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A hybridized discontinuous Galerkin method for 3d time-harmonic Maxwell’s equations

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Abstract — Hybridized discontinuous Galerkin methods preserve the advantages of classical discontinuous Galerkin methods and in addition enable to circumvent the issue of the number of degrees of freedom. The principles of these numerical methods are summed up for 3d time-harmonic Maxwell’s equations and basic examples are proposed to assess their efficiency.

I. INTRODUCTION

Discontinuous Galerkin (DG) methods are currently widespread for the discretization of time-transient Maxwell’s equations [1, 2]. This success can be explained by several advantages compared to other approaches such as the flexibility for hp-adaptivity or a natural parallelism [3]. Nonetheless for stationary problems as time-harmonic Maxwell’s equations, the number of degrees of freedom is several times larger than when using a classical conforming finite element method and this aspect is cumbersome. Hybridized DG (HDG) methods have been proposed recently to circumvent this problem [4]: preserving the discontinuous approximations of both electric and magnetic fields, the only global linear system to solve determines the degrees of freedom of a new hybrid variable “living” on the interface of each element of the mesh. Thus, a HDG method leads to the solution of a reduced-size linear system without altering the nice properties of a DG method. We consider here such a HDG method for solving the 3d time-harmonic Maxwell’s equations and illustrate the accuracy and the computing cost of the approach.

II. HDG DISCRETIZATION

Electric E and magnetic H fields are approximated by discontinuous vector fields $E_h$ and $H_h$ whose components are polynomials in each element, i.e. $E_h$ and $H_h$ belong to $\{P_k(K)\}^3$ in each element $K$ with $P_k(K)$ the space of polynomials of degree $k$ on $K$. The fields $E_h$ and $H_h$ are computed by weakly verifying time-harmonic Maxwell’s equations in each element $K$ of the mesh; it means that for all $\mathbf{v}$ in $\{P_k(K)\}^3$, $E_h$ and $H_h$, satisfy

\[
\begin{align*}
(\omega \varepsilon_r \mathbf{E}_h, \mathbf{v}) - (\mathbf{H}_h, \nabla \times \mathbf{v}) + (|H_h|^2, \mathbf{n} \times \mathbf{v}) &= 0, \\
(i \mu \mu_{r} \mathbf{H}_h, \mathbf{v}) + (\mathbf{E}_h, \nabla \times \mathbf{v}) - (|E_h|^2, \mathbf{n} \times \mathbf{v}) &= 0.
\end{align*}
\]

Here $i$ is the imaginary unit, $\omega$ the angular frequency, $\varepsilon_r$ and $\mu_r$ the relative permittivity and permeability. Symbols $(\cdot, \cdot)$ and $(\cdot, \cdot, \cdot)$ denote the scalar product of complex-valued vectors integrated over $K$ and $\partial K$; the boundary of $K$; $\mathbf{n}$ is the outward unit normal vector on $\partial K$ with $\times$ the cross-product. Quantities $\mathbf{H}_h$ and $\mathbf{E}_h$ are only defined on $\partial K$ and are approximated as follows:

\[
\mathbf{E}_h = \mathbf{E}_h^0 + \tau n \times (\mathbf{H}_h - \mathbf{H}_h^0) \text{ on } \partial K, \quad (\mathbf{E}_h^0, \mathbf{n}) = (\mathbf{E}_0, \mathbf{n}), \quad (\mathbf{H}_h^0, \mathbf{n}) = (\mathbf{H}_0, \mathbf{n}),
\]

where $\tau$ is a (often strictly) positive parameter. And $\mathbf{H}_h$, $\mathbf{E}_h$, $\mathbf{E}_h^0$ has a priori two distinct values on each side of an interface between two elements. We thus need to weakly enforce that $\mathbf{E}_h$ is single-valued on each interface and $\mathbf{H}_h$ on $\partial K$.

$(\mathbf{E}_h, \eta)_{\mathbf{v}} = 0$, $\forall \mathbf{v} 
\mathbf{E}_h ∈ \{P_k(K)\}^3$,

and that the scheme is consistent with the boundary conditions (Silver-Muller conditions for instance)

\[
\mathbf{E}_h \times \mathbf{n} + \mathbf{H}_h - \mathbf{G}, \eta \mathbf{v} = 0, \forall \mathbf{v} ∈ \mathbf{F}_h.
\]

Here $\mathbf{F}_h$ and $\mathbf{F}_h^0$ are the set of the interior and boundary faces of the mesh $(\mathbf{F}_h = \mathbf{F}_h^0 \cup \mathbf{F}_h^b)$, $[\mathbf{v}]$ denotes the jump of $\mathbf{v}$ on any face, $\mathbf{G}$ is a source term provided by an incident electromagnetic wave on the boundary and the meaning of $(\cdot, \cdot, \cdot)_{\mathbf{v}}$ comes from (1) with $F$ replacing $\partial K$.

\[
\mathbf{M}_h^0 = \{\eta ∈ \{L^2(\mathbf{F}_h)\}^3 | (\eta, \mathbf{v})_{\mathbf{v}} = 0, \forall \mathbf{v} ∈ \mathbf{F}_h^b\},
\]

with $L^2(\mathbf{F}_h)$ the space of square-integrable functions on $\mathbf{F}_h$. The hybrid variable $\mathbf{H}_h^0$ belongs to this space $\mathbf{M}_h^0$.

Thus, from (1), we obtain disconnected local problems in each element and a reduced-size linear system can be assembled only for the hybrid variable in a classical finite element way starting from the interface conditions (3); see for instance [5] for more details. Once the linear system for the unknowns associated to the hybrid variable has been solved, the local electromagnetic field within each $K$ can be recovered locally by solving (1).

III. PRELIMINARY NUMERICAL RESULTS

A. Accuracy of the method

We consider the propagation of a plane wave in vacuum. The computational domain is chosen to be the unit cube...
Ω = (−0.5, 0.5)³, and a Silver-Müller absorbing boundary condition is imposed on the whole boundary. Parameters , , , and are set to 1, and .

A sequence of regular tetrahedral meshes is employed. Table I displays the error in L²-norm and the convergence orders estimated between two consecutive meshes, i.e. an estimate of the exponent in the asymptotic convergence rate h", for the HDG method and p = 1. Table II yields the same characteristics for p = 2. In Table I, "mesh size" denotes the edge length of the tetrahedrons on the edge of the unit cube and the number of degrees of freedom for the hybrid variable. We observe that the asymptotic convergence orders of the approximate solutions for both HDG and H are optimal, i.e. of order p + 1 for both H and HDG when using polynomial order p. This convergence rate has been proved for the Helmholtz equation in 2d [6] but no theoretical result is currently available for Maxwell’s equations in 3d.

### TABLE I: CONVERGENCE OF THE HDG METHOD WITH p = 1

<table>
<thead>
<tr>
<th>Mesh</th>
<th>N₀</th>
<th>Error [H₁ₐ]</th>
<th>Error [H₂₀]</th>
</tr>
</thead>
<tbody>
<tr>
<td>1/4</td>
<td>720</td>
<td>2.86e-1</td>
<td>2.86e-1</td>
</tr>
<tr>
<td>1/4</td>
<td>5184</td>
<td>6.02e-2</td>
<td>6.68e-2</td>
</tr>
<tr>
<td>1/4</td>
<td>39168</td>
<td>1.54e-2</td>
<td>1.76e-2</td>
</tr>
</tbody>
</table>

### TABLE II: CONVERGENCE OF THE HDG METHOD WITH p = 2

<table>
<thead>
<tr>
<th>Mesh</th>
<th>N₀</th>
<th>Error [H₁ₐ]</th>
<th>Error [H₂₀]</th>
</tr>
</thead>
<tbody>
<tr>
<td>1/4</td>
<td>1490</td>
<td>3.15e-2</td>
<td>3.36e-2</td>
</tr>
<tr>
<td>1/4</td>
<td>10368</td>
<td>4.00e-3</td>
<td>4.44e-3</td>
</tr>
<tr>
<td>1/4</td>
<td>78336</td>
<td>4.96e-4</td>
<td>5.53e-4</td>
</tr>
</tbody>
</table>

### IV. CONCLUSION AND PERSPECTIVES

HDG methods lead to the same optimal accuracy as some classical DG methods but greatly reduce the computational burden in comparison. They make DG methods more attractive for stationary problems. We are currently working on efficient strategies for solving the linear systems for the hybrid variable unknowns, as well as on the further development and application of the proposed HDG method to more challenging propagation problems. Other questions need to be addressed, especially the use of a variable order p and non-conforming meshes, as it can be naturally dealt with the HDG strategy [4]. Moreover, it should also facilitate the rigorous coupling with other discretizations because the HDG approach clearly separates the local volume problems (1) and the interface conditions (3).

### REFERENCES


