a speed of 500 rpm (round per minute). The ball-to-powder ratio (BPR) was 15/1, with stainless balls of 2 mm in diameter. A first batch was incrementally ball-milled up to $t_{BM} = 360$ min by sampling 20 mg of powder from the vial inside a Ar-filled glove box at $t = 2, 8, 20, 45, 90, 180$ and $360$ min. This sampled powder was used to record a PXRD pattern (with a sample holder protecting the powder from the air), to carry out a scanning electron microscopy (SEM) investigation and for DSC measurements with a $10 \degree C \text{ min}^{-1}$ heating rate. Three other batches were ball-milled without interruption during $t_{BM} = 20, 45, 130$ min. They were also characterized by PXRD (with a low noise single-crystalline Si sample holder) and by DSC. These two different ball milling conditions (with and without interruptions) may slightly influence the temperature inside the vial and consequently the crystallization kinetics, even if large majority of the energy is provided by the chocks.

2.2. Spark Plasma Sintering (SPS)

The elaboration conditions of the different Cu$_{15}$As$_{30}$Te$_{55}$ pellets are summarized in Table 1 and are described below. The Cu$_{15}$As$_{30}$Te$_{55}$ powders milled for different durations ($t_{BM} = 0, 20, 45, 130$ and $360$ min) were sintered by SPS using a Dr. Sinter 505 Syntex apparatus to obtain glass-ceramic materials. It is noteworthy that for this specific SPS equipment, there is no way to bridle the intensity delivered by the generator (up to 1500 A). For each ball-milling time $t_{BM}$, 3 samples were elaborated for reproducibility concerns. 0.7 g of powder was introduced in an 8 mm diameter graphite die. The sintered temperature was checked by a thermocouple located in the wall of the die. In order to follow the evolution of the crystallization, 3 dwell times were applied (0, 10 and 20 min) at 423 K. An initial pressure of 50 MPa was implemented for $t_{BM} = 0, 130$ and $360$ min. Given the larger number of crystallized phases obtained as compared to a previous study 6, the following ball-milled
samples (t_{BM} = 20 and 45 min) and another reference sample (t_{BM} = 0 min) were sintered under a higher pressure of 80 MPa. However, as it will be discussed below (see t_{BM} = 0 min in Table 1), the applied pressure is not a significant parameter in this study compared to other parameters. In the case of t_{BM} = 0 min, an extra sample was synthesized with a 30 min dwell time under 80 MPa. In the following, the powder obtained from hand-ground Cu_{15}As_{30}Te_{55} glass (t_{BM} = 0 min) will be considered as a reference in this work. Figure 1 shows a representative SPS cycle used to elaborate glass-ceramics. The temperature rises at a heating rate of 90 K min\(^{-1}\) up to T = 390 K. At this temperature, the heating rate is decreased to 45 K min\(^{-1}\) to reach the dwell temperature of 423 K. As shown in Fig. 1 (a), an overshoot (T_{max}) can be seen before the set-temperature plateau. At the end of the dwell time, the sample undergoes a natural cooling. The density d of the samples was measured using the Archimedes’s method and is presented in Table 1. An average density of \(d = 6.28 \pm 0.05 \text{ g.cm}^{-3}\) over all the samples can be deduced. The standard deviation of 0.05 g.cm\(^{-3}\) is equal to the uncertainty of the measurement. Thus, there is no significant variation in the density with the duration of the initial ball-milling or the SPS parameters.

2.3. Transport properties of SPS glass-ceramics

Electrical resistivity measurements were directly performed on the SPS pellets with parallel faces using the van der Pauw technique\(^{14}\). The thermal diffusivity \(a\) was measured by the laser flash technique (Netzsch LFA 427) between 300 and 380 K. A piece of the sample was used to measure the heat capacity \(C_p\) from room temperature up to 400 K with a Netzsch Pegasus DSC apparatus, in order to calculate the thermal conductivity \(\lambda = a \times C_p \times d\) where \(d\) is the experimental density.

The Seebeck coefficient was measured between room temperature and 400 K on a parallelepiped-shaped (about 1 \(\times\) 1 \(\times\) 8 mm\(^3\)) sample cut with a diamond wire saw. These
measurements were carried out using a homemade apparatus based on the differential method. Measurement uncertainties are estimated to be 3% for Seebeck coefficient, 5% for electrical resistivity, 10% for thermal conductivity and 15% for ZT.

Finally, the samples used for the Seebeck coefficient measurements were hand-ground and PXRD was performed in the 2θ range 20-60° in order to determine the nature of the crystalline phases stabilized in the glass-ceramics. The amount of residual glass was estimated by integrating the area under the main XRD peak around 2θ = 29.58°, normalized to the high angle background. Moreover, the amount of α-As2Te3, β-As2Te3, AsTe and Cu1.4Te crystalline phases in the SPS samples was evaluated from Rietveld refinements.

3. Results and discussion

3.1 Structural, microstructural and thermal characterizations of ball milled Cu15As30Te55 glass

Figure 2 shows SEM micrographs corresponding to the glass particle size reduction after ball-milling. After 2 minutes, the particle size is reduced from hundreds of microns (Fig.2a) to about 10-20 µm (Fig.2b). After 20 and 360 minutes (Fig.2c and 2d), the SEM picture shows very small particles typically less than 1 µm. However, some larger particles, up to 5 µm, are still present. Larger magnification (Fig. 2e-g) indicates that these particles are aggregates of grains smaller than 100 nm.

Ball-milling effects in the Cu15As30Te55 glass samples were also followed by PXRD. In Fig.3, a + 0.6° shift in 2-theta for the main broad peak is noticeable around 29°, between the non-ball-milled glass and the 20 minutes ball-milled glass. This shift is too large to be ascribed to an artifact such as a sample displacement (max 0.1°). It can be noticed that this shifted peak is located at the same angle as a diffraction line of the crystallized phase. It could be due i) to a structural transformation of the initial glass into another polyamorph as already
observed in other chalcogenide glasses \textsuperscript{16, 17}, ii) to the emergence of nanometric crystalline domains, considering the low stability ($T_c - T_g$) of the glass ($\sim 30$K) where $T_c$ and $T_g$ are the crystallization and glass transition temperatures, respectively. Other evidences for this glass-transformation will be further discussed below. PXRD also highlights crystallization phenomenon which occurs during ball milling. After 45 min of ball-milling, some diffraction lines appear at 29.58° and 44.89° (Fig.3). While they cannot be observed at $t_{BM} = 90$ min, they appear again at $t_{BM} = 130$ min and become more intense for longer $t_{BM}$. Thus, the onset of crystallization is ambiguous when considering only the XRD data. By combining the former data with the DSC data (see Fig. 6), the beginning of crystallization can be better determined to occur between $t_{BM} = 45$ min and $t_{BM} = 90$ min. These diffraction lines at 29.58° and 44.89° are characteristic of one of the As$_2$Te$_3$ crystalline phases. Due to their very low intensity, it is difficult to distinguish the $\alpha$-As$_2$Te$_3$ from the $\beta$-As$_2$Te$_3$ crystalline polymorphs based solely on PXRD data. Nevertheless, the presence of the diffraction line near 45° is in favor of the metastable $\beta$-As$_2$Te$_3$ form, which is likely to appear under out-of-equilibrium ball-milling conditions in this specific glass composition \textsuperscript{6}.

The DSC curves, evidencing the temperatures of the glass transition, crystallization and phase transformation, are plotted in Figure 4 for all the incremented $t_{BM}$. The DSC curve of the reference Cu$_{15}$As$_{30}$Te$_{55}$ sample ($t_{BM} = 0$ min) shows an endothermic glass transition at $T_g = 403$ K and a first crystallization peak at $T_{c1} = 440$ K. This crystallization peak is shifted to higher temperature in comparison to our previous study \textsuperscript{6}. This enlightens the importance of the particle size on the crystallization phenomenon which occurs in the present case at the surface of the particles as already discussed in \textsuperscript{18}. Upon ball milling, $T_g$ decreases to 399 K for $t_{BM} = 90$ min (Table 2). For larger $t_{BM}$ (samples are partially crystallized), the glass transition is hidden by the first crystallization peak.
The first exothermic crystallization peak which extends from 440 K to 465 K in the reference sample \( (t_{BM} = 0 \text{ min}) \), is shifted to lower temperatures with \( t_{BM} \) (Table 2) in accordance with the decrease of the particle size \(^{19}\). A second weak peak is also observed between 465 K and 480 K in the reference sample \( (t_{BM} = 0 \text{ min}) \). Upon increasing the ball milling durations, this peak separates into two components, which both shift towards lower crystallization temperatures. In order to assess the crystalline composition, three 90 min interrupted ball-milled \( \text{Cu}_{15}\text{As}_{30}\text{Te}_{55} \) samples were heat-treated at a temperature higher than the maximum of each three DSC peaks \( i.e. \) at 438, 459 and 478 K respectively, \( \text{(black dots in Figure 4)} \) and their PXRD patterns were recorded. The results are displayed in Fig. 5. The diffractogram recorded on the 438 K heat-treated sample can be indexed with metastable \( \beta\)-\( \text{As}_2\text{Te}_3 \) as the major phase and \( \text{AsTe} \) as a secondary phase. The first crystallization peak observed in the DSC thermograms corresponds to the crystallization of these two phases, in good agreement with a previous work \(^6\). The second diffractogram at 459 K shows appearance of diffraction lines corresponding to the \( \text{Cu}_{1.4}\text{Te} \) rickardite phase. Thus, the second crystallization peak can be ascribed to the crystallization of the rickardite phase. Finally, the third diffractogram (Fig. 5) obtained after a heat-treatment at 478 K, shows diffraction peaks related to the \( \alpha\)-\( \text{As}_2\text{Te}_3 \) and \( \text{Cu}_{1.4}\text{Te} \) phases indicating that the third DSC peak corresponds to a \( \beta\)- to \( \alpha\)-\( \text{As}_2\text{Te}_3 \) phase transition. For \( t_{BM} = 180 \) and 360 min, the intensity of the second and third crystallization peaks hardly varies while the intensity of the first peak dramatically decreases as seen in Fig. 4. This reduction in intensity could be related to the partial crystallization of the glass particles into the \( \beta\)-\( \text{As}_2\text{Te}_3 \) crystalline phase upon ball-milling as evidenced by the PXRD patterns shown in Fig. 3.

Based on the area \( A \) of the first crystallization peak, a glass transformation degree \( f \) was defined and plotted in Fig. 6:
$f = 100 \times \frac{A(t_{BM=0}) - A(t_{BM})}{A(t_{BM=0})}$ (1)

A sharp increase in the glass transformation degree for small $t_{BM}$ is likely to arise from a transformation of the initial glass. As previously mentioned, this could be related to a nanostructuration of the grains. Nonetheless, the presence of a plateau following this sharp increase gives a highly non monotonous character to these variations. These are in favor of a glass-to-glass transformation (polyamorph), in agreement with both the shift of the main PXRD broad peak at $2\theta \sim 29^\circ$ and the change in the glass transition temperature $T_g$. For $t_{BM} > 45 - 90$ min, the increase in the glass transformation degree corresponds to the onset of the crystallization process during ball-milling and the presence of pre-existing crystallites, in good agreement with PXRD patterns (Fig. 3).

3.2 Glass-ceramics elaborated by SPS

In order to evaluate the effect of the ball-milling time on both the devitrification process and the thermoelectric properties, four series of Cu$_{15}$As$_{30}$Te$_{55}$ samples, ball-milled during various times (20, 45, 130 and 360 min) were sintered by SPS and compared to the reference sample sintered from hand-ground glass powder ($t_{BM} = 0$ min). During the sintering, a competition may occur between the glass particle sintering through a viscous flow mechanism and the crystallization phenomenon that occurs at the same time since the thermal stability of this specific glass composition is rather low ($\Delta T = 38$K), as discussed previously $^6$.

3.2.1. Microstructures

The microstructures of three characteristic sintered glass-ceramics obtained for $t_{BM} =$ 0, 20 and 360 min are presented in Fig. 7. The grain size after the SPS process on hand-ground samples is around 10 $\mu$m, whereas the grain size after the SPS process on the 20 minutes ball-milled sample is around 200 nm. The same grain size is observed for $t_{BM} = 360$
Ball-milled samples display the same microstructure and thus, cannot be differentiated by this feature. EDX maps reveal some differences in the chemical homogeneity of the non-ball-milled and ball-milled samples ($t_{\text{BM}} = 0$ and 20 minutes). While Te and As elements are well distributed within the matrix for both samples, Cu element is found as aggregates with micrometric size in the sample $t_{\text{BM}} = 0$ min (Fig. 7g). These aggregates, corresponding mostly to the crystallization of the $\text{Cu}_{1.4}\text{Te}$ rickardite phase, are more homogeneously distributed in the sample $t_{\text{BM}} = 20$ minutes for comparison (Fig. 7h). As will be further discussed, the fraction of crystalline phases is the critical parameter that best characterizes the samples.

### 3.2.2. Volume fractions of crystals and “glass rate”

Figure 8 shows a typical PXRD pattern of a sintered glass-ceramic sample ($t_{\text{BM}} = 130$ min, SPS dwell time $= 0$ min, 80MPa, 423K). A remaining quantity of glass is accompanied by four crystalline phases: $\alpha\text{-As}_2\text{Te}_3$, $\beta\text{-As}_2\text{Te}_3$, $\text{AsTe}$, and $\text{Cu}_{1.4}\text{Te}$. Rietveld refinements were performed to evaluate the relative crystalline phase fractions, which are presented in Table 1. A glass rate was defined by integrating the area under the main XRD peak around $2\theta = 29.58^\circ$, normalized to the high angle background around $2\theta = 59^\circ$.

Figure 9 presents the evolution of the crystalline fractions and the glass rate as a function of the ball-milling time. It is obvious that these fractions do not vary monotonously. The fraction of $\beta\text{-As}_2\text{Te}_3$ goes through a sharp maximum for short ball-milling times (until 20 and 45 min) and then slowly decreases for longer ball-milling times. On the contrary, the $\alpha\text{-As}_2\text{Te}_3$ crystalline phase initially displays the largest fraction in the non-ball-milled sample,
strongly decreases for short ball-milling times, and then slowly increases for longer durations. Meanwhile, the average crystalline fractions of Cu₁.₄Te and AsTe phases remain under 20 wt% and their variance is relatively large.

These evolutions of the crystalline phase fractions are in good agreement with the evolution of the glass transformation degree plotted in Fig.6. For short \( t_{BM} \) (< 45 min), the large fraction of \( \beta\)-As\₂Te\₃ phase is correlated to the first glass-to-glass transformation. Indeed, it appears that this glass-to-glass transformation yields a glass reorganization which favors the germination of the metastable \( \beta\)-As\₂Te\₃ phase at the sintering stage. On the other hand, for longer \( t_{BM} \), the crystallization, which occurs during ball-milling, favors larger fractions of the \( \alpha\)-As\₂Te\₃ phase at the SPS stage. The larger amount of energy provided by long \( t_{BM} \) finally favors the crystallization of the most thermodynamically stable phase, \( \alpha\)-As\₂Te\₃. Figure 9 gathers all the densified samples regardless of the sintering pressure (50 or 80 MPa). Since no distinction can be made between the samples, this parameter has no or a negligible influence on the phase fractions. Regarding the glass rate in the glass-ceramics, no trend can be noticed in the plot of the normalized glass rate versus \( t_{BM} \).

### 3.2.3 Thermal and electrical transport properties

Figure 10 shows the Seebeck coefficient, the electrical resistivity, the thermal conductivity, and \( ZT \) of the SPS glass-ceramics as a function of \( t_{BM} \). All the data shown here were measured at 400 K and are also summarized in Table 1. On the one hand, the evolution of the Seebeck coefficient with the ball-milling time mimics the evolution of the \( \beta\)-As\₂Te\₃ crystalline fraction up to \( t_{BM} = 45 \) min. For longer \( t_{BM} \), the Seebeck coefficient weakly varies despite the decrease of the \( \beta\)-As\₂Te\₃ fraction. This arises from the increase of the \( \alpha\)-As\₂Te\₃ phase fraction, which also presents good thermoelectric properties. Although the values of the thermal conductivity present a high variability in the non-ball-milled sintered samples,
Similarly to the Seebeck coefficient, they are correlated to \( t_{BM} \) and to the crystalline phase fractions, e.g. they decrease in the ball-milled samples. For long ball-milling times, the values weakly varies since the \( \alpha\text{-As}_2\text{Te}_3 \) and \( \beta\text{-As}_2\text{Te}_3 \) phases display similar values of thermal conductivities (1.0 W m\(^{-1}\) K\(^{-1}\) and 1.4 W m\(^{-1}\) K\(^{-1}\), respectively)\(^7,24\). An effect of the numerous interfaces created by the fine microstructure (~200 nm) on the phonon transport cannot as well be excluded to explain this trend. Note that (i) the variability for \( t_{BM} = 0 \) min can probably be related to the extremely variable remaining glass fraction for these samples (Fig. 9); (ii) the fact that the thermal conductivity reaches 1.8 W.m\(^{-1}\).K\(^{-1}\), while the thermal conductivity of the Cu\(_{15}\)As\(_{30}\)Te\(_{55}\) parent glass is about 0.22 W.m\(^{-1}\).K\(^{-1}\), or that of the \( \alpha \) or \( \beta\text{-As}_2\text{Te}_3 \) crystalline phase does not exceed 1.0 W m\(^{-1}\) K\(^{-1}\)\(^24\) and 1.4 W m\(^{-1}\) K\(^{-1}\), respectively\(^7\) could be explained by the presence of additional metallic phases, namely the Cu\(_{1.4}\)Te rickardite\(^25\) and AsTe\(^22\).

On the other hand, the resistivity values are subjected to a larger sample-to-sample variability, which probably masks the evolution of this property. The values for all the samples vary in an interval (±0.60 mΩ cm) around the average value of 1.5 mΩ cm but weakly influence the \( ZT \) values. This variability stems from: (i) the differences in crystalline phases fractions; (ii) the likely different carrier concentrations of the As\(_2\)Te\(_3\) phases, known for severely impacting their electrical properties, especially in the case of \( \alpha\text{-As}_2\text{Te}_3 \)\(^26\).

The figure of merit \( ZT \) follows the variations in the Seebeck coefficient: it goes through a maximum for short \( t_{BM} \), and weakly changes for longer \( t_{BM} \). Again, this evolution mimics the evolution of the \( \beta\text{-As}_2\text{Te}_3 \) phase fraction, which improves the thermoelectric properties of the glass-ceramics. This result is in full agreement with our previous work on the partial crystallization of the Cu\(_{15}\)As\(_{30}\)Te\(_{55}\) glass composition which enabled an increase in the \( ZT \) values from 0.08 to 0.14 at 365K by crystallizing the metastable \( \beta\text{-As}_2\text{Te}_3 \) phase\(^6\). Indeed, this phase was shown to display good thermoelectric properties at 400 K with \( ZT = 0.4 \)\(^7\). In
the present work, the effect of short ball-milling time leads to an increase in the $ZT_{\text{max}}$ value, from 0.15 for the best non-ball-milled sample to 0.3 at 400K for $t_{\text{BM}} = 20$ min.

4. Conclusions

In this study, the effect of ball-milling and crystallization from the Cu$_{15}$As$_{30}$Te$_{55}$ parent glass on the thermoelectric properties of the corresponding glass-ceramic materials are presented. The partial crystallization of glass remains an interesting and original way to stabilize new crystalline phases. Thanks to ball-milling, the grain size is reduced to less than 100 nm. PXRD and DSC measurements indicate that the first stages ($t_{\text{BM}} < 45$ min) induce a glass-to-glass transformation followed by the crystallization of the $\beta$-As$_2$Te$_3$ phase which occurs after $t_{\text{BM}} = 130$ min for interrupted ball-milling. Upon sintering, the grain size increases up to 200 nm in every ball-milled sample. Four crystalline phases, namely $\alpha$-As$_2$Te$_3$, $\beta$-As$_2$Te$_3$, AsTe and Cu$_{1.4}$Te, were evidenced in the different samples in various proportions. The fraction of the metastable $\beta$-As$_2$Te$_3$ phase is maximum for short $t_{\text{BM}}$ (20 - 45 min), i.e. for samples which underwent the glass-to-glass transformation (in comparison to non-ball-milled samples). These samples exhibit the largest thermoelectric figure of merit $ZT$ up to 0.3 at 400K, a value nearly doubled when compared to the non-ball-milled glass-ceramics.

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Figures captions

Figure 1: (a) SPS cycle of the reference sample (423 K, 80 MPa, dwell time = 30 min). Open squares are related to the temperature T inside the die and closed squares to load applied to the sample. (b) Corresponding shrinkage (mm) as a function of time during the same SPS cycle.

Figure 2: SEM micrographs of Cu$_{15}$As$_{30}$Te$_{55}$ powder obtained after different duration of ball milling $t_{BM} = 0$ (a), 2 (b), 20 (c) and 360 min (d) at low magnification and for $t_{BM} = 20$ (e) and 360 min (f, g) at high magnification.

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Figure 4: DSC thermograms of the Cu$_{15}$As$_{30}$Te$_{55}$ powders for various ball milling times. Black dots correspond to different annealing temperatures (438, 459 and 478 K, see text).

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Figure 6: Glass transformation degree $\mu$ as a function of ball-milling time $t_{BM}$. Squares are related to the incrementally ball-milled batch. Triangles are related to three other batches ball-milled without interruption during $t_{BM} = 20$, 45, 130 min. The continuous line is a spline through the entire dataset.

Figure 7: SEM micrographs of SPS glass-ceramics at two magnifications ($\times$ 1k and $\times$ 10k); a & b: the hand-ground reference ($t_{BM} = 0$ min); c & d: $t_{BM} = 20$ min ball-milled sample; e & f: $t_{BM} = 360$ min. SPS was performed at 423 K with 10 min dwell time, under a pressure of 50 MPa (reference and $t_{BM} = 360$ min) or 80 MPa ($t_{BM} = 20$ min). Elemental EDX maps of these last two samples: $t_{BM} = 0$ min (g) and $t_{BM}$=20 min (h).

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Table captions

Table 1: Ball-milling times $t_{\text{BM}}$, experimental SPS conditions ($T = 423\text{K}$, dwell time, pressure and maximal temperature in SPS $T_{\text{max}}$), density, crystallized phase fractions (extracted from Rietveld refinement), normalized glass rate, and thermoelectric properties measured at 400K.

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<table>
<thead>
<tr>
<th>$t_{BM}$ (min)</th>
<th>SPS Conditions:</th>
<th>$T_{max}$ (K)</th>
<th>Density</th>
<th>Crystallized phase fractions from Rietveld refinements (%mass)</th>
<th>Glass rate normalized toward background*</th>
<th>Seebeck coefficient ($\mu$V.K$^{-1}$)</th>
<th>Resistivity (mΩ.cm)</th>
<th>Thermal conductivity (W.m$^{-1}$.K$^{-1}$)</th>
<th>ZT</th>
</tr>
</thead>
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<tr>
<td>0</td>
<td>0 min / 80 MPa</td>
<td>425.4</td>
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<tr>
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<td>10 min / 80 MPa</td>
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<td>0.13</td>
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<td>0.08</td>
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<tr>
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<td>0.20</td>
<td>86</td>
<td>1.6</td>
<td>1.0</td>
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<td>20 min / 50 MPa</td>
<td>425.9</td>
<td>6.25</td>
<td>24 43 15 18</td>
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<td>86</td>
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<td>0.17</td>
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<td>91</td>
<td>1.2</td>
<td>1.2</td>
<td>0.24</td>
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*The values have been normalized to the 100% glassy compound.
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<table>
<thead>
<tr>
<th>$t_{BM}$ (min)</th>
<th>Incremental ball-milling (= interrupted BM)</th>
<th>Continuous BM</th>
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<tbody>
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<td>0</td>
<td>402.9(5)</td>
<td>397.0(5)</td>
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<td>400.3(5)</td>
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<td>90</td>
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<td>130</td>
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</tbody>
</table>
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115x169mm (150 x 150 DPI)
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174x235mm (150 x 150 DPI)
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272x208mm (300 x 300 DPI)
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307x290mm (150 x 150 DPI)
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108x237mm (150 x 150 DPI)
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306x104mm (150 x 150 DPI)