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Non-adiabaticity in mantle convection

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Abstract. Seismic observations indicate that Earth's lower mantle is homogeneous as revealed by smooth depth variations of the bulk sound velocity and Bullen's inhomogeneity parameter η being close to one. Here we show with 3D spherical convection simulations that it should also be non-adiabatic because a significant proportion of mantle heat sources is internal. The computer simulations predict non-adiabaticity of 100-300 degrees for the lower mantle geotherm, with $\eta \sim 1.01$ implying that the temperature drop across the Core Mantle Boundary exceeds 1000 degrees.

1. Introduction

Seismic studies of free oscillations and body waves provide by far the most detailed information on the large-scale structure of the Earth [Giardini & Woodhouse, 1987, van der Hilst *et al.*, 1997]. Through much of the lower mantle they reveal smooth depth variations of seismic sound velocities, consistent with a homogeneous material under adiabatic self compression. Assuming adiabaticity Birch [1952] combined the Adams-Williamson equation with bulk sound velocity measurements to calculate a lower mantle density profile. Lower mantle density can also be obtained directly, with no assumption of adiabaticity, from the study of spheroidal low frequency normal modes. The agreement between body wave and normal data suggests that the lower mantle is nearly adiabatic [Masters, 1979; Dziewonski & Anderson, 1981] and modeled geotherms have been found in good agreement with the adiabat [Shankland & Brown, 1985], supporting a long-held view that vigorous convection of the deep mantle results in adiabaticity.

Geotherms derived from vigorously convecting fluid models of the mantle are also adiabatic provided that both internal and viscous heating can be neglected, with strongly superadiabatic temperature gradients concentrated into narrow thermal boundary layers near the surface and the core mantle boundary (CMB) [Jarvis & McKenzie, 1980]. Strong internal heating, however, affects the geotherm giving rise to subadiabaticity in the interior of mantle convection cells even as the Rayleigh number approaches the convective

vigor of the mantle [Tackley *et al.*, 1994; Parmentier *et al.*, 1994; Bunge *et al.*, 1997]. The effect is important in mantle convection, where observational evidence suggests about 10% core heating in the overall heating budget of the mantle [Davies, 1988]. The physical reason for subadiabaticity is that material is slowly reheated by radioactivity while being entrained toward the surface by diffuse upwellings. Jeanloz & Morris [1987] estimate the impact of internal heating on the mantle geotherm and here we further investigate it in fully dynamic models. We demonstrate that subadiabaticity of 100-300 degrees may be expected in most of the lower mantle, while superadiabaticity in the bottom boundary layer should be confined to a layer of ~ 500 km.

2. Convection models

We model highly time-dependent convection in 3D spherical geometry [Bunge & Baumgardner, 1995]. Figure 1a shows a snapshot from an isoviscous reference calculation obtained assuming the mantle is compressible and heated purely from within. Material properties are identical to Bunge *et al.* [1997] and the Rayleigh number is 1.1×10^8 based on internal heating. The near surface and interior (cut-away) temperature fields are shown and the reference adiabat has been subtracted. Temperatures are subadiabatic as evidenced by colder fluid being concentrated at the mantle base. Shown in Figure 1e is the non-adiabatic radial temperature profile. Below the superadiabatic upper thermal boundary layer, temperatures drop subadiabatically by about 300 degrees from the upper mantle to the CMB. Figure 1b shows a model identical to the reference case, except that we have stiffened the lower mantle by a factor of 40 relative to the upper mantle viscosity (2×10^{22} Pa s), as suggested by studies of the geoid [Ricard *et al.*, 1993]. Subadiabaticity (Figure 1f) is comparable to the reference case, except for a region near the CMB where temperatures increase somewhat steeper than the adiabat. An isoviscous convection model with 50% core heating is shown in Figure 1c. Note that this core heat flux is significantly larger than estimates for the Earth [Davies 1988], and that it is in addition to the internal heating rate which remains unchanged from the reference case. The temperature field shows numerous hot upwellings from the CMB in addition to cold downwellings from the surface. The radial temperature profile (Figure 1g) reveals that mantle subadiabaticity is reduced to half its value in the reference case, or about 150 degrees. A model combining a lower mantle viscosity in-

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Figure 1. a) Non-adiabatic temperatures for the isoviscous, purely internally heated reference case. Blue is cold and red is hot. The top 100 km of the mantle is removed. Cold fluid is concentrated near the mantle base. b) Same as (a) except that lower mantle viscosity has been increased by a factor of 40. c/d) Same as (a/b) except for the addition of 50% core heating. (e)-(h) Non-adiabatic geotherms for cases (a)-(d). Superadiabaticity is concentrated into narrow thermal boundary layers. The geotherm is subadiabatic for the intervening region.

crease (factor 40) and 50% core heating with the internal heat production of the reference calculation is shown in Figure 1d. In this hybrid model hot mantle upwellings from the CMB are swept into a few major plumes due the long-wavelength planform induced by viscosity layering [Bunge *et al.*, 1996], while most of the lower mantle away from plumes appears to be relatively cold. This observation is borne out in the radial temperature profile (Figure 1h), which shows a non-adiabatic temperature gradient of -0.1 degree/km before reaching a broad thermal boundary layer at the CMB.

3. Density profiles and seismic inhomogeneity

The modeled geotherms may be converted into density profiles to estimate the subadiabatic effect on mantle density. In Figure 2a we calculate the lower mantle excess density relative to the adiabatic reference density profile in a

mineralogic model of pyrolite composition [Matas, 1999]. As expected, the cold geotherms steepen the density increase with depth for all four models shown in Figure 1. However, the anomalous excess density remains below 0.2% due to the small value of the thermal expansivity under lower mantle conditions ($1.5 \times 10^{-5} \text{ K}^{-1}$) [Chopelas & Boehler, 1992; Duffy & Ahrens, 1993]. The subadiabatic effect on the isentropic bulk modulus K_s (Figure 2b) is somewhat larger, with excess values of K_s approaching one percent. To go one step further in comparing our results with seismic observations, we compute Bullen's stratification parameter η [Bullen, 1963] in Figure 3. η measures departures from the Adams-Williamson relation, and is exactly one for adiabatic and chemically homogeneous regions. From Figure 3 we identify three distinct regimes for η . In the thermal

Figure 2. a) Lower mantle excess density in a pyrolite model relative to the adiabat for the four geotherms in Figure 1. Density variations remain below 0.2% due to the small value of lower mantle thermal expansivity. b) Same as a), but for the isentropic bulk modulus K_s .

boundary layers near the top and bottom of the mantle η is substantially smaller than one, as expected. There is also a sharp increase in the upper mantle for models with layered viscosity due to the strongly subadiabatic geotherm in the low viscosity zone. For the rapidly convecting interior of the lower mantle, we find η is about 1.01 in agreement with results obtained by Matyska & Yuen [2000] for 2D convection.

4. Discussion

The magnitude of mantle non-adiabaticity implied by the convection models is significant. Observational evidence [Davies, 1988] suggests predominantly internal mantle heat generation with an additional component arising possibly from secular cooling [Davies, 1980], which would act as a further source of internal heat. We explore the consequences in Figure 4, where we show absolute lower mantle temperatures tied to a temperature of 1800 K at the 660 km depth ringwoodite/perovskite+oxide phase change for two end-member calculations. The geotherm computed from the reference model (Figure 1a) increases by about 400 degrees through the lower mantle, with most of the lower mantle being stably stratified in agreement with earlier results obtained by Jarvis & McKenzie [1980]. The temperature inferred for the CMB is 2200 K in this case. Setting the core temperature to a geophysically more plausible value (3800 K) [Boehler, 1996] results in 50% heating from the core even for the more realistic convection model with a high viscosity lower mantle (Figure 1d) due to the sharp temperature drop of nearly 1600 K across the thermal boundary layer. This modeled heat flux is substantially higher than the core

Figure 4. Lower mantle geotherms for the isoviscous, bottom insulated reference model (Fig. 1a), and the stratified viscosity case with 50% heating from the core (Fig. 1d). A geothermal gradient of about 0.2 degree/km characterizes the lower mantle away from the thermal boundary layer at the CMB.

Figure 3. (a)-(d) Seismic inhomogeneity parameter η (see text) for the four models in Fig. 1a-d. η is defined as $\eta = -\frac{\phi}{\rho g} \frac{d\rho}{dr}$, where ρ is density, r is radial distance from the Earth's center, g is the acceleration due to gravity, and $\phi (= \frac{K_A}{\rho})$ is the seismic parameter.

heat flux inferred from mantle hotspots [Davies, 1988]. A chemically distinct layer at the bottom of the mantle would lower the core heat flux by distributing the temperature drop across an additional thermal boundary layer. It would also help to reduce the high excess temperature of mantle plumes [Farnetani, 1997] suggested by our isochemical calculations. The existence of this layer appears likely from geochemical [Hofmann, 1997] and seismic considerations [van der Hilst & Karason, 1999], and has been explored in dynamic models [Christensen & Hofmann, 1994; Kellogg et al., 1999; Coltice & Ricard, 2000]. A seismological consequence of our result is that homogeneous regions in the lower mantle should appear otherwise, with $\eta \sim 1.01$, and that η should be as low as $\eta \sim 0.92$ for a broad thermal boundary layer region in the bottom 500 km of the mantle (Figure 3). Seismic observations of η are resolved to about 2% in the lower mantle [Masters, 1979], which suggests that constraints on η should be improved to infer the mantle geotherm.

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