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Explicit computation of the electrostatic energy for an elliptical charged disc

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\section*{Abstract}
This letter describes a method for obtaining an explicit expression for the electrostatic energy of a charged elliptical infinitely thin disc. The charge distribution is assumed to be polynomial. Such explicit values for this energy are fundamental for assessing the accuracy of boundary element codes. The main tools used are an extension of Copson’s method and a diagonalization, given by Leppington and Levine, of the single-layer potential operator associated with the electrostatic potential created by a distribution of charges on the elliptical disc.

\section{1. Introduction}

In recent years, integral equations have become an essential tool for solving both industrial and scientific problems in electromagnetism and acoustics. The assessment of the accuracy delivered by such codes, in particular in their handling of the singular integrals involved, is a major issue. Here, we present a method for deriving an analytical expression for the electrostatic energy of a charged elliptical infinitely thin plate, providing a means for the validation of these codes.

Let us denote by \( A = \{(x_1, x_2) \in \mathbb{R}^2 \text{ with } \frac{x_1^2}{a^2} + \frac{x_2^2}{b^2} - 1 < 0 \text{ and } a > b\} \) the ellipse with major and minor semi-axes \( a \) and \( b \). Let \( f \) be the electrostatic potential generated by a density of charges \( \sigma \) distributed over \( A \):

\[
f(x) = \frac{1}{4\pi} \int_A \frac{\sigma(y)}{|x - y|} dy \quad \text{for all } x \in A,
\]

with \( x = (x_1, x_2) \). The units have been chosen such that the electric permittivity of air is 1. The electrostatic energy \( I \) can be expressed in either the two following forms:

\[
I_\sigma = \int_A f(x) \sigma(x) dx = \int_A \frac{1}{4\pi} \int_A \frac{\sigma(x) \sigma(y)}{|x - y|} ds_x ds_y.
\]
We aim in this letter at proving and numerically illustrating the following theorem, where \( \varepsilon \) is the eccentricity of the ellipse \( A \) given by \( \varepsilon = \sqrt{1 - b^2/a^2} \).

**Theorem 1.1.** Let \( \sigma(x) = \alpha_0 + \alpha_1x_1/a + \alpha_2x_2/b \), with \( \alpha \in \mathbb{R}^3 \), be the distribution of charges over \( A \). The corresponding electrostatic energy is given by

\[
I_\sigma = \frac{8ab^2}{15\pi} \left[ (5\alpha_0^2 + \alpha_1^2 + \alpha_2^2) K(\varepsilon) + (\alpha_1^2 - \alpha_2^2) \frac{K(\varepsilon) - E(\varepsilon)}{\varepsilon^2} \right],
\]

with \( K(\varepsilon) \) and \( E(\varepsilon) \) the complete elliptic integrals of the first and second kind

\[
K(\varepsilon) = \int_0^{\pi/2} \frac{d\phi}{\sqrt{1 - \varepsilon^2 \sin^2 \phi}} \quad \text{and} \quad E(\varepsilon) = \int_0^{\pi/2} \sqrt{1 - \varepsilon^2 \sin^2 \phi} d\phi. \tag{3}
\]

2. **Diagonalization of the electrostatic energy**

Following [1], we consider the spheroidal coordinate system \((\theta, \varphi)\) giving a parametrization of \( A \) in terms of the unit half-sphere

\[
x_1 = a \sin \theta \cos \varphi \quad \text{and} \quad x_2 = b \sin \theta \sin \varphi, \quad \text{with} \quad \theta \in [0, \pi/2], \ \varphi \in [0, 2\pi]. \tag{4}
\]

The elemental area associated with the new variables is \( ab \cos \theta \sin \theta \) d\theta d\varphi. In these spheroidal coordinates, the electrostatic potential \( f \) defined in (1) and the electrostatic energy can also be written in terms of \( \theta \) and \( \varphi \) as

\[
\begin{align*}
\left\{ \begin{array}{l}
f(\theta, \varphi) = \frac{ab}{4\pi} \int_0^{\pi/2} \int_0^{2\pi} g(\theta', \varphi') \sin \theta' \, d\theta' \, d\varphi', \\
I_\sigma = ab \int_0^{\pi/2} \int_0^{2\pi} f(\theta, \varphi) g(\theta, \varphi) \sin \theta \, d\theta d\varphi,
\end{array} \right.
\end{align*}
\]

with \( g(\theta, \varphi) = \sigma(\theta, \varphi) \cos \theta \) and \( d(\theta, \varphi, \theta', \varphi') \) the distance separating \( x \) from \( y \) expressed in the spheroidal coordinates (4).

The next step consists in introducing a well chosen spectral basis for the half-sphere involving the even Legendre functions \( Q_n^m \) normalized by

\[
\int_0^{\pi} Q_n^m(\cos \theta) Q_{n'}^{m'}(\cos \theta) \sin \theta d\theta = \delta_{n,n'}. \tag{5}
\]

This basis yields a block diagonalization of the convolution operator (see [1])

\[
\frac{1}{d} = \frac{1}{\sqrt{ab}} \sum_{n=0}^{\infty} \sum_{m=-n}^{n} \sum_{m'=-n}^{n} d_{nmn'} Q_n^m(\cos \theta) Q_{n'}^{m'}(\cos \theta') e^{i(m\varphi - m'\varphi')},
\]

with

\[
d_{nmn'} = \frac{Q_n^m(0) Q_{n'}^{m'}(0)}{2n+1} \int_0^{2\pi} \frac{e^{i(m-m')\varphi}}{\sqrt{a^2 \cos^2 \varphi + b^2 \sin^2 \varphi}} d\varphi. \tag{6}
\]
The functions $f$ and $g$ can be expanded in this basis as

$$u(\theta, \phi) = \sum_{n,m=-n}^{\infty} u_n^m Q_n^m(\cos \theta) e^{im\phi} \quad \text{with } u = f \text{ or } g$$  \hspace{1cm} (7)

with

$$u_n^m = \frac{1}{\pi} \int_0^{\pi/2} \int_0^{2\pi} u(\theta, \phi) Q_n^m(\cos \theta) e^{-im\phi} \sin(\theta) d\theta d\phi.$$  \hspace{1cm} (8)

Due to (6), coefficients $f_n^m$ are related to $g_n^m$ by

$$f_n^m = \sqrt{ab} \frac{4}{n-m} \sum_{m'=-n}^{n} d_{nmn'}^m g_{m'n'}^m.$$  \hspace{1cm} (9)

Moreover, the orthogonal properties of the spectral basis yield

$$I_{\sigma} = \frac{\pi}{4} (ab)^{3/2} \sum_{n,m,m'} d_{nmn'}^m g_{n'm'}^m,$$

where we have lightened the notation by making the range of the summation index implicit. Indices $n, n'$ are varying from 0 to $\infty$, and $m, m'$ are such that $|m| \leq n$, $|m'| \leq n'$, with $n-m$ and $n-m'$ even. Substituting expression (6) for $d_{nmn'}^m$ and introducing the eccentricity of the ellipse $\varepsilon$, we get

$$I_{\sigma} = \pi ab^2 \sum_{n,m,m'} g_{n'm'}^m g_{n'm'}^m \frac{Q_n^m(0) Q_{n'}^m(0)}{2n+1} \int_0^{\pi/2} \cos(m-m')\varphi \sqrt{1-\varepsilon^2 \cos^2 \varphi} d\varphi.$$  \hspace{1cm} (10)

3. Electrostatic energy for an affine distribution of charges

This section is dedicated to the calculation of the electrostatic energy generated by an affine density of charges.

**Proof of Theorem 1.1.** Let $\sigma_0(x) = 1$, $\sigma_1(x) = x_1/a$ and $\sigma_2(x) = x_2/b$. In view of the symmetry of $A$ with respect to $x_1$ and $x_2$, we have

$$\int_A \frac{1}{4\pi} \int_A \sigma_i(x) \sigma_j(y) \frac{ds_x ds_y}{|x-y|} = 0 \quad \text{for } i \neq j.$$  \hspace{1cm} (11)

Consequently, the electrostatic energy $I_{\sigma}$ can be expanded as

$$I_{\sigma} = \alpha_0^3 I_{\sigma_0} + \alpha_1^2 I_{\sigma_1} + \alpha_2^2 I_{\sigma_2}.$$  \hspace{1cm} (12)

The result will follow from the computation of $I_{\sigma_0}$, $I_{\sigma_1}$ and $I_{\sigma_2}$.

3.1. Computation of $I_{\sigma_0}$

For $\sigma(x) = \sigma_0(x) = 1$, the function $g(\theta, \varphi) = \cos \theta$ does not depend on $\varphi$. The $g_n^m$ coefficients are independent from $a$ and $b$, and since $g_n^m = 0$ for all
$m \neq 0$, the function $g$ can be expanded as $g(\theta) = \sum_{n=0}^{+\infty} g_n^0 Q_n^0 (\cos \theta)$. Due to (10), the electrostatic energy $I_{\sigma_0}$ depends only on $a$ and $b$ and is given by

$$I_{\sigma_0}(a, b) = \pi ab^2 \sum_{n=0}^{+\infty} \frac{|g_n|^2} {2n + 1} \int_0^{\pi/2} \frac{d\varphi} {\sqrt{1 - \varepsilon^2 \cos^2 \varphi}} = \kappa ab^2 K(\varepsilon),$$

(11)

with $\varepsilon = \sqrt{1-b^2/a^2}$ and $\kappa$ a constant depending neither on $a$ nor on $b$. The constant $\kappa$ is deduced from the classical case of an unit circle which has been detailed for example in [2]

$$I_{\sigma_0}(1, 1) = \int_C \int_C \frac{1}{|x - y|} ds_x ds_y = 4/3.$$  
(12)

Comparing (11) and (12), this yields to $\kappa = 8/3\pi$ since $K(0) = \pi/2$. Therefore,

$$I_{\sigma_0} = \frac{8}{3\pi} ab^2 K(\varepsilon).$$
(13)

3.2. Computation of $I_{\sigma_1}$

For $\sigma(x) = \sigma_1(x) = x_1/a$, the function $g$ is given by $g(\theta, \varphi) = \sin \theta \cos \theta \cos \varphi$. In that case, the $g_n^m$ are zero except for $|m| = 1$. By definition of the Legendre functions, we have $Q_n^{-1} = -Q_n^1$. As the function $g$ is even, it emerges that $g_n^{-1} = -g_n^1$, and thus

$$I_{\sigma_1} = \pi ab^2 \sum_n \frac{(g_n^1 Q_n^1 (0))^2} {2n + 1} \left[ 2 \int_0^{\pi/2} \frac{d\varphi} {\sqrt{1 - \varepsilon^2 \cos^2 \varphi}} + 2 \int_0^{\pi/2} \frac{\cos 2\varphi} {\sqrt{1 - \varepsilon^2 \cos^2 \varphi}} d\varphi \right].$$

Due to (3), it emerges that

$$I_{\sigma_1}(a, b) = \kappa ab^2 \frac{K(\varepsilon) - E(\varepsilon)} {\varepsilon^2}.$$  
(14)

To determine the constant $\kappa$, we consider again the case of an unit circle. In this case, $I_{\sigma_1}$ can be explicitly computed (see the Appendix A), and is given by $I_{\sigma_1}(1, 1) = 2/15$. Evaluating (14) at $a = b = 1$ we get

$$I_{\sigma_1}(1, 1) = \frac{\pi} {4} \text{ since } \lim_{\varepsilon \to 0} \frac{K(\varepsilon) - E(\varepsilon)} {\varepsilon^2} = \frac{3\pi} {4}.$$  
(15)

It follows that $\kappa = 8/15\pi$ and therefore we have

$$I_{\sigma_1} = \frac{8}{15\pi} ab^2 \frac{K(\varepsilon) - E(\varepsilon)} {\varepsilon^2}.$$  
(16)
3.3. Computation of $I_{\sigma_2}$

For $I_{\sigma_2}$, we consider $\sigma(x) = \sigma_2(x) = x^2/b$, meaning that $g(\theta, \varphi) = \sin \theta \cos \theta \sin \varphi$. We still have $g_m^m = 0$ except for $|m| = 1$, but in that case, $g_m^{-1} = g_m^1$. Thus

$$I_{\sigma_2} = \pi ab^2 \sum_{n=0}^{+\infty} \left( \frac{g_n^1 Q_n^1(0)}{2n+1} \right)^2 \left[ 2 \int_0^{\pi/2} d\varphi \frac{\cos \varphi}{\sqrt{1 - \varepsilon^2 \cos^2 \varphi}} - 2 \int_0^{\pi/2} \cos 2\varphi \ d\varphi \right]$$

Moreover, both integrals $I_{\sigma_1}$ and $I_{\sigma_2}$ are equal on $C$ by symmetry. We obtain

$$I_{\sigma_2} = \frac{8}{15\pi} ab^2 \left( K(\varepsilon) - \frac{E(\varepsilon)}{\varepsilon^2} \right). \quad (17)$$

4. Numerical tests and conclusion

Tables 1 and 2 give a comparison of the exact values given by an analytical expression with numerical approximate values obtained by the boundary element code CESC of CERFACS with $P^1$ continuous elements. It can be observed that the two values coincide at least up to the fourth decimal digit. Table 3 shows the maximum relative error for each of the cases of Tables 1 and 2, which is less than 0.35 per mil.

![Table 1](image1.png)

![Table 2](image2.png)

![Table 3](image3.png)
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Appendix A. The case of an unit circle disc

Let C be the circle with radius 1. We aim here at computing the integral

\[ I_C = \frac{1}{4\pi} \int_C \frac{x_1 y_1}{|x - y|} \, ds_x ds_y \]

This integral is rewritten in polar coordinates (r, \( \phi \) for x and \( \rho, \phi' \) for y) as

\[ I_C = \frac{1}{4\pi} \int_0^1 \int_0^{2\pi} \int_0^1 \int_0^{2\pi} \frac{r \cos \phi \rho \cos \phi'}{\sqrt{r^2 + \rho^2 - 2r \rho \cos(\phi - \phi')}} r dr d\phi \rho d\rho' d\phi' \]

and evaluated using the formula (3.4.5) of [3, p70]

\[ \int_0^{2\pi} \cos \phi \, d\phi = 4 \cos \phi' \rho \int_0^{\min(\rho, r)} t^2 \, dt \int_0^{2\pi} \cos^2 \phi' \, d\phi'. \]

This leads to

\[ I_C = \frac{1}{\pi} \int_0^1 \int_0^1 \int_0^{\min(\rho, r)} \frac{\rho t^2}{\sqrt{\rho^2 - t^2} \sqrt{r^2 - t^2}} \, d\rho \, dt \, dr \int_0^{2\pi} \cos^2 \phi' \, d\phi'. \]

This integral is symmetric in \( \rho \) and \( r \), and since \( \int_0^{2\pi} \cos^2 \phi' \, d\phi' = \pi \), we have

\[ I_C = 2 \int_{t=0}^1 t^2 \int_{\rho=t}^1 \frac{\rho}{\sqrt{\rho^2 - t^2}} \int_{r=\rho}^1 \frac{r}{\sqrt{r^2 - t^2}} \, dr \, d\rho \, dt = \frac{2}{15}. \]

References

