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FROM DYNAMICAL SYSTEMS TO RENORMALIZATION

FRÉDÉRIC MENOUS

ABSTRACT. We study in this paper logarithmic derivatives associated to derivations on graded complete Lie algebra, as well as the existence of inverses. These logarithmic derivatives, when invertible, generalize the exp–log correspondence between a Lie algebra and its Lie group. Such correspondences occur naturally in the study of dynamical systems when dealing with the linearization of vector fields and the non-linearizability of a resonant vector fields corresponds to the non-invertibility of a logarithmic derivative and to the existence of normal forms. These concepts, stemming from the theory of dynamical systems, can be rephrased in the abstract setting of Lie algebra and the same difficulties as in perturbative quantum field theory (pQFT) arise here. Surprisingly, one can adopt the same ideas as in pQFT with fruitful results such as new constructions of normal forms with the help of the Birkhoff decomposition. The analogy goes even further (locality of counter terms, choice of a renormalization scheme) and shall lead to more interactions between dynamical systems and quantum field theory.

1. INTRODUCTION.

Since the work of A. Connes and D.K Kreimer (see [3], [4]) in perturbative quantum field theory (pQFT), it has been possible to have a purely algebraic interpretation of some renormalization schemes, as the Birkhoff decomposition of regularized characters (with values in a commutative algebra of Laurent series), that is of elements of the group of algebra morphisms from a graded commutative Hopf algebra to the algebra of Laurent series (see [17], [5]). Many graded commutative Hopf algebras (shuffle, quasi shuffle, tree) do appear in this framework and, surprisingly, the same objects (groups of characters on Hopf algebras) arise in the study of analytic dynamical systems (see for example [10]), where changes of coordinates can be computed as elements either of the group of formal identity-tangent diffeomorphisms (namely characters over the Faà di Bruno Hopf algebra) or of some subgroup that corresponds to characters over a combinatorial Hopf algebra of trees or words (using mould calculus, see [10], [14]).

Such interactions should allow to enrich each domain with the ideas of the other and some steps have already been done in this direction (see [21] or [20]).

The origin of this paper comes from the observation that, in pQFT, the need for renormalization comes from the ill-definedness of some character on a Hopf algebra, which, after dimensional regularization, is replaced by a character with values in an algebra of Laurent series. The attempted character is then obtained with the help of the Birkhoff decomposition. The same phenomenon occur in dynamical systems, in the study of normal forms of vector fields (when the change of coordinates that should linearize the vector field is ill-defined). Apart from some remarks in the literature (see for example [15]), there does not seem that renormalization schemes have been used to compute, for example, normal forms.

We shall explain in this paper in an abstract algebraic context – a common framework to pQFT and dynamical systems – how the ideas developed in pQFT (dimensional regularization, Birkhoff decomposition ...) provide a way to compute normal forms as well as other objects that play a crude role in dynamics.

In the next section, we briefly recall the definition of complete graded Lie algebras and their link with commutative Hopf algebras. We define then in section 3 graded derivations on Lie algebras and their associated logarithmic derivative. Under some invertibility condition on such derivations, we define the inverse of such logarithmic derivative (see also section 4) that generalizes the log–exp correspondence between a Lie algebra and its Lie group. Whenever such a derivation is not invertible, the theory of dynamical systems (see appendix 9), provide, in the framework of Lie algebra, a way to study these

logarithmic derivatives, using d -conjugacy and "normal" forms (see section 5). There are surprising similarities with the need for renormalization in pQFT, the same ideas –regularization and Birkhoff decomposition – provide a new method to compute normal forms (section 6) and the analogy goes even further, see section 7, since such notions as the locality of counter terms or the choice among several "renormalization schemes" also appear here.

2. LIE AND HOPF ALGEBRAS.

In the sequel we consider complete graded Lie algebras L^k (as well as complete vector spaces or algebras) over a field k of characteristic 0. In other words, L^k is the completion $\prod_{n \geq 1} L_n^k$ of a graded Lie algebra $\bigoplus_{n \geq 1} L_n^k$ (with finite dimensional graded components L_n^k and $[L_n^k, L_m^k] \subset L_{n+m}^k$) with respect to the graduation. We write for short $x = \sum_{n \geq 1} x_n$ ($x_n \in L_n^k$) and the Lie bracket extends to the completion L^k : if $x = \sum_{n \geq 1} x_n$ and $y = \sum_{n \geq 1} y_n$ then $[x, y] = \sum_{n \geq 1} z_n$ with

$$z_n = \sum_{k=1}^{n-1} [x_k, y_{n-k}]$$

Under these hypotheses, the universal enveloping algebra $\mathcal{U}(L^k)$ of L^k is a unital, graded associative algebra with a graduation 0 component $k.1$. This algebra turns into a graded co-commutative Hopf algebra for the graded co-product defined on L^k by

$$\forall x \in L^k, \quad \Delta(x) = 1 \otimes x + x \otimes 1$$

that extends, thanks to the universal property of Lie algebras, to $\mathcal{U}(L^k)$ (see [2]). Using once again the graduation, we note $\mathcal{U}^k = \prod_{n \geq 0} \mathcal{U}_n^k$ the completion of $\mathcal{U}(L^k)$ and, as a graded map the co-product extends to a map from \mathcal{U}^k to the completion of $\mathcal{U}(L^k) \otimes \mathcal{U}(L^k)$ (that may only contain $\mathcal{U}^k \otimes \mathcal{U}^k$). \mathcal{U}^k is not necessarily a Hopf algebra but it contains L^k as the Lie algebra of *primitive elements* of \mathcal{U}^k , that is the elements x of \mathcal{U}^k such that $\Delta(x) = 1 \otimes x + x \otimes 1$, but also the Lie group G^k of group-like elements of \mathcal{U}^k , that is the elements x of \mathcal{U}^k such that $\Delta(x) = x \otimes x$.

Let us also remind that there is a one to one correspondence between L^k and G^k , using the well-defined exponential and logarithm:

$$\begin{array}{ll} \exp : \mathcal{U}_{\geq 1}^k \rightarrow 1 + \mathcal{U}_{\geq 1}^k & \log : 1 + \mathcal{U}_{\geq 1}^k \rightarrow \mathcal{U}_{\geq 1}^k \\ x \mapsto 1 + \sum_{n \geq 1} \frac{x^n}{n!} & 1 + y \mapsto \sum_{n \geq 1} (-1)^{n-1} \frac{y^n}{n} \end{array}$$

where $\mathcal{U}_{\geq 1}^k = \prod_{n \geq 1} \mathcal{U}_n^k$.

For a given commutative k -algebra \mathcal{A} , one can also consider the Lie algebra $L^{\mathcal{A}}$, which inherits grading of L^k and is defined as the completion, for the same graduation as in L^k , of the Lie algebra $\bigoplus_{n \geq 1} \mathcal{A} \otimes L_n^k$ with the Lie bracket

$$[a \otimes x, b \otimes y] = (ab) \otimes [x, y]$$

Note that $L^{\mathcal{A}}$ is both a vector space over k (whose graded component may be of infinite dimension) and a module over \mathcal{A} (with graded components of finite dimension) and one can as well define $\mathcal{U}^{\mathcal{A}}$ and $G^{\mathcal{A}}$.

There is no difficulty to extend iteratively this construction. For example, if $k \subset \mathcal{A} \subset \mathcal{B}$, then $L^k \subset L^{\mathcal{A}} \subset L^{\mathcal{B}}$ and the same inclusions hold for the completed enveloping algebras as well as for the groups. On the same way, we could define $(L^{\mathcal{A}})^{\mathcal{B}}$ as the completion of $\bigoplus_{n \geq 1} \mathcal{B} \otimes L_n^{\mathcal{A}}$. In the sequel, we will adopt this Lie algebra point of view but there is no difficulty to identify such Lie algebra $L^{\mathcal{A}}$ (resp. Lie groups $G^{\mathcal{A}}$) with the Lie algebra (resp. group) of infinitesimal characters (resp. characters) on a graded commutative algebra, with values in \mathcal{A} . The reader can refer to [6] on this matter. This simply means that the ideas stemming from the study of characters on Hopf algebras can be easily translated in the context of such Lie algebras.

Let us give, as an example, a "toy model" stemming from the theory of dynamical systems. When dealing with the simple differential equation:

$$y' = a(x)y^2 \quad \text{with } a(x) \in \mathbb{C}\{x\}$$

the Lie algebra (over \mathbb{C}) of differential operator (or vector fields),

$$(1) \quad L = \bigoplus_{n \geq 1} \text{Vect}_{\mathbb{C}}\{x^{n-1}y^2\partial_y\}$$

appears with a natural action on series of $\mathbb{C}[[x, y]]$ and the elements of the completion of L are formal operators $a(x)y^2\partial_y$ with $a(x) \in \mathbb{C}[[x]]$. For short we write $y^2\mathbb{C}[[x]]\partial_y$ this Lie algebra. Note that it coincides, for a different graduation, with the Lie algebra $L^{\mathcal{A}}$ with $\mathcal{A} = \mathbb{C}[[x]]$ and $L^{\mathbb{C}} = L_1^{\mathbb{C}} = \text{Vect}_{\mathbb{C}}\{y^2\partial_y\}$.

As another example related to dynamical systems, consider the Lie algebra $L^{k[t]}$, whose elements can be written:

$$x = \sum_{n \geq 1} \sum_{l=0}^{N_n} t^l x_{n,l} \quad \text{with } N_n \in \mathbb{N}, x_{n,l} \in L_n^k.$$

The element $\varphi^t = \exp(tx) \in G^{k[t]}$ ($x \in L^k$) is the unique solution of the differential equation

$$\partial_t \varphi^t = \varphi^t . x$$

with the initial condition $\varphi^0 = 1$. This suggests to explore similar "differential" equations that may lead to new correspondences between Lie algebras and Lie groups, since such equations appear naturally in the framework of differential equations (see section 9).

In the case of our toy model, the differential operator $\exp(a(x)y^2\partial_y)$ is a substitution automorphism:

$$\forall f(x, y) \in \mathbb{C}[[x, y]], \quad \exp(a(x)y^2\partial_y).f(x, y) = f(x, \exp(a(x)y^2\partial_y).y)$$

with $\exp(a(x)y^2\partial_y).y = \frac{y}{1-a(x)y}$ and the map $\exp(a(x)y^2\partial_y) \mapsto \frac{y}{1-a(x)y}$ is an anti-isomorphism between the Lie group of $y^2\mathbb{C}[[x]]\partial_y$ and the group (for the composition) of formal identity-tangent diffeomorphisms

$$\left\{ (x, y) \mapsto \left(x, \frac{y}{1-a(x)y} \right) \right\}.$$

The reader can check that $\exp(ta(x)y^2\partial_y) = \frac{y}{1-ta(x)y} = \alpha^t(x, y)$ is the flow of the differential equation $y' = a(x)y^2$ (see section 9 for details).

For short, in the sequel L (resp. \mathcal{U} , G) will denote $L^{\mathcal{A}}$ (resp. $\mathcal{U}^{\mathcal{A}}$, $G^{\mathcal{A}}$). Note also that any k -linear endomorphism u of \mathcal{B} (resp. graded k -linear endomorphism of L) can be extended to a k -linear graded endomorphism of $L^{\mathcal{B}}$, still noted u when there is no ambiguity, setting $u(b \otimes x) = u(b) \otimes x$ (resp. $u(b \otimes x) = b \otimes u(x)$). In order to avoid ambiguities, whenever such an endomorphism is defined on possibly many k -vector spaces (or algebras) V , we will write the image $u(V)$ and the corresponding kernel

$$\ker_V u = \ker u \cap V \subset V.$$

3. DERIVATIONS AND d -LOGARITHMS IN LIE ALGEBRAS.

3.a. Definitions and first properties.

Definition 1. A derivation d is a graded k -linear endomorphism on L (L^k or $L^{\mathcal{A}}$) such that

$$\forall x, y \in L, \quad d([x, y]) = [d(x), y] + [x, d(y)]$$

From the universal property of Lie algebras, as d is graded ($d(L_n) \subset L_n$) it extends to a graded derivation on \mathcal{U} :

- i. $d(1) = 0$,
- ii. $d(\mathcal{U}_n) \subset \mathcal{U}_n$ and
- iii. $\forall x, y \in \mathcal{U}, \quad d(x.y) = d(x).y + x.d(y)$.

For such a derivation, we can define and study the “differential” equation

$$(2) \quad d(\varphi) = \varphi.x$$

Proposition 1. *Any such derivation defines a map:*

$$(3) \quad \begin{array}{ccc} \log_d & : & 1 + \mathcal{U}_{\geq 1} \rightarrow \mathcal{U}_{\geq 1} \\ & & \varphi \mapsto \varphi^{-1}.d(\varphi) \end{array}$$

that sends G on L .

The map is well defined since any $\varphi = 1 + \sum_{n \geq 1} \varphi_n \in 1 + \mathcal{U}_{\geq 1}$ is invertible for the product in \mathcal{U} with an inverse

$$\varphi^{-1} = 1 + \sum_{s \geq 1} (-1)^s \sum_{n_1, \dots, n_s \geq 1} \varphi_{n_1} \dots \varphi_{n_s}$$

and, thanks to the completion, \log_d is well-defined in the complete algebra \mathcal{U} . It remains to prove that it maps G on L and this results comes from the Magnus formula (see [8]):

Lemma 1. *Let $\varphi = \exp(\alpha) \in G$ ($\alpha \in L$), then*

$$(4) \quad \log_d(\varphi) = \sum_{s \geq 0} \frac{(-1)^s}{(s+1)!} \text{ad}_\alpha^s(d(\alpha)) = \frac{e^{-\text{ad}_\alpha} - 1}{-\text{ad}_\alpha}(d(\alpha))$$

where $\text{ad}_\alpha(x) = [\alpha, x] = \alpha x - x\alpha$ is the adjoint action of α .

Since \mathcal{U} is complete, the above series converges in $\mathcal{U}_{\geq 1}$ when α is in $\mathcal{U}_{\geq 1}$. Moreover, since α is in L ($\varphi \in G$), $d(\alpha)$ and $\text{ad}_\alpha^s(d(\alpha))$ belong to L thus $\log_d(\varphi)$ is in L .

Note that we would have similar formulas for $(d(\varphi)).\varphi^{-1}$ ($(d(\varphi)).\varphi^{-1} = \varphi.\log_d(\varphi).\varphi^{-1}$) and, since $\varphi^{-1}.\varphi = 1$,

$$\begin{cases} \log_d(\varphi^{-1}) = \varphi.d(\varphi^{-1}) = -d(\varphi).\varphi^{-1} = -\varphi.\log_d(\varphi).\varphi^{-1} \\ \log_d(\varphi_1.\varphi_2) = \log_d(\varphi_2) + \varphi_2^{-1}\log_d(\varphi_1)\varphi_2 \end{cases}.$$

This logarithmic derivative is a kind of generalization of the logarithm and such d -logarithm do appear in the study of dynamical systems (see section 9).

For a given d -logarithm from G to L it is natural to ask if there exists an inverse.

3.b. The invertible case. Whenever the derivation d has a graded right inverse on L (L^k or L^A), there is a positive answer to the previous question:

Theorem 1. *If d admits a graded right inverse I on L then for any $u \in L$ there exists $\varphi \in G$ such that $\log_d(\varphi) = u$.*

Proof. Let $u \in L$. If we look for $\varphi = \exp(v)$ and set, thanks to the graduation,

$$v = \sum_{n \geq 1} v_n, \quad v_n \in L_n.$$

Thanks to lemma 1,

$$(5) \quad d(v_n) + \sum_{k=1}^{n-1} \sum_{i=1}^{n-k} \frac{(-1)^i}{(i+1)!} \sum_{n_1+\dots+n_i=n-k} \text{ad}_{v_{n_1}}(\text{ad}_{v_{n_2}} \dots (\text{ad}_{v_{n_i}}(d(v_k)))) = u_n$$

In L_1 we have,

$$d(v_1) = u_1$$

and we define $v_1 = I(u_1) \in L$. For $n = 2$,

$$d(v_2) - \frac{1}{2}[v_1, d(v_1)] = u_2$$

and we have the solution

$$v_2 = I(u_2) - \frac{1}{2}I([I(u_1), u_1])$$

and the proof follows recursively: $v = \sum_{n \geq 1} v_n$ is in L and $\varphi = \exp(v)$ is a solution of $\log_d(\varphi) = u \in G$. ■

If d is invertible on L then the solution φ of $\log_d \varphi = u$ is unique and noted

$$\varphi = \exp_d(u).$$

The reader can notice that in the case of L^k , since we assumed that the graded components L_n^k are of finite dimension, a derivation has a right inverse if and only if it is invertible. Otherwise,

Proposition 2. *For any derivation d , $\ker_L d$ is a sub-Lie algebra of L ($\{0\}$ when d is invertible) that gives rise to a subgroup of G :*

$$G^0 = \{\varphi \in G \mid d(\varphi) = 0\}$$

and φ_1 and φ_2 are two solutions of $\log_d \varphi = u$ if and only if $\varphi_2 \varphi_1^{-1} \in G^0$

Proof. Consider $\psi = \varphi_2 \varphi_1^{-1}$, that is $\psi \varphi_1 = \varphi_2$, then

$$\log_d(\psi \varphi_1) = \log_d(\varphi_1) + \varphi_1^{-1} \log_d(\psi) \varphi_1 = u + \varphi_1^{-1} \log_d(\psi) \varphi_1 = u$$

thus

$$\psi^{-1} d(\psi) = \log_d(\psi) = 0$$

and $d(\psi) = 0$. ■

Note that, using once again lemma 1, we get:

$$(6) \quad G^0 = \{\varphi = \exp(v) \in G \mid v \in \ker_L d\} = \exp(\ker_L d)$$

3.c. The Dynkin operator, the logarithm and other examples. A classical (and universal) example of invertible derivation in graded Lie algebras is the derivation associated to the graduation:

$$(7) \quad \forall n \geq 1, \forall x \in L_n, \quad Y(x) = nx.$$

In this case, \log_Y is the Dynkin operator, which is a bijection between L and G (see [6]).

Let us just point out another example that relates the Dynkin operator and the usual logarithm. For a given complete graded Lie algebra $L = L^A$ over a commutative k -algebra \mathcal{A} , consider the Lie algebra $L^{\mu^A[\mu]}$. This algebra inherits the graduation of L and the graded derivation $\mu \partial_\mu$ is invertible. If $x = \sum x_n \in L$, then μx and $\mu^Y(x) = \sum \mu^n x_n$ are in $L^{\mu^A[\mu]}$ and:

$$\begin{aligned} \exp_{\mu \partial_\mu}(\mu x) &= \exp(\mu x) \\ \exp_{\mu \partial_\mu}(\mu^Y(x)) &= \mu^Y(\exp_Y(x)) \end{aligned}$$

so that the classical exponential as well as \exp_Y are related to the same differential equation in $L^{\mu^A[\mu]}$:

$$\mu \partial_\mu \varphi = \varphi \cdot x$$

We will come back to the interaction between the logarithm and the Dynkin operator in the following section but let us end by the definition of a wide class of derivations that appear in dynamical systems. Suppose that L is a sub-algebra of a Lie algebra $L_0 \oplus L$ where L_0 is a Lie algebra of graduation 0, then for any $x_0 \in L_0$, the adjoint action of x_0 on L defines a graded derivation on L . Whenever this adjoint action is invertible, we have a unique solution to the equation

$$ad_{x_0}(\varphi) = [x_0, \varphi] = \varphi \cdot x$$

which is fundamental is the study of vector fields (see section 9). But conversely, the study of vector fields provides many ideas to deal with the case of non invertible derivations. Let us illustrate this on the toy model $L = y^2 \mathbb{C}[[x]] \partial_y$ with $L_0 = \{\alpha x \partial_x + \beta y \partial_y, \alpha, \beta \in \mathbb{C}\}$. The associated derivation $d = ad_{X_{\alpha, \beta}}$ (with $X_{\alpha, \beta} = \alpha x \partial_x + \beta y \partial_y$) gives for $n \geq 0$

$$\begin{aligned} d(x^n y^2 \partial_y) &= [\alpha x \partial_x + \beta y \partial_y, x^n y^2 \partial_y] = (\alpha n + \beta) x^n y^2 \partial_y \\ d(a(x) y^2 \partial_y) &= (\alpha a'(x) + \beta a(x)) y^2 \partial_y \end{aligned}$$

and the derivation is invertible if $\alpha n + \beta$ does not vanish for $n \geq 0$ or, equivalently, when, for any $b(x) \in \mathbb{C}[[x]]$, there exists a unique formal solution $a(x)$ to the equation

$$\alpha x a'(x) + \beta a(x) = b(x).$$

More precisely, for a fixed $X_{\alpha,\beta} = \alpha x \partial_x + \beta y \partial_y \in L_0$ and $y^2 b(x) \partial_y$, if we look for an element $\exp(a(x)y^2 \partial_y)$ in the Lie group such that

$$d(\exp(a(x)y^2 \partial_y)) = \exp(a(x)y^2 \partial_y) \cdot y^2 b(x) \partial_y$$

then, since any Lie bracket in $y^2 \mathbb{C}[[x]] \partial_y$ vanishes, Theorem 1 ensures that $\exp(a(x)y^2 \partial_y)$ is given by the solution of

$$\alpha x a'(x) + \beta a(x) = b(x)$$

From the dynamical systems point of view, the reader can check that, when identifying $\exp(a(x)y^2 \partial_y)$ to the formal diffeomorphism

$$\psi(x, y) = (x, z) = \left(x, \frac{y}{1 - a(x)y} \right),$$

the (formal) solutions of the differential system

$$\begin{cases} \frac{dx}{dt} = \alpha x \\ \frac{dy}{dt} = \beta y \end{cases}$$

are in one-to-one correspondence with the solutions of

$$\begin{cases} \frac{dx}{dt} = \alpha x \\ \frac{dz}{dt} = \beta z + b(x)z^2 \end{cases}$$

since

$$\begin{aligned} \frac{dz}{dt} &= \frac{d}{dt} \left(\frac{y}{1 - a(x)y} \right) \\ &= \frac{dy}{dt} \cdot \frac{1}{1 - a(x)y} + \left(a'(x)y \frac{dx}{dt} + a(x) \frac{dy}{dt} \right) \cdot \frac{y}{(1 - a(x)y)^2} \\ &= \frac{\beta y}{1 - a(x)y} + (\alpha x a'(x) + \beta a(x)) \frac{y^2}{(1 - a(x)y)^2} \\ &= \beta z + b(x)z^2. \end{aligned}$$

Depending on the choice of α and β , the associated derivation could, in the best case, be invertible on the complete enveloping algebra and, in the worst case (for example $\beta = 0$), be non-invertible on the Lie algebra.

4. THE “OVERALL” INVERTIBLE CASE AND ROTA-BAXTER ALGEBRAS.

If d is invertible as well on L as on $\mathcal{U}_{\geq 1}$, we will say that d is “overall” invertible and it is natural to transform the equation $d\varphi = \varphi x$ into the so-called Atkinson’s recursion:

$$(8) \quad \varphi = 1 + I(\varphi x)$$

where I is the inverse of d on $\mathcal{U}_{\geq 1}$. In this case,

$$\forall x, y \in \mathcal{U}_{\geq 1}, \quad I(x)I(y) = I(I(x)y + xI(y)).$$

Namely, I is a Rota-Baxter operator of weight 0:

Definition 2. Let A an associative algebra and R a linear operator on A . R is a Rota-Baxter operator of weight θ if:

$$\forall a, b \in A, \quad R(a)R(b) = R(R(a)b + bR(a)) + \theta R(ab)$$

In this case (see [8]), the equation (with perturbative parameter α):

$$(9) \quad \varphi_\alpha = 1 + \alpha I(\varphi_\alpha x)$$

admits a perturbative solution:

$$(10) \quad \varphi_\alpha = 1 + \sum_{n \geq 1} \alpha^n R^{[n]}(x)$$

with $R^{[1]}(x) = I(x)$ and $R^{[n+1]}(x) = I(R^{[n]}(x).x)$. This solution is in G and its logarithm, as well as its image by the Dynkin map \log_Y , are given by combinatorial formulas (see [8]). We proved in [19]:

Theorem 2. *Let δ a graded derivation that commutes with d (thus with I), the above perturbative solution satisfies:*

$$(11) \quad \log_\delta(\varphi_\alpha) = \varphi_\alpha^{-1} \cdot (\delta(\varphi_\alpha)) = \sum_{n \geq 1} \alpha^n \mathcal{L}_\delta^{[n]}(x)$$

with $\mathcal{L}_\delta^{[1]}(x) = I(\delta(x))$ et $\mathcal{L}_\delta^{[n+1]}(x) = I([\mathcal{L}_\delta^{[n]}(x), x])$.

The proof of this theorem is based on the fact that d is overall invertible. This is necessary to get the perturbative expansion of the solution $d\varphi = \varphi.x$ and then the formula 11. But this formula only involves the inverse of d on L and shall still hold if d is only invertible on L . This idea could certainly lead to a new proof of theorem 1, when d is invertible on L but not necessary on $\mathcal{U}_{\geq 1}$.

For example, in the case $L = L^k$, since the vector spaces (over k) L_n and thus \mathcal{U}_n are finite dimensional, for a given parameter ε , the regularized operator $d + \varepsilon Y$ is a derivation on the Lie algebra $L^{k[[\varepsilon]][[\varepsilon^{-1}]}$ where $k[[\varepsilon]][[\varepsilon^{-1}]}$ is the algebra of Laurent series in ε with coefficients in k .

On one hand, on each $\mathcal{U}_n^{k[[\varepsilon]][[\varepsilon^{-1}]}$ the restriction of $d + \varepsilon Y$ is invertible: $\mathcal{U}_n^{k[[\varepsilon]][[\varepsilon^{-1}]}$ is a finite-dimensional $k[[\varepsilon]][[\varepsilon^{-1}]$ -module and the determinant of the restriction of $d + \varepsilon Y$ to this module is a non zero Laurent series. The derivation $d + \varepsilon Y$ is thus invertible, of inverse I_ε , on $\mathcal{U}_{\geq 1}^{k[[\varepsilon]][[\varepsilon^{-1}]}$.

On the other hand, on each $L_n \subset L_n^{k[[\varepsilon]][[\varepsilon^{-1}]}$, since d is invertible, I_ε still makes sense when $\varepsilon = 0$ and coincides, for this value, with the inverse of d (say I_0).

The perturbative solution of $d\varphi = \varphi.x$ makes no sense since d is not necessarily overall invertible but, after the regularization, $d + \varepsilon Y$ is overall invertible and the perturbative expansion of the solution of $(d + \varepsilon Y)\varphi_\varepsilon = \varphi_\varepsilon.x$ makes sense, especially when x is in L^k . The formula 11, with δ commuting with $d + \varepsilon Y$ still hold. If we take $\delta = Y$ then the right part of this formula make sense when $\varepsilon = 0$ thus $\log_Y \varphi_\varepsilon$ and then φ_ε can be evaluated at $\varepsilon = 0$ and gives the solution in the invertible case! There is no contradiction here, this simply means that the expansion giving the perturbative solution may have poles, but once rewritten, for example in a basis of \mathcal{U} , poles in ε cancel out.

This " ε -regularization" turn a derivation in an overall invertible regularized derivation. Such regularizations appear in quantum field theory and will provide a way to deal with non-invertible derivations, see section 6.

5. THE NON-INVERTIBLE CASE.

5.a. **d -conjugacy.** Inspired by the situation in dynamical systems, whether d is invertible or not, it is relevant to study d -conjugate elements in L

Proposition 3. *Two elements in L are d -conjugate if there exists $\varphi \in G$ such that*

$$d(\varphi) + v.\varphi = \varphi.u$$

we say that φ d -conjugates u to v . This is an equivalence relation and we note $u \sim_d v$.

If φ d -conjugates u to v , then φ^{-1} conjugates v to u since:

$$d(\varphi^{-1}) = -\varphi^{-1}d(\varphi)\varphi^{-1} = -\varphi^{-1}(\varphi u - v\varphi)\varphi^{-1} = -u\varphi^{-1} + \varphi^{-1}v$$

On the same way, if φ d -conjugates u to v and ψ d -conjugates v to w , then $\psi\varphi$ d -conjugates u to w :

$$\begin{aligned} d(\psi\varphi) + w\psi\varphi &= d(\psi)\varphi + \psi d(\varphi) + w\psi\varphi \\ &= (\psi v - w\psi)\varphi + \psi(\varphi u - v\varphi) + w\psi\varphi \\ &= \psi\varphi u \end{aligned}$$

The idea of d -conjugacy comes from dynamical systems and is motivated by the following fact: On one hand, if, for example d is invertible on L , then there is only one d -conjugacy class since any u in L is conjugated to 0 by $\varphi = \exp_d(u)$. On the other hand, whenever d is not invertible, these equivalence classes are not trivial and their classification or description is of great interest, especially in the framework of dynamical systems.

The complete classification, up to d -conjugacy, shall, in many cases, be out of reach. For example, if $d = 0$. the equation reads

$$v\varphi = \varphi u.$$

Using the graduation and $\varphi = \exp(\alpha)$ ($\alpha \in L$), in the first graded components we get:

$$v_1 = u_1, \quad v_2 + v_1\alpha_1 = u_2 + \alpha_1 u_1$$

thus we must find α_1 such that

$$[\alpha_1, v_1] = v_2 - u_2$$

and the existence of a solution highly depends on the structure and the relations in L . This last equation suggest, in order to prove that two elements u and v are d -conjugate, to look for a solution $\varphi = \exp(\alpha)$ of the equation:

$$(12) \quad \log_d(\varphi) + \varphi^{-1}v\varphi = u.$$

If $\alpha = \sum_{n \geq 1} \alpha_n$ and $\beta = \log_d(\varphi) = \sum_{n \geq 1} \beta_n$, let us recall that we already have (see proposition 1 and equation 5):

$$(13) \quad \beta_n = d(\alpha_n) + \sum_{k=1}^{n-1} \sum_{i=1}^{n-k} \frac{(-1)^i}{(i+1)!} \sum_{n_1+\dots+n_i=n-k} ad_{\alpha_{n_1}}(ad_{\alpha_{n_2}}\dots(ad_{\alpha_{n_i}}(d(\alpha_k)))) = d(\alpha_n) + P_{n-1}(\alpha, d(\alpha))$$

where P_{n-1} is a Lie polynomial in α_k and $d(\alpha_k)$ with $k < n$. On the same way,

Lemma 2. *If v and α are in L , then*

$$(14) \quad \exp(-\alpha).v.\exp(\alpha) = \left(\sum_{i \geq 0} \frac{(-1)^i}{i!} ad_{\alpha}^i \right) (v).$$

The straightforward proof follows the same lines as for lemma 4 and if $\alpha = \sum_{n \geq 1} \alpha_n$, $v = \sum_{n \geq 1} v_n$ and $w = \exp(-\alpha).v.\exp(\alpha) = \sum_{n \geq 1} w_n$:

$$(15) \quad w_n = v_n + \sum_{k=1}^{n-1} \sum_{i=1}^{n-k} \frac{(-1)^i}{i!} \sum_{n_1+\dots+n_i=n-k} ad_{\alpha_{n_1}}(ad_{\alpha_{n_2}}\dots(ad_{\alpha_{n_i}}(v_k))) = v_n + Q_{n-1}(\alpha, v)$$

where Q_{n-1} is a Lie polynomial in α_k and v_k with $k < n$.

Thanks to equations 13 and 15, one can already give a particular element in each conjugacy class:

Theorem 3. *Let d a graded derivation on L and $F = \prod_{n \geq 1} F_n$ a supplementary vector space of $d(L)$ in L . For any $u \in L$, there exists an element $u_F \in F$ which is a d -conjugate of u .*

Proof. If $u = \sum_{n \geq 1} u_n$ we must find $u_F = v = \sum_{n \geq 1} v_n \in F$ and $\varphi = \exp(\alpha) \in G$ ($\alpha = \sum_{n \geq 1} \alpha_n$) such that

$$\log_d(\exp(\alpha)) + \exp(-\alpha)v\exp(\alpha) = u$$

that is to say, for all $n > 0$:

$$(16) \quad d(\alpha_n) + P_{n-1}(\alpha, d(\alpha)) + v_n + Q_{n-1}(\alpha, v) = u_n.$$

For $n = 1$, this read $d(\alpha_1) = u_1 - v_1$. Let p_F the projection on F , parallel to $d(L)$. This is a graded linear operator and if $v_1 = p_F(u_1) \in F_1$ then $u_1 - v_1$ is in $d(L_1)$ and one can find a solution to the equation $d(\alpha_1) = u_1 - v_1$. Suppose for a given $n > 0$, that we have found $\alpha_1, \dots, \alpha_n$ in L and v_1, \dots, v_n in F that solve equation 16. The polynomials $P_n(\alpha, d(\alpha))$ and $Q_n(\alpha, v)$ are well defined and, at order $n + 1$, we must find $\alpha_{n+1} \in L_{n+1}$ and $v_{n+1} \in F_{n+1}$ such that

$$d(\alpha_{n+1}) = u_{n+1} - P_n(\alpha, d(\alpha)) - Q_n(\alpha, v) - v_{n+1}$$

As in the case $n = 1$, we take $v_{n+1} = p_F(u_{n+1} - P_n(\alpha, d(\alpha)) - Q_n(\alpha, v))$ and we can then find an element α_{n+1} that solves the equation. This ends the recursive construction of α and $u_F = v$. ■

Note that the element u_F is not unique since there is a latitude in the determination of α . There would be more to say if one could chose F as a sub-Lie algebra, especially $\ker_L d$. This is the case if we assume that

$$(17) \quad L = \ker_L d \oplus d(L).$$

This hypothesis is natural dynamical systems and, in the sequel, we will deal with derivations satisfying equation (17). Under this assumption, the restriction of d from $d(L)$ to itself is invertible and we note in the sequel I this inverse, defined from $d(L)$ to itself.

From now on, we can choose $F = \ker_L d$ in the previous theorem and we note p the projector on $\ker_L d$, parallel to $d(L)$. This choice allow to define and study "normal forms".

Following the example of $L = y^2\mathbb{C}[[x]]\partial_y$ if $d = ad_{x\partial_x}$ then we have

$$d(L) = xy^2\mathbb{C}[[x]]\partial_y, \quad \ker_L d = \text{Vect}\{y^2\partial_y\}$$

so that $L = \ker_L d \oplus d(L)$. In this case, with $F = \ker_L d = L_1$, one can easily see, using equation 16 and the very simple Lie bracket in $y^2\mathbb{C}[[x]]\partial_y$, that $a(x)y^2\partial_y$ is d -conjugate to $a(0)y^2\partial_y$. As we shall see now, the choice $F = \ker_L d$ provides d -normal forms that play a crude role in dynamical systems.

5.b. Normalization. Let us look closer to the classification of L , up to d -conjugacy, when $L = \ker_L d \oplus d(L)$.

Theorem 4. *For any $u \in L$, there exists $v \in \ker_L d$ such that*

$$u \sim_d v.$$

Moreover, $v, w \in \ker_L d$ are d -conjugated to u if and only if there exists $\varphi \in \exp(\ker_L d)$ such that

$$v\varphi = \varphi w$$

Such elements are called d -normal forms of u .

Proof. The existence of such normal forms follows from theorem 3 with $F = \ker_L d$. For the second part of this theorem, let us assume that v is a d -normal form of u and let $\varphi \in \exp(\ker_L d)$. If $v = \varphi^{-1}w\varphi$, that is to say $\varphi v = w\varphi$, then, thanks to lemma 2, w is in the sub-Lie algebra $\ker_L d$ and, since $d(\varphi) = 0$,

$$d(\varphi) + w\varphi = w\varphi = \varphi v$$

thus w is d -conjugated to v and to u : w is a d -normal form.

Conversely, if v and w are two normal forms, then they are d -conjugated: there exists $\varphi \in G$ such that

$$\log_d(\varphi) + \varphi^{-1}w\varphi = v$$

and it remains to prove that $\varphi \in \exp(\ker_L d)$.

As in the proof of theorem 3, if $\varphi = \exp \alpha$, we have for $n \geq 1$:

$$d(\alpha_n) + P_{n-1}(\alpha, d(\alpha)) + w_n + Q_{n-1}(\alpha, w) = v_n.$$

In graduation 1 this gives $d(\alpha_1) + w_1 = v_1$ or $d(\alpha_1) = v_1 - w_1 \in \ker_L d \cap d(L) = 0$. Thus $w_1 = v_1$ and $\alpha_1 \in \ker_L d$. In graduation 2 the equation

$$d(\alpha_2) - \frac{1}{2}[\alpha_1, d(\alpha_1)] + w_2 - [\alpha_1, w_1] = v_2$$

reduces to

$$d(\alpha_2) = v_2 - w_2 - [\alpha_1, w_1] \in \ker_L d \cap d(L)$$

since $[\ker_L d, \ker_L d] \in \ker_L d$. Once again $d(\alpha_2) = 0$ and $\alpha_2 \in \ker_L d$. We get recursively that $\varphi \in \exp(\ker_L d)$, using the fact that the polynomials $Q_{n-1}(\alpha, w)$ are Lie polynomials of elements of the sub-Lie algebra $\ker_L d$. ■

This theorem gives a description of the d -conjugacy classes : the d -conjugacy classes are exactly the classes of $\ker_L d$ for the classic conjugacy by the group $\exp(\ker_L d)$. As this will be used in section 7.b, this implies that if u is d -conjugated to 0, then 0 is its unique normal form because $\varphi 0 \varphi^{-1} = 0$!

The second part of this theorem implies that there can be more than one normal form. Let us mention that, in the framework of dynamical systems, J. Ecalle and B. Vallet (see [13], [12]), provide a way to define for each $u \in L$ a unique normal form, with the help of a supplementary condition:

Proposition 4. *Let δ be an invertible derivation such that $\delta(\ker_L d) = \ker_L d$, there exist a unique solution to the system:*

$$\begin{cases} d(\varphi) + v\varphi &= \varphi u \\ d(v) &= 0 \\ p((\delta(\varphi))\varphi^{-1}) &= 0 \end{cases}$$

Let us give an idea of the proof. Once a solution (v, φ) of

$$\begin{cases} d(\varphi) &= \varphi u - v\varphi \\ d(v) &= 0 \end{cases}$$

is given the other solutions are given by $(\psi v \psi^{-1}, \psi \varphi)$ with $\psi \in \exp(\ker_L d)$. In order to prove this proposition we have to show that the condition

$$p((\delta(\psi \varphi))(\psi \varphi)^{-1}) = 0$$

determines a unique $\psi \in \exp(\ker_L d)$. But,

$$(\delta(\psi \varphi))(\psi \varphi)^{-1} = \psi(\delta \varphi)\varphi^{-1}\psi^{-1} + \delta(\psi)\psi^{-1} = \psi(\delta \varphi)\varphi^{-1}\psi^{-1} - \psi\delta(\psi^{-1})$$

Once again, using the graduation, if $\delta(\varphi)\varphi^{-1} = \beta \in L$ and $\psi^{-1} = \exp(\alpha) \in \exp(\ker_L d)$, we get for $n = 1$,

$$p(-\delta(\alpha_1) + \beta_1) = 0$$

thus $\delta(\alpha_1) = p(\beta_1)$ that has a unique solution in $\ker_L d$ since δ is invertible and $\ker_L d$ is stable by δ . The remainder of the proof would follow recursively, as in the previous theorems. For any such δ we can define a map N_δ from L to $\ker d$ (note that the conjugating map is also unique).

The reader can notice that in the case of $d = ad_{x\partial_x}$ on $y^2\mathbb{C}[[x]]\partial_y$ there is exactly one normal form of $a(x)y^2\partial_y$, namely $a(0)y^2\partial_y$. This is mainly due to the extreme simplicity of this toy model but it does not reflect the complexity of the general case, for which the choice of a normal form and its computation are not so easy to deal with.

6. NORMALIZATION AND RENORMALIZATION.

6.a. Main theorem. The aim of this section is to point out that similar procedures as those used in perturbative quantum field theory can be used to compute normal forms in the case of non-invertible derivations. In the Connes-Kreimer picture of renormalization (see [3] and [4]), the computation of Feynman integrals shall deliver a character φ on the the graded Hopf algebra of 1PI Feynman graphs, that is an element of the Lie group associated to the complete graded Lie algebra of infinitesimal characters on the Hopf algebra of Feynman graphs L_{FG}^k (with $k = \mathbb{C}$). This character is defined by Feynman integrals that could be divergent but can be renormalized, using dimensional regularization and Birkhoff decomposition.

More precisely, after dimensional regularization, the Feynman integrals define a character $\varphi(\varepsilon)$ with values in the algebra of Laurent series $\mathcal{A} = k[[\varepsilon]][[\varepsilon^{-1}]]$, that is to say an element of the Lie group $G^{\mathcal{A}}$ of the Lie algebra $L^{\mathcal{A}}$. The presence of poles at 0 reflects the divergence of the initial character. The usual renormalization of such a group-like element, is then based on the Birkhoff decomposition: since

$\mathcal{A} = \mathcal{A}^- \oplus \mathcal{A}^+$ where $\mathcal{A}^- = \varepsilon^{-1}k[\varepsilon^{-1}]$ and $\mathcal{A}^+ = k[[\varepsilon]]$ are two sub-algebras, it is possible to factorize $\varphi(\varepsilon)$, that is

$$(18) \quad \varphi(\varepsilon) = \varphi^-(\varepsilon)\varphi^+(\varepsilon)$$

where $\varphi^-(\varepsilon)$ (resp. $\varphi^+(\varepsilon)$) is in $G^{\mathcal{A}^-}$ (resp. $G^{\mathcal{A}^+}$) and the evaluation of $\varphi^+(\varepsilon)$ at $\varepsilon = 0$ gives the renormalized character φ^{ren} . We refer to [7] for further details on the Birkhoff decomposition.

Such ideas can be applied in order to find normal forms. As in perturbative quantum field theory, the solutions of

$$(19) \quad d\varphi = \varphi u$$

are ill-defined whenever the derivation d is not invertible. One can then chose to "regularize" d as follows. Let d a derivation and δ a derivation on $L = L^k$ such that $\ker_L d$ is stable by δ and the restriction of δ from $\ker_L d$ to itself is invertible. These operators d and δ can be extended to $L^{\mathcal{A}}$ by \mathcal{A} -linearity and we get a " δ -renormalization" scheme that delivers a normal form:

Theorem 5. *Let d and δ two derivations on $L = L^k$ such that:*

- (1) $L = \ker_L d \oplus d(L)$,
- (2) $\ker_L d$ is stable by δ ,
- (3) The restriction of δ from $\ker_L d$ to $\ker_L d$ is invertible.

For $u \in L = L^k \subset L^{\mathcal{A}}$ ($\mathcal{A} = k[[\varepsilon]][\varepsilon^{-1}]$), the equation

$$(20) \quad (d + \varepsilon\delta)\varphi = \varphi u$$

has a solution $\varphi = \varphi(\varepsilon) \in G^{\mathcal{A}}$ and, after Birkhoff decomposition,

$$\varphi = \varphi^- \varphi^+ \quad (\varphi^-, \varphi^+) \in G^{\mathcal{A}^-} \times G^{\mathcal{A}^+},$$

the group-like element $\varphi^{\text{ren}} = \varphi^+(\varepsilon)|_{\varepsilon=0}$ conjugates u to a normal form $\beta \in \ker_L d$. Moreover, we have $\delta\varphi_\varepsilon^- = \varphi_\varepsilon^-(\varepsilon^{-1}\beta)$.

Note that this result still hold on a Lie algebra $L^{\mathcal{B}}$ by considering Laurent series in $\mathcal{A} = \mathcal{B}[[\varepsilon]][\varepsilon^{-1}]$. As we shall see in section 7.a, the normal form β is an analog to the β function that appears in perturbative quantum field theory (see [4], [6]).

In the case $d = ad_{x\partial_x}$ on $L = y^2\mathbb{C}[[x]]\partial_y$, the reader can check that one can choose $\delta = ad_{x\partial_x + y\partial_y}$, that coincides with the grading Y . In this case, $d + \varepsilon\delta = ad_{(1+\varepsilon)x\partial_x + \varepsilon y\partial_y}$ and for a given element

$$u = a(x)y^2\partial_y, \quad a(x) = \sum_{n \geq 0} a_n x^n$$

A solution of $(d + \varepsilon\delta)\varphi = \varphi u$ is given by $\varphi(\varepsilon) = \exp(b(x)y^2\partial_y)$ with

$$b(x) = \sum_{n \geq 0} \frac{a_n}{n(1+\varepsilon) + \varepsilon} x^n = \frac{a_0}{\varepsilon} + \sum_{n \geq 1} \frac{a_n}{n(1+\varepsilon) + \varepsilon} x^n$$

where fractions in ε identify with Laurent series and, thanks to the simplicity of the Lie algebra,

$$\varphi^- = \exp\left(\frac{a_0}{\varepsilon}y^2\partial_y\right), \quad \varphi^+ = \exp\left(\sum_{n \geq 1} \frac{a_n}{n(1+\varepsilon) + \varepsilon} x^n y^2 \partial_y\right)$$

so that a normal form of $u = a(x)y^2\partial_y$ is

$$\beta = (\varphi_\varepsilon^-)^{-1}(\varepsilon\delta\varphi_\varepsilon^-) = \varepsilon \frac{a_0}{\varepsilon} y^2 \partial_y = a_0 y^2 \partial_y.$$

6.b. **Proof of theorem 5.** Let d a graded derivation on $L = L^k$ such that $L = \ker_L d \oplus d(L)$ and $p = p_d$, $q = \text{id} - p$ the respective projections on the kernel of d on L and its image. It is clear that d , p and q extend to L^A with $L^A = \ker_{L^A} d \oplus d(L^A)$. For example, if $u = \sum_{n \geq 1} u_n \in L^A$,

$$(21) \quad u_n = \sum_{m \geq M_n} \varepsilon^m u_{n,m}, \quad M_n \in \mathbb{Z}, \quad u_{n,m} \in L_n^k$$

and $d(u) = \sum_{n \geq 1} d(u_n)$ with

$$d(u_n) = \sum_{m \geq M_n} \varepsilon^m d(u_{n,m})$$

The identity $L^A = \ker_{L^A} d \oplus d(L^A)$ implies that the restriction of d from $d(L^k)$ (or $d(L^A)$) to itself is invertible and we note I this graded inverse. On the same way δ can be extended to L^A , its restriction from \ker_L (or \ker_{L^A}) to itself is invertible and we note abusively δ^{-1} its inverse on \ker_L (or \ker_{L^A}). Note that all these operators are graded.

The first step of the proof is based on the following lemma:

Lemma 3. *The endomorphism $d + \varepsilon\delta$ is an invertible derivation on L^A and its (graded) inverse is the linear map I_ε defined for $u \in L^A$ by*

$$(22) \quad I_\varepsilon(u) = \varepsilon^{-1}\delta^{-1}(p(u)) + (\text{id} - \delta^{-1} \circ p \circ \delta) \circ I \left(\sum_{k \geq 0} (-1)^k \varepsilon^k (q \circ \delta \circ I)^{\circ k} (q(u)) \right).$$

Proof. It is a matter of fact to check that $d + \varepsilon\delta$ and I_ε are well-defined: as all the operators do respect the graduation, we can work on graded components L_n^A with elements as in equation (21). It remains to prove that $(d + \varepsilon\delta) \circ I_\varepsilon = I_\varepsilon \circ (d + \varepsilon\delta) = \text{id}_{L^A} = \text{id}$. Since $\ker_{L^A} d$ is stable by δ , $d \circ \delta^{-1} \circ p = 0$ thus

$$(d + \varepsilon\delta)(\varepsilon^{-1}\delta^{-1}(p(u))) = p(u)$$

and, when restricted to $d(L^A)$,

$$(d + \varepsilon\delta) \circ (\text{id} - \delta^{-1} \circ p \circ \delta) \circ I = \text{id} + \varepsilon\delta \circ I - \varepsilon p \circ \delta \circ I = \text{id} + \varepsilon q \circ \delta \circ I.$$

If we apply to

$$\sum_{k \geq 0} (-1)^k \varepsilon^k (q \circ \delta \circ I)^{\circ k} (q(u)) \in d(L^A)$$

it is now clear that

$$(d + \varepsilon\delta)(I_\varepsilon(u)) = p(u) + q(u) = u$$

thus $(d + \varepsilon\delta) \circ I_\varepsilon = \text{id}$. In order to prove that $I_\varepsilon \circ (d + \varepsilon\delta) = \text{id}$, let us notice that

$$\varepsilon^{-1}\delta^{-1} \circ p \circ (d + \varepsilon\delta) = \delta^{-1} \circ p \circ \delta,$$

$I \circ q \circ d = q$ and $q \circ \delta \circ p = 0$. We get

$$\begin{aligned}
\sum_{k \geq 0} (-1)^k \varepsilon^k (q \circ \delta \circ I)^{\circ k} \circ q \circ (d + \varepsilon \delta) &= \sum_{k \geq 0} (-1)^k \varepsilon^k (q \circ \delta \circ I)^{\circ k} \circ q \circ d \\
&\quad - \sum_{k \geq 1} (-1)^k \varepsilon^k (q \circ \delta \circ I)^{\circ k-1} \circ q \circ \delta \\
&= q \circ d \\
&\quad + \sum_{k \geq 1} (-1)^k \varepsilon^k (q \circ \delta \circ I)^{\circ k-1} \circ (q \circ \delta \circ I \circ q \circ d - q \circ \delta) \\
&= q \circ d \\
&\quad + \sum_{k \geq 1} (-1)^k \varepsilon^k (q \circ \delta \circ I)^{\circ k-1} \circ (q \circ \delta \circ q - q \circ \delta) \\
&= q \circ d \\
&\quad - \sum_{k \geq 1} (-1)^k \varepsilon^k (q \circ \delta \circ I)^{\circ k-1} \circ (q \circ \delta \circ p) \\
&= q \circ d
\end{aligned}$$

and finally

$$\begin{aligned}
I_\varepsilon \circ (d + \varepsilon \delta) &= \delta^{-1} \circ p \circ \delta + (\text{id} - \delta^{-1} \circ p \circ \delta) \circ I \circ q \circ d \\
&= \delta^{-1} \circ p \circ \delta + (\text{id} - \delta^{-1} \circ p \circ \delta) \circ q \\
&= \delta^{-1} \circ p \circ \delta \circ p + q = p + q = \text{id}
\end{aligned}$$

■

Note that the expression of this inverse can be simplified if we assume that d and δ commute. Since $d + \varepsilon \delta$ is invertible on L^A , thanks to theorem 1, the equation $(d + \varepsilon \delta)\varphi = \varphi u$ has a solution $\varphi = \varphi(\varepsilon) \in G^A$ for any u in L^A or L . Thanks to the Birkhoff decomposition (see [7]), we have

$$\varphi(\varepsilon) = \varphi^-(\varepsilon)\varphi^+(\varepsilon), \quad \varphi^\pm(\varepsilon) = \exp(\alpha^\pm(\varepsilon)), \quad \alpha^\pm(\varepsilon) \in L^{A^\pm}$$

If $u \in L$, since $(d + \varepsilon \delta)\varphi = \varphi u$ we get

$$\begin{aligned}
(d + \varepsilon \delta)\varphi &= (d + \varepsilon \delta)(\varphi^- \varphi^+) \\
&= \varphi^- ((d + \varepsilon \delta)\varphi^+) + ((d + \varepsilon \delta)\varphi^-) \varphi^+ \\
&= \varphi^- \varphi^+ u
\end{aligned}$$

and then

$$(23) \quad ((d + \varepsilon \delta)\varphi^+) (\varphi^+)^{-1} + (\varphi^-)^{-1} ((d + \varepsilon \delta)\varphi^-) = \varphi^+ u (\varphi^+)^{-1}$$

But $(\varphi^-)^{-1}((d + \varepsilon \delta)\varphi^-)$ is in $L^{k[\varepsilon^{-1}]}$ whereas the other terms are in $L^{k[[\varepsilon]]}$ thus these two parts of the identity do not depend on ε . Let $\beta = (\varphi^-)^{-1}((d + \varepsilon \delta)\varphi^-) \in L^k$ (not in $L^A!$), then

$$((d + \varepsilon \delta)\varphi^+ + \beta\varphi^+ = \varphi^+ u$$

and for $\varepsilon = 0$ the element $\varphi^{\text{ren}} = \varphi^+(0)$ conjugates u to β . It remains to prove that β is a normal form, that is $d(\beta) = 0$. Note that if $\varphi^- = \exp(\alpha)$ with $d(\alpha) = 0$, then $\delta\varphi^- = \varphi^-(\varepsilon^{-1}\beta)$ with $d(\beta) = 0$ (see lemma 1). The proof of this last property is based on the following lemma:

Lemma 4. *Let $\psi \in G^{A^-}$. If*

$$\log_{d+\varepsilon\delta}(\psi) = \psi^{-1} \cdot (d + \varepsilon \delta)(\psi) \in L^k$$

then $\psi \in \exp(\ker_{L^{A^-}} d)$

Proof. The proof of the lemma is based on the following observation: If, for a given $n \geq 1$,

$$x_n = \sum_{M \leq k \leq -1} \varepsilon^k x_{n,k} \in L_n^{A^-}$$

is such that $(d + \varepsilon\delta)(x_n) = y_n + z_n \in L_n^A$ with $y_n \in d(L_n^k)$ and

$$z_n = \sum_{M \leq k \leq 0} \varepsilon^k z_{n,k} \in \ker_{L^{k[\varepsilon^{-1}]}} d$$

then $y_n = 0$ and $x_n = I_\varepsilon(z_n) \in \ker_{L^{A^-}} d$. One can check that $z_{n,k} \in \ker_{L^k} d$ and

$$I_\varepsilon(z_n) = \varepsilon^{-1} \delta^{-1}(z_n) = \sum_{M-1 \leq k \leq -1} \varepsilon^k \delta^{-1}(z_{n,k+1}) \in \ker_{L^{A^-}} d \cap L_n^{A^-}$$

On the other hand, if $y_n \in d(L_n^k)$,

$$I_\varepsilon(y_n) = (\text{id} - \delta^{-1} \circ p \circ \delta) \circ I \left(\sum_{k \geq 0} (-1)^k \varepsilon^k (q \circ \delta \circ I)^{\circ k} (y_n) \right) = x - I_\varepsilon(z_n) \in L_n^{A^+} \cap L_n^{A^-}.$$

It follows that necessarily $I_\varepsilon(y_n) = 0$ (thus $y_n = 0$) and $x_n = I_\varepsilon(z_n) \in \ker_{L^{A^-}} d$.

This is the key to prove the lemma. If $\beta = (\psi)^{-1} ((d + \varepsilon\delta)\psi) = \log_{d+\varepsilon\delta}(\psi) \in L^k$ with $\psi = \exp(\alpha)$ ($\alpha \in L^{A^-}$) then, using lemma 1 and the graduation (see equation (13)), we get:

$$\beta_n = (d + \varepsilon\delta)(\alpha_n) + P_{n-1}(\alpha, (d + \varepsilon\delta)(\alpha)) \in L_n^k$$

where P_{n-1} is a Lie polynomial in α_k and $(d + \varepsilon\delta)(\alpha_k)$ with $k < n$.

In graduation 1:

$$(d + \varepsilon\delta)\alpha_1 = \beta_1$$

and, thanks to the previous result (with $y_1 = q(\beta_1) \in d(L_1^k)$ and $z_1 = p(\beta_1) \in \ker_{L^{k[\varepsilon^{-1}]}} d$), $q(\beta_1) = 0$ and finally $\beta_1 \in \ker_{L^k} d$, $\alpha_1 \in \ker_{L^{A^-}} d$ and $\delta\alpha_1 = \varepsilon^{-1}\beta_1$. Let us suppose now that

$$d(\alpha_1) = \dots = d(\alpha_{n-1}) = 0 = d(\beta_1) = \dots = d(\beta_{n-1})$$

we have

$$(d + \varepsilon\delta)(\alpha_n) = \beta_n - P_{n-1}(\alpha, (d + \varepsilon\delta)(\alpha)) = \beta_n - P_{n-1}(\alpha, \varepsilon\delta(\alpha))$$

with

$$P_{n-1}(\alpha, \varepsilon\delta(\alpha)) = \sum_{k=1}^{n-1} \sum_{i=1}^{n-k} \frac{(-1)^i}{(i+1)!} \sum_{n_1+\dots+n_i=n-k} \text{ad}_{\alpha_{n_1}}(\text{ad}_{\alpha_{n_2}} \dots (\text{ad}_{\alpha_{n_i}}(\varepsilon\delta(\alpha_k))) \dots)$$

Since, for any k , $\alpha_k \in L^{A^-}$, it is clear that we can set

$$\begin{aligned} x_n &= \alpha_n \in L_n^{A^-} \\ z_n &= p(\beta_n) - P_{n-1}(\alpha, \varepsilon\delta(\alpha)) \in \ker_{L^{k[\varepsilon^{-1}]}} d \\ y_n &= q(\beta_n) \in d(L_n^k) \end{aligned}$$

and we already proved that $y_n = 0$ (thus $d(\beta_n) = 0$) and $x_n = \alpha_n \in \ker_{L^{A^-}} d$

It follow recursively that $d(\alpha_k) = 0$, $\beta_k \in \ker d$ for $k \geq 1$ ■

Using this lemma, together with the Birkhoff decomposition, $\beta = \log_{d+\varepsilon\delta} \varphi^- \in L^k$ is a normal form related to φ^- by

$$(24) \quad \delta\varphi^- = \varphi^-(\varepsilon^{-1}\beta)$$

7. FURTHER DEVELOPMENTS.

We enclose in this section many ideas, developments that are suggested by the strong similarity between our results, stemming from dynamical systems, and the algebraic machinery developed in the context of perturbative quantum field theory.

7.a. **Locality and Residues.** The last identity (24) has to be related to residues. If we write

$$\varphi^- = 1 + \sum_{k \geq 1} \varepsilon^{-k} \varphi_k^-, \quad \varphi_k^- \in \mathcal{U}^k$$

then, the residue $\text{Res } \varphi^- = \varphi_1^-$ is such that

$$\delta \text{Res } \varphi^- = \beta$$

and it looks very close to the beta function in perturbative quantum field theory, see [4], [6] and [17]. This concept is related to the "locality" of counter terms (that is φ^-). In our framework one can define on $L^{\mathcal{A}[[\tau]]}$ (and then on $\mathcal{U}^{\mathcal{A}[[\tau]]}$) the automorphism θ_τ :

$$\theta_\tau(x) = e^{\tau \varepsilon \delta} x$$

Following section 7 in [6] (in the case $\delta = Y$), $\varphi \in G^{\mathcal{A}}$ is δ -local if for any τ , the Birkhoff decomposition of $\varphi^\tau = \theta_\tau(\varphi)$ is such that

$$\partial_\tau(\varphi^\tau)^- = 0$$

that is the counter terms do not depend on τ . Under the assumption that $[d, \delta] = 0$ (thus $[d, \theta_\tau] = 0$), the same results as in [6] could be obtained here and one can check that our renormalization procedure delivers a δ -local element. In fact, we have $\partial_\tau \varphi^\tau = \theta_\tau(\varepsilon \delta \varphi)$ and

$$(d + \partial_\tau) \varphi^\tau = \theta_\tau((d + \varepsilon \delta)(\varphi)) = \varphi^\tau u^\tau$$

and after Birkhoff decomposition (as in equation (23)):

$$((d + \partial_\tau)(\varphi^\tau)^+) ((\varphi^\tau)^+)^{-1} + ((\varphi^\tau)^-)^{-1} ((d + \partial_\tau)(\varphi^\tau)^-) = (\varphi^\tau)^+ u^\tau ((\varphi^\tau)^+)^{-1}$$

but as in the proof of theorem 5, this identity shows that $((\varphi^\tau)^-)^{-1} ((d + \partial_\tau)(\varphi^\tau)^-) \in L^{\mathcal{A}[[\tau]]^-} \cap L^{\mathcal{A}[[\tau]]^+} = \{0\}$ and $(d + \partial_\tau)(\varphi^\tau)^- = 0$. If we expand in powers of τ , since $(\varphi^0)^- = \varphi^-$ with $d(\varphi^-) = 0$, we get

$$\partial_\tau(\varphi^\tau)^- = 0$$

and this proves the δ -locality of φ .

As we can see here, this renormalization is very similar to the one used in physics. Note also that it could be of some interest to consider "perturbations" of d such as $d + \varepsilon \delta_1 + \varepsilon^2 \delta_2 + \dots$. On the same way, the following section suggests that other renormalization procedure appear in the framework of dynamical systems.

7.b. **Another "renormalization" : the correction.** As pointed out in [9], in pQFT, the attempted (but ill-defined) group-like element φ is associated to a Lagrangian \mathcal{L} and the renormalization procedure can be interpreted as an iterative process, based on the graduation of the considered Hopf algebra, that consists in modifying the Lagrangian:

$$\mathcal{L} \rightarrow \mathcal{L} - \mathcal{L}_1 \rightarrow \mathcal{L} - \mathcal{L}_1 - \mathcal{L}_2 \rightarrow \dots \rightarrow \mathcal{L}^{ren}$$

so that the renormalized group-like element φ^{ren} corresponds to the modified Lagrangian \mathcal{L}^{ren} .

The same principle appear in the framework of dynamical systems but the final object differs from a normal form, this is the correction:

Theorem 6. *Let d a derivation such that $L = \ker d \oplus d(L)$. For any $u \in L$ there exists a unique $u^c \in \ker d$ such that $u - u^c$ is in the d -conjugacy class of 0. u^c is called the correction of u . Moreover, if $u \in \ker d$, then $u^c = u$.*

As in the Lagrangian interpretation of renormalization in pQFT, the idea of the proof is to modify u , graded component by graded component, so that we can solve the equation:

$$d\varphi = \varphi(u - u_1^c - u_2^c - \dots)$$

Proof. Whenever there exists φ such that $\log_d(\varphi) = u$, then we can take $u^c = 0$, but, when d is not invertible, such φ does not necessarily exist. Let us consider

$$u = \sum u_n, \quad u^c = \sum u_n^c, \quad u - u^c = v = \sum v_n, \quad \varphi = \exp\left(\sum \alpha_n\right)$$

The equation $\log_d(\varphi) = v = u - u^c$ reads (see lemma 1 and equation (13)):

$$d(\alpha_n) = v_n - P_{n-1}(\alpha, d(\alpha))$$

where

$$P_{n-1}(\alpha, d(\alpha)) = \sum_{k=1}^{n-1} \sum_{i=1}^{n-k} \frac{(-1)^i}{i+1} \sum_{n_1+\dots+n_i=n-k} ad_{\alpha_{n_1}}(ad_{\alpha_{n_2}}\dots(ad_{\alpha_{n_i}}(d(\alpha_k))))\dots$$

depends only on $\alpha_1, \dots, \alpha_{n-1}, d(\alpha_1), \dots, d(\alpha_{n-1})$ (with $P_0 = 0$). For $n = 1$, we must solve:

$$d(\alpha_1) = v_1 = u_1 - u_1^c$$

thus $u_1^c = p(u_1) \in \ker d$ (unique), $v_1 = q(u_1)$ and then we can choose α_1 up to an element of $\ker d$: $\alpha_1 = I(q(u_1)) + z_1$, $z_1 \in \ker d$. Now suppose that for $1 \leq k \leq n$ we can find u_1^c, \dots, u_n^c so that there exists solutions (say $\alpha_1, \dots, \alpha_n$) to the equations:

$$\forall 1 \leq k \leq n, \quad d(\alpha_k) = u_k - u_k^c - P_{k-1}(\alpha, d(\alpha))$$

for the next equation:

$$d(\alpha_{n+1}) = u_{n+1} - u_{n+1}^c - P_n(\alpha, d(\alpha)),$$

as $L = \ker d \oplus d(L)$, we must define

$$u_{n+1}^c = p(u_{n+1} - P_n(\alpha, d(\alpha)))$$

and this ends the (recursive) proof of the existence of u^c . But in this recursive definition u^c seems to depend on the choice of α . Indeed u^c is unique : suppose we have u^1 and u^2 in $\ker d$ such that $u - u^1 \sim_d 0 \sim_d u - u^2$, then there exists φ, ψ such that:

$$d(\varphi) = \varphi(u - u^1), \quad d(\psi) = \psi(u - u^2)$$

In G , let $\varphi\psi^{-1} = \phi$, then

$$d(\varphi) = \varphi(u - u^1) = \phi\psi(u - u^1) = d(\phi)\psi + \phi d(\psi) = d(\phi)\psi + \phi\psi(u - u^2)$$

thus

$$\phi\psi(u^2 - u^1) = d(\phi)\psi$$

or rather

$$\psi(u^2 - u^1)\psi^{-1} = \log_d(\phi)$$

Using lemmas 1 and 2 with $\phi = \exp \alpha$ and $\psi^{-1} = \exp \beta$:

$$\sum_{i \geq 0} \frac{1}{i!} ad_{\beta}^i(u^2 - u^1) = \sum_{i \geq 0} \frac{(-1)^i}{(i+1)!} ad_{\alpha}^i(d(\alpha)).$$

Thanks to the graduation we get first $d(\alpha_1) = u_1^2 - u_1^1$: since it is in the kernel and the image of d , we get $d(\alpha_1) = 0$ and $u_1^1 - u_1^2 = 0$. One can then prove recursively on n that, for $1 \leq k \leq n$:

$$d(\alpha_k) = 0, \quad u_k^1 = u_k^2.$$

If this is true then the equation reduces, in graduation $n + 1$ to

$$d(\alpha_{n+1}) = u_{n+1}^2 - u_{n+1}^1$$

thus $d(\alpha_{n+1}) = u_{n+1}^2 - u_{n+1}^1 = 0$ and this ends the recursion. ■

Note that, φ is unique if we assume that $\varphi = \exp(\alpha)$ with $\alpha \in d(L)$, using the inverse I of the restriction of d from $d(L)$ to $d(L)$.

We don't know if this correction can be interpreted in the framework pQFT but let us just end by two remarks on the link between normal forms obtained by the "renormalization scheme" described in theorem 5 and the correction. For a given δ as in theorem 5, we get a function N_δ from L to $\ker_L d$ such that, for all $u \in L$, $u \sim_d N_\delta(u)$ and the correction associates to any $u \in L$ a unique $u^c = C(u)$, independent of δ , such that $u - C(u) \sim_d 0$. As pointed out in section 5.b, there is no other normal form than 0 which is conjugated to 0:

$$N_\delta(u - C(u)) = 0$$

thus $u - C(u)$ is a fixed point for the operator $\text{id}_L - N_\delta$. One could even go further: if we introduce a parameter t then,

$$\lim_{t \rightarrow 0} t^{-1} C(tu) = \lim_{t \rightarrow 0} t^{-1} N_\delta(tu) = p(u)$$

so that, in a first approximation, the normal form and the correction coincide. This looks like two renormalizations, with different prescriptions but with a common "minimal substraction" $u \mapsto u - p(u)$.

Under the same hypothesis as in theorem 5, since $d + \varepsilon\delta$ is invertible on $L^{\mathcal{A}}$ with $\mathcal{A} = k[[\varepsilon]][[\varepsilon^{-1}]]$, the equation $\log_{d+\varepsilon\delta} \varphi = u \in L^k$ has a unique solution $\varphi = \exp(\alpha) \in G^{\mathcal{A}}$. As for the Birkhoff decomposition, the reader can check that there exist $(\alpha^+, \alpha^-) \in L^{\mathcal{A}^+} \times L^{\mathcal{A}^-}$ such that, if $\varphi^\pm = \exp(\alpha^\pm)$,

$$(25) \quad \log_{d+\varepsilon\delta} \varphi = \log_{d+\varepsilon\delta} \varphi^+ + \log_{d+\varepsilon\delta} \varphi^- = u.$$

As before, the recursive proof is based on equation (13): For $n \geq 1$,

$$\begin{aligned} (d + \varepsilon\delta)(\alpha_n) + P_{n-1}(\alpha, (d + \varepsilon\delta)(\alpha)) &= (d + \varepsilon\delta)(\alpha_n^+) + P_{n-1}(\alpha^+, (d + \varepsilon\delta)(\alpha^+)) \\ &\quad + (d + \varepsilon\delta)(\alpha_n^-) + P_{n-1}(\alpha^-, (d + \varepsilon\delta)(\alpha^-)) \end{aligned}$$

thus

$$\alpha_n^+ + \alpha_n^- = \alpha_n + I_\varepsilon (P_{n-1}(\alpha, (d + \varepsilon\delta)(\alpha)) - P_{n-1}(\alpha^+, (d + \varepsilon\delta)(\alpha^+)) - P_{n-1}(\alpha^-, (d + \varepsilon\delta)(\alpha^-))).$$

But $\log_{d+\varepsilon\delta} \varphi^- \in L^{k[[\varepsilon^{-1}]}}$, $\log_{d+\varepsilon\delta} \varphi^+ \in L^{k[[\varepsilon]]}$ and their sum is $u \in L^k$. This means that $\gamma = \log_{d+\varepsilon\delta} \varphi^- \in L^k$ and, thanks to lemma 4, $d(\gamma) = 0$ thus

$$\log_{d+\varepsilon\delta} \varphi^+ = u - \gamma \text{ with } d(\gamma) = 0$$

and this identity makes sense at $\varepsilon = 0$. We recover the correction $\gamma = u_c$ by this "additive" Birkhoff decomposition.

8. CONCLUSION.

This use of "renormalization schemes" is new in the framework of dynamical systems and shall open new perspectives in this area. It shall interact with J. Ecalle's mould calculus that provides a way to "lift" the computation of diffeomorphisms to computations in the Lie group of the Lie algebra of infinitesimal characters on a shuffle Hopf algebra (see [10], [14]). If mould calculus is sufficient to lead to explicit computations, this does not tackle the difficult and crucial question of the analyticity of the computed diffeomorphisms.

Surprisingly, the analytic regularity of such diffeomorphism can sometimes be easily obtained by computing in a Hopf algebras of trees, which is strongly related to the PreLie structure of vector fields (see [14]). There shall be no major difficulty to adapt the results of this paper to the case of derivations on PreLie algebras.

This may even be the right theoretical framework to deal with the similar questions in "discrete dynamical systems" that are related to group equations

$$\theta\varphi = \psi\varphi, \quad \psi, \varphi \in G$$

where θ is a graded automorphism on L : $\theta([x, y]) = [\theta(x), \theta(y)]$, with an "invertibility" condition on $\theta - \text{id}_L$. This is the right framework to deal with Hurwitz multizetas (see [1]) and their regularization.

Let us end this paper with a remark on the choice of an ε -regularization. In section 6.a, we illustrated theorem 5 in the case $d = ad_{x\partial_x}$ on $L = y^2\mathbb{C}[[x]]\partial_y$ with $\delta = ad_{x\partial_x + y\partial_y}$ and the solution of

$$\log_{d+\varepsilon\delta} \varphi = a(x)y^2$$

was $\varphi(\varepsilon) = \exp(b(x)y^2\partial_y)$ where

$$(1 + \varepsilon)xb'(x) + \varepsilon b(x) = a(x).$$

Theorem 5 gave then the normal form $a(0)y^2\partial_y$. Following [21], we could have chosen the regularization $d_\varepsilon = ad_{x^{1-\varepsilon}\partial_x}$, assuming that one can now work on $\mathbb{C}[[x, x^\varepsilon]]$. After expansion $x^\varepsilon = \sum_{n \geq 0} \frac{\varepsilon^n \log^n x}{n!}$ and Birkhoff decomposition, it appear that there exists solutions to the equation

$$\log_d \psi = a(x)y^2, \quad \psi = \exp(c(x)y^2\partial_y)$$

but with a series $c(x)$ in x and $\log x$. On one hand, this result is very close to the results of J. Ecalle on "ramified linearization" (see [11]). On the other hand, the apparition of logarithmic terms shall be familiar to the experts of pQFT and shall be related to the toy model described in [9] (section 4.2).

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9. APPENDIX: DYNAMICAL SYSTEMS

Let us first recall some results on autonomous analytic local differential equations (see [16] for proofs and details), that is a system (in dimension ν)

$$(26) \quad \dot{x} = f(x)$$

where $x = (x_1, \dots, x_\nu)$ and $f = (f_1, \dots, f_\nu)$ is an analytic map defined in a neighborhood of the origin of \mathbb{C}^ν , with values in \mathbb{C}^ν such that $f(0) = 0$. There are several strategies to study such a local dynamical system. The first one is to compute the flow (whose existence, uniqueness and analyticity is due to the Cauchy-Lipschitz theorem), that is the local biholomorphism from $\mathbb{C}^{1+\nu}$ to $\mathbb{C}^{1+\nu}$:

$$\varphi(t, x) = (t, \varphi^t(x))$$

that describes the local solution, of the system $(t, y(t))$ such that $y(0) = x$. Note that, since the system is autonomous (f does not depend on t), for small values t and s , we have $\varphi^t \circ \varphi^s = \varphi^{t+s}$. Although this flow encodes all the properties of the solutions, its computation does not give many information on the geometry of the solutions. Since Poincaré, in order to get more geometrical information, a better approach is to study this solutions up to a change of coordinates. Let ψ be a local ($\psi(0) = 0$) biholomorphism of \mathbb{C}^ν , then if $y(t) = \psi(x(t))$ where x is a solution of the initial system, then :

$$\dot{y} = \sum_{i=1}^{\nu} \dot{x}_i \frac{\partial \psi}{\partial x_i}(x) = \sum_{i=1}^{\nu} f_i(x) \frac{\partial \psi}{\partial x_i}(x) = g(y) = g(\psi(x))$$

and the vector field f is transformed into the vector field g by the equation

$$(27) \quad \sum_{i=1}^{\nu} f_i(x) \frac{\partial \psi}{\partial x_i}(x) = g(\psi(x))$$

This defines an equivalence relation of vector field up to local biholomorphism and, in order to grab the geometry of the solutions, it seem reasonable to try to find the most simple vector field g which is conjugated to f . For the sake of simplicity we will suppose in the sequel that the *linear part* of f is diagonal:

$$f(x) = \Lambda x + u(x)$$

where Λ is a diagonal matrix whose diagonal $\lambda = (\lambda_1, \dots, \lambda_\nu)$ is called the *spectrum*, and the components of $u = (u_1, \dots, u_\nu)$ are of valuation at least 2 (for the usual valuation in $\mathbb{C}\{x\}$). In this context (as in the context of perturbative quantum field theory), it seem reasonable to consider f as a *perturbation* of

its linear part and to ask if it is conjugated to its linear part: Does there exist a local biholomorphism of C^ν , tangent to the identity (in order to preserve the linear part) such that:

$$\sum_{i=1}^{\nu} f_i(x) \frac{\partial \psi}{\partial x_i}(x) = \sum_{i=1}^{\nu} \lambda_i x_i \frac{\partial \psi}{\partial x_i}(x) + \sum_{i=1}^{\nu} u_i(x) \frac{\partial \psi}{\partial x_i}(x) = \Lambda \cdot \psi(x)$$

This equation is called the homological equation and it provides recursive equations on the coefficients of ψ . Unfortunately (see for example [18] or [16]) many difficulties arise in the resolution of such equation (see also below):

1. If, for any $1 \leq j \leq \nu$ and any sequence of non-negative integers $n = (n_1, \dots, n_\nu)$ with $|n| = n_1 + \dots + n_\nu \geq 2$, we have:

$$(28) \quad \langle \lambda, n \rangle - \lambda_j \neq 0 \quad (\langle \lambda, n \rangle = \lambda_1 n_1 + \dots + \lambda_\nu n_\nu)$$

then the vector field is *non-resonant* and one can solve recursively the homological equation. Unfortunately, the series ψ is not necessarily convergent (i.e. analytic), unless the spectrum satisfies some *diophantine condition* (see [18]).

2. Otherwise, when there are some cancellations of the values $\langle \lambda, n \rangle - \lambda_j$, then the homological equation cannot *even formally* be solved. In this case, the case of resonant vector fields, solving recursively the homological equation involves some divisions by 0. This also means that, in this case, the vector field cannot be *linearized*: at best, one can try to conjugate f to a simpler vector field (but not its linear part) which is called a normal (or pre-normal) vector field $g(x) = \Lambda \cdot x + v(x)$ such that the components v_j of v contains only monomials $x^n = x_1^{n_1} \dots x_\nu^{n_\nu}$ such that $\langle \lambda, n \rangle - \lambda_j = 0$. This is indeed possible but nothing ensures once again that we get an analytic solution, either for ψ nor for g .

This small survey of the situation calls for some remarks. We shall leave aside the analytic difficulties (for which there still exists open problems) and focus on the algebraic difficulties, that is computing formal series rather than analytic series. Even formally, there are some obstruction to the resolution of the homological equation: we started with the naive guess that, since f can be seen as perturbation of its linear part, it shall be conjugated to it. This look like the situation in perturbative quantum field theory (pQFT), where we see the action as a perturbation of the free action, and as in pQFT, this gives a ill-defined solution (here the map ψ). The analogy between resonant vector field and pQFT goes even further: in the resonant case, the computation of a normal form can be obtained by a renormalization procedure (as in pQFT).

As pointed out in [3] and [4], pQFT leads to the computation of a character (maybe ill-defined) on a graded commutative Hopf algebra, that is to say to the computation of an element of the Lie group of a complete graded Lie algebra. The situation is indeed the same in the study of *formal* differential equation, when considering vector fields as *operators*: let f is a formal vector field without constant term ($f \in C_{\geq 1}[[x_1, \dots, x_\nu]]^\nu = C_{\geq 1}[[x]]^\nu$), and $A \in C[[x]]$ a formal series, if $\dot{x} = f(x)$ and $y = A(x)$ then:

$$\dot{y} = A(x) = \sum_{i=1}^{\nu} \dot{x}_i \frac{\partial A}{\partial x_i}(x) = \left(\sum_{i=1}^{\nu} f_i(x) \frac{\partial}{\partial x_i} \right) \cdot A(x)$$

so that we can identify, in a unique way, a vector field f with a differential operator (even a derivation) $X = X_f$ acting on $C[[x]]$:

$$\sum_{i=1}^{\nu} f_i(x) \frac{\partial}{\partial x_i} = X(x)$$

On the same way, if ψ is a local formal, invertible diffeomorphism of C^ν , that is to say

$$\psi \in C_{\geq 1}[[x_1, \dots, x_\nu]]^\nu = C_{\geq 1}[[x]]^\nu$$

with an invertible linear part $\left(\left(\frac{\partial \psi_i}{\partial x_j}(0) \right) \right)$, then one can associate an invertible operator (called a substitution automorphism) F_ψ acting on $C[[x]]$:

$$(F_\psi \cdot A)(x) = A(\psi(x))$$

Moreover, it is a matter of fact to check that a linear operator F on $\mathbb{C}[[x]]$ such that:

- i. $\psi(x) = (F.x_1, \dots, F.x_\nu)$ is a local formal, invertible diffeomorphism of \mathbb{C}^ν
- ii. For any series A, B in $\mathbb{C}[[x]]$, $F.(AB) = (F.A)(F.B)$

is the substitution automorphism associated to ψ .

The link with Lie algebras can be put this way (see [16]):

Proposition 5. *Let \mathcal{D} (resp. $\mathcal{D}_{\geq 1}$) the set of derivations $\sum_{i=1}^{\nu} f_i(x) \frac{\partial}{\partial x_i} = X(x)$ such that $f(0) = 0$ (resp. f has no linear part). \mathcal{D} (resp. $\mathcal{D}_{\geq 1}$) is a complete graded Lie algebra (for the Lie bracket associated to the composition of operators). Moreover, if \mathcal{G} (resp. $\mathcal{G}_{\geq 1}$) is the group of substitution automorphisms associated to local formal, invertible diffeomorphisms (resp. identity-tangent formal diffeomorphisms), then the exponential map:*

$$\exp(X) = \text{Id}_{\mathbb{C}[[x]]} + \sum_{s \geq 1} \frac{1}{s!} X^s$$

defines an injective (resp. bijective) map from \mathcal{D} (resp. $\mathcal{D}_{\geq 1}$) to \mathcal{G} (resp. $\mathcal{G}_{\geq 1}$) whose inverse on $\mathcal{G}_{\geq 1}$ is given by the logarithm:

$$\log(F) = \sum_{s \geq 1} \frac{(-1)^{s-1}}{s} (F - \text{Id}_{\mathbb{C}[[x]]})$$

A proof can be found in [16] and we refer to [22] and [21] for some applications of this operator formalism in dimension $\nu = 1$. There is no difficulty in proving such a proposition. Let us just notice that the graduation is related to the action of such derivation on monomials : A vector field in \mathcal{D} (resp. $\mathcal{D}_{\geq 1}$) is given by a series of operators such as $cx_1^{n_1} \dots x_\nu^{n_\nu} \frac{\partial}{\partial x_i}$ with $n_1 + \dots + n_\nu \geq 1$ (resp. $n_1 + \dots + n_\nu \geq 2$) that acts on monomials $x_1^{m_1} \dots x_\nu^{m_\nu}$:

$$\left(cx_1^{n_1} \dots x_\nu^{n_\nu} \frac{\partial}{\partial x_i} \right) . x_1^{m_1} \dots x_\nu^{m_\nu} = x_1^{m_1+n_1} \dots x_i^{m_i+n_i-1} \dots x_\nu^{m_\nu+n_\nu}$$

so that the total degree goes from $m_1 + \dots + m_\nu$ to $m_1 + \dots + m_\nu + n_1 + \dots + n_\nu - 1$ and the graduation for such an operator is then $n_1 + \dots + n_\nu - 1$.

Note also that this proposition implies that any substitution automorphism in $\mathcal{G}_{\geq 1}$ is a differential operator. This could have been proved directly, using the Taylor formula: If $\psi(x) = (x_1 + u_1(x), \dots, x_\nu + u_\nu(x)) = x + u(x)$, then

$$F_\psi . A(x) = A(\psi(x)) = A(x + u(x)) = A(x) + \sum_{s \geq 1} \frac{1}{s!} \sum_{1 \leq i_1, \dots, i_s \leq \nu} u_{i_1}(x) \dots u_{i_s}(x) \frac{\partial^s A(x)}{\partial x_{i_1} \dots \partial x_{i_s}}$$

If we look back at the (formal) flow of a vector field (or rather derivation) X in \mathcal{D} (resp. $\mathcal{D}_{\geq 1}$), the reader can check that the associated one parameter family of substitution automorphisms F_X^t in \mathcal{G} (resp. $\mathcal{G}_{\geq 1}$) is given by $F_X^t = \exp(tX)$ and satisfies the operator equation

$$\frac{\partial}{\partial t} . F_X^t = F_X^t . X$$

On the same way the homological equation can be formulated as follows : Let $X = X_0 + B$ be a vector field of \mathcal{D} , with linear part X_0 and $B \in \mathcal{D}_{\geq 1}$, if this vector field is conjugated by ψ to its linear part and $F = F_\psi$, then:

$$F.X_0 = X.F = (X_0 + B).F$$

and if G is the inverse of F , this reads

$$\text{ad}_{X_0}(G) = [X_0, G] = X_0.G - G.X_0 = G.B$$

so that the homological equation is an equation relating an element of $\mathcal{D}_{\geq 1}$ and an element of $\mathcal{G}_{\geq 1}$, through the adjoint action of X_0 : B is a logarithmic derivative of G for the derivation $d = \text{ad}_{X_0}$ (that respects the graduation). This is precisely the kind of equation we studied in the abstract context of complete graded Lie algebras, equipped with a *graded derivation* d that has essentially the same

properties as in the case of the adjoint action for vector fields. For example, if $X_0 = \sum_{i=1}^{\nu} \lambda_i x_i \frac{\partial}{\partial x_i}$, then

$$d \left(cx_1^{n_1} \dots x_{\nu}^{n_{\nu}} \frac{\partial}{\partial x_i} \right) = \text{ad}_{X_0} \left(cx_1^{n_1} \dots x_{\nu}^{n_{\nu}} \frac{\partial}{\partial x_i} \right) = (\langle \lambda, n \rangle - \lambda_i) cx_1^{n_1} \dots x_{\nu}^{n_{\nu}} \frac{\partial}{\partial x_i}$$

so that a vector field $X_0 + B$ is resonant if d is not invertible on $\mathcal{D}_{\geq 1}$. One can also notice in this case that, on restriction to $\mathcal{D}_{\geq 1}$, we have $\mathcal{D}_{\geq 1} = \ker d \oplus \text{Im } d$ and $\ker d$ is a sub-Lie algebra.

REFERENCES

- [1] Olivier Bouillot. *Invariants analytiques des difféomorphismes et Multizetas*. PhD thesis, Paris-Sud University, 2011.
- [2] Nicolas Bourbaki. *Lie groups and Lie algebras. Chapters 1–3*. Elements of Mathematics (Berlin). Springer-Verlag, Berlin, 1998. Translated from the French, Reprint of the 1989 English translation.
- [3] Alain Connes and Dirk Kreimer. Renormalization in quantum field theory and the Riemann-Hilbert problem. I. The Hopf algebra structure of graphs and the main theorem. *Comm. Math. Phys.*, 210(1):249–273, 2000.
- [4] Alain Connes and Dirk Kreimer. Renormalization in quantum field theory and the Riemann-Hilbert problem. II. The β -function, diffeomorphisms and the renormalization group. *Comm. Math. Phys.*, 216(1):215–241, 2001.
- [5] Alain Connes and Matilde Marcolli. *Noncommutative geometry, quantum fields and motives*, volume 55 of *American Mathematical Society Colloquium Publications*. American Mathematical Society, Providence, RI, 2008.
- [6] Kurusch Ebrahimi-Fard, José M. Gracia-Bondía, and Frédéric Patras. A Lie theoretic approach to renormalization. *Comm. Math. Phys.*, 276(2):519–549, 2007.
- [7] Kurusch Ebrahimi-Fard, Li Guo, and Dominique Manchon. Birkhoff type decompositions and the Baker-Campbell-Hausdorff recursion. *Comm. Math. Phys.*, 267(3):821–845, 2006.
- [8] Kurusch Ebrahimi-Fard, Dominique Manchon, and Frédéric Patras. A noncommutative Bohnenblust-Spitzer identity for Rota-Baxter algebras solves Bogoliubov’s recursion. *J. Noncommut. Geom.*, 3(2):181–222, 2009.
- [9] Kurusch Ebrahimi-Fard and Frédéric Patras. Exponential renormalization. *Ann. Henri Poincaré*, 11(5):943–971, 2010.
- [10] Jean Ecalle. Singularités non abordables par la géométrie. *Ann. Inst. Fourier (Grenoble)*, 42(1-2):73–164, 1992.
- [11] Jean Ecalle. Compensation of small denominators and ramified linearisation of local objects. *Astérisque*, (222):4, 135–199, 1994. Complex analytic methods in dynamical systems (Rio de Janeiro, 1992).
- [12] Jean Ecalle and Bruno Vallet. Passive and active resonance. Non-linear resurgence and iso-resurgent deformations. In *The Stokes phenomenon and Hilbert’s 16th problem (Groningen, 1995)*, pages 103–138. World Sci. Publ., River Edge, NJ, 1996.
- [13] Jean Ecalle and Bruno Vallet. Correction and linearization of resonant vector fields and diffeomorphisms. *Math. Z.*, 229(2):249–318, 1998.
- [14] Frédéric Fauvet and Frédéric Menous. Ecalle’s arborification–coarborification transforms and connes–kreimer hopf algebra. arXiv:1212.4740.
- [15] Giovanni Gallavotti. A criterion of integrability for perturbed nonresonant harmonic oscillators. “Wick ordering” of the perturbations in classical mechanics and invariance of the frequency spectrum. *Comm. Math. Phys.*, 87(3):365–383, 1982/83.
- [16] Yulij Ilyashenko and Sergei Yakovenko. *Lectures on analytic differential equations*, volume 86 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, 2008.
- [17] Dominique Manchon. Bogota lectures on hopf algebras, from basics to applications to renormalization. In *Comptes-rendus des Rencontres mathématiques de Glanon 2001*.
- [18] Jean Martinet. Normalisation des champs de vecteurs holomorphes (d’après A.-D. Brjuno). In *Bourbaki Seminar, Vol. 1980/81*, volume 901 of *Lecture Notes in Math.*, pages 55–70. Springer, Berlin, 1981.
- [19] F. Menous and F. Patras. Logarithmic Derivatives and Generalized Dynkin Operators. *ArXiv e-prints*, June 2012.
- [20] Frédéric Menous. On the stability of some groups of formal diffeomorphisms by the Birkhoff decomposition. *Adv. Math.*, 216(1):1–28, 2007.
- [21] Frédéric Menous. Formal differential equations and renormalization. In *Renormalization and Galois theories*, volume 15 of *IRMA Lect. Math. Theor. Phys.*, pages 229–246. Eur. Math. Soc., Zürich, 2009.
- [22] Frédéric Menous. Formulas for the Connes-Moscovici Hopf algebra. In *Combinatorics and physics*, volume 539 of *Contemp. Math.*, pages 269–285. Amer. Math. Soc., Providence, RI, 2011.