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Determination and modelling of the cure shrinkage of epoxy vinylester resin and associated composites by considering thermal gradients

Yasir Nawaba,b, Xavier Tardif a, Nicolas Boyard a,1, Vincent Sobotka a, Pascal Casari b, Frédéric Jacquemin b

a Université de Nantes, CNRS, Laboratoire de Thermocinétique de Nantes, UMR 6607, La Chantrerie, rue Christian Pauc, BP 50609, 44306 Nantes cedex 3, France
b Université de Nantes, Institut de Recherche en Génie Civil et Mécanique, UMR CNRS 6183, 37 boulevard de l'Université, BP 406, 44600 Saint-Nazaire, France

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Quantification and understanding of the evolution of chemical shrinkage of thermoset polymers is of crucial importance for modelling of residual strains and stresses. Thermal properties of resin and the strong coupling between chemical reactions and thermal fields lead to non-negligible thermal and curing degree gradients in the part. In this paper, modelling of the volume chemical shrinkages of an unsaturated epoxy vinylester resin and associated glass fibre composites is proposed, by taking into account the coupling between volume variation and thermal gradients. Modelling is also compared with the measurements done with a home-made instrument (PVT–a). Results suggest that chemical shrinkage is non linear as a function of degree of cure. Moreover, for an equal mass of resin, chemical shrinkage of resin carrying fibres is lesser than the neat resin.

1. Introduction

During the curing cycle, thermoset composites undergo the residual strains and stresses. These latter depend mostly on the shrinkage behaviour of the polymer matrix, from the point where curing stresses cannot be relaxed anymore. The matrix shrinkage may have a thermal and/or chemical origin. The chemical shrinkage is a direct consequence of crosslinking of the thermosetting polymer and its characterisation is thus of crucial interest for modelling the residual stresses [1,2] and shape evolution of composite part during curing [3].

Several studies have been carried out to determine the curing based dimensional changes in thermoset resins using the various dilatometric techniques. These techniques include capillary-type dilatometer [4,5], plunger type dilatometer [6–8], gravimetric method [9,10], rheometer [11], pycnometer [11,12], fibre Bragg grating [13–17], thermomechanical analyser [18], dynamic mechanical analyser [19,20], ultrasonic waves [21], using optical fibres, etc.

Since chemical shrinkage is expressed as a function of degree of cure, another instrument (DSC in most cases) was used to determine degree of cure. This tends to be prejudicial against a wholly accurate interpretation, since the experimental conditions may not exactly be the same. Furthermore, all the above mentioned instruments (except plunger type) work in low pressure, which does not match the industrial conditions.

Keeping in view the limitations of these instruments, a plunger type device called PVT–a mould was developed in our lab, which works in industrial like conditions (i.e., adjustable pressure up to 10 MPa in moulding cavity and temperature up to 200 °C), and is able to measure simultaneously the flux density hence the degree of cure, volume and temperature of mould surface for a given pressure. This instrument has already been used to determine the chemical shrinkage of thermoset resins [22–25].

In most of above mentioned studies, the chemical shrinkage was found by neglecting the thermal gradients, and evolution of shrinkage curve versus degree of cure was found linear. But in reality, thermal gradients exist in the thermoset resins due to the strong coupling between heat transfers and reaction kinetics, and due to low thermal diffusivity of resin. These gradients become more significant in case of thick parts and may change considerably the evolution of shrinkage curve. Therefore, there is a real need to quantify these gradients, and then determine the chemical shrinkage by considering them.

Studies on determination of volume chemical shrinkage of composites are found rare. For example, Russell [26] used GNONEX PVT device for determination of global volume chemical shrinkage of UD laminated (carbon/epoxy) composite. Problems faced were the toxicity of sample confining fluid (mercury), voids in the...
sample and inaccurate determination of degree of cure. Same device was also used to study the epoxy moulding compound with and without fillers [27]. Holst et al. [28] determined the volume chemical shrinkage of epoxy system containing fillers using a laser scanning dilatometer. Some authors also carried out some significant studies to find the in-plane chemical shrinkage of composites, such as Daniel et al. [29] which followed the warpage of 16 plies unidirectional glass/epoxy laminated plate. In a comprehensive study, Olivier [30] proposed a setup to determine the chemical shrinkage and thermo-mechanical mechanical properties of lami-nated composite samples.

In the present article, volume chemical shrinkage of vinylester resin was determined under high pressure and heating ramp conditions by considering the thermal gradients. Moreover, volume chemical shrinkage of glass/vinylester laminated composite was determined (under pressure, and heating ramp conditions) to study the effect of reinforcement on the chemical shrinkage while considering the thermal gradients. Some results on thermokinetic characterisation of the resin are also presented.

The volume measurements of resin and composite during curing were taken using PVT–α mould. The contributions of thermal expansion and chemical shrinkage to the experimental displacement curve during the curing cycle were then deduced from thermal and curing fields. These fields were computed by solving heat transfer and volume variation equation coupled to crosslinking kinetics, and by assuming a shrinkage model. The evolution of the cure shrinkage was then compared to experimental curves for which temperature gradients were neglected. The calculated shrinkage at the end of reaction is found very close to the experimental values.

2. Materials and methods

Vinylester resin (DERAKANE MOMENTUM 411-350 epoxy vinyl ester) was used with 0.05 wt% of initiator (Cobalt Naphthenate) and 1 wt% of Catalyst (NOROX MEKP-925H Methylmethyketone peroxide (MEKP)). These proportions were selected to get maximum gel time at room temperature (30 ± 10 min at 25 °C). Cross-ply glass/vinylester composite having two fibre volume fractions (32 vol% and 49 vol%) were selected to study the effect of the fibres on chemical shrinkage. Following thermal cycle was used for the curing of samples in PVT–α mould at the pressure of 0.65 MPa.

1. Equilibrating the system at room temperature and maintaining for 2 min.
2. Heating to 180 °C at the rate of 4 °C/min, and maintaining at 180 °C for 10 min.
3. Cooling to room temperature at 4 °C/min.

Studied samples have diameter equal to 40.5 ± 1 mm and thickness about 6 mm.

2.1. Description of PVT–α mould

PVT–α mould (Fig. 1) consists of a piston which is able to move in a cylindrical stainless steel cavity having a 61 mm internal diameter. The mould is designed so that heat transfers are 1D through the sample thickness. The heating is ensured by three heaters located at the top and bottom to heat the piston and cavity respectively. The mould is placed between two platens of an electric press. The position of the piston is adjusted in real time to keep the pressure constant. More details about this device are given in [31,32]. In this instrument, the experiments can be done on bulk samples.

Before starting PVT–α experiment, resin and the associated composites were sealed in a low modulus elastomer capsule, which is selected to avoid the jamming of the mould. Infusion of resin in capsule was done using vacuum assistance to remove air bubbles. Capsule was then placed in PVT–α cavity under the piston, which moves following the variations of the sample volume to get a constant pressure. These variations were recorded by LVDT-type displacement sensor with a precision of 1 μm and a limit of 10 mm. The heat flux was measured by two heat flux sensors [33] placed in the bottom of cavity and in the piston. Since the composite sample was placed in the deformable but incompressible elastomer cavity, the pressure is assumed to be hydrostatic and the surface of the elastomer, on which the pressure was applied, does not vary. As a consequence, the measured thickness is directly related to the sample volume and its variations.

![General view of PVT–α device.](image)
3. Thermo-physical characterisation of the resin

3.1. Determination and modelling of glass transition temperature

Glass transition temperature $T_g$ has a strong effect on thermo-physical properties and kinetics of reaction during isothermal curing or when the sample temperature reaches close to $T_g$ during curing. Therefore, it is important to determine and to model it versus the cure degree ($\alpha$). The glass transition temperature of raw resin ($T_g(0)$) and of fully cured resin ($T_g(\infty)$) were found equal to $-55.5\, ^\circ C$ and $118.9\, ^\circ C$ respectively by using DSC. The classical model developed by Pascalut and Williams [34] was used to describe the unique relationship (1) between $T_g$ and $\alpha$.

$$\frac{T_g(0) - T_g(\infty)}{1 - (1 - k) \cdot \alpha} = \frac{\alpha \cdot k}{(1 - (1 - k) \cdot \alpha)}$$  \hspace{1cm} (1)

where $'k'$ is a structure-dependent parameter and was found equal to 0.332 to fit the experimental data.

3.2. Characterisation of thermal conductivity

In the present study, guarded hot plate method [35] was used to measure thermal conductivity ($\lambda$) of cured resin between 25 $^\circ$C and 75 $^\circ$C. A second technique based on an optimisation approach, was also applied to determine the thermal conductivity of resin in both raw and cured resin states [36]. It consists in identifying the thermal conductivity from temperature measurements done at several locations through the sample thickness. The sample was placed between two heating plates, assuming that heat transfers are 1D. Temperatures recorded closest to the plates were used as boundary conditions, whereas the others were used for estimation. Identification was done by minimising a least square criterion using the conjugate gradient algorithm. The gradient of the criterion was determined by solving a set of adjoint equations. Thermal conductivity versus temperature results were fitted using linear regressions (2)-(4). The effect of the rubbery state ($T > T_g$) on thermal conductivity can be noted.

$$\lambda = 0.165 \text{ (W m}^{-1} \text{ K}^{-1}) \quad \text{(W m K}^{-1} \text{ raw resin, } 20 \text{ C < T < 70 C})$$ \hspace{1cm} (2)

$$\lambda = 1.41 \times 10^{-4} + 0.187 \quad \text{(W m K}^{-1} \text{ cured resin, with T in C, } T < T_g(\infty))$$ \hspace{1cm} (3)

$$\lambda = -1.63 \times 10^{-4} + 0.216 \quad \text{(W m K}^{-1} \text{ cured resin, with T in C, } T > T_g(\infty))$$ \hspace{1cm} (4)

3.3. Determination of heat of cure and specific heat

DSC (TA Q200 from TA Instruments$^\text{a}$) was used to measure these characteristic parameters of the resin. Classical dynamic scans (heating rate: 5 $^\circ$C/min, across a temperature range from -10 to 170 $^\circ$C) were performed under nitrogen purge. Sample weights ranged between 7 and 9 mg.

The total heat of cure was determined by integration of the exothermal peak area using a sigmoid integration line, resulting into average value (five experiments) of 324.9 $\pm$ 3.6 J/g. During the curing cycle, the resin passes from the raw (uncured liquid) to the cured rubbery and ultimately, to the cured glassy state (during cooling). $C_p$ in each of these three states was determined from DSC experiments using a sapphire calibration (5), according to the DIN standard 51007:

$$C_p = \frac{\Delta H}{\Delta T} = \frac{m \cdot dH}{m \cdot dT} = \frac{E}{T}$$  \hspace{1cm} (5)

Here $'E'$ is the calibration coefficient, whose values were determined from experiment on synthetic sapphire disc with known $C_p$. The variations of $C_p$ with temperature ($T$ in $^\circ$C) for all resin states were found linear and fitted accordingly

$$C_{p,\text{raw}} = 0.00073T + 1.8544 \quad \text{(J g}^{-1} \text{ K}^{-1}, \text{with } T \text{ in } ^\circ C)$$ \hspace{1cm} (6)

$$C_{p,\text{cured, glassy}} = 0.00497 + 1.1805 \quad \text{(J g}^{-1} \text{ K}^{-1}, \text{with } T \text{ in } ^\circ C), \text{if } T < T_g(\infty)$$ \hspace{1cm} (7)

$$C_{p,\text{cured, rubbery}} = 0.00177 + 1.8292 \quad \text{(J g}^{-1} \text{ K}^{-1}, \text{with } T \text{ in } ^\circ C), \text{if } T > T_g(\infty)$$ \hspace{1cm} (8)

The $C_p$ of solid cured resin during the glass transition was modeled from the values of $C_p$ in cured glassy state and that in cured rubbery state using:

$$C_{p,\text{cured}}(T) = C_{p,\text{glassy}}(1 - I) + I \cdot C_{p,\text{rubbery}}$$  \hspace{1cm} (9)

where $d$ is the half temperature range of glass transition ($17.5 \, ^\circ C$ in our case), and $\delta$ is a temperature factor to determine (found equal to 17 $^\circ C$ in our case).

Knowing the values of $\rho$ and $C_p$ in raw and cured states, a mixing law [37] was used to get $(\rho C_p)$ during the curing for any temperature and degree of cure.

$$(\rho C_p)_{\text{resin}}(T, \alpha) = \alpha \cdot (\rho C_p)_{\text{cured, rubbery}}(T) + (1 - \alpha) \cdot (\rho C_p)_{\text{raw}}(T)$$  \hspace{1cm} (10)

3.4. Characterisation of kinetics of reaction

Within the aim to compute the cure degree field through the sample during the curing process, the knowledge of reaction kinetics is of crucial importance due to its strong coupling with heat transfers. The determination of this field will be required later onto model the evolution of shrinkage.

A model of cure reaction should be able to describe the kinetics of reaction in material as well as the physical phenomenon like vitrification. A number of models have been reported for thermoset resins, which can be categorised as phenomenological and mechanistic. Phenomenological models can accurately describe the main features of the kinetic reaction ignoring the elementary reactions. Among these models, one can cite the catalytic models [38,39]. Prime [40] proposed a simply model, relating the kinetics of crosslinking with the degree of cure. Kamal and Sourour [41] proposed the first comprehensive model, used later on in several studies. Adolf [42,43] considered the effect of glass transition in his model. Bailleul [44] proposed a model stating the kinetic rate as a product of temperature function and degree of cure function. This model was also be extended to include the effects of diffusion associated to network formation and thus $T_g$ [1].

In our case, the resin cure is modelled by an empirical autocatalytic kinetic equation that describes polymerisation rate through a combination of Arrhenius and polynomial functions proposed by Bailleul [44]. It is also assumed that the Arrhenius relationship and the diffusion rate act in parallel, leading to

$$\frac{dz}{dt} = W(v(T)) \cdot K(T) \cdot G(z) \cdot (F_{\text{diffusion}}(z))$$  \hspace{1cm} (11)

$W(v(T))$ is a function corresponding to the induction of reaction. Note that it can also be used to consider an inhibition period when inhibitors are added to the resin. It is defined as $W(v(T)) = 0$ if $v(T) \geq 0$, $W(v(T)) = 1$ if $v(T) < 0$, where $v(T)$ is expressed as follows:

$$v(T) = t_{\text{ref}} - \int_{0}^{t_{\text{ref}}} \exp \left[ -B \cdot \left( \frac{T - T_{\text{ref}}}{T_{\text{ref}} - 1} \right) \right] dt$$  \hspace{1cm} (12)

$K(T)$ is an Arrhenius function that describes the dependence of rate of reaction on temperature.
and Poehlein found in a satisfactory agreement. Calculated and recorded temperatures were located in the middle of a resin sample and the calculated one parameters (Table 1) were determined by using CureKinetics model and model curves, plotted for five heating rates. The model a

diffusion control is sensitive to due to network formation. The diffusion becomes the controlling factor when a

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cure and the associated degree of cure a

determination became too slow to observe any effect on the heat flux curve). However as the tempera-

ture increased, the reaction started resulting in exothermal peak. At about 118 °C, the resin seemed to be fully cured (as a

1.9107 × 10⁻², G4 = 6.6887 × 10⁻³

Fdiffusion a = 0.99, E1 = 8.75

F(T) = Kref exp \[ -A \left( \frac{T_{\text{ref}}}{T} - 1 \right) \] \tag{13}

G(x) is a polynomial function (14) that determines the dependence of reaction rate on the degree of cure.

\[ G(x) = \sum_{i=0}^{n} G_i x^i \] \tag{14}

Finally, the diffusion-controlled effects are included in the model using the semi-empirical relationship (15) proposed by Chern and Poehlein [45].

\[ F_{\text{diffusion}}(x) = \frac{1}{1 + \exp(E_1(x - \alpha_c))} \] \tag{15}

Fdiffusion is the function used to describe the diffusion control in cure reaction associated to the decrease in mobility of molecules due to network formation. The diffusion control is sensitive to the absolute value |T - Tg(x)| [34]. In Eq. (15), it is considered that diffusion becomes the controlling factor when a approaches \( \alpha_c \). Here \( \alpha_c \) and \( E_1 \) are parameters to be determined.

Fig. 2 shows the reaction rates versus time graph of experimental and model curves, plotted for five heating rates. The model parameters (Table 1) were determined by using CureKinetics® software. This model has been validated by comparing the temperature recorded by a 80 µm-diameter thermocouple (K-type) located in the middle of a resin sample and the calculated one using (17a)–(17d). Calculated and recorded temperatures were found in a satisfactory agreement.

### 4.1. Interpretation of PVT–a curves

Vinylester resin used in the study is a low-temperature polymerisation resin so care was taken to start PVT–a experiment quickly after the resin preparation, but still it took 4–6 min. In order to take into account the possible heat losses of the mould, thermal expansions of the LVDT sensor, the elastomer and the mould, an aluminium disc with known properties was placed in the elastomer capsule to be submitted to the same thermal cycle as the sample. This data was then used as “reference” experiment. After data treatment, the total heat flux density (\( \phi \)) and the thickness variation purely associated to resin were obtained.

In Fig. 3, a typical evolution of the resin thickness (representative of volume) during the curing cycle, sample surface temperature and the associated degree of cure a (calculated from heat flux density data) are plotted versus time. During the heating cycle, this thickness variation is the result of thermal expansion and chemical shrinkage, while the only prominent phenomenon during cooling cycle was the thermal contraction.

Linear (A–B) part of the thickness curve represents the thermal expansion of neat resin. As the reaction starts at the low temperature, this region is very small and it is difficult to determine CTE (coefficient of thermal expansion) of raw resin accurately. Between the points B and C (curing zone), the chemical shrinkage of resin also appeared in addition to thermal expansion, and tended to decrease the sample thickness up to the point C where the resin was close to be fully cured (as \( a = 1 \)). A linear thickness curve from point C to D shows that the only prominent effect was again the thermal expansion.

When the thickness is plotted versus temperature, slopes of part AB and CD of that curve gives the CTE of raw and cured rubbery resin respectively.

In order to get the value of thermal expansion coefficient CTE at each time step in the crosslinking zone (B–C), a mixing law [23] (16) was used.

\[ \text{CTE}(a) = \text{CTE}_{\text{raw}} \cdot (1 - a) + a \cdot \text{CTE}_{\text{rubbery}} \text{ (K}^{-1} \text{, during crosslinking phase) } \tag{16} \]

In the beginning, when the reaction was not started, an endothermic effect was observed on the heat flux curve (Not shown for brevity) due to the heating of the sample. However as the temperature increased, the reaction started resulting in exothermal peak. At about 118 °C, the resin seemed to be fully cured (or reaction became too slow to observe any effect on the heat flux curve). In the isothermal step, flux curve returned to zero level showing that there was no more thermal or chemical phenomenon. Total heat of cure (327 ± 4.1 J/g) was determined by integrating the heat flux curve using sigmoid base line, which is found in agreement with DSC results (324.9 ± 3.6 J/g).

### 4.2. Effect of thermal gradients on macroscopic shrinkage during the crosslinking

To deduce the exact contribution of thermal expansion and chemical shrinkage to the experimental displacement curve, the temperature and cure degree fields (which depend on the sample thickness variation and coupled to reaction kinetics) have to be computed simultaneously during the curing cycle. Within this context, we described heat transfer and volume variation (17a) coupled to crosslinking kinetics (11). The used domain is cylindrical. Heat transfer is through the thickness, and the sample is assumed to be isothermal along the radius at a given thickness. Since \( \mu, C_p, \lambda \) depend on the temperature and cure degree, the problem to be solved is non-linear. Initial temperature field (17d) is known, and boundary conditions are of the third kind (17b)–(17c). The problem is given by the following set of equations, where \( h \) is the inverse of the thermal resistance induced by the silicon capsule and thermal contact, and \( T_{\text{mould}}, T_{\text{mouldy}} \) are the surface temperatures of the mould cavity and piston respectively.
\[ \rho C_p(T, x) \frac{\partial T}{\partial t} = \frac{\partial}{\partial x} \left( \lambda(T, x) \frac{\partial T}{\partial x} \right) + \rho(T, x) \cdot \Delta H \cdot \frac{\partial \alpha(T)}{\partial t} \cdot x \]
\[ \in [0, L], \forall t > 0 \]  
(17a)

\[ -\frac{\partial T}{\partial x} \bigg|_{x=0} = h(T_{mould 1} - T), \quad x = 0, \quad \forall t > 0 \]  
(17b)

\[ -\frac{\partial T}{\partial x} \bigg|_{x=L} = h(T - T_{mould 2}), \quad x = L, \quad \forall t > 0 \]  
(17c)

\[ T(0, x) = T_{ini}(x) \]  
(17d)

In order to calculate the temperature at each step of thickness, the Crank–Nicholson method was used to discretize 1-D heat conduction equation. The sample was supposed to compose of a number of slices (meshes). The volume of the mesh (an element of discretization) changes due to chemical shrinkage and thermal expansion. Therefore, the problem is a moving boundary one with constant mesh mass. Since temperature is varying through the thickness, all the meshes may not necessarily have the same volume. The volume variation of an elementary volume mesh \( dV_i \) can be expressed by:

\[ \frac{1}{d\rho_i(T, x)} = dV_i(T, x) = dV_0(1 + \text{CTE}(x) \cdot \Delta T + \text{CCS} \cdot \Delta \alpha) \]  
(18)

With \( dV_0 \) initial volume, and \( \Delta T \) and \( \Delta \alpha \) the temperature and degree of cure variations between two time steps. Since the resin was in the liquid/gel state during curing and a hydrostatic pressure was applied over the surface, all the changes in volume appear along the thickness, and radius was assumed to be constant.

Assuming that the cure shrinkage is a linear function of \( \alpha \) (observed when shrinkage was calculated by neglecting thermal gradients and plotted versus degree of cure, as shown in Fig. 7) in each mesh of the neat resin and composite samples, we thus define a coefficient of chemical shrinkage (CCS) to describe this linear relationship. Both coefficients (CCS and CTE) were then used to model the total thickness variation of each mesh such that the difference between modelled and experimental thickness was minimum. By iterative calculation, thickness of each mesh was calculated according to its temperature and degree of cure, and then summed to get final thickness at each time step.

In Fig. 4, experimental displacement curve and the modelled one for the neat resin are plotted. Thermal expansion curve represents the thickness of sample including the thermal expansion. Chemical shrinkage curve corresponds to the thickness evolution induced by resin crosslinking.

According to our calculations, the temperature difference between the surface and the core of the sample during curing reached up to 20°C (for 6 mm thick sample), which is significant. This difference will become higher for the thicker sample. The peak observed on the “thermal expansion” curve corresponds to exothermy of reaction and shows that present study is closer to reality as compared to literature, since it has not been presented before. The thermal gradient induces the gradient of the degree of cure as well, which will ultimately induce the gradients of cure dependent properties. Therefore, surface and core degree of cure versus time are also plotted on the same graph. The difference between the surface and the core degree of cure reached up to 15% of total degree of cure \( \Delta \alpha = 0.15 \), as this parameter is strongly coupled to
the temperature field. Therefore, evolution of shrinkage versus degree of cure will be different at different points in the thickness.

The final value of computed volume shrinkage for resin was found equal to −7.14 vol%, which is in agreement with literature shrinkage value [25] found by neglecting the gradients.

The composite experiments were also treated in the same way, and the results are presented in Figs. 5 and 6 for 32 vol% and 49 vol% fibre volume fraction respectively. The volume shrinkage induced by the crosslinking was found equal to −4.40% for the composite with 32 vol% of fibres, and −2.76% for that containing 49 vol% of fibres. The lower volume fraction of resin leads to the lower temperature and degree of cure gradients in composite samples, as expected. The proposed shrinkage model, coupled to thermal effects can thus be used to simulate dilatometric behaviour of both resin and associated composites.

Classically, the shrinkage is plotted versus degree of cure [23,25,28]. Therefore, it was also interesting for us to analyse the evolution of the volume shrinkage versus the degree of cure. It is previously demonstrated that the degrees of cure as a function of time are different through the sample thickness (Fig. 4–6). Therefore, we plotted the chemical shrinkage (for each sample) versus the mean cure degree (solid curves), surface cure degree (big dotted curves), and cure degree in the core of the part (small dotted curves) as shown in Fig. 7.

Regarding resin, a significant difference between the evolution of shrinkage curves plotted versus surface and core degree of cure can be observed. The difference of magnitude of shrinkage between these curves reaches 1% (14% of total chemical shrinkage) at $a = 0.65$, which is significant. For comparison with literature results [23,25], a curve for resin, calculated by neglecting the thermal gradients, is also plotted in Fig. 7 (dashed curve). This curve evolves in the similar way to the shrinkage curve versus core degree of cure (calculated considering the gradients). The results show that the shrinkage curve versus degree of cure is not linear when thermal gradients are taken into account (for any of three cases). Moreover, it can be noted that dotted curves act as upper and lower bounds of the solid line curve. As thermal gradients become lower in the presence of fibres (amount of resin decreases), these bounds become narrower for composites. This behaviour continues with increasing the fibre volume fraction.

Taking into account thermal gradients through the thickness of the composite sample, we can observe that the volume shrinkage of composites can be simply modelled assuming a linear behaviour of the volume variation of the matrix shrinkage in each element of the discretization scheme. Compared to the resin sample, the curves associated to composites have a more linear evolution which can be attributed to the fact that the temperature gradients are rather low when the volume fibre fraction increases. Even if anisotropic shrinkage exists, it is neither highlighted from the experimental nor from the modelled volume variations.

To study the effect of presence of fibres on the chemical shrinkage, shrinkage of each composite sample (found experimentally) was converted to shrinkage of an equivalent neat resin sample. This equivalent shrinkage of composite with 32% and 49% fibres were found equal to −6.55% and −5.51% respectively. It can be noted that for the same mass of resin, shrinkage of resin without fibres is higher than the shrinkage of resin with the fibres. It means that presence of fibres hindered the resin to shrink chemically during curing. The effect becomes prominent by increasing the fibre fraction.

Using the chemical shrinkage of resin and fibre fraction, volume chemical shrinkage of composite can be found by using:

$$\text{Shrinkage}_\text{composite} = \left( \text{Shrinkage}_\text{resin} - \text{fibre vol}\% \right) \times H \quad (19)$$

where $H$ is a hindrance factor, which can be found experimentally. Its value depends on type of reinforcement and resin. In the present study, its value was found equal to 0.03.

5. Conclusion

In this study, mould PVT–$\alpha$ was used to determine the volume chemical shrinkage of vinylester resin and glass/vinylester composites [0/90] with two fibre volume fractions, during the curing by taking into account the thermal gradients. Firstly, the thermophysical properties and the kinetics of the resin were determined and validated. Thermal and cure degree gradients were then computed to calculate the exact thermal expansion and chemical shrinkage contributions (which are coupled) to the sample volume variation during the curing cycle (heating step). It was done by assuming a linear model to describe the shrinkage evolution and by solving a heat conduction equation coupled to crosslinking kinetics model and a moving boundary problem. The large amount of heat of reaction, which was released in a short time range, led to a thermal expansion peak and temperature and cure degree gradients through the sample thickness for neat resin. The final resin chemical shrinkage value was found to be equal to −7.1 vol%. These observations were attenuated for composites, since the volume fraction of resin was lower. Moreover, the graph of the calculated shrinkage versus the mean cure degree showed that curves are not fully linear. This seems consistent when the thermal and cure degree gradients exist in the sample.
Finally, it was found that for an equal mass, chemical shrinkage of resin carrying fibres is lesser than the shrinkage of neat resin. This result can lead to the assumption that a part of the shrinkage has been hindered by the fibres which will reside as residual stresses in the composite part.

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