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Experimental study and modelling approaches for the thermal conductivity evolution of hydrating cement paste

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Thermal stresses due to hydration heat release within concrete structures can lead to early age cracking particularly in cases of massive structures. In order to prevent, minimise and avoid this, a good knowledge of the evolution of cement-based material heat transfer properties at early stages of hydration is essential. In order to address the problem of the evolution of the thermal conductivity of fresh cement paste, two models developed by the authors are discussed: first, the 'apparent' thermal conductivity is calculated, assuming the fresh cement paste to be a multiphase material, which makes it possible to consider the material chemical evolution. The second modelling approach is based on the theory of heat and mass transfer within non-saturated porous media. The second modelling results appear to fit the experimental results better.

Notation

a	air volume content	T	temperature: °C
A	parameter of Hashin and Shtrikman's model	V_i	volume fraction of compound i
c	mass of cement: kg	w_n	mass of non-evaporable water: kg
C	volumetric heat capacity: J m ⁻³ K ⁻¹	λ_0	thermal conductivity at normal state: W m ⁻¹ K ⁻¹
D_{va}	diffusion coefficient of water vapour in air: m ² s ⁻¹	λ	thermal conductivity: W m ⁻¹ K ⁻¹
D_{Tv}	diffusivity of vapour under temperature gradient effect: m ² s ⁻¹ K ⁻¹	∇T	temperature gradient: K m ⁻¹
f	dimensionless factor	λ_i	thermal conductivity of compound i : W m ⁻¹ K ⁻¹
J_v	vapour mass flow density: kg m ⁻² s ⁻¹	α	hydration degree
L_v	latent heat of water vaporisation: J kg ⁻¹	∇P	pressure gradient: N m ⁻³
m	mass of sample: kg	ρ	density: kg m ⁻³
M	molar mass: kg mol ⁻¹	λ^*	'apparent' thermal conductivity: W m ⁻¹ K ⁻¹
P	total pressure: N m ⁻²	ΔT_{cor}	corrected temperature increase: K
P_v	partial vapour pressure: N m ⁻²	ΔT_{meas}	measured temperature increase: K
P_{vs}	saturation vapour pressure: N m ⁻²	ΔT_{hydr}	temperature increase due to hydration: K
Q	heat flow density: W m ⁻²	φ	relative humidity
q	linear heat flow density: W m ⁻¹	λ_{diff}	thermal conductivity of vapour diffusion: W m ⁻¹ K ⁻¹
R	perfect gas constant: J mol ⁻¹ K ⁻¹	θ	water volume content
t	time: s	β	dimensionless factor

Superscripts

c	consumed
i	initial

Subscripts

a	air
c	cement

h	hydrates
m	hydrate gel micropores
w	water

Introduction

The hydration process of cement-based materials generates considerable heat release. Both exothermic and thermo-activated characteristics of the hydration reaction can cause cracking at early ages of the material and alter the durability of concrete structures, especially large ones.

The prevention of cracks requires a good understanding of the thermal behaviour of the material during hydration. The determination of concrete thermal properties in hardened state is extensively documented.^{1–5} In comparison, very limited studies have been conducted to examine fresh material, when heat release occurs.^{6–8} These studies, however, reveal contradictory results.

The choice in this paper to study pure cement paste makes it possible to link the evolution of the thermal properties directly to both chemical and structural evolutions of the material. The study presented in this paper focuses on the evolution of the thermal conductivity of fresh ordinary cement paste (OPC).

The hydration processes are complex because anhydrous cement is a heterogeneous material. All the chemical phases have specific hydration kinetics and interact with one another.⁹ From a structural point of view, the negative volume balance, between reaction products (mainly calcium silicate hydrates, noted C–S–H) and initial reactants, causes internal porosity and, therefore, introduces an internal water vapour phase.

The model of hydration, proposed by Powers and quoted by Neville¹⁰ and Taylor,¹¹ is a simplification of this process, which assumes a global cement hydration degree and a linear production of C–S–H products according to this degree. Considering cement paste as a multiphase material and assigning, to each phase, the values of its own thermal properties, it is possible to estimate the thermal conductivity evolution of the material as a function of the chemical reaction progress (classical Maxwell's problem). The second approach consists in taking the microstructure of the material (formation of non-saturated internal porosity) into account and in assuming evaporation–condensation mechanisms. Philip and De Vries' analysis,^{12–14} is used to quantify the impact of these mechanisms on the thermal conductivity variations of wet porous materials.

Both modelling approaches have been applied to the case of OPC. The computed results are compared with the experimental results previously obtained using the 'hot wire method'.

Experimental set-up

Material

The cement used is an ordinary type I Portland cement (CPA CEM I 52.5 HTS), containing 63.3% of C₃S, 17.9% of C₂S, 4.7% of C₃A and 5.6% C₄AF. The Blaine specific surface is 3350 cm² g^{−1}. The water-to-cement (w/c) ratio is 0.348. A 3 min mixing guarantees the homogeneity of the cement–water mixture. Immediately after mixing, the cement paste is poured into a PVC tube containing the two parallel wires.

Thermal conductivity measurements

'Hot wire method' and 'two wire method'. The method used here to determine the thermal conductivity of OPC is the 'two wire method'. It is an adaptation of the 'hot wire method'. The 'hot wire method' is a classical transient-type method.¹⁵ It consists in enforcing a heating rectangular pulse in a wire dived into the studied material. From temperature evolution versus time, thermal conductivity is determined using (Δ being some linear variations)

$$\lambda = \frac{q}{4\pi} \frac{\Delta \ln \left(\frac{t}{t - t_{\text{heating}}} \right)}{\Delta T} \quad (1)$$

where λ is the thermal conductivity of the material (in W m^{−1} K^{−1}); q is the heat flow provided to the 'hot wire' (in W m^{−1}); t is the time (in s); t_{heating} is the heating time (in s); and T is the temperature (in K).

The physical parameter measured here will be more properly called 'apparent' thermal conductivity of the material (λ^*) since, as demonstrated by Moyne *et al.*,¹⁶ this is the only value that transient-type methods can determine. The notion of 'apparent' thermal conductivity λ^* will be clarified below.

Fitting specially building materials, variants of the 'hot wire method' were achieved by Bastian and Khelidj;⁶ the same probe was used simultaneously as a heat transmitter and as a temperature detector.

In the 'two wire method', a second wire is embedded, parallel to the first. From the temperature evolution of this second (non-heated) wire versus time, it becomes possible to determine the volumetric heat capacity C (in J m^{−3} K^{−1}) of the material

$$C = \frac{q}{e\pi r^2} \left(\frac{\Delta t}{\Delta T} \right)_I \quad (2)$$

where r is the distance between the two wires (in m); and I the inflexion point of the 'temperature versus time' evolution of the non-heated wire.

Test device. The test device is schematically illustrated in Fig. 1. Two stainless steel wires ($\varnothing 0.5$ mm), insulated with a thermo-retractable sheath, are used as temperature sensors for measuring temperature variations. One of the wires is also used as a heat line source: every half hour, a 4 A current passes

Fig. 1. Schematic diagram of the test device used for the measurement of fresh cement paste thermal conductivity

through the wire. A transformer linked to a potentiometer supplies the electric current. The duration of heating steps is 60 s; between heating steps (and throughout the whole experiment for the non-heated wire), a much lower current (80 mA) is injected in both wires in order to generate the potential difference necessary to make voltage measurements. The evolution of the wire electrical resistance is determined by measuring voltage with a data logging system during heating and non-heating steps. From the previously determined 'resistance versus temperature' law of the wires, the evolution of their temperature and, therefore, of the 'apparent' thermal conductivity and the volumetric heat capacity of the material can be deduced. The theoretical development leading to the calculation of these thermal properties has been detailed by Bastian and Khelidj⁶ and by Carslaw and Jaeger.¹⁷

The exterior mould containing the fresh cement paste is a PVC tube 1200 mm long with a diameter of 160 mm. A wire tensioning system is placed on both ends of the pipe to maintain an identical mechanical state (50 N) of both wires, which are placed symmetrically in relation to the cylinder axis. The distance between both wires is 15 mm. Four other sensors are used

- (a) three thermocouples, placed respectively at 200, 400 and 600 mm from the top of the PVC tube
- (b) a relative humidity probe, which is inserted 25 h after the casting of the cement paste, inside an orifice bored with a steel rod ($\varnothing 5$ mm). This probe measures the decrease in the internal relative humidity due to the mixing water consumption inherent to the hydration process.

Once the paste has been cast, the mould is closed and sealed with silicone grease to prevent any mass exchanges between specimen and outside. The experiment is conducted in a temperature and relative

humidity controlled room ($T = 20 \pm 2^\circ\text{C}$, relative humidity = $50 \pm 5\%$).

Hydration degree measurements

The evolution of the most important properties of cement paste (mechanical resistance,¹⁸ chemical shrinkage,¹⁹ hydration heat,²⁰ etc.) are directly related to its degree of hydration. In order to establish a correlation between the evolution of the thermal properties of cement paste and the chemical changes induced by the hydration process the evolution of the hydration degree of cement paste is measured at the same time. For this purpose, prismatic specimens ($40 \times 40 \times 60$ mm) are prepared from the same batch as the cylindrical specimen. The method adopted here for the sealing of the prismatic samples consists in, first, wrapping the moulds with a plastic film in a damp cloth for 24 h and, then, in sealing the specimens with a resin layer and an aluminium adhesive tape after removal from the mould at 24 h.

The hydration degree is determined on the basis that the hydration of 1 g of anhydrous cement produces 0.23 g of non-evaporable water, a ratio calculated from Bogue's equation.²¹ The non-evaporable water corresponds to the loss on ignition (LOI) of the hydrating cement paste sample between 105 and 1050°C.²² Finally, we obtain the following equation for the calculation of the hydration degree

$$\alpha = \frac{w_n(t)}{0.23 \cdot c} \times 100 = \frac{w_n(t)}{0.23 \cdot m} \cdot (1 + w/c) \times 100 \quad (3)$$

where α is the hydration degree of cement paste (in %); $w_n(t)$ is the mass (in kg) of non-evaporable water measured at time t ; c is the initial mass (in kg) of anhydrous cement in the sample of cement paste; and m the mass of this sample (in kg).

The furnace used for the sample ignition is a type 6000 Thermolyne. For each measurement, three prismatic specimens are cut and two pieces (1 cm³) taken from each of these specimens to determine the hydration degree.²³

Figure 2 shows the evolution of the cement paste hydration degree as a function of time.

Fig. 2. Evolution of the hydration degree of cement paste (measured by loss on ignition) plotted against time

Results and discussion

Correction of the experimental results

As mentioned in the introduction, cement paste is characterised, at early age (between 5 and 10 h after mixing), by an important heat release. This exothermic phenomenon disturbs measurement because the real temperature evolution of the wires is a combination of the Joule heat and the heat generated by the chemical reactions. Since the equations for the determination of the thermal conductivity only take the power supplied by the electric current into account, a correction is necessary to eliminate the influence of the hydration heat. A pure additive effect between the power generated by the hydration process and the one supplied by the electric current is considered here. Consequently, correcting the experimental results simply consists in deducting the temperature increase ($\Delta T_{hydr}(t)$) due to the hydration heat from the measured wire temperature increase ($\Delta T_{meas}(t)$)

$$\Delta T_{cor}(t) = \Delta T_{meas}(t) - \Delta T_{hydr}(t) \quad (4)$$

where $\Delta T_{cor}(t)$ is the corrected temperature increase of the ‘hot wire’ (in K). $\Delta T_{hydr}(t)$ evolution is measured using the thermocouples embedded in the cement paste (see Fig. 3).

All the calculations, presented below, take this correction into account.

Evolution of the ‘apparent’ thermal conductivity

Two series of tests are carried out. The first series, conducted over a relatively long period (28 days), shows a variation of the ‘apparent’ thermal conductivity at early ages. In order to examine this variation, a second test series is performed over a much shorter period (the first 24 h after OPC casting). Figs 4(a) and 4(b) present the evolution of the ‘apparent’ thermal conductivity of cement paste. Three main phases are observed.

- (a) From 0 to 6 h after mixing, the ‘apparent’ thermal conductivity remains nearly constant, at about $1.0 \text{ W m}^{-1} \text{ K}^{-1}$.

Fig. 3. Evolution of the internal temperature of fresh cement paste plotted against time

Fig. 4. Evolution of the measured ‘apparent’ thermal conductivity of cement paste (corrected results) plotted against time: (a) during 28 days; (b) during the first 24 h after casting

- (b) At about 6 h, the ‘apparent’ thermal conductivity shows a maximum value, of about $1.4 \text{ W m}^{-1} \text{ K}^{-1}$.
- (c) After the peak, the ‘apparent’ thermal conductivity value decreases and finally stabilises at about $1.1 \text{ W m}^{-1} \text{ K}^{-1}$.

The final value of the ‘apparent’ thermal conductivity satisfactorily agrees with that of a previous study.²⁴ However, nowhere in the literature can it be found any information about a peak in the evolution of cement paste thermal conductivity at very early age.

In order to explain this phenomenon, two different modelling approaches have been developed and compared with the experimental results.

Modelling approaches

Linear chemical evolution model and classical Maxwell’s problem

This modelling approach consists in linking the evolution of the ‘apparent’ thermal conductivity directly to the chemical evolution of cement paste. The modelling principle is based on the classical Maxwell’s problem: on the assumption of only pure conductive heat transfers, the ‘apparent’ thermal conductivity λ^* of cement paste is estimated to be the equivalent thermal conductivity of a multiphase material

$$\lambda^* = g(\lambda_i(T), V_i(a)) \quad (5)$$

where λ_i and V_i are the thermal conductivity (in $\text{W m}^{-1} \text{K}^{-1}$) and the volume fraction of phase i , respectively, T is the temperature (in $^{\circ}\text{C}$); and g the different types of multiphase models (series/parallel model, Hashin and Shtrikman's model, for example).

Considering the reactive behaviour of OPC, the volume fraction evolution of each phase has to be quantified, first. In order to express this evolution as a function of hydration degree α , the Powers' model, quoted by Taylor,¹¹ is used. This model integrates several hypotheses.

- (a) The various and complex hydration reactions are represented by a global reactive balance (anhydrous cement + water $\rightarrow$ hydrates).
- (b) The hydration of 1 g of anhydrous cement consumes 0.23 g of water.
- (c) The absolute volume V_h of hydrates, formed by reactions, is lower than the sum of anhydrous cement, V_c^c and water, V_w^c , volumes consumed by hydration (Le Chatelier's contraction²⁵)

$$V_h = (V_c^c + V_w^c) - 0.254 \cdot V_w^c = V_c^c + 0.746 \cdot V_w^c \quad (6)$$

- (d) The hydrate structure contains micro-pores, whose volume V_m is expressed as

$$\frac{V_m}{V_m + V_h} = 0.28 \quad (7)$$

Finally, the following equation system is obtained, for one volume unit of cement paste

$$\begin{cases} V_c = \frac{1}{1 + \rho_c \cdot \frac{w}{c}} (1 - \alpha) \\ V_h = 1.539 \frac{1}{1 + \rho_c \cdot \frac{w}{c}} \cdot \alpha \\ V_w = V_w^i - V_w^c = \frac{\rho_c \cdot \frac{w}{c}}{1 + \rho_c \cdot \frac{w}{c}} - 0.722 \frac{1}{1 + \rho_c \cdot \frac{w}{c}} \cdot \alpha \\ V_a = 1 - (V_c + V_h + V_w) \end{cases} \quad (8)$$

where V_c , V_w , V_h , V_a are the anhydrous cement volume fraction, the water volume fraction, the hydrate volume fraction and the gaseous volume fraction, respectively, present in cement paste at $\alpha\%$ of hydration and ρ_c is the density of anhydrous cement

($\rho_c = 3.15 \text{ g mm}^{-3}$). Superscripts c and i are used for 'consumed' and 'initial', respectively.

Figure 5 shows the evolution of the volume fractions of hydration products and reactants calculated from equations (4–6) as a function of hydration degree, for a cement paste with a $w/c = 0.348$.

The thermal conductivity values of the different phases (anhydrous cement + hydrates, liquid water and air) are given in Table 1.

From equation (8) and the values of Table 1, it is possible to propose a framing of λ^* value. Two approaches are considered.

- (a) First, with disregard for the material structure, we have used the series/parallel models (Wiener bounds) whose lower limit is given by

$$\lambda_{\perp} = \left(\sum_{i=1}^n \frac{V_i}{\lambda_i} \right)^{-1}, \text{ (series model)} \quad (9a)$$

and upper limit by

$$\lambda_{\parallel} = \sum_{i=1}^n V_i \cdot \lambda_i, \text{ (parallel model)} \quad (9b)$$

- (b) Second, the structure of fresh cement paste is assumed isotropic and the more restrictive Hashin and Shtrikman's model is applied. The thermal conductivity λ^* of the multiphase material is bounded by

Fig. 5. Evolution of the volume fractions of hydration reactants and products computed using the model of Powers as a function of hydration degree cement paste with $w/c = 0.348$ between 0 and 50% of hydration

Table 1. Thermal conductivity values of the different phases of fresh cement paste

Phase	Thermal conductivity: $\text{W m}^{-1} \text{K}^{-1}$	Sources
Anhydrous cement and hydrates	2.8	Bouguerra ²⁷ (on crushed hydrated cement paste)
Water	0.6	Farouki ³⁰
Air	0.025	Farouki ³⁰

$$\lambda_l + \frac{A_l}{1 - \frac{A_l}{3\lambda_l}} \leq \lambda^* \leq \lambda_n + \frac{A_n}{1 - \frac{A_n}{3\lambda_n}} \quad (10a)$$

in which $\lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_n$ are the thermal conductivities of phases 1, 2, ..., n of the material, respectively

$$A_j = \sum_{i \neq j}^n \frac{V_i}{\frac{1}{\lambda_i - \lambda_j} + \frac{1}{3\lambda_j}} \quad (10b)$$

where λ_i and V_i are the thermal conductivity (in $\text{W m}^{-1} \text{K}^{-1}$) and the volume fraction of phase i of the material, respectively.

The results of the simulation, calculated from equations (9) and (10), are presented in Fig. 6. The comparison between simulated and experimental results reveals that

- (a) On the one hand, both models give a correct bounding of the linear part of the experimental curve. The Hashin and Shtrickman's model, furthermore, makes it possible to obtain a more precise interval, which confirms that fresh cement paste structure can be considered as isotropic.
- (b) On the other, both models do not account for the 'apparent' thermal conductivity peak observed at about 30% of hydration.

Philip and De Vries' analysis (PDVA)

As demonstrated by Krischer and Kröll,²⁶ the vapour diffusive transport contributes to heat transfer processes within wet porous media through the evaporation–condensation mechanism (water evaporates from the pore hot surfaces, then, migrates by gaseous diffusion process and condenses on the cold surfaces, thus transferring its latent heat vaporisation).

Considering the case of fresh cement paste, it is

Fig. 6. Determination of the interval to which the measured 'apparent' thermal conductivity (corrected results) belongs, using different multiphase models (series/parallel models and Hashin and Shtrickman's model) and comparison with the "apparent" thermal conductivity computed using PDVA ($f = 0.17$)

known that the hydration process affects the material microstructure doubly: first, through the creating of internal porosity; and second through the consumption of water by the chemical reactions. The combination of both phenomena leads to some phase changes (liquid water – vapour) and to the creation of interfaces (meniscuses) between liquid and water vapour. The decrease in the internal relative humidity measured during cement paste hydration is the result of the formation of these meniscuses (Fig. 7).

The classical Philip and De Vries' analysis (PDVA)^{12–14} allows for the consideration of the evaporation–condensation mechanism, which occurs within non-saturated porous media submitted to relatively high temperatures. This analysis has been already applied successfully on hardened civil engineering materials (wood concrete,²⁷ autoclaved aerated concrete,²⁸ for example). Originally, the Philip and De Vries' analysis can only be applied to non-reactive bodies.¹² In spite of this, we have chosen to use it with a view to examining whether it could give a correct order of magnitude for the experimental results obtained on hardening cement paste.

In the theoretical framework proposed by Philip and De Vries, the heat flow density, Q , inside wet porous media is written as

$$Q = -\lambda_0 \nabla T + L_v J_v \quad (11)$$

in which λ_0 is the 'pure' thermal conductivity (Fourier's law) in $\text{W m}^{-1} \text{K}^{-1}$; ∇T is the temperature gradient (in K m^{-1}); L_v is the water latent heat of vaporisation (in J kg^{-1}); and J_v is the mass flow of water vapour (in $\text{kg s}^{-1} \text{m}^{-2}$). The term $L_v J_v$ represents the evaporation–condensation effects. The amount of sensible heat transferred is negligible because the vapour density value remains low and the latent heat vaporisation value is high compared to the specific heat of water. The vapour flow density J_v can be expressed as

$$J_v = -f D_{va} \frac{P}{P - P_v} \frac{M}{RT} \nabla P_v \quad (12)$$

a specific application of Fick's law, in which P is the total pressure (in N m^{-2}), P_v the partial pressure of

Fig. 7. Evolution of the internal relative humidity of fresh cement paste

water vapour (in N m^{-2}); M its molar mass (in kg mol^{-1}); R the perfect gas constant (in $\text{J mol}^{-1} \text{K}^{-1}$); D_{va} the diffusion coefficient of water vapour in air (in $\text{m}^2 \text{s}^{-1}$); T the temperature (in K) and f a dimensionless factor taking the tortuosity effect of the porous medium into account. Assuming water vapour as a perfect gas and considering a very high moisture content (relative humidity $\varphi \approx 100\%$ during the first hours of hydration), then

$$P_v = \varphi \cdot P_{vs} \approx P_{vs} \quad (13)$$

where P_{vs} is the saturation vapour pressure (in N m^{-2}).

Therefore, when neglecting in first approximation the moisture rate gradient, equation (12) becomes

$$J_v = -f D_{va} \frac{P}{P - P_{vs}} \frac{M}{RT} \frac{dP_{vs}}{dT} \nabla T \quad (14)$$

When equation (14) is introduced into equation (11), we obtain

$$Q = - \left[\lambda_0 + f L_v D_{va} \frac{P}{P - P_{vs}} \frac{M}{RT} \frac{dP_{vs}}{dT} \right] \nabla T = -\lambda^* \nabla T \quad (15)$$

which is the classical form of Fourier's law, if the concept of 'apparent' thermal conductivity λ^* of the material is assumed.

$$L_v D_{va} \frac{P}{P - P_{vs}} \frac{M}{RT} \frac{dP_{vs}}{dT}$$

Dimensions are homogeneous with thermal conductivity ones. This term, which is generally noted λ_{diff} , represents the 'apparent' contribution of water vapour diffusion to heat transfer. Its consideration for the calculation of the 'apparent' thermal conductivity of wet porous materials can take different forms. De Vries¹⁴ advises the use of $(\lambda_{diff} + \lambda_a)$ for the effective thermal conductivity of the air component, where λ_a is the normal thermal conductivity of air whereas Bouguerra²⁷ and Laurent *et al.*²⁸ use the relationship presented in equation (15) and obtain the following expression for λ^*

$$\lambda^* = \lambda_0 + f \lambda_{diff} \quad (16)$$

The expression of J_v proposed by Crausse *et al.*²⁹ is slightly different since the authors explicitly consider the volume content a of the gaseous phase. They obtain the expression of the 'apparent' thermal conductivity λ^* from

$$\lambda^* = \lambda_0 + \rho_l L_v D_{Tv} \quad (17)$$

where

$$D_{Tv} = \beta a D_{va} \frac{P}{P - P_v} \frac{M}{RT} \frac{1}{\rho_l} \left[\varphi \frac{dP_{vs}}{dT} + P_{vs} \left(\frac{\partial \varphi}{\partial T} \right)_{\theta} \right] \quad (18)$$

in which β is a dimensionless factor taking the

geometrical influence of the porous media on vapour diffusion into account; and θ is the volume water content of the material. Considering, like previously, a constant relative humidity $\varphi \approx 100\%$, equation (18) becomes

$$D_{Tv} = \beta a D_{va} \frac{P}{P - P_v} \frac{M}{RT} \frac{1}{\rho_l} \frac{dP_{vs}}{dT} \quad (19)$$

Introducing equation (19) into equation (17), λ^* becomes

$$\lambda^* = \lambda_0 + \beta a L_v D_{va} \frac{P}{P - P_v} \frac{M}{RT} \frac{dP_{vs}}{dT} \quad (20)$$

The identification of the different parameters shows that factor f , found in equation (16), corresponds to factor βa used in equation (20). Considering the high moisture content rate, the only difference between both expressions of λ^* is the explicit introduction, by Crausse *et al.*,²⁹ of the value of volume air content a .

Experimental results^{16,27,28} show that f can be considered as a unique function of the material saturation rate S . In practice, f is determined by measuring λ^* at different temperatures and by assuming that, at low temperature (from 0 to 20°C), evaporation-condensation effects can be neglected (then equation (16) gives $\lambda^* = \lambda_0$). Therefore, factor f can be expressed as

$$f(S) = \frac{\lambda^*(S, T)_{measured} - \lambda_0(S)}{\lambda_{diff}(T)} \quad (21)$$

Fitting this to our case, we consider that

- (a) The value of λ_0 changes during the hydration process, since, in autogenous conditions, the saturation rate is a direct function of both hydration degree and initial water content of the cement paste. In first approximation, the evolution of λ_0 can be considered linear within the interval constituted by the 'apparent' thermal conductivity measured at $t = 1$ h ($T = 25.65^\circ\text{C}$; $\alpha = 12.74\%$), denoted λ_1^* , and the same parameter measured at $t = 73$ h ($T = 21.33^\circ\text{C}$; $\alpha = 58.88\%$), denoted λ_2^* . Therefore, λ_0 becomes

$$\lambda_0(\alpha) = \frac{\lambda_2^* - \lambda_1^*}{46.14} \cdot (\alpha - 58.88) + \lambda_2^* \quad (\alpha \text{ in } \%) \quad (22)$$

- (b) For a given material, the λ_{diff} value depends on the internal pressure of the gaseous phase, the temperature, and the relative humidity of the material.³⁰ At atmospheric pressure and for a relative humidity value close to 100%, the following expression, already used by Krischer and Kröll,²⁶ can be applied

$$\lambda_{diff} = 0.02213 \times \exp(0.054198 \cdot T) \quad (T \text{ in } ^\circ\text{C}) \quad (23)$$

where T is the temperature of the cement paste measured by the thermocouples embedded in the material (see Fig. 3).

- (c) Our experimental results cannot be used to

determine factor f . However, we know that, at an early age, the saturation rate of cement paste is high. Furthermore, the experiments carried out on hardened wood concrete and autoclaved concrete have shown that, within the range of high saturation rate, the value of f does not exceed 0.2. So, in order to achieve some simulations and compare them with our experimental results, we consider, in first approximation, that f remains constant. Different f -values are, then, simulated.

The experimental and simulated results, obtained from equations (22) and (23), are presented in Fig. 6.

It reveals that, for $f = 0.17$, PDVA gives a good order of magnitude of the experimental results.

However, a small difference between the experimental peak of the 'apparent' thermal conductivity evolution of fresh cement paste and the theoretical peak calculated by mean of PDVA appears.

Different reasons may be put forward to explain this difference. PDVA does not consider reactive and shrinking materials (it is the case of fresh cement paste). In this type of material, complex physico-chemical processes take place locally. Two phenomena in competition are in progress: first the hydration reactions, which correspond to a consumption of water; and second, the increasing of temperature, which leads to the evaporation of water in the internal porosity. Therefore, the quantity of water available for evaporation–condensation transfers is continuously modified. At the present time, it is not possible to take this chemical evolution of the material into account through the analysis developed by De Vries and Philip.

Moreover, moisture displacements inside the material have been neglected. These displacements, combined with the temperature gradient induced by the exothermic feature of cement hydration, may affect the 'apparent' thermal properties of the material.

Conclusion

This paper examined the thermal behaviour of fresh cement paste maintained in autogenous conditions shows a variation in its 'apparent' thermal conductivity at very early ages. From the comparison between the experimental results and the simulations presented in this paper, the following main conclusions can be drawn.

- (a) The models simply based on the linear chemical evolution approach (series and parallel model and Hashin and Shtrikman's model) cannot simulate the 'apparent' thermal conductivity peak observed at very early ages of the cement paste.
- (b) Considering the existence of evaporation–condensation phenomena due to the structural changes of the material, the Philip and De Vries' analysis^{12–14} has

been carried out and has made it possible to achieve a better agreement with the experimental results. However, a limited time lag between modelled and experimental results is observed.

- (c) Further experimental work is necessary, especially as regards isothermal conditions, to improve our knowledge of the phenomenon observed and get closer to previous experimental and theoretical studies carried out in isothermal conditions.¹⁶ For this purpose, a specific device is currently being developed within the laboratory. This device will be used to study smaller fresh cement paste samples under different selected isothermal curing temperatures. From a theoretical point of view, the integration of deformation and chemical reactions into the equation set may be achieved, by considering, for example, the larger framework of the thermodynamics field of non-saturated reactive porous media.³¹

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