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# DENSITY OF THE SPAN OF POWERS OF A FUNCTION À LA MÜNTZ-SZÁSZ 

PHILIPPE JAMING \& ILONA SIMON


#### Abstract

The aim of this paper is to establish density properties in $L^{p}$ spaces of the span of powers of functions $\left\{\psi^{\lambda}: \lambda \in \Lambda\right\}, \Lambda \subset \mathbb{N}$ in the spirit of the Müntz-Szász Theorem. As density is almost never achieved, we further investigate the density of powers and a modulation of powers $\left\{\psi^{\lambda}, \psi^{\lambda} e^{i \alpha t}: \lambda \in \Lambda\right\}$. Finally, we establish a Müntz-Szász Theorem for density of translates of powers of cosines $\left\{\cos ^{\lambda}\left(t-\theta_{1}\right), \cos ^{\lambda}\left(t-\theta_{2}\right): \lambda \in \Lambda\right\}$. Under some arithmetic restrictions on $\theta_{1}-\theta_{2}$, we show that density is equivalent to a Müntz-Szász condition on $\Lambda$ and we conjecture that those arithmetic restrictions are not needed. Some links are also established with the recently introduced concept of Heisenberg Uniqueness Pairs.


## 1. Introduction

The aim of this paper is to establish density properties in $L^{p}$ spaces of the span of powers of a single or a pair of functions in the spirit of the Müntz-Szász Theorem.

Representing a generic function of some function space in terms of a family of simple functions is one of the main tasks in analysis. For instance, complex analysis deals with functions that can be expressed as power series, that is, the span of the functions $\left\{x^{k}, k \in \mathbb{N}\right\}$. Fourier analysis deals with the representation of functions in terms of the simple functions $\{\cos 2 k \pi t, \sin 2 k \pi t\}_{k \in \mathbb{Z}}$ or alternatively $\left\{e^{2 i k \pi t}=\left(e^{2 i \pi t}\right)^{k}\right\}_{k \in \mathbb{Z}}$. Exploring the spanning properties (basis, minimal set, $\ldots$ ) of the restricted trigonometric system $\left\{e^{2 i k \pi t}\right\}_{k \in \Lambda}, \Lambda \subset \mathbb{Z}$ in various function spaces has lead to a considerable bulk of Literature (see e.g. [Ru] as a starting point). In order to establish good spanning properties of the restricted trigonometric system, the first step consists in knowing if this system is total (that is, if its span is dense) in a given function space. Our aim here is to set a basic stone for similar properties when the basic brick $e^{2 i \pi t}$ is replaced by some other functions.

When considering the power functions $\left\{t^{\lambda}, \lambda \in \Lambda\right\}$, the problem dates back to the early 20 th century. This problem leads to one of the most intriguing results, the Müntz-Szász Theorem $[\mathrm{Mu}, \mathrm{Sz}]$ which relates the density of powers $\left\{x^{\lambda}: \lambda \in \Lambda\right\}$ in $\mathcal{C}([0,1])$ with an arithmetic property of $\Lambda$, namely the divergence of the series $\sum_{\lambda \in \Lambda \backslash\{0\}} \frac{1}{\lambda}$. This theorem has been extended in many ways, in particular to $L^{p}$ spaces, see e.g. $[\mathrm{BE}, \mathrm{CE}, \mathrm{Er}, \mathrm{EJ}, \mathrm{S}]$ and the nice survey $[\mathrm{Al}]$ for more on the subject. We will recall precise statements needed here in the next section.

The question we are asking here is of the same nature but we want to allow powers of more general functions than the identity. More precisely, we want to investigate the density of systems of the form $\left\{\psi^{\lambda}: \lambda \in \Lambda\right\}$ in $L^{p}([a, b])$ or $\mathcal{C}([a, b])$ when $\psi:[a, b] \rightarrow \mathbb{R}$ is a smooth function and $\Lambda$ is a set of integers ( $\psi$ may change sign). It is rather easy to notice that such a density can only occur when $\psi$ is monotonic (see Proposition 3.1 below). On the other hand, if $\psi$ has a local extrema then $\psi$ has some symmetry and this symmetry will also occur in the entire closed span of $\left\{\psi^{\lambda}: \lambda \in \Lambda\right\}$. Therefore, density can not be achieved for such functions. The question then arises on how to complete this system in order to obtain density.

[^0]One idea is to add translations of $\psi$. For instance, for a given $f \in \mathcal{C}([0,1])$ (here seen as the space of 1-periodic functions), we can consider the space

$$
\mathcal{T}(f)=\operatorname{span}\left\{f^{n}(t-\tau), n \in \mathbb{N}, \tau \in[0,1]\right\}
$$

As $\mathcal{T}(\cos 2 \pi t)$ is an algebra under pointwise multiplication, then, according to the Stone-Weierstrass Theorem, it is dense in $\mathcal{C}([0,1])$. This has been further investigated by Kerman and Weit [KW] who gave a characterization of the $f$ 's for which $\mathcal{T}(f)$ is dense in $\mathcal{C}([0,1])$. Further generalizations can be found e.g. in [RSW]. We address here a similar question for $f(t)=\cos 2 \pi t$ and we show that the set of powers and translates can then be substantially reduced. This should call for more research on density of

$$
\mathcal{T}_{\Lambda, T}(f)=\operatorname{span}\left\{f^{\lambda}(t-\tau), \lambda \in \Lambda, \tau \in T\right\} .
$$

A second option consists in adding modulations, instead of translate. In other words we are now looking for density criteria for

$$
\mathcal{M}_{\Lambda, \Omega}(f)=\operatorname{span}\left\{f^{\lambda}(t) e^{i \omega t}, \lambda \in \Lambda, \omega \in \Omega\right\}
$$

Here we show that if $\Lambda$ satisfies a Müntz-Szász type condition, two modulations suffice when $f$ has a single local maximum.

More precisely, our main results can be stated as follows (the general statement is more precise):
Theorem. Let $\Lambda$ be a set of non-negative integers containing zero and write $\Lambda=\{0\} \cup \Lambda_{e} \cup \Lambda_{o}$ where $\Lambda_{e}$ (resp. $\Lambda_{o}$ ) are the non-zero even (resp. odd) integers in $\Lambda$. Let $\theta_{1}, \theta_{2} \in \mathbb{R}$ such that $\theta_{1}-\theta_{2}$ is an irrational algebraic number and let $T=\left\{\theta_{1}, \theta_{2}\right\}$ and $\Omega=\{0, \omega\}$ with $|\omega|<1 / 2$. Then the following are equivalent:
(1) $\sum_{\lambda \in \Lambda_{e}} \frac{1}{\lambda}=+\infty$ and $\sum_{\lambda \in \Lambda_{o}} \frac{1}{\lambda}=+\infty$;
(2) $\mathcal{M}_{\Lambda, \Omega}(\cos \pi t)$ is dense in $L^{p}([0,1]), 1<p<+\infty$;
(3) $\mathcal{T}_{\Lambda, T}(\cos 2 \pi t)$ is dense in $L^{p}([0,1]), 1<p<+\infty$.

Moreover, the result stays true if $L^{p}([0,1])$ is replaced by $\mathcal{C}([0,1])$.
For $\mathcal{M}_{\Lambda, \Omega}$, the function $\cos \pi t$ can be replaced by any $\mathcal{C}^{2}$ smooth function $\psi:[0,1] \rightarrow \mathbb{R}$ such that $\psi^{\prime}$ vanishes at a single point $t_{0} \in(0,1)$ and $\psi^{\prime \prime}\left(t_{0}\right) \neq 0$.

We conjecture that the density of $\mathcal{T}_{\Lambda, T}(\cos 2 \pi t)$ is valid as soon as $\theta_{1}-\theta_{2}$ is irrational, while we prove that it is not valid when $\theta_{1}-\theta_{2}$ is rational.

The remaining of the paper is organized as follows. In the next section, we present some background on the Müntz-Szász Theorem. We then devote a section to our results on modulations while in Section 4 we prove our result concerning density of translates of the cosine function. In the last section we conclude by establishing some links with the recently introduced concept of Heisenberg Uniqueness Pairs.

## 2. Background and notations

Definition 1. Let $\Lambda \subset \mathbb{N}:=\{0,1,2, \ldots\}$ and $I=[a, b], a<b$ be a bounded interval. We will denote by $\Lambda_{e}=\Lambda \cap(2 \mathbb{N} \backslash\{0\})$ and $\Lambda_{o}=\{0\} \cup(\Lambda \cap(2 \mathbb{N}+1))$.

Let us define an ( $I-\mathrm{MS}$ ) sequence in the following way:

- when either $a=0$ or $b=0$, then we call $\Lambda$ an $I$-Müntz-Szász sequence, if $0 \in \Lambda$ and

$$
\sum_{\lambda \in \Lambda \backslash\{0\}} \frac{1}{\lambda}=+\infty ;
$$

- when either $a>0$ or $b<0$, then we call $\Lambda$ an $I$-Müntz-Szász sequence, if

$$
\sum_{\lambda \in \Lambda \backslash\{0\}} \frac{1}{\lambda}=+\infty ;
$$

- when $a<0<b$, then we call $\Lambda$ an $I$-Müntz-Szász sequence, if $0 \in \Lambda$,

$$
\sum_{\lambda \in \Lambda_{e}} \frac{1}{\lambda}=+\infty \quad \text { and } \quad \sum_{\lambda \in \Lambda_{o}} \frac{1}{\lambda}=+\infty
$$

We will further use the following notation: for $p \in[1, \infty]$, we write $X_{p}(I)=L^{p}(I)$ if $1 \leq p<+\infty$ and $X_{\infty}(I)=\mathcal{C}(I)$. We then define $p^{\prime}$ to be the usual dual index, $\frac{1}{p}+\frac{1}{p^{\prime}}=1$ with the convention that $1 / \infty=0$. Finally, we write $X_{p}^{\prime}=X_{p^{\prime}}$.

The classical Müntz-Szász Theorem [S, page 23], see also [BE, Section 6], states that
Theorem 2.1 (Müntz-Szász). Let $\Lambda \subset \mathbb{N}, 1 \leq p \leq+\infty, \frac{1}{p}+\frac{1}{p^{\prime}}=1$ and $I \subset \mathbb{R}$ be a bounded interval. The following conditions are equivalent
(i) The set $\left\{x^{\lambda}, \lambda \in \Lambda\right\}$ is total in $X_{p}(I)$.
(ii) If $f \in X_{p}^{\prime}(I)$ is such that $\int_{I} f(s) s^{\lambda} d s=0$ for every $\lambda \in \Lambda$, then $f=0$.
(iii) $\Lambda$ is an $I$-Müntz-Szász sequence.

Moreover,

- if $I \subset \mathbb{R}^{+}$or $\mathbb{R}^{-}$and $\sum_{\lambda \in \Lambda \backslash\{0\}} \frac{1}{\lambda}<+\infty$, then every function in the closed linear span of $\left\{x^{\lambda}, \lambda \in \Lambda\right\}$ is analytic in the interior of $I$;
- if $I=[a, b]$ with $a<0<b$ and $\sum_{\lambda \in \Lambda_{e}} \frac{1}{\lambda}<+\infty$ (resp. $\sum_{\lambda \in \Lambda_{o}} \frac{1}{\lambda}<+\infty$ ) then the even (resp. odd) part of each function in the closed linear span of $\left\{x^{\lambda}, \lambda \in \Lambda\right\}$ is analytic on $(a, b) \backslash\{0\}$.

Of course, the equivalence of (i) and (ii) is a direct consequence of the Hahn-Banach Theorem. The classical Müntz-Szász Theorem covers only the case $I=[0,1]$ (and thus $I=[a, b]$ with $a b=0$ ), the more general case $I=[a, b], a b \neq 0$ is due to Clarkson-Erdős and Schwartz. The case where $I$ is no longer included in a half-line is an easy consequence of the classical Müntz-Szász Theorem by writing $f$ in (ii) as a sum of an even and odd function (after extending $f$ by 0 so that it is defined on a symmetric interval). Also, this theorem is usually stated for density in $\mathcal{C}(I)$ but the statement is the same for $L^{p}(I)$ when $\Lambda \subset \mathbb{N}$, see e.g. [BE, Section 6].

Note that when $I$ intersects both $\mathbb{R}_{+}$and $\mathbb{R}_{-}$the statement can be reformulated in terms of the Fourier transform that we normalize as

$$
\widehat{f}(\zeta)=\mathcal{F}[f](\zeta):=\int_{\mathbb{R}} f(s) e^{-i s \zeta} \mathrm{~d} s
$$

if $f \in L^{1}(\mathbb{R})$ and extended to $L^{2}(\mathbb{R})$ in the usual way. In this case, if $\zeta_{0} \in \mathbb{R}$, then (i),(ii),(iii) are equivalent to
(iv) $\frac{\mathrm{d}^{\lambda}}{\mathrm{d} x^{\lambda}} \widehat{f}\left(\zeta_{0}\right)=0$, for every $\lambda \in \Lambda$ implies $f=0$.

## 3. Density of powers of a fixed function and modulation

In this section $I$ will still be a fixed bounded closed interval and $\psi: I \rightarrow \mathbb{R}$ a $\mathcal{C}^{1}$-smooth function (one may slightly weaken this condition). We will first prove the following result:

Proposition 3.1. Let $a, b \in \mathbb{R}$ and $\psi:[a, b] \rightarrow \mathbb{R}$ be a $\mathcal{C}^{2}$ function such that $\psi^{\prime}$ and $\psi^{\prime \prime}$ do not vanish simultaneously. Let $p \in[1,+\infty]$. Let $J=\psi([a, b])$ and let $\Lambda \subset \mathbb{N}$. The following are equivalent:
(i) $\left\{\psi^{\lambda}: \lambda \in \Lambda\right\}$ is total in $X_{p}(a, b)$.
(ii) $\psi$ is one-to-one and $\Lambda$ is a J-Müntz-Szász sequence.

Proof. Let us first assume that $\psi$ is not one-to-one. Then $\psi$ has a local extremum at a point $x_{0}$ in the interior of $[a, b]$. Therefore, there exists $a \leq a^{\prime}<x_{0}<b^{\prime} \leq b$ and a map $\varphi:\left[a^{\prime}, x_{0}\right] \rightarrow\left[x_{0}, b^{\prime}\right]$
such that $\varphi$ is one-to-one and onto and $\psi \circ \varphi=\psi$ on $\left[a^{\prime}, x_{0}\right]$. Let $f$ be any non-zero $\mathcal{C}^{1}$ function on [ $\left.x_{0}, b^{\prime}\right]$ and extend $f$ to $\left[a^{\prime}, x_{0}\right]$ by setting

$$
f(x)=-\varphi^{\prime}(x) f(\varphi(x))
$$

and then extend $f$ further to $[a, b] \backslash\left[a^{\prime}, b^{\prime}\right]$ by setting $f(x)=0$. Then $f \in L^{p^{\prime}}(a, b)\left(1 / p+1 / p^{\prime}=1\right)$ is non-zero and

$$
\int_{a}^{b} f(x) \psi^{\lambda}(x) \mathrm{d} x=\int_{a^{\prime}}^{b^{\prime}} f(x) \psi^{\lambda}(x) \mathrm{d} x=\int_{a^{\prime}}^{x_{0}}+\int_{x_{0}}^{b^{\prime}} f(x) \psi^{\lambda}(x) \mathrm{d} x
$$

Changing variable $x=\varphi(t)$ in the first integral we obtain

$$
\int_{a}^{b} f(x) \psi^{\lambda}(x) \mathrm{d} x=\int_{b^{\prime}}^{x_{0}} f(\varphi(x)) \psi^{\lambda}(\varphi(x)) \varphi^{\prime}(x) \mathrm{d} x+\int_{x_{0}}^{b^{\prime}} f(x) \psi^{\lambda}(x) \mathrm{d} x=0
$$

It follows that $\left\{\psi^{\lambda}: \lambda \in \Lambda\right\}$ is not total in $X_{p}(a, b)$.
Let us now assume that $\psi$ is one-to-one so that $\psi^{\prime}$ does not vanish (otherwise, if $\psi^{\prime}\left(x_{0}\right)=0$ then, by assumption, $\psi^{\prime \prime}\left(x_{0}\right) \neq 0$ so that $\psi^{\prime}$ changes sign at $x_{0}$ and $\psi$ would not be one-to-one). In particular, $\left|\psi^{\prime}\right|$ is bounded below. For a function $f$ on $(a, b)$ we define the function $\psi_{*} f$ on $J$ by $\psi_{*} f(t)=\frac{f\left(\psi^{-1}(t)\right)}{\psi^{\prime}\left(\psi^{-1}(t)\right)}$. Then, as $\left|\psi^{\prime}\right|$ is bounded below, $f \in X_{p}^{\prime}(a, b)$ if and only if $\psi_{*} f \in X_{p}^{\prime}(J)$.

Further, changing variable $t=\psi(x)$ we get

$$
\int_{a}^{b} f(x) \psi^{\lambda}(x) \mathrm{d} x=\int_{J} \frac{f\left(\psi^{-1}(t)\right)}{\psi^{\prime}\left(\psi^{-1}(t)\right)} t^{\lambda} \mathrm{d} t
$$

Applying the Müntz-Szász Theorem, the above proposition follows.
The question now arises on how to modify the set $\left\{\psi^{\lambda}, \lambda \in \Lambda\right\}$ in order to obtain a total set when $\psi$ is not one-to-one. In our opinion, there are two natural ways to do so, if one considers the Müntz-Szász theorem as a statement about the cancellation of the Fourier transform of a compactly supported function in a point. The first one consists of adding modulations the second one consists of adding translations. The following result deals with modulations and shows the equivalence $(1) \Leftrightarrow(2)$ of the theorem stated in the introduction.
Theorem 3.2. Let $a, b \in \mathbb{R}, \psi$ be a $\mathcal{C}^{2}$ function $[a, b] \rightarrow \mathbb{R}$ such that $\psi^{\prime}$ changes sign in a single point $x_{0} \in(a, b)$. Let $-\frac{1}{b-a}<\alpha<\frac{1}{b-a}$ and define $e_{\alpha}(t)=e^{i \alpha t}$. Let $\Lambda, \Lambda^{\prime} \subset \mathbb{N}$ and $p \in(1,+\infty]$. The following are equivalent.
(i) $\left\{\psi^{\lambda}, \lambda \in \Lambda\right\} \cup\left\{\psi^{\lambda} e_{\alpha}, \lambda \in \Lambda^{\prime}\right\}$ is total in $X_{p}$.
(ii) $\Lambda$ and $\Lambda^{\prime}$ are both $\psi([a, b])-M u ̈ n t z-S z a ́ s z ~ s e q u e n c e s . ~$

Example 1. Typical examples we have in mind are the functions $\psi(t)=\cos \pi t$ and $\psi(t)=1-\cos \pi t$ on $[-1 / 2,1 / 2]$ or equivalently $\psi(t)=\sin \pi t$ and $\psi(v)=1-\sin \pi t$ on $[0,1]$.

Further examples are $\psi(t)=t^{2}$ on $[-1,1], \psi(t)=1-t^{2}$ on $[0,1]$. A translation and dilation then gives a density criteria for the family $\left\{[t(1-t)]^{\lambda},[t(1-t)]^{\lambda} e^{i t}: \lambda \in \Lambda\right\}$ in $L^{p}([0,1])$.
Remark 1. The function $e_{\alpha}$ can be replaced by any function of the form $e^{i \varphi(t)}$ where the real valued function $\varphi$ is chosen such that, if $\psi(v)=\psi\left(v^{\prime}\right)$ with $v \neq v^{\prime}$, then $e^{i\left(\varphi(v)-\varphi\left(v^{\prime}\right)\right)} \neq 1$.

If one chooses $\varphi$ such that, if $\psi(v)=\psi\left(v^{\prime}\right)$ with $v \neq v^{\prime}$ implies $e^{i\left(\varphi(v)-\varphi\left(v^{\prime}\right)\right)} \neq \pm 1$, then the same result stays true for the system $\left\{\psi^{\lambda} \cos \varphi, \lambda \in \Lambda\right\} \cup\left\{\psi^{\lambda} \sin \varphi, \lambda \in \Lambda^{\prime}\right\}$. We leave the necessary adaptation of the proof below to the reader.

Proof. Let $p^{\prime}$ be given by $\frac{1}{p}+\frac{1}{p^{\prime}}=1$. We will only prove the theorem for $p \in(1, \infty)$ as no change is needed for $p=+\infty$.

We will use the following notation: set $J=\psi([a, b]), J_{+}=\psi\left(\left[x_{0}, b\right]\right)$ and $J_{-}=\psi\left(\left[a, x_{0}\right]\right)$ and $\psi_{+}^{-1}: J_{+} \rightarrow\left[x_{0}, b\right]$ be the inverse of $\psi$ on $\left[x_{0}, b\right]$ while $\psi_{-}^{-1}: J_{-} \rightarrow\left[x_{0}, b\right]$ is the inverse of $\psi$ on [ $a, x_{0}$ ]. Then

$$
\begin{aligned}
\int_{a}^{b} f(x) \psi(x)^{\lambda} \mathrm{d} x & =\int_{a}^{x_{0}} f(x) \psi(x)^{\lambda} \mathrm{d} x+\int_{x_{0}}^{b} f(x) \psi(x)^{\lambda} \mathrm{d} x \\
& =\int_{J_{-}} \frac{f\left(\psi_{-}^{-1}(y)\right)}{\psi^{\prime}\left(\psi_{-}^{-1}(y)\right)} y^{\lambda} \mathrm{d} y+\int_{J_{+}} \frac{f\left(\psi_{+}^{-1}(y)\right)}{\psi^{\prime}\left(\psi_{+}^{-1}(y)\right)} y^{\lambda} \mathrm{d} y \\
& =\int_{J}\left(\mathbf{1}_{J_{-}}(y) \frac{f\left(\psi_{-}^{-1}(y)\right)}{\psi^{\prime}\left(\psi_{-}^{-1}(y)\right)}+\mathbf{1}_{J_{+}}(y) \frac{f\left(\psi_{+}^{-1}(y)\right)}{\psi^{\prime}\left(\psi_{+}^{-1}(y)\right)}\right) y^{\lambda} \mathrm{d} y
\end{aligned}
$$

But then if we set*

$$
g(y)=\mathbf{1}_{J_{-}}(y) \frac{f\left(\psi_{-}^{-1}(y)\right)}{\psi^{\prime}\left(\psi_{-}^{-1}(y)\right)}+\mathbf{1}_{J_{+}}(y) \frac{f\left(\psi_{+}^{-1}(y)\right)}{\psi^{\prime}\left(\psi_{+}^{-1}(y)\right)}
$$

we get

$$
\int_{a}^{b} f(x) \psi(x)^{\lambda} \mathrm{d} x=\int_{J} g(y) y^{\lambda} \mathrm{d} y
$$

Similarly, if we set

$$
\tilde{g}(x)=\mathbf{1}_{J_{-}}(y) \frac{f\left(\psi_{-}^{-1}(y)\right) e^{i \alpha \psi_{-}^{-1}(y)}}{\psi^{\prime}\left(\psi_{-}^{-1}(y)\right)}+\mathbf{1}_{J_{+}}(y) \frac{f\left(\psi_{+}^{-1}(y)\right) e^{i \alpha \psi_{+}^{-1}(y)}}{\psi^{\prime}\left(\psi_{+}^{-1}(y)\right)}
$$

we get

$$
\int_{a}^{b} f(x) e^{i \alpha x} \psi(x)^{\lambda} \mathrm{d} x=\int_{J} \tilde{g}(y) y^{\lambda} \mathrm{d} y
$$

Let us now prove $(\mathrm{ii}) \Rightarrow(\mathrm{i})$. Assume that 0 is not in the interior of $J$ and that $\Lambda$ and $\Lambda^{\prime}$ are both $J$-Müntz-Szász sequences (the proof when 0 is in the interior of $J$ and $\Lambda_{e}, \Lambda_{o}$ and $\Lambda_{e}^{\prime}, \Lambda_{o}^{\prime}$ are all Müntz-Szász sequences is similar). Notice that if $f \in L^{p}(a, b)$, then $g, \tilde{g} \in L^{1}(J)$. According to the Müntz-Szász Theorem, if

$$
\int_{a}^{b} f(x) \psi(x)^{\lambda} \mathrm{d} x=\int_{a}^{b} f(x) \psi(x)^{\lambda^{\prime}} e^{i \alpha x} \mathrm{~d} x=0
$$

for every $\lambda \in \Lambda, \lambda^{\prime} \in \Lambda^{\prime}$, then $g=\tilde{g}=0$. But, writing

$$
f_{-}(y)=\mathbf{1}_{J_{-}}(y) \frac{f\left(\psi_{-}^{-1}(y)\right)}{\psi^{\prime}\left(\psi_{-}^{-1}(y)\right)} \quad \text { and } \quad f_{+}(y)=\mathbf{1}_{J_{+}}(y) \frac{f\left(\psi_{+}^{-1}(y)\right)}{\psi^{\prime}\left(\psi_{+}^{-1}(y)\right)}
$$

and $u_{ \pm}=e^{i \alpha \psi_{ \pm}^{-1}(y)}, g=\tilde{g}=0$ is equivalent to

$$
\left\{\begin{array}{ccc}
f_{-}(y) & +f_{+}(y) & =0 \\
u_{-} f_{-}(y) & + & u_{+} f_{+}(y)
\end{array}=0 .\right.
$$

As $u_{-} \neq u_{+}$, this implies $f_{+}=f_{-}=0$ thus $f=0$.
Conversely, for (i) $\Rightarrow$ (ii), assume that one of $\Lambda, \Lambda^{\prime}$ is not a $J$-Müntz-Szász sequence. Let $\tilde{p}=\frac{3 p^{\prime}}{3 p^{\prime}-1}$ so that $\frac{1}{\tilde{p}}+\frac{1}{3 p^{\prime}}=1$. Applying the Müntz-Szász Theorem in $L^{\tilde{p}}(J)$, there exist $g, \tilde{g} \in L^{3 p^{\prime}}$, one of them non zero and the other 0 , such that

$$
\int_{J} g(y) y^{\lambda} \mathrm{d} y=\int_{J} \tilde{g}(y) y^{\lambda^{\prime}} \mathrm{d} y=0
$$

[^1]for every $\lambda \in \Lambda, \lambda^{\prime} \in \Lambda^{\prime}$. If we find $f \in L^{p^{\prime}}([a, b])$ such that the associated $f_{ \pm}$are solution of
\[

\left\{$$
\begin{array}{ccc}
f_{-}(y) & +f_{+}(y) & =g  \tag{3.1}\\
u_{-} f_{-}(y) & +u_{+} f_{+}(y) & =\tilde{g}
\end{array}
$$\right.
\]

then also

$$
\int_{a}^{b} f(x) \psi(x)^{\lambda} \mathrm{d} x=\int_{a}^{b} f(x) \psi(x)^{\lambda^{\prime}} e^{i \alpha x} \mathrm{~d} x=0
$$

But, the system (3.1) has as solution

$$
f_{+}(y)=\frac{u_{-} g-\tilde{g}}{u_{-}-u_{+}} \quad \text { and } \quad f_{-}(y)=-\frac{u_{+} g-\tilde{g}}{u_{-}-u_{+}}
$$

that is

$$
\mathbf{1}_{J_{+}}(y) f\left(\psi_{+}^{-1}(y)\right)=\psi^{\prime}\left(\psi_{+}^{-1}(y)\right) \frac{e^{i \alpha \psi_{-}^{-1}(y)} g(y)-\tilde{g}(y)}{e^{i \alpha \psi_{-}^{-1}(y)}-e^{i \alpha \psi_{+}^{-1}(y)}}
$$

and

$$
\mathbf{1}_{J_{-}}(y) f\left(\psi_{-}^{-1}(y)\right)=-\psi^{\prime}\left(\psi_{-}^{-1}(y)\right) \frac{e^{i \alpha \psi_{+}^{-1}(y)} g(y)-\tilde{g}(y)}{e^{i \alpha \psi_{-}^{-1}(y)}-e^{i \alpha \psi_{+}^{-1}(y)}} .
$$

From this, we get

$$
f(x)=\left\{\begin{array}{ll}
\psi^{\prime}(x) \frac{e^{i \alpha \psi \psi_{+}^{-1} \circ \psi(x)} g \circ \psi(x)-\tilde{g} \circ \psi(x)}{e^{i \alpha \psi_{+}^{-1} \circ \psi(x)}-e^{i \alpha x}} & \text { for } x \in\left[a, x_{0}\right] \\
\psi^{\prime}(x) \frac{e^{i \alpha \psi_{-}^{-1} \circ \psi(x)} g \circ \psi(x)-\tilde{g} \circ \psi(x)}{e^{i \alpha \psi \psi_{-}^{-1} \circ \psi(x)}-e^{i \alpha x}} & \text { for } x \in\left[x_{0}, b\right]
\end{array} .\right.
$$

But now, as exactly one of $g, \tilde{g}$ is zero, $f$ is not the zero function. It remains to prove that $f \in$ $L^{p^{\prime}}(a, b)$. For this, define $f_{-}$on $\left[a, x_{0}\right]$ and $f_{+}$on $\left[x_{0}, b\right]$ by

$$
f_{ \pm}(x)=\psi^{\prime}(x) \frac{e^{i \alpha \psi_{\mp}^{-1} \circ \psi(x)} g \circ \psi(x)}{e^{i \alpha \psi_{\mp}^{-1} \circ \psi(x)}-e^{i \alpha x}}
$$

and $\tilde{f}_{ \pm}=f-f_{ \pm}$. It is enough to show that $f_{-}, \tilde{f}_{-} \in L^{p^{\prime}}\left(a, x_{0}\right)$ and $f_{+}, \tilde{f}_{+} \in L^{p^{\prime}}\left(x_{0}, b\right)$.
Next, changing variable $x=\psi_{-}^{-1}(t)$, we get

$$
\begin{aligned}
\int_{a}^{x_{0}}\left|f_{-}(x)\right|^{p^{\prime}} \mathrm{d} x & =\int_{a}^{x_{0}}\left|\psi^{\prime}(x) \frac{e^{i \alpha \psi_{+}^{-1} \circ \psi(x)} g \circ \psi(x)}{e^{i \alpha \psi_{+}^{-1} \circ \psi(x)}-e^{i \alpha x}}\right|^{p^{p^{\prime}}} \mathrm{d} x \\
& =\int_{J_{-}} \frac{\left|\psi^{\prime}\left(\psi_{-}^{-1}(t)\right)\right|^{p^{\prime}-1}}{\left|e^{i \alpha \psi_{+}^{-1}(t)}-e^{i \alpha \psi_{-}^{-1}(t)}\right|^{p^{p^{\prime}}}|g(t)|^{p^{\prime}} \mathrm{d} t} \\
& =\frac{1}{2^{p^{\prime}}} \int_{J_{-}} \frac{\left|\psi^{\prime}\left(\psi_{-}^{-1}(t)\right)\right|^{p^{\prime}-1}}{\left|\sin \frac{\alpha}{2}\left(\psi_{+}^{-1}(t)-\psi_{-}^{-1}(t)\right)\right|^{p^{2}}}|g(t)|^{p^{\prime}} \mathrm{d} t
\end{aligned}
$$

But now, as $\psi^{\prime \prime}\left(x_{0}\right) \neq 0, \psi(x)=\psi\left(x_{0}\right)+\frac{\psi^{\prime \prime}\left(x_{0}\right)}{2}\left(x-x_{0}\right)^{2}+o\left(\left(x-x_{0}\right)^{2}\right)$. From this, one immediately gets that

$$
\Phi(t)=\frac{\left|\psi^{\prime}\left(\psi_{-}^{-1}(t)\right)\right|^{p^{\prime}-1}}{\left|\sin \frac{\alpha}{2}\left(\psi_{+}^{-1}(t)-\psi_{-}^{-1}(t)\right)\right|^{p^{\prime}}} \approx C\left(t-\psi\left(x_{0}\right)\right)^{-1 / 2}
$$

when $t \rightarrow \psi\left(x_{0}\right)$ one of the end points of $J_{-}$. Further, the assumption on $\psi$ implies that $\Phi$ is $\mathcal{C}^{2}$ smooth on $J_{-} \backslash\left\{\psi\left(x_{0}\right)\right\}$. In particular, $\Phi \in L^{3 / 2}\left(J_{-}\right)$(say). Thus, from Hölder's inequality,

$$
\int_{a}^{x_{0}}\left|f_{-}(x)\right|^{p^{\prime}} \mathrm{d} x \leq \frac{1}{2^{p^{\prime}}}\|\Phi\|_{L^{3 / 2}\left(J_{-}\right)}\|g\|_{L^{3 p^{\prime}\left(J_{-}\right)}}^{p^{\prime}}<+\infty .
$$

The proof for $f_{+}$and $\tilde{f}_{ \pm}$is similar.

## 4. Density of translates of powers of the cosine function

In this section, functions on $[0,1]$ will be identified with 1-periodic functions, so that even and odd functions on $[0,1]$ make sense. We are interested in the density of translates of powers of the cosine (or sine) function, $\left\{\cos ^{\lambda} 2 \pi(t-\theta), \lambda \in \Lambda\right\}$. According to Proposition 3.1, this system is never dense in $L^{p}(0,1)$. Actually, for this function, it is easy to describe the "orthogonal":

Lemma 4.1. Let $p \in[1,+\infty], \Lambda \subset \mathbb{N}, \theta \in[0,1)$ and

$$
\mathcal{T}_{p, \Lambda, \theta}=\overline{\operatorname{span}}\left\{\cos ^{\lambda} 2 \pi(t-\theta), \lambda \in \Lambda\right\}
$$

be the closed subspace of $X_{p}(0,1)$ spanned by $\left\{\cos ^{\lambda} 2 \pi(t-\theta), \lambda \in \Lambda\right\}$. Let

$$
\mathcal{T}_{p, \Lambda, \theta}^{\perp}=\left\{f \in X_{p}^{\prime}(0,1): \int_{0}^{1} f(t) \cos ^{\lambda} 2 \pi(t-\theta) d t=0 \quad \forall \lambda \in \Lambda\right\}
$$

If $\Lambda$ is a $[-1,1]$-Müntz-Szász sequence, then

$$
\mathcal{T}_{p, \Lambda, \theta}=\left\{f \in X_{p}(0,1): f(\theta+t)-f(\theta-t)=0 \text { a.e. on }[0,1]\right\}
$$

and

$$
\mathcal{T}_{p, \Lambda, \theta}^{\perp}=\left\{f \in X_{p}^{\prime}(0,1): f(\theta+t)+f(\theta-t)=0 \text { a.e. on }[0,1]\right\} .
$$

Proof of Lemma 4.1. Up to translating by $\theta$, we may assume that $\theta=0$. We thus want to prove that $\mathcal{T}_{p, \Lambda, \theta}^{\perp}$ is the space of odd functions in $X_{p}^{\prime}(0,1)$. Once this is established, it is obvious that $\mathcal{T}_{p, \Lambda, \theta}$ is the space of even functions in $X_{p}(0,1)$.

Note that

$$
\int_{0}^{1} f(t) \cos ^{\ell} 2 \pi t \mathrm{~d} t=\int_{-1 / 2}^{1 / 2} f(t) \cos ^{\ell} 2 \pi t \mathrm{~d} t=\int_{0}^{1 / 2}(f(t)+f(-t)) \cos ^{\ell} 2 \pi t \mathrm{~d} t
$$

We can now change variable $u=\cos 2 \pi t$ to get

$$
\begin{equation*}
\int_{0}^{1} f(t) \cos ^{\ell} 2 \pi(t) \mathrm{d} t=\int_{-1}^{1} g(u) u^{\ell} \mathrm{d} u \tag{4.2}
\end{equation*}
$$

where

$$
\begin{equation*}
g(u)=\frac{1}{2 \pi \sqrt{1-u^{2}}}\left[f\left(\frac{\arccos u}{2 \pi}\right)+f\left(-\frac{\arccos u}{2 \pi}\right)\right] . \tag{4.3}
\end{equation*}
$$

Now $g \in L^{1}(-1,1)$ since

$$
\begin{aligned}
\int_{-1}^{1}|g(u)| \mathrm{d} u & \leq \frac{1}{2 \pi} \int_{-1}^{1}\left|f\left(\frac{\arccos u}{2 \pi}\right)\right|+\left|f\left(-\frac{\arccos u}{2 \pi}\right)\right| \frac{\mathrm{d} u}{\left(1-u^{2}\right)^{1 / 2}} \\
& =\int_{0}^{1 / 2}(|f(t)|+|f(-t)|) \mathrm{d} t<+\infty
\end{aligned}
$$

as $f \in X_{p}(0,1) \subset L^{1}(0,1)$. If $\Lambda$ is a $[-1,1]$-Müntz-Szász sequence we deduce that $g=0$ which is equivalent to $f(t)+f(-t)=0$ i.e. $f$ is odd.

Lemma 4.2. Let $p \in[1,+\infty], \Lambda \subset \mathbb{N}, \theta \in[0,1)$. Assume that $\Lambda$ is not $a[-1,1]$-Müntz-Szász sequence. Let $f \in \mathcal{T}_{p, \Lambda, \theta}$, that is, $f$ is even with respect to $\theta, f(\theta+t)=f(\theta-t)$ a.e., and assume further that

- if $\sum_{\lambda \in \Lambda_{e}} \frac{1}{\lambda}=+\infty$, then $f(\theta+1 / 2-t)=f(\theta+t)$
- if $\sum_{\lambda \in \Lambda_{o}} \frac{1}{\lambda}=+\infty$, then $f(\theta+1 / 2-t)=-f(\theta+t)$.

Then $f$ is analytic on $[0,1)$ except possibly at two points.

Remark 2. Write $c_{k}(f)$ for the $k$-th Fourier coefficient of $f$ and note that $f(\theta-t)=f(\theta+t)$ is equivalent to $c_{-k}(f) e^{-2 i \pi k \theta}=c_{k}(f) e^{2 i \pi k \theta}$ for all $k$. On the other hand, $f(\theta+1 / 2-t)= \pm f(\theta+t)$ is equivalent to $c_{-k}(f) e^{-2 i \pi k \theta}= \pm(-1)^{k} c_{k}(f) e^{2 i \pi k \theta}$ so that $c_{k}=0$ when $k$ is odd - resp. $c_{k}=0$ when $k$ is even.

Proof. Up to translating by $\theta$, we may assume that $\theta=0$ i.e., $f$ is even. Let $f \in \mathcal{T}_{p, \Lambda, 0}$ be even and define $h$ on $[-1,1]$ by

$$
h(x)=f\left(\frac{\arccos x}{2 \pi}\right)
$$

so that, when $f(1 / 2-t)=f(t), h$ is even, while $f(1 / 2-t)=-f(t)$ implies that $h$ is odd.
Let $F \subset \Lambda_{e}$ when $h$ is even (resp. $F \subset \Lambda_{o}$ when $h$ is odd) be a finite set.
Changing variable $t=\frac{\arccos x}{2 \pi}$, we get

$$
\begin{aligned}
\int_{-1 / 2}^{1 / 2}\left|f(t)-\sum_{\lambda \in F} c_{\lambda} \cos ^{\lambda} 2 \pi t\right|^{p} \mathrm{~d} t & =2 \int_{0}^{1 / 2}\left|f(t)-\sum_{\lambda \in F} c_{\lambda} \cos ^{\lambda} 2 \pi t\right|^{p} \mathrm{~d} t \\
& =\frac{1}{\pi} \int_{-1}^{1}\left|h(x)-\sum_{\lambda \in F} c_{\lambda} x^{\lambda}\right|^{p} \frac{\mathrm{~d} x}{\sqrt{1-x^{2}}} \\
& =\frac{2}{\pi} \int_{0}^{1}\left|h(x)-\sum_{\lambda \in F} c_{\lambda} x^{\lambda}\right|^{p} \frac{\mathrm{~d} x}{\sqrt{1-x^{2}}} \\
& \geq \frac{1}{\pi} \int_{0}^{1}\left|h(x)-\sum_{\lambda \in F} c_{\lambda} x^{\lambda}\right|^{p} \mathrm{~d} x
\end{aligned}
$$

so that $h \in \mathcal{P}_{p, \Lambda_{e}}$ (resp. $h \in \mathcal{P}_{p, \Lambda_{o}}$ ) where we denote by

$$
\mathcal{P}_{p, A}=\overline{\operatorname{span}}\left\{t^{a}, a \in A\right\}
$$

the closed subspace of $X_{p}(-1,1)$ spanned by $\left\{x^{a}, a \in A\right\}$.
As $\Lambda_{e}$ (resp. $\Lambda_{o}$ ) does not satisfy the Müntz-Szász condition, $h$ is analytic on $(0,1)$ so that $h$ is analytic in $(-1,1) \backslash\{0\}$. As $f(t)=h(\cos 2 \pi t)$, the result follows.

As one system $\left\{\cos ^{\lambda} 2 \pi(t-\theta), \lambda \in \Lambda\right\}$ is not total in $L^{p}(0,1)$, we may ask if adding a second system of this kind improves the situation. Let us first show that the second system can not be arbitrary:

Lemma 4.3. Let $p \in[1, \infty], \theta_{1}, \theta_{2} \in[0,1)$ and $\Lambda, \Lambda^{\prime} \subset \mathbb{N}$. Assume that the system

$$
\left\{\cos ^{\lambda} 2 \pi\left(t-\theta_{1}\right), \lambda \in \Lambda\right\} \cup\left\{\cos ^{\lambda^{\prime}} 2 \pi\left(t-\theta_{2}\right), \lambda^{\prime} \in \Lambda^{\prime}\right\}
$$

is total in $X_{p}(0,1)$, then $\theta_{1}-\theta_{2} \notin \mathbb{Q}$.
Proof. Let us write $\theta_{1}-\theta_{2}=\frac{m}{n} \in \mathbb{Q}$. Let $\varphi$ be a continuous, odd and $1 / n$-periodic function and $f(t)=\varphi\left(t-\theta_{2}\right)$. Then, for every $\ell \in \mathbb{N}$, as $f$ is also 1-periodic,

$$
\int_{0}^{1} f(t) \cos ^{\ell} 2 \pi\left(t-\theta_{2}\right) \mathrm{d} t=\int_{-1 / 2}^{1 / 2} \varphi(t) \cos ^{\ell} 2 \pi t \mathrm{~d} t=0
$$

since $\varphi$ is odd. Next, write $\theta_{1}=\theta_{2}+\frac{m}{n}$ then, as $\varphi$ is $1 / n$-periodic,

$$
f(t)=\varphi\left(t-\theta_{2}\right)=\varphi\left(t-\theta_{2}-m / n\right)=\varphi\left(t-\theta_{1}\right)
$$

therefore, for each $\ell \in \mathbb{N}$,

$$
\int_{0}^{1} f(t) \cos ^{\ell} 2 \pi\left(t-\theta_{1}\right) \mathrm{d} t=\int_{0}^{1} \varphi\left(t-\theta_{1}\right) \cos ^{\ell} 2 \pi\left(t-\theta_{1}\right) \mathrm{d} t=\int_{-1 / 2}^{1 / 2} \varphi(t) \cos ^{\ell} 2 \pi t \mathrm{~d} t=0
$$

using 1-periodicity again. It follows that $\left\{\cos ^{\lambda} 2 \pi\left(t-\theta_{1}\right), \lambda \in \Lambda\right\} \cup\left\{\cos ^{\lambda^{\prime}} 2 \pi\left(t-\theta_{2}\right), \lambda^{\prime} \in \Lambda^{\prime}\right\}$ is never total in $X_{p}(0,1)$.

We will also need the following lemma which provides a (non-orthogonal) decomposition trigonometric polynomials into a sum of two trigonometric polynomials with specific parity. In a sense, this generalizes the decomposition of a function into an even and an odd function. Unfortunately, it is only valid in full generality for trigonometric polynomials.
Lemma 4.4. Let $P$ be a 1-periodic trigonometric polynomial with zero-mean and $\theta_{1}, \theta_{2} \in \mathbb{R}$ such that $\theta_{1}-\theta_{2} \notin \mathbb{Q}$. Then there exists a unique pair $\left(P_{1}, P_{2}\right)$ of 1-periodic trigonometric polynomials with zero-mean such that $P_{1}\left(\theta_{1}-t\right)=P_{1}\left(\theta_{1}+t\right), P_{2}\left(\theta_{2}-t\right)=P_{2}\left(\theta_{2}+t\right)$ and $P=P_{1}+P_{2}$.
Proof. By expanding $P, P_{1}, P_{2}$ in Fourier series, one sees that the existence of the desired decomposition is equivalent to the systems

$$
\left\{\begin{array}{cllll}
c_{k}(P) & = & c_{k}\left(P_{1}\right) & + & c_{k}\left(P_{2}\right) \\
c_{-k}(P) & =e^{4 i \pi k \theta_{1}} & c_{k}\left(P_{1}\right) & +e^{4 i \pi k \theta_{2}} & c_{k}\left(P_{2}\right)
\end{array} \quad, \forall k \in \mathbb{Z} \backslash\{0\}\right.
$$

As $\theta_{1}-\theta_{2} \notin \mathbb{Q}$, the determinant of this system is non zero for every $k \neq 0$ and its solutions are given by

$$
\begin{equation*}
c_{k}\left(P_{1}\right)=\frac{c_{-k}(P)-c_{k}(P) e^{4 i \pi k \theta_{2}}}{e^{4 i \pi k \theta_{1}}-e^{4 i \pi k \theta_{2}}} \quad, \quad c_{k}\left(P_{2}\right)=\frac{c_{-k}(P)-c_{k}(P) e^{4 i \pi k \theta_{1}}}{e^{4 i \pi k \theta_{2}}-e^{4 i \pi k \theta_{1}}} . \tag{4.4}
\end{equation*}
$$

which gives both existence and uniqueness.
Remark 3. When $\theta_{1}-\theta_{2}=\frac{m}{n} \in \mathbb{Q}$, the lemma stays true with the same proof if we impose $P, P_{1}, P_{2}$ to have degree $<n / 2$ for even $n$ and $<n$ for odd $n$.

Note also that the zero mean assumption is only used to guarantee uniqueness as constant functions satisfy both parities $P_{1}\left(\theta_{1}-t\right)=P_{1}\left(\theta_{1}+t\right), P_{2}\left(\theta_{2}-t\right)=P_{2}\left(\theta_{2}+t\right)$. Actually, the proof shows that the $P_{1}, P_{2}$ 's we obtain have both zero mean. If $P$ has non-zero mean, we apply the lemma to $P-c_{0}(P)$, which has zero mean. We may thus write $P=c_{0}(P)+P_{1}+P_{2}=$ $\left(P_{1}+\lambda c_{0}(P)\right)+\left(P_{2}+(1-\lambda) c_{0}(P)\right), \lambda \in \mathbb{R}$. Uniqueness would still be guaranteed if we ask for, say, $P_{2}$ to have zero mean which would imply that we take $\lambda=1$ in the above decomposition.

Note that if $P$ is no longer a trigonometric polynomial but a function in $L^{1}$, the formula (4.4) will in general lead to sequences that do not go to zero, so that they are not sequences of Fourier coefficients of $L^{1}$ functions.

Before going on with our main subject, let us elaborate a bit on this topic:
Lemma 4.5. Let $\theta_{1}, \theta_{2} \in \mathbb{R}$ such that $\theta:=\theta_{1}-\theta_{2} \in \mathbb{R} \backslash \mathbb{Q}$. Let $a>0$ and assume that $\theta$ is a-approximable by rational numbers in the sense that there is a constant $C_{\theta}>0$ such that the set

$$
\left\{(m, n) \in \mathbb{Z}:|m-n \theta|<C_{\theta} n^{-a}\right\}
$$

is finite.
Let $s \geq a$ and $f \in H^{s}(0,1)$ with mean zero, then there exists a unique pair $f_{1}, f_{2} \in L^{2}(0,1)$ such that $f=f_{1}+f_{2}, f_{1}\left(\theta_{1}-t\right)=f_{1}\left(\theta_{1}+t\right)$ and $f_{2}\left(\theta_{2}-t\right)=f_{2}\left(\theta_{2}+t\right)$.

Moreover, if $s>a+1 / 2+j$ for some integer $j$, then $f_{1}, f_{2}$ are of class $\mathcal{C}^{j}$.
Remark 4. From the Dirichlet Theorem, no irrational number is 1-approximable. However, according to Khinchin's theorem that for $a>1$, almost every real number is $a$-approximable. Further, from the Thue-Siegel-Roth Theorem, every algebraic number is $1+\varepsilon$-approximable for every $\varepsilon>0$. On the other hand, Liouville numbers are not $a$-approximable for any number. See e.g. [HS, Sc, Wa] and references therein for more on the subject.

Proof. As previously, if the decomposition exists then the Fourier coefficients of $f_{1}, f_{2}$ are given by

$$
c_{k}\left(f_{1}\right)=\frac{c_{-k}(f)-c_{k}(f) e^{4 i \pi k \theta_{2}}}{e^{4 i \pi k \theta_{1}}-e^{4 i \pi k \theta_{2}}} \quad, \quad c_{k}\left(f_{2}\right)=\frac{c_{-k}(f)-c_{k}(f) e^{4 i \pi k \theta_{1}}}{e^{4 i \pi k \theta_{2}}-e^{4 i \pi k \theta_{1}}} .
$$

Conversely, if these two sequences are in $\ell^{2}$ (resp. if $k^{j} c_{k}\left(f_{1}\right), k^{j} c_{k}\left(f_{2}\right) \in \ell^{1}$ ) then the corresponding Fourier series define $f_{1}, f_{2}$ as functions in $L^{2}(0,1)$ (resp. in $\mathcal{C}^{j}(0,1)$ ).

Since $f \in H^{s},(1+|k|)^{s} c_{k}(f) \in \ell^{2}$. On the other hand

$$
\left|e^{4 i \pi k \theta_{1}}-e^{4 i \pi k \theta_{2}}\right|=2|\sin \pi k 2 \theta|=2|\sin \pi \operatorname{dist}(2 k \theta, \mathbb{Z})| \geq 4 \operatorname{dist}(2 k \theta, \mathbb{Z})
$$

But, $\operatorname{dist}(2 k \theta, \mathbb{Z})>2^{-a} C_{\theta} k^{-a}$ for all but finitely many $k$ 's so that $\left|e^{4 i \pi k \theta_{1}}-e^{4 i \pi k \theta_{2}}\right|>2^{2-a} C_{\theta}|k|^{-a}$ for all but finitely many $k$ 's, thus for all $|k| \geq k_{\theta}$ for some $k_{\theta} \in \mathbb{N}$.

Now, if $f \in H^{s}$, for $|k| \geq k_{\theta}$,

$$
\frac{\left|c_{k}(f)\right|}{\left|e^{4 i \pi k \theta_{1}}-e^{4 i \pi k \theta_{2}}\right|} \leq \frac{2^{a-2}}{C_{\theta}}|k|^{a}\left|c_{k}(f)\right| \in \ell^{2}(\mathbb{Z})
$$

if $s \geq a$. Further

$$
\begin{aligned}
\sum_{|k| \geq k_{\theta}} \frac{|k|^{j}\left|c_{k}(f)\right|}{\left|e^{4 i \pi k \theta_{1}}-e^{4 i \pi k \theta_{2}}\right|} & \leq \sum_{|k| \geq k_{\theta}} \frac{|k|^{s}\left|c_{k}(f)\right|}{2^{2-a} C_{\theta}|k|^{s-j-a}} \\
& \leq \frac{2^{a-2}}{C_{\theta}}\left(\sum_{|k| \geq k_{\theta}} \frac{1}{|k|^{2(s-j-a)}}\right)^{1 / 2}\left(\sum_{|k| \geq k_{\theta}}|k|^{2 s}\left|c_{k}(f)\right|^{2}\right)^{1 / 2}
\end{aligned}
$$

which is finite as soon as $s>j+a+1 / 2$. The result follows immediately.
Remark 5. Note that if $f$ is such that $c_{k}(f)=0$ when $k$ is odd - resp. $c_{k}(f)=0$ when $k$ is even- then the same is true for $f_{1}, f_{2}$. Thus, for $j=1,2, f_{j}(\theta+1 / 2-t)=f_{j}(\theta+t)$-resp. $f_{j}(\theta+1 / 2-t)=-f_{j}(\theta+t)$.

We are now in position to prove the following result:
Theorem 4.6. Let $p \in[1,+\infty]$. Let $\theta_{1}, \theta_{2} \in[0,1)$ be such that $\theta_{1}-\theta_{2} \notin \mathbb{Q}$ and $\Lambda, \Lambda^{\prime} \subset \mathbb{N}$. If $\Lambda, \Lambda^{\prime}$ are $[-1,1]$-Müntz-Szász sequences, then the system $\left\{\cos ^{\lambda} 2 \pi\left(t-\theta_{1}\right), \lambda \in \Lambda\right\} \cup\left\{\cos ^{\lambda^{\prime}} 2 \pi\left(t-\theta_{2}\right), \lambda^{\prime} \in \Lambda^{\prime}\right\}$ is total in $X_{p}(0,1)$.

We will give two proofs of this theorem. The first one is "constructive" while the second one relies on the Hahn-Banach theorem. The advantage of the second one is that it is more illustrative for our conjecture.

Direct proof. Let $f \in X_{p}(0,1)$ and $\varepsilon>0$. There exists a trigonometric polynomial such that $\|f-P\|_{p}<\varepsilon$ (such polynomials can be given explicitly via Fejér sums). Write $P=P_{1}+P_{2}$ with $P_{1}, P_{2}$ given by Lemma 4.4 (those are again explicit).

Finally, as $\Lambda, \Lambda^{\prime}$ are $[-1,1]$-Müntz-Szász sequences, $P_{1} \in \mathcal{T}_{\infty, \Lambda, \theta_{1}}$ and $P_{2} \in \mathcal{T}_{\infty, \Lambda^{\prime}, \theta_{2}}$. Therefore, there exists $Q_{1}, Q_{2}$ two polynomials such that, if we set $\pi_{j}=Q_{j}\left(\cos 2 \pi\left(t-\theta_{j}\right)\right)$, then $\left\|P_{j}-\pi_{j}\right\|_{p}<\varepsilon$ ( $Q_{1}, Q_{2}$ can be explicitly given with the help of the constructive proofs of the Müntz-Szász theorem). But then

$$
\left\|f-\pi_{1}-\pi_{2}\right\|_{p} \leq\|f-P\|_{p}+\left\|P_{1}-\pi_{1}\right\|_{p}+\left\|P_{2}-\pi_{2}\right\|_{p}<3 \varepsilon
$$

as expected.
Indirect proof. If $\Lambda, \Lambda^{\prime}$ are $[-1,1]$-Müntz-Szász sequences and $f \in\left(\mathcal{T}_{p, \Lambda, \theta_{1}}+\mathcal{T}_{p, \Lambda^{\prime}, \theta_{2}}\right)^{\perp}$ then $f \in$ $\mathcal{T}_{p, \Lambda, \theta_{1}}^{\perp} \cap \mathcal{T}_{p, \Lambda^{\prime}, \theta_{2}}^{\perp}$. Lemma 4.1 then implies that

$$
\left\{\begin{array}{l}
f\left(\theta_{1}+t\right)+f\left(\theta_{1}-t\right)=0 \\
f\left(\theta_{2}+t\right)+f\left(\theta_{2}-t\right)=0
\end{array} .\right.
$$

This implies that $f=0$ (see e.g. [Sj, Le]). For sake of completeness, here is a simple proof: looking at Fourier coefficients, this system is equivalent to

$$
\left\{\begin{array}{l}
e^{2 i \pi k \theta_{1}} c_{k}(f)+e^{-2 i \pi k i \theta_{1}} c_{-k}(f)=0 \\
e^{2 i \pi k \theta_{2}} c_{k}(f)+e^{-2 i \pi k i \theta_{2}} c_{-k}(f)=0
\end{array} \quad, \forall k \in \mathbb{Z}\right.
$$

When $k=0$ we get $c_{0}(f)=0$. For $k \neq 0$, the system has determinant $2 i \sin 2 k \pi\left(\theta_{1}-\theta_{2}\right) \neq 0$ since $\theta_{1}-\theta_{2} \notin \mathbb{Q}$. Therefore $c_{k}(f)=0$ for every $k \in \mathbb{Z}$ and thus $f=0$.

We conjecture that the reverse of this theorem is true as well:
Conjecture 1. Let $p \in[1,+\infty]$. Let $\theta_{1}, \theta_{2} \in[0,1)$ be such that $\theta_{1}-\theta_{2} \notin \mathbb{Q}$ and $\Lambda, \Lambda^{\prime} \subset \mathbb{N}$. If the system $\left\{\cos ^{\lambda} 2 \pi\left(t-\theta_{1}\right): \lambda \in \Lambda\right\} \cup\left\{\cos ^{\lambda^{\prime}} 2 \pi\left(t-\theta_{2}\right): \lambda^{\prime} \in \Lambda^{\prime}\right\}$ is total in $X_{p}(0,1)$, then $\Lambda, \Lambda^{\prime}$ are $[-1,1]$-Müntz-Szász sequences.

We can now prove the following partial version of the conjecture:
Theorem 4.7. Let $p \in[1,+\infty]$ and $\Lambda \subset \mathbb{N}$. Let $\theta_{1}, \theta_{2} \in[0,1)$ be such that $\theta_{1}-\theta_{2}$ is irrational and a-approximable for some $a>0$. If the system $\left\{\cos ^{\lambda} 2 \pi\left(t-\theta_{1}\right), \cos ^{\lambda} 2 \pi\left(t-\theta_{2}\right): \lambda \in \Lambda\right\}$ is total in $X_{p}(0,1)$, then $\Lambda$ is $a[-1,1]$-Müntz-Szász sequence.

Proof. If $\Lambda$ is not a $[-1,1]$-Müntz-Szász sequence, then at least one of the series $\sum_{\lambda \in \Lambda_{e}} \lambda^{-1}$ or $\sum_{\lambda \in \Lambda_{o}} \lambda^{-1}$ diverges, say the first one.

Let $j>a+1 / 2$ and $f$ be a $\mathcal{C}^{j}$ 1-periodic function with zero-mean and assume that $f$ is not analytic in at least 5 points. Assume further that $c_{k}(f)=0$ when $k$ is odd.

Write $f=f_{1}+f_{2}$ be the decomposition given by Lemma 4.5. Note that, from the remark following its proof, $f_{1}, f_{2}$ satisfy $f_{\ell}\left(\theta_{\ell}+1 / 2-t\right)=f_{\ell}\left(\theta_{\ell}+t\right), \ell=1,2$.

Now, as functions in $\mathcal{T}_{p, \Lambda, \theta_{2}}$ satisfy $\varphi\left(\theta_{2}-t\right)=\varphi\left(\theta_{2}+t\right)$, and $f_{1}\left(\theta_{1}-t\right)=f_{1}\left(\theta_{1}+t\right)$, necessarily, $f_{1} \in \mathcal{T}_{p, \Lambda, \theta_{1}}$ since otherwise $f_{1}$ would be constant. Similarly, $f_{2} \in \mathcal{T}_{p, \Lambda, \theta_{2}}$. From Lemma 4.2 we therefore know that $f_{1}, f_{2}$ are analytic except possibly at two points each. Finally $f=f_{1}+f_{2}$ is analytic except at 4 points, a contradiction.

## 5. A link with Heisenberg Uniqueness Pairs

The original idea behind this paper comes from an other problem, namely the notion of Heisenberg Uniqueness Pairs recently introduced by Hedenmalm and Montes-Rodríguez [HMR] and further investigated e.g. in [JK, Le, Sj]:

Definition 2. Let $\Lambda \subset \mathbb{R}^{2}$ and $\Gamma$ a smooth curve. Then $(\Gamma, \Lambda)$ is a Heisenberg Uniqueness Pair if the only finite measure $\mu$ that is supported on $\Gamma$, absolutely continuous with respect to arc length on $\Gamma$ and such that $\left.\widehat{\mu}\right|_{\Lambda}=0$ is the measure $\mu=0$.

Take for instance $\Gamma=\{(\cos 2 \pi t, \sin 2 \pi t), t \in[0,1)\}$ to be the unit circle and $\Lambda$ to be a set of two lines through the origin, $\Lambda=\left\{\left(t \cos \theta_{1}, t \sin \theta_{1}\right), t \in \mathbb{R}\right\} \cup\left\{\left(t \cos 2 \pi \theta_{1}, t \sin 2 \pi \theta_{1}\right), t \in \mathbb{R}\right\}$, $\theta_{1} \neq \theta_{2} \in[0,1)$.

Then Lev [Le] and Sjölin $[\mathrm{Sj}]$ independently proved that $(\Gamma, \Lambda)$ is a Heisenberg Uniqueness Pair if and only if $\theta_{1}-\theta_{2}$ is irrational.

But, a measure $\mu$ that is supported on $\Gamma$ and absolutely continuous with respect to arc length on $\Gamma$ is determined by a density $f \in L^{1}(0,1)$ in the following way

$$
\langle\mu, \varphi\rangle=\int_{0}^{1} \varphi(\cos 2 \pi t, \sin 2 \pi t) f(t) \mathrm{d} t \quad f \in \mathcal{C}\left(\mathbb{R}^{2}\right)
$$

In particular, $\widehat{\mu}(\eta, \xi)=\int_{0}^{1} f(t) e^{2 i \pi(\eta \cos 2 \pi t+\xi \sin 2 \pi t)} \mathrm{d} t$ so that $\widehat{\mu}=0$ on $\Lambda$ means that

$$
\int_{0}^{1} f(t) e^{2 i \pi r \cos 2 \pi\left(t-\theta_{1}\right)} \mathrm{d} t=\int_{0}^{1} f(t) e^{2 i \pi r \cos 2 \pi\left(t-\theta_{2}\right)} \mathrm{d} t=0, \quad r \in \mathbb{R}
$$

In particular, differentiating $\lambda$ times with respect to $r$ at $r=0$ leads to

$$
\begin{equation*}
\int_{0}^{1} f(t) \cos ^{\lambda} 2 \pi\left(t-\theta_{1}\right) \mathrm{d} t=\int_{0}^{1} f(t) \cos ^{\lambda} 2 \pi\left(t-\theta_{2}\right) \mathrm{d} t=0 \tag{5.5}
\end{equation*}
$$

for each $\lambda \in \Lambda$. From Theorem 4.6, if (5.5) holds for every $\lambda$ in some $[-1,1]$-Müntz-Szász set, then $f=0$. Note that the fact that $f \in L^{1}$ and the fact that the circle $\Gamma$ is compact makes it easy to justify all computations.

The first author's paper [JK] contains many more examples of compact curves $\Gamma$ and pairs of lines $\Lambda$ such that $(\Gamma, \Lambda)$ is a Heisenberg Uniqueness Pair. Each such example leads to new pairs of functions for which the sufficient part of our Müntz-Szász Theorem 4.6 holds. However, the converse seems much more difficult to establish.

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[^1]:    *with the obvious abuse of notation when $y \notin J_{-}$or $y \notin J_{+}$.

