A PREDATOR-PREY SYSTEM WITH HOLLING-TYPE FUNCTIONAL RESPONSE

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ABSTRACT. We consider the Gause-type predator-prey system with a large class of growth and response functions, in the case where the response function is not smooth at the origin. We discuss the conditions under which this system has exactly one stable limit cycle or has a positive stable equilibrium point and we describe the basin of attraction of the stable limit cycle and the stable equilibrium point, respectively. Our results correct previous results of the existing literature obtained for the Holling response function $x^p/(a + x^p)$, in the case where $0 < p < 1$.

1. INTRODUCTION

Consider a predator-prey system of the so-called Gause type in the following general form:

$$\begin{aligned} \dot{x} &= xf(x) - yg(x), \\ \dot{y} &= y(\mu g(x) - D), \end{aligned} \tag{1.1}$$

where $\dot{} = d/dt$; x and y represent the prey population and the predator population, respectively; μ and D are positive parameters. Here $xf(x)$ is the growth function of the prey species in the absence of the predators, and $g(x)$ represents the functional response of predators to the growth of prey. For the background on this model, see Freedman [4].

In mathematical models for biology, the focus is on the global behavior of the system in the positive cone of the phase plane. The global behavior of the system is typically determined by the existence of a positive global attractor, either an equilibrium or a limit cycle. For this reason, the existence of a positive equilibrium of (1.1), its global asymptotic stability, the existence of a limit cycle (in case where the positive equilibrium is unstable), the uniqueness of the limit cycle, its global orbital attractivity, have attracted much interest in recent years, see [6, 7, 8, 9, 11, 12, 13, 14, 15] and the references therein.

The functions f and g in (1.1) are assumed to satisfy the following hypotheses.

(H1) f is smooth for $x \geq 0$ and there exists $k > 0$ such that $f(k) = 0$ and $(x - k)f(x) < 0$ for $x \geq 0$ and $x \neq k$.

(H2) g is continuous for $x \geq 0$, $g(0) = 0$ and $\mu g(k) > D$. It is smooth for $x > 0$ and $g'(x) > 0$ for $x > 0$.

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In (H1) the value k is interpreted as the carrying capacity of the prey species in the absence of the predators. The system (1.1) admits two boundary equilibrium points, $E_0 = (0, 0)$ and $E_1 = (k, 0)$. It follows from (H2) that there is a unique x^* such that $0 < x^* < k$ and $g(x^*) = D/\mu$. The predator isocline in the positive cone, i.e., the curve determined by the equation $\mu g(x) - D = 0$, is a vertical line $x = x^*$. Meanwhile, the prey isocline is the graph of the function $y = \varphi(x)$ where

$$(1.2) \quad \varphi(x) = \frac{xf(x)}{g(x)}.$$

It is easy to see that $\varphi(k) = 0$ and $\varphi(x) > 0$ for $0 < x < k$. The value of $\varphi(0)$ depends on additional properties on g near the origin. The intersection point of the isoclines, $E^* = (x^*, y^*)$ with $y^* = \varphi(x^*)$, is the only positive equilibrium point of system (1.1).

Since g is not assumed to be smooth at 0, E_0 is not a hyperbolic equilibrium. Using linear stability analysis (Jacobian) it is easy to see that E_1 is always a saddle point, while E^* can be a stable or unstable focus or node depending on the sign of $\varphi'(x^*)$, where φ is defined by (1.2). For $\varphi'(x^*) < 0$, E^* is an attractor. For $\varphi'(x^*) > 0$, E^* is a repeller.

A prototype of $f(x)$ and $g(x)$ in (1.1) is given by the logistic growth and the Holling functional response, see [12, 13],

$$(1.3) \quad f(x) = r \left(1 - \frac{x}{k}\right), \quad g(x) = \frac{x^p}{a + x^p},$$

where r , k , a and p are positive parameters (p is not always an integer). In this case, if $0 < p < 1$, the function $x \mapsto g(x)$ is not smooth at $x = 0$. Using numerical simulations, it was observed in [2] that, if $0 < p < 1$, some results of [12, 13] on the existence of a limit cycle or on the global stability of the positive equilibrium, are not true. We show in this paper why the classical arguments of dynamical systems do not apply, and why the results on existence and uniqueness of limit cycle, together with the result on global attractivity of the positive equilibrium or the limit cycle obtained in [12, 13] are not true when $0 < p < 1$.

Liu [9] considered the system (1.1). This author did not required the smoothness of $g(x)$ at $x = 0$, because (see the lines following Hypothesis (H2) in [9]), in some models it may happen that

$$(1.4) \quad g'(0^+) := \lim_{x \rightarrow 0^+} \frac{g(x)}{x} = +\infty.$$

For instance, for Holling functions (1.3) with $0 < p < 1$, (1.4) holds.

The purpose of this paper is to consider the case where $g(x)$ is not smooth at $x = 0$ and to study the global asymptotic stability of the positive equilibrium point (when it is locally exponentially stable), and also to study the existence of limit cycles (when the positive equilibrium is unstable) and the orbital asymptotic stability, and hence uniqueness, of the limit cycle (when the limit cycle exists).

The paper is organized as follows. In the next section we give a condition on the response function g , which is satisfied by the Holling function (1.3), with $0 < p < 1$, such that all points of the positive y -axis are points of non-uniqueness of the solutions of (1.1), that is to say the Cauchy problem (1.1) and the conditions $x(0) = 0$, $y(0) > 0$ has not only the trivial solution $x(t) \equiv 0$, but also has solutions satisfying $x(t) > 0$ for some $t < 0$. In Section 3 we discuss the global results on (1.1) and we describe the basin of attraction of the limit cycle, if it exists, and the

basin of attraction of the positive equilibrium point, when it is a local attractor. In Section 4 we apply our results to the system (1.1) where f and g are given by (1.3) and we correct some of the results in the existing literature which asserted that the system has a unique limit cycle (which is hence globally attracting), if and only if the positive equilibrium is a repeller, and that the positive equilibrium is globally attracting if and only if it is a local attractor. Some complements and technical proofs are given in the appendix.

2. NON UNIQUENESS OF SOLUTIONS

The function $g(x)$ is assumed to satisfy the following hypothesis.

(H3) $g(x)$ is not smooth for $x = 0$, (1.4) holds and the integral $\int_{\varepsilon}^{\alpha} \frac{dx}{g(x)}$ converges as $\varepsilon \rightarrow 0$, where $\alpha > 0$ is fixed.

It follows from (1.4) that $\varphi(0) = 0$. The condition on the convergence of the integral in (H3) is a technical condition that will be useful for the proofs presented in this paper. For instance, for Holling functions (1.3) with $0 < p < 1$, (H2) and (H3) hold. However, it should be noticed that (1.4) does not imply that the improper integral $\int_{\varepsilon}^{\alpha} \frac{dx}{g(x)}$ converges. For instance if function g is equivalent to $-x \ln(x)$ near 0, then (1.4) holds and the improper integral does not converge.

Let x_0 be such that $0 < x_0 < x^*$. Let $(X(t, x_0), Y(t, x_0))$ be the solution of (1.1) with initial condition $(x_0, \varphi(x_0))$. As long as $x^* > x > 0$ and $y > \varphi(x)$, the solution $(X(t, x_0), Y(t, x_0))$ can be parametrized by x , that is to say its orbit is the graph of a function of x :

$$(2.1) \quad y = \eta(x, x_0), \text{ for } x_0 \leq x \leq x^*.$$

Our aim is to show that the limit of these orbits (2.1) exists when $x_0 \rightarrow 0$ and is a positive orbit. If the equilibrium point E_0 was hyperbolic, then the limit of orbits (2.1) is simply the y -axis, which is the stable separatrix of E_0 . However, under hypothesis (H3) all points $(0, y_0)$, with $y_0 > 0$ are non uniqueness points and some solutions with initial conditions $x(0) > 0$, $y(0) > 0$ can reach in finite time the y -axis at some point $(0, y_0)$ after with y decreases to 0. Therefore the orbits (2.1) are bounded above, so that their limit exists and is a positive orbit.

More precisely, we consider the initial value problem

$$(2.2) \quad \begin{aligned} \dot{x} &= xf(x) - yg(x), & x(0) &= 0, \\ \dot{y} &= y(\mu g(x) - D), & y(0) &= y_0. \end{aligned}$$

Besides the trivial solution $(x(t) = 0, y(t) = y_0 e^{-Dt})$, the initial value problem (2.2) admits nontrivial positive solutions $(x_1(t), y_1(t))$, where $x_1(t)$ is not identically 0, which tends also to E_0 as $t \rightarrow +\infty$, as shown in the following result.

Theorem 2.1. *For any $y_0 > 0$, $(0, y_0)$ is a non uniqueness point for (2.2). The system admits a maximal solution $(X_1(t, y_0), Y_1(t, y_0))$ satisfying*

- (1) $X_1(t, y_0) > 0$ and $Y_1(t, y_0) > 0$ for $t < 0$.
- (2) $X_1(t, y_0) = 0$ and $Y_1(t, y_0) = y_0 e^{-Dt}$ for $t \geq 0$.

Proof. The concept of maximal solution used in this theorem should not be confused with the usual concept of *maximal solution*, defined on a maximal interval of definition. Here the solution is said to be *maximal* in the sense that it is the solution $(X_1(t), Y_1(t))$ of the initial value problem (2.2), which is non trivial (i.e.

106 $X_1(t) > 0$), on the maximal possible interval, that is to say for all $t < 0$. The
 107 details of the proof is given in Appendix B. $\square$

108 As long as $0 < x < x^*$ and $y > \varphi(x)$, the solution $(X_1(t, y_0), Y_1(t, y_0))$ can be
 109 parametrized by x , that is to say, its orbit is the graph of a function of x :

$$(2.3) \quad y = \eta_1(x, y_0), \text{ for } 0 \leq x \leq x^*.$$

Hence, we have the following inequalities, see Fig. 1(a):

$$\varphi(x) < \eta(x, x_0) < \eta_1(x, y_0), \text{ for } x_0 > 0, y_0 > 0 \text{ and } x_0 \leq x \leq x^*.$$

Therefore the following limits

$$w(x) = \lim_{x_0 \rightarrow 0} \eta(x, x_0) = \sup_{0 < x_0 < x^*} \eta(x, x_0),$$

$$w_1(x) = \lim_{y_0 \rightarrow 0} \eta_1(x, y_0) = \inf_{y_0 > 0} \eta_1(x, y_0),$$

110 exit and satisfy $\varphi(x) < w(x) \leq w_1(x)$, for $0 < x < x^*$, see Fig. 1(a). Since the
 111 graphs of functions $y = w(x)$ and $y = w_1(x)$ are limits, when $x_0 \rightarrow 0$ and $y_0 \rightarrow 0$
 112 respectively, of orbits of (1.1), they themselves are orbits of (1.1). We denote by
 113 $W^s(E_0)$ and $W_1^s(E_0)$ the global orbits corresponding to the local orbits $y = w(x)$
 114 and $y = w_1(x)$ respectively, see Fig. 1(b).

FIGURE 1. (a): The definitions of the local orbits $y = w(x)$ (colored in blue) and $y = w_1(x)$ (colored in red). (b): The global corresponding orbits $W^s(E_0)$ (colored in blue) and $W_1^s(E_0)$ (colored in red) and the elliptic sector, above $W^s(E_0)$, and hyperbolic sector, below $W^s(E_0)$ (colored in yellow), of E_0 .

115 *Remark 2.2.* Although the origin E_0 is not a hyperbolic singular point, we use the
 116 notations $W^s(E_0)$ and $W_1^s(E_0)$ since the solutions corresponding to these orbits
 117 tend to E_0 as $t \rightarrow +\infty$. The description of the behavior of the system near E_0 is as
 118 follows: The region between $W^s(E_0)$ and the x -axis is a hyperbolic sector and the
 119 region between $W^s(E_0)$ and the y -axis is an elliptic sector, see Fig. 1(b). Orbits
 120 who are between $W^s(E_0)$ and $W_1^s(E_0)$ tend toward E_0 when $t \rightarrow +\infty$ and those
 121 who are between the y axis and $W_1^s(E_0)$ reach the y axis in finite time and then,
 122 go toward the origin when $t \rightarrow +\infty$.

3. GLOBAL BEHAVIOUR

123

124 Assume that E^* is a repeller, that is to say $\varphi'(x^*) > 0$. If g were smooth at 0, then the system would necessarily have a limit cycle surrounding E^* . The existence
 125 of a limit cycle results from the Poincaré-Bendixson Theorem and the fact that E_0 is
 126 a saddle point, see Appendix A. Is this limit cycle globally orbitally asymptotically
 127 stable (GOAS) ? If the system has at most one cycle, then this cycle is necessarily
 128 GOAS.
 129

130 Assume that E^* is an attractor, that is to say $\varphi'(x^*) < 0$. Is E^* globally
 131 asymptotically stable (GAS) ? In the case where g is smooth at 0, if a solution with
 132 positive initial condition does not converge toward E^* , then, from the Poincaré-
 133 Bendixson Theorem and the fact that E_0 is a saddle point, it must converge toward
 134 a limit cycle, see Appendix A. If the system cannot have a cycle, then E^* is
 135 necessarily GAS.

136 In both cases the argument that the system has a limit cycle that must surround
 137 E^* uses the fact that E_0 is a saddle point, which does not hold when g is not
 138 smooth at 0, as shown in the preceding section, see Remark 2.2. In the case where
 139 hypothesis (H3) holds, E_0 is not a saddle point and the Poincaré-Bendixson theorem
 140 is inapplicable. In this case, the description of the global behavior of the system is
 141 obtained by considering the relative position of the orbit $W^s(E_0)$ with respect to
 142 the unstable separatrix $W^u(E_1)$ of the saddle point E_1 .

FIGURE 2. $W^s(E_0)$ (in blue) is above $W^u(E_1)$ (in red). The isoclines $\dot{x} = 0$ and $\dot{y} = 0$ are represented in green. (a): E^* is a repeller; (b): E^* is an attractor.

143 **3.1. The case where $W^s(E_0)$ is above $W^u(E_1)$.** We have the following result

144 **Proposition 3.1.** *Assume that $W^s(E_0)$ is above $W^u(E_1)$. If E^* is a repeller,*
 145 *then it is surrounded by at least one limit cycle. If the system can have at most one*
 146 *cycle, then E^* is surrounded by a unique limit cycle which is orbitally asymptotically*
 147 *stable. This limit cycle is not GOAS, even if it is unique. If E^* is an attractor and*
 148 *if the system has no cycles, then all orbits below $W^s(E_0)$ converge towards E^* . E^**
 149 *is not GAS, even if it is not surrounded by any unstable limit cycle.*

150 *Proof.* Assume that E^* is a repeller. The orbit $W^u(E_1)$ lies in the hyperbolic
 151 sector. Therefore, its ω -limit set is a limit cycle Γ surrounding E^* . This limit
 152 cycle is attracting at least on its exterior, see Fig. 2(a). Since E^* is repelling, it is
 153 surrounded also by a limit cycle Γ_1 which is attracting at least on its interior. If
 154 the system can have at most one cycle then $\Gamma_1 = \Gamma$. Therefore a limit cycle exists
 155 but it is not GOAS: all orbits above $W^s(E_0)$ converge toward E_0 and all positive
 156 orbits below $W^s(E_0)$, excepted for E^* , converge toward the limit cycle Γ .

157 Assume that E^* is an attractor. The orbit $W^u(E_1)$ lies in the hyperbolic
 158 sector. Therefore, its ω -limit set, if it does not contain E^* , is necessarily a limit cycle
 159 surrounding E^* . If the system cannot have a cycle then the limit set of $W^u(E_1)$
 160 is equal to E^* . Therefore E^* is not GAS since: all orbits above $W^s(E_0)$ converge
 161 toward E_0 and all positive orbits below $W^s(E_0)$ converge toward E^* , see Fig.
 162 2(b). $\square$

FIGURE 3. $W^s(E_0)$ (in blue) is below $W^u(E_1)$ (in red). The iso-
 clines $\dot{x} = 0$ and $\dot{y} = 0$ are represented in green. (a): E^* is a
 repeller; (b): E^* is an attractor.

163 **3.2. The case where $W^s(E_0)$ is below $W^u(E_1)$.** We have the following result

164 **Proposition 3.2.** *Assume that $W^s(E_0)$ is below $W^u(E_1)$. If E^* is a repeller and*
 165 *assuming that the system can have at most one cycle and, if it exists, this cycle*
 166 *cannot be semi-stable (i.e stable on one side and unstable in another side), then all*
 167 *positive orbits, excepted for E^* , converge toward E_0 . If E^* is an attractor then it*
 168 *is surrounded by an unstable limit cycle.*

169 *Proof.* Assume that E^* is a repeller. The orbit $W^s(E_0)$ lies in the hyperbolic sector.
 170 Therefore, its α -limit set, if it does not contain E^* , is necessarily a limit cycle Γ
 171 surrounding E^* . This limit cycle is repelling on its exterior. Since E^* is repelling,
 172 it is surrounded also by a limit cycle Γ_1 which is attracting on its interior. Since
 173 the system can have at most one cycle, $\Gamma_1 = \Gamma$. Therefore Γ is attracting on one
 174 side and repelling on the other side. Since this situation is impossible, the α -limit
 175 set of $W^s(E_0)$ is equal to E^* . Therefore no limit cycle can exist and all positive
 176 solutions, excepted for E^* , converge toward E_0 .

177 Assume that E^* is an attractor. The orbit $W^s(E_0)$ lies in the hyperbolic sector.
 178 Therefore, its α -limit set is necessarily a limit cycle Γ surrounding E^* . This limit
 179 cycle is repelling on its exterior. $\square$

180 The first case of the proposition (E^* is a repeller) is depicted in Fig.3(a). This
 181 case presents an unstable positive equilibrium, which is not surrounded by a cycle.
 182 The second case of the proposition (E^* is an attractor) is depicted in Fig. 3(b). If
 183 the system cannot have a cycle then the situation where E^* is an attractor and in
 184 the same time $W^s(E_0)$ is below $W^u(E_1)$ cannot hold.

185 4. APPLICATIONS

186 The purpose of this section is to clarify the necessary and sufficient conditions
 187 for the existence and uniqueness of limit cycles of the predator-prey system (1.1),
 188 where f and g are given by (1.3), with $0 < p < 1$. The system is written

$$(4.1) \quad \begin{aligned} \dot{x} &= rx \left(1 - \frac{x}{k}\right) - \frac{x^p y}{a + x^p}, \\ \dot{y} &= y \left(\frac{\mu x^p}{a + x^p} - D\right). \end{aligned}$$

189 The condition $\mu g(k) > D$ in (H2) is equivalent to $k > x^*$, where x^* is defined by
 190 $g(x^*) = D/\mu$, that is to say

$$(4.2) \quad \mu > D, \quad k > x^* := \left(\frac{aD}{\mu - D}\right)^{\frac{1}{p}}.$$

191 Hence functions f and g satisfy the hypotheses (H1), (H2) and (H3) if and only if
 192 $0 < p < 1$ and (4.2) is satisfied. Therefore the results in Propositions 3.1 and 3.2,
 193 have the corollary that if E^* is a repeller, the system does not always have a limit
 194 cycle, and if E^* is a attractor, it is not GAS.

195 **4.1. Existence of a limit cycle.** Sugie, Kohno and Miyazaki have stated (see
 196 Theorem 2.1 in [13]) that if

$$(4.3) \quad (pD - (p-1)\mu)k > (pD - (p-2)\mu)x^*,$$

197 then the system (4.1) has a unique limit cycle. Owing the the non uniqueness of
 198 the solution, this claim is not true. We have the following result:

199 **Proposition 4.1.** *Assume that (4.2) and (4.3) hold. Then, if $W^s(E_0)$ is above*
 200 *$W^u(E_1)$, the system (4.1) has a unique limit cycle whose basin of attraction is the*
 201 *positive region of the plane located below $W^s(E_0)$. If $W^s(E_0)$ is below $W^u(E_1)$, the*
 202 *system (4.1) has no limit cycle and all positive solutions, excepted for E^* converge*
 203 *toward E_0 .*

204 *Proof.* One can check that the condition (4.3) is equivalent to $\varphi'(x^*) > 0$. Therefore
 205 E^* is a repeller and it can be surrounded by a limit cycle. However, this limit cycle
 206 is not GOAS since it cannot attract the orbits above $W^s(E_0)$. Actually, if it exists,
 207 the limit cycle attracts all the orbits below $W^s(E_0)$. Using the result of [13] that,
 208 in this case, the system can have at most one limit cycle, we obtain as a corollary
 209 of Proposition 3.1 that if $W^s(E_0)$ is above $W^u(E_1)$, then the system has a unique
 210 limit cycle. On the other hand, using Proposition 3.2 we obtain that if $W^s(E_0)$ is
 211 below $W^u(E_1)$, then the system has no cycle and all positive solutions are attracted
 212 by E_0 . $\square$

FIGURE 4. The relative positions of $W^s(E_0)$ and $W^u(E_1)$ for (4.1) when the parameter values are $k = 9.5$, $a = 2$, $p = 0.6$, $\mu = 2$, $D = 1$ and (a): $r = 0.6$, (b): The heterocline bifurcation $W^s(E_0) = W^u(E_1)$ holds for $1.4026 < r < 1.4027$, (c): $r = 2$.

These results are illustrated in Fig. 4. The numerical simulations show that for the given values of the parameters, there exists a bifurcation value r_0 for which the heterocline bifurcation $W^s(E_0) = W^u(E_1)$ occurs: If $r < r_0$ then $W^s(E_0)$ is above $W^u(E_1)$ and the system has a unique limit cycle which is not GOAS; If $r > r_0$ then $W^s(E_0)$ is below $W^u(E_1)$ and the system has no limit cycle, and the positive solutions are attracted by the origin.

4.2. Global asymptotic stability of E^* . Sugie and Katayama have stated (see Theorem 1.1 in [12]) that the positive equilibrium E^* is GAS if and only if

$$(4.4) \quad (pD - (p - 1)\mu)k \leq (pD - (p - 2)\mu)x^*.$$

Owing to the non uniqueness of the solution, this claim is not true. We have the following result:

Proposition 4.2. *Assume that (4.2) and (4.4) hold. Then $W^s(E_0)$ is above $W^u(E_1)$ and the basin of attraction of E^* is the positive region of the plane located below $W^s(E_0)$. Therefore E^* is not GAS.*

Proof. One can check that the condition (4.4) is equivalent to $\varphi'(x^*) \leq 0$. Therefore E^* is an attractor. However, it is not GAS since it cannot attract the orbits above $W^s(E_0)$. Actually, it attracts all the orbits below $W^s(E_0)$. Using the result of [9, 12] that, in this case the system has no cycle, we obtain as a corollary of Proposition 3.2 that $W^s(E_0)$ is always above $W^u(E_1)$, that is to say, the situation depicted in Fig. 3 never occurs. Therefore, using Propositions 3.1 we obtain that E^* attracts all the orbits below $W^s(E_0)$. $\square$

These results are illustrated in Fig. 5, where it is shown that E^* is never GAS.

5. DISCUSSION

We have shown in this work that the differentiability of g at $x = 0$ could drastically affect the dynamics of (1.1). For instance, when $0 < p < 1$, and E^* is a repeller, the system (4.1) does not necessarily have a limit cycle. When $W^s(E_0)$ is above $W^u(E_1)$, the limit cycle exists. However, it is not GOAS. All solutions with initial conditions above $W^s(E_0)$ reach the y -axis in finite time. All solutions with initial conditions below $W^s(E_0)$ excepted the equilibrium point E^* converge to the

FIGURE 5. When E^* is an attractor, $W^s(E_0)$ is always above $W^u(E_1)$. The parameter values are the same as in Fig. 4, excepted that $D = 1.1$ and $r = 1.6$.

limit cycle. When $W^s(E_0)$ is below $W^u(E_1)$, the limit cycle does not exist. All solutions, excepted the equilibrium point E^* , reach the y -axis in finite time.

An heterocline bifurcation occurs when $W^s(E_0) = W^u(E_1)$ giving rise to a poly-cycle: all solutions with initial conditions outside the poly-cycle reach the y -axis in finite time. All solution with initial conditions inside the poly-cycle excepted the equilibrium E^* converge to the poly-cycle.

On the other hand, when $0 < p < 1$ and E^* is an attractor, it is not GAS. In this case $W^s(E_0)$ is always above $W^u(E_1)$, and the basin of attraction of E^* is the positive region below $W^s(E_0)$.

Our findings clarify some misunderstandings on the behaviour of this prey predator model that appeared in the existing literature [12, 13]. It is not true that when $0 < p < 1$ the system admits a limit cycle as long as its positive equilibrium E^* is a repeller, nor that this limit cycle is globally attracting whenever it exists, nor that E^* is GAS, whenever it is an attractor.

In [12], the authors consider that the function $g(x) = \frac{x^p}{a+x^p}$ in (4.1) is of Holling type II if $p \leq 1$ and of Holling type III if $p > 1$. Actually the functional response belongs to Holling type III if $p > 1$, since its has an inflexion point and the functional response curve is sigmoid [5]. When $0 < p \leq 1$, the function g is strictly increasing, bounded and convex, although the case $0 < p < 1$ deserves special interest because of the singularity of g in 0. From a biological and ecological point of view, the functional response g with $0 < p < 1$, defines a class of prey and predator populations characterized by predators with a very high attack rate for very small densities prey and a need for a large amount of prey to reach satiety. This response function has the property that the prey can go extinct in finite time, which never happens when $p \geq 1$.

APPENDIX A. EXISTENCE OF LIMIT CYCLES WHEN g IS SMOOTH AT 0

Most of the classical works on (1.1) or (4.1) have employed the transformation that converts the system into an equivalent generalized Liénard system. This method requires that $f(x)$ and $g(x)$ are sufficiently smooth so that existence, uniqueness, and continuation for all positive t are satisfied for initial-value problems. For instance, the proof of the existence of a cycle, when g is smooth at 0, and the positive equilibrium is a repeller, needs the Poincaré-Bendixson Theorem

and the fact that E_0 is hyperbolic. However, since this proof is standard, it is often omitted in the existing literature by authors who attempted to extend the results to functions g which are not smooth at 0. This is precisely the point where these proofs are not correct.

For the convenience of the reader, we give here the details on the proof of the existence of a limit cycle, when the positive equilibrium point is a repeller and the function g is smooth at 0. The hypothesis of hyperbolicity of E_0 is crucial to applying the so-called Butler-MacGehee Theorem which we recall here and whose proof is given in [10].

Theorem A.1. (*Butler-McGehee*). *If an omega (or alpha)-limit set ω contains a hyperbolic equilibrium point a , but it is not equal to $\{a\}$, then ω contains a point $b \neq a$ of the stable manifold $W^s(a)$, as well as a point $c \neq a$ of the unstable manifold $W^u(a)$.*

We have the following result

Theorem A.2. *Assume that (H1) and (H2) hold and in addition g is smooth at 0. If E^* is a repeller, then (1.1) has a limit cycle that contains E^* in its interior.*

Proof. It is known that under the hypotheses (H1) and (H2), all solutions of (1.1) are positively bounded. Therefore the omega-limit set ω of a positive trajectory is nonempty, invariant and compact. If g is smooth at $x = 0$ then E_0 is hyperbolic (a saddle point).

The limit set ω does not contain the saddle points E_0 and E_1 . Indeed, if ω contains E_0 , and since it is different from E_0 , because E_0 is the omega-limit set of only the trajectories lying in the y -axis, then by the Butler-MacGehee Theorem, it must contain a point $b = (0, y_0)$ with $y_0 > 0$ of $W^s(E_0)$ which is the positive y -axis. Since ω is invariant it must contain the whole trajectory of b , which is the positive y axis. This contradicts the compactness of ω .

In the same manner if ω contains E_1 , and since it is different from E_1 , because E_1 is the omega-limit set of only trajectories lying in the x -axis, then by the Butler-MacGehee Theorem, it must contain a point $b = (x_0, 0)$ with $x_0 > 0$ of $W^s(E_1)$ which is the positive x -axis. Two cases must be distinguished: If $x_0 > k$ then ω contains the whole trajectory of b which is unbounded. This contradicts the compactness of ω ; If $x_0 < k$ then ω contains the whole trajectory of b which is the segment $(0, k)$ of the x -axis. Since ω is closed, it must contain $[0, k]$, the closure of $(0, k)$. Therefore ω contains E_0 which is impossible as we have previously seen.

Assuming that E^* is a repeller, then it cannot belong to the omega-limit set ω . Therefore ω contains no equilibrium points. By the Poincaré-Bendixson Theorem, ω is a cycle. Since in the plane, a cycle must contain an equilibrium point in its interior, ω surrounds E^* . $\square$

If one have to our disposal the result that (1.1) admits at most one cycle, then the limit cycle whose existence was proved in the previous theorem is necessarily GOAS. Therefore, all solutions with positive initial conditions, excepted for E^* converge toward this limit cycle. The solutions with initial conditions on the y -axis converge toward E_0 , and the solutions with initial conditions on the positive x -axis converge toward E_1 .

When E^* is an attractor, we have the following result

Theorem A.3. Assume that (H1) and (H2) hold and in addition g is smooth at 0. If E^* is an attractor and (1.1) has no cycles then E^* is GAS.

Proof. Assume that E^* is an attractor. If a solution with positive initial condition does not converge toward E^* , then its omega-limit set ω cannot contain E^* . By the Butler-McGehee Theorem, it cannot contain E_0 nor E_1 as shown in the proof of the previous theorem. Therefore ω contain no equilibrium points. From the Poincaré-Bendixson Theorem, ω is a limit cycle, which contradicts the hypothesis that (1.1) has no cycles. Therefore $\omega = \{E^*\}$, which proves that E^* is GAS. $\square$

In the case where $g(x)$ is not smooth at $x = 0$, E_0 is not a hyperbolic saddle point and the arguments used in the proofs of Theorem A.2 and Theorem A.3 do not apply. For instance, in his proof of Theorem 1 in [9], the author claimed that $(0, 0)$ is a saddle, which is not true when $g(x)$ is not smooth at $x = 0$. Therefore, the system does not have necessarily a cycle when (x^*, y^*) is a repeller. Moreover, the claim of the author following the proof of his Theorem 1, that “local asymptotic stability of (x^*, y^*) implies its global asymptotic stability” is not true. Indeed Fig. 5 provides a counterexample to this claim.

APPENDIX B. PROOF OF THEOREM 2.1

The proof needs the following result, known as the Tamarkine’s non uniqueness theorem (see [1], p. 98).

Theorem B.1. Let $F(u, x)$ be a continuous function on $S = [-a, a] \times [-b, b]$, with $a > 0$ and $b > 0$. Assume that for all $(u, x) \in S$

$$|F(u, x) - F(u, x_0(u))| \geq G(|x - x_0(u)|),$$

where $u \mapsto x_0(u)$ is a solution of the initial value problem

$$(B.1) \quad \frac{dx}{du} = F(u, x), \quad x(0) = 0.$$

Assume that $z \mapsto G(z)$ is an increasing function for $z \geq 0$, $G(0) = 0$ and the integral $\int_{\varepsilon}^{\alpha} \frac{dz}{G(z)}$, ($\alpha > \varepsilon$) converges as $\varepsilon \rightarrow 0$. Then, the initial value problem (B.1) has at least two solutions in $[-a, a]$.

Proof. (of Theorem 2.1) The change of variables $y = y_0 e^u$, transforms (2.2) into the system

$$(B.2) \quad \begin{aligned} \dot{x} &= x f(x) - y_0 e^u g(x), & x(0) &= 0, \\ \dot{u} &= \mu g(x) - D, & u(0) &= 0. \end{aligned}$$

In the region

$$D = \left\{ (x, u) : 0 \leq x < x^*, u > \ln \frac{\varphi(x)}{y_0} \right\},$$

where $\varphi(x)$ is defined by (1.2), we have $x f(x) - y_0 e^u g(x) < 0$ and $\mu g(x) - D < 0$. Therefore (B.2) is equivalent to the scalar initial value problem

$$(B.3) \quad \frac{dx}{du} = F(u, x, y_0), \quad x(0) = 0,$$

with

$$F(u, x, y_0) = \frac{x f(x) - y_0 e^u g(x)}{\mu g(x) - D},$$

where y_0 is considered as a parameter. Using (1.2), $F(u, x, y_0)$ can be written

$$F(u, x, y_0) = g(x) \frac{y_0 e^u - \varphi(x)}{D - \mu g(x)}.$$

347 Since $g(0) = 0$, $F(u, 0, y_0) = 0$. Therefore $x_0(u) = 0$ is a solution of (B.3). Actually
 348 the initial value problem (B.3) has also non trivial solutions.

If $u \geq 0$ then $y_0 e^u \geq y_0$ and hence, for any $u \geq 0$ and $x \in [0, x^*)$, we have

$$|F(u, x, y_0) - F(u, x_0(u), y_0)| = F(u, x, y_0) > G(x),$$

where

$$G(z) = g(z) \frac{y_0 - \varphi(z)}{D - \mu g(z)}.$$

Let us denote by

$$h(z) = \frac{y_0 - \varphi(z)}{D - \mu g(z)}.$$

The function G is written $G(z) = g(z)h(z)$. It satisfies that $G(0) = 0$ and, for z small enough, G is increasing. Indeed, from $h(0) = y_0/D > 0$ we deduce that $h(x) > 0$ for x small enough, and from $g(0) = 0$ we deduce that $g(x)$ is small for x small enough, so that

$$G'(z) = g'(z)h(z) + g(z)h'(z) \approx g'(z)h(z) > 0.$$

For z small enough

$$\frac{1}{G(z)} = \frac{D - \mu g(z)}{g(z)(y_0 - \varphi(z))} \sim \frac{D}{y_0 g(z)}.$$

349 Therefore, using (H3), $\int_\varepsilon^\alpha \frac{dz}{G(z)}$ converges as $\varepsilon \rightarrow 0$. Using Theorem B.1 the initial
 350 value problem (B.3) has a positive nontrivial solution $x_1(u)$ such that $x_1(u)$ is not
 351 identically 0.

Let us prove now that the initial value problem (B.3) has a solution $\xi(u, y_0)$ satisfying the following conditions:

$$\xi(u, y_0) > 0, \text{ for } 0 < u < u_0 \text{ and } \xi(u, y_0) = 0, \text{ for } u \leq 0.$$

The set $\Phi(y_0)$ of all solutions $u \mapsto x(u)$ of (B.3) defined on a open interval I containing 0 contains at least two distinct solutions, the trivial solution x_0 for which $x_0(u)$ is identically 0 and a non trivial solution x_1 for which $x_1(u)$ is not identically 0. For each $u \in I$ let

$$\xi(u, y_0) = \sup_{x \in \Phi(y_0)} x(u).$$

It is well known (see [3], section 10.5, problem 7) that ξ belongs to $\Phi(y_0)$. It is called the *maximal solution* of (B.3) on I , corresponding to the initial condition $x(0) = 0$. Let $u_0 \in I$, such that $u_0 > 0$. Using the definition of ξ , we have $\xi(u, y_0) > 0$ for $0 < u < u_0$. The trivial solution $x_0(u) = 0$ of (B.3) corresponds to the trivial solution $(x(t) = 0, y(0) = y_0 e^{-Dt})$ of (2.2) and a nontrivial solution $x_1(u)$ of (B.3) corresponds to a nontrivial positive solution $(x_1(t), y_1(t))$ of (2.2). For instance, to the maximal nontrivial solution $\xi(u, y_0)$ corresponds the maximal solution $(X_1(t, y_0), Y_1(t, y_0))$ of (2.2) constructed as follows. First we consider the second equation in (B.2), where x is replaced by $\xi(u, y_0)$:

$$\dot{u} = \mu g(\xi(u, y_0)) - D, \quad u(0) = 0.$$

Let $U(t, y_0)$ be the solution of this initial value problem. Since $\xi(u, y_0) = 0$ for $u \leq 0$, one has $U(t, y_0) = -Dt$ for $u \leq 0$. Using the change of variables $y = y_0 e^u$, the solution $(X_1(t, y_0), Y_1(t, y_0))$ of (2.2) is defined by

$$X_1(t, y_0) = \xi(U(t, y_0), y_0), \quad Y_1(t, y_0) = y_0 e^{U(t, y_0)}.$$

This solution satisfies $X_1(t, y_0) = 0$, $Y_1(t, y_0) = y_0 e^{-Dt}$ for $t \leq 0$ and $X_1(t, y_0) > 0$, $Y_1(t, y_0) > 0$ for $t > 0$, which are the conditions in the theorem. $\square$

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