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TURBULENT FUNCTIONS AND SOLVING THE NAVIER-STOKES EQUATION
BY FOURIER SERIES

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Abstract : I give a resolution of the Navier-Stokes [2] equation by using the series of Fourier.

Résumé: Je donne une résolution de l'équation de Navier-Stokes [2] par les séries de Fourier.

Keywords : Navier-Stokes, Fourier, Séries de Fourier.

I . Introduction

The Navier–Stokes equations is considered to be the first step to understanding the elusive phenomenon of turbulence, the Clay Mathematics Institute in May 2000 made this problem [2] one of its seven Millennium Prize problems in mathematics.

In this article I will prove that the Navier-Stokes equation have a solutions and I will give techniques to resolve this beautiful equation.

The Navier-Stokes equation, established in the nineteenth century by the French Navier and the British Stokes. It is an equation that describes the velocity field of a fluid. More specifically, it is a differential equation whose velocity field is unknown.

The Navier-Stokes equation is also used to predict the weather, the oceans simulate, optimize aircraft wings ...

Knowing that a link between the Boltzmann equation and the Navier-Stokes equation was established, by studying the latter problem, I found that for to solve it we can reduce the problem of the heat-equation which is known can be solved by several methods : one of the first methods of solving the heat-equation was proposed by Joseph Fourier in his treatise analytical Theory of heat [1] in 1822 .

After giving a specific solutions to the Navier-Stokes equation, I will demonstrate how to find all solutions of this equation if they exist, and I give the necessary and sufficient conditions for their existence.

It will be seen in a remark that if the turbulence function is negligible, then the fluid will tend to behave like an ideal gas.

II. Recall, notations and definitions

Here are the Navier-Stokes equation:

$$\rho \left(\frac{\partial u}{\partial t} + (u \cdot \nabla) u \right) = -\nabla p + \mu \nabla^2 u$$

$$\operatorname{div} u = 0$$

Where u is the velocity field, p is the pressure, ρ the density of the fluid, and μ its viscosity.

And :

$$\nabla = \left(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n} \right)$$

$$(u \cdot \nabla) u = \sum_{i=1}^n u_i \frac{\partial u}{\partial x_i}$$

$$\nabla^2 u = \sum_{i=1}^n \frac{\partial^2 u}{\partial x_i^2}$$

$$\operatorname{div} u = \sum_{i=1}^n \frac{\partial u}{\partial x_i}$$

In the following, by dividing by ρ , the Navier-Stokes equations is of the form :

$$\frac{\partial u}{\partial t} + (u \cdot \nabla) u = \alpha \nabla p + \beta \nabla^2 u$$

$$\operatorname{div} u = 0$$

III- Existence of the solutions for the Navier-Stokes equation :

On each axis i , try to find the solutions of the form :

$$\frac{\partial u_i}{\partial t} + \left(u_i \frac{\partial u_i}{\partial x_i} \right) = \alpha \frac{\partial p_i}{\partial x_i} + \beta \frac{\partial^2 u_i}{\partial x_i^2}$$

This is equivalent to:

$$\frac{\partial u_i}{\partial t} + \frac{\partial \left(\frac{1}{2} u_i^2 - \alpha p_i \right)}{\partial x_i} = \beta \frac{\partial^2 u_i}{\partial x_i^2}$$

If u_i is a solution of the equation :

$$\frac{\partial u_i}{\partial t} = \beta \frac{\partial^2 u_i}{\partial x_i^2}$$

Such solutions u_i exist because the equation is analogous to the heat-equation which is resolvable by the Fourier series [1].

If p_i is such $\frac{1}{2} u_i^2 - \alpha p_i = f_i(t)$, then $\frac{\partial \left(\frac{1}{2} u_i^2 - \alpha p_i \right)}{\partial x_i} = 0$, and the equation is solved.

We do the same for all axes i until $i = n - 1$.

For the axe $i = n$:

Let be $u_n = - \sum_{i=1}^{n-1} u_i$.

We have : $\frac{\partial u_n}{\partial t} = \beta \frac{\partial^2 u_n}{\partial x_i^2}$, and if p_n is such $\frac{1}{2} u_n^2 - \alpha p_n = f_n(t)$, then

$$\frac{\partial \left(\frac{1}{2} u_n^2 - \alpha p_n \right)}{\partial x_n} = 0, \text{ and the equation is solved for the axe } n.$$

It is clear that if e_i is the vector for the axes i , then $u = \sum_1^n u_i e_i$ is one solution of the Naviers-Stokes equation if $\operatorname{div} u = 0$.

Else, to have $\operatorname{div} u = 0$, we take u_i of the form :

$$u_i = e^{\beta t + \left(\sum_1^n x_i \right)}, \forall i \in \{1, \dots, n-1\}$$

So we have solutions of the Navier-Stokes equation.

IV- Necessary conditions :

Any solution (u, p) of the Navier-Stokes equation verifies that :

$$u_i^2 - \alpha p_i = f_i(t), \forall i \in \{1, \dots, n\}$$

Indeed :

If (u, p) is a solution of the Navier-Stokes equation, we must have :

$$\frac{\partial u_i}{\partial t} + \frac{\partial \left(\frac{1}{2} u_i^2 - \alpha p_i \right)}{\partial x_i} = \beta \frac{\partial^2 u_i}{\partial x_i^2}$$

Therefore :

$$-\frac{\partial u_i}{\partial t} = \frac{\partial \left(\frac{1}{2} u_i^2 - \alpha p_i - \beta \frac{\partial u_i}{\partial x_i} \right)}{\partial x_i}$$

And :

$$-\partial x_i \frac{\partial u_i}{\partial t} = \partial \left(\frac{1}{2} u_i^2 - \alpha p_i - \beta \frac{\partial u_i}{\partial x_i} \right)$$

When the fluid flows in one direction, then the space-time flows in the opposite direction with the same speed value:

We deduce that :

$$-u_i \partial u_i = \partial \left(\frac{1}{2} u_i^2 - \alpha p_i - \beta \frac{\partial u_i}{\partial x_i} \right)$$

Therefore :

$$0 = \partial \left(u_i^2 - \alpha p_i - \beta \frac{\partial u_i}{\partial x_i} \right)$$

And :

$$u_i^2 - \alpha p_i - \beta \frac{\partial u_i}{\partial x_i} = f_i(t)$$

So :

$$\sum_1^n u_i^2 - \alpha p_i - \beta \frac{\partial u_i}{\partial x_i} = f(t)$$

And consequently $\sum_1^n u_i^2 - \alpha p_i = g(t)$ because $\operatorname{div} u = 0$.

We deduce therefore that :

$u_i^2 - \alpha p_i = h_i(t), \forall i \in \{1, \dots, n\}$ because : $\forall i \in \{1, \dots, n\}, u_i^2 - \alpha p_i = l_i(x_i, t)$, and we must have $\frac{\partial(u_i^2 - \alpha p_i)}{\partial x_j} = 0 \forall j \in \{1, \dots, n\}$.

V- Conclusion :

Theorem:

The Navier-stokes equation have a solution, Moreover, any solution (u, p) must check:

$$u_i^2 - \alpha p_i = f_i(t) \text{ and } \frac{\partial u_i}{\partial t} = \beta \frac{\partial^2 u_i}{\partial x_i^2}, \quad \forall i \in \{1, \dots, n\} \text{ where } u = \sum_1^n u_i e_i \text{ and}$$

$$p = \sum_1^n p_i e_i.$$

Conversely any pair (n, p) satisfying these conditions with $\operatorname{div} u = 0$ is solution of the Navier-Stokes equation.

Remarks :

1- We note in the above equations the dependence between pressure, density, speed vector fields and viscosity

2- By dividing by α in the equation $u_i^2 - \alpha p_i = f_i(t)$ we deduce that :

$$\rho u_i^2 - p_i = \rho f_i(t) \text{ where } \frac{1}{\alpha} = \rho \text{ is the density of the fluid.}$$

Let's $\rho = \frac{m}{V}$, we will have : $mu_i^2 - Vp_i = mf_i(t)$.

And the equation $mu_i^2 - Vp_i = mf_i(t)$ is linking energy, mass, pressure, temperature, volume and time ... This may not be surprising since a link between the Boltzmann equation and the

Navier-Stokes has been established.

When $f_i(t)$ tends to 0, we will have $mu_i^2 \simeq Vp_i$, there is therefore a tendency towards the law of an ideal gas, and the function $f_i(t)$ can be regarded as a turbulent function.

References

[1] Joseph Fourier, Théorie analytique de la chaleur, Firmin Didot Père et Fils (Paris-1822).
Réédition Jacques Gabay, 1988 ([ISBN 2-87647-046-2](#))

[2] <http://www.claymath.org/sites/default/files/navierstokes.pdf>