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Reality-Based Real-Time Cell Indentation Simulator

Hamid Ladjal, Jean-Luc Hanus, Anand Pillarisetti, Carol Keefer, Antoine Ferreira, Member, IEEE, and Jaydev P. Desai, Senior Member, IEEE

Abstract—Training simulators that provide realistic visual and haptic feedback during cell indentation tasks are currently investigated. Complex cell geometry inherent to biological cells and intricate mechanical properties drive the need for precise mechanical and numerical modeling to assure accurate cell deformation and force calculations. Advances in alternative finite-element formulation, such as the mass–tensor approach, have reached a state, where they are applicable to model soft-cell deformation in real time. The geometrical characteristics and the mechanical properties of different cells are determined with atomic force microscopy (AFM) indentation. A real-time, haptics-enabled simulator for cell centered indentation has been developed, which utilizes the AFM data (mechanical and geometrical properties of embryonic stem cells) to accurately replicate the indentation task and predict the cell deformation during indentation in real time. This tool can be used as a mechanical marker to characterize the biological state of the cell. The operator is able to feel the change in the stiffness during cell deformation between fixed and live cells in real time. A comparative study with finite-element simulations using a commercial software and the experimental data demonstrate the effectiveness of the proposed physically based model.

Index Terms—Atomic force microscope (AFM), finite-element modeling (FEM), haptics, modeling, real-time interaction, stem cell.

I. INTRODUCTION

MECHANICAL manipulation and characterization of biological cells is currently one of the most promising research areas in the field of biorobotics applied to cellular level interactions. The mechanical properties, such as elasticity, membrane tension, cell shape, and adhesion strength, may play an important role in cell development and differentiation [1]–[3]. In contrast to adult stem cells, embryonic stem cells differentiate into every cell type in the body, and thus, have a huge therapeutic potential. Stem-cell-based systems offer a very promising and innovative solution for obtaining large number of cells for early efficacy and toxicity screening. Stem cell technology provides a new tool for drug development and a better insight to understand mammalian gene function (see Fig. 1). The mechanical properties of biological cells have been studied with different techniques [4], the most popular being optical tweezers [5], magnetic beads [6], and micropipette aspiration [7]. Different microrobotics systems with various end effectors have been proposed [33], [34]. However, these methods cannot compete with the precision that can be obtained with the atomic force microscopy (AFM) method [8], [32]. AFMs have been widely used in the study of micro- and nanostructures including living cells. Modern AFM techniques allow solving several problems of cell biomechanics by simultaneous evaluation of the local mechanical properties and the topography of living cells, at a high spatial resolution and force sensitivity [9]. In these experiments, an AFM cantilever serves as a microindenter to probe the cell, and further analysis of force-indentation data yields the local Young’s modulus. In addition, the AFM indentation technique can be used to characterize the viscoelastic behavior of the cell cytoskeleton [10], including viscosity [11], loss and storage moduli [12], and stress relaxation times. Mechanical phenotyping prove actually to be a valuable tool in the development of improved methods of targeted cellular differentiation of embryonic, induced pluripotent (iPS), and adult stem cells for therapeutic and diagnostic purposes [30], [31]. In this study, we propose to characterize the mechanical behavior of an individual mouse embryonic stem cells (mESC) in undifferentiated state...
Using AFM indentation [13]. As we hypothesize that the cell stiffness of mESC may reflect the nuclear physical changes of the cell, it should be possible to detect modifications in stiffness by probing the cell directly. However, different limitations are inherent in AFM-based indentation techniques. First, deformations induced by the AFM probe and cell positions cannot be seen by the user. Second, the force sensing of the AFM cannot resolve the force interaction between the AFM tip and the cell’s surface independently in all three dimensions to avoid cytoskeleton-disrupted experiments. Therefore, the user performs extremely challenging manipulations requiring advanced manual skills. Modeling the physics of the interaction between tip and cell is necessary to realize complicated indentation tasks like cancer cell probing [14].

We developed finite-element models (FEMs) that include the topological information of adherent stem cells (shape and dimensions), AFM tips (conical and spherical), and biological structure (cell models with cytoplasm layers, cytoskeletons, and nucleus) [15]. To validate the viscoelastic and nonlinear dynamics of the cell in real time using finite-element procedures, we report herein the mechanical characterization of single mESC indentation-relaxation for different colonies of fixed and live undifferentiated mESCs (termed, undiff mESC). The proposed indentation tests use microsphere-modified AFM probes in order to estimate the global elastic modulus of the cell to reflect the true global response of a mESC [14], [15]. This paper is a subsequent contribution to [27], where various analytical models are tested and compared to estimate the mechanical properties of the biological cells. Based on the proposed physical model, we have developed a computer-based training system using force feedback to simulate cell indentation procedures in virtual environments (VEs) for biologist training. The paper consists of five sections. In Section II, we present the biomechanical finite-element approach dedicated to real-time indentation. In Section III, we present the methodology of the VE system for cell indentation. In Section IV, we present the experimental setup using the force microscopy. In Section V, we make a comparative study between the physics-based model and experimental data. In Section VI, we present our real-time cell indentation simulator using a spherical tip. Finally, we give some concluding remarks and the directions for future work.

II. PHYSICS-BASED DEFORMABLE MODELS FOR CELL NANOINDENTATION SIMULATION

In this section, we present further developments to the linear elastic mass–tensor model introduced by [16], [17] and extended by Schwartz et al. [18]. In our paper, the cell is modeled by a volume object discretized into a conformal tetrahedral mesh as defined by finite-element theory.

Using classical notations, inside each tetrahedron $T^k$, the displacement field, defined by a linear interpolation $[N^k]$ of the nodal displacement vector $\{u^k\}$ of the four vertices of tetrahedron, is written as follows:

\[
\{U(x)^k\} = [N^k(x)] \{u^k\},
\]

(1)

For a linear elastic material, the relation between the Cauchy stress tensor and the linearized strain tensor is written with Lame’s coefficient in condensed vector notation as follows:

\[
\{\sigma\} = \lambda ([\varepsilon])_1 + ([\varepsilon])_2 + ([\varepsilon])_3 [I] + 2\mu ([\varepsilon])
\]

(2)

where $[I]$ is the identity matrix.

The principle of virtual work applied to a single tetrahedron $T^k$ leads to the elementary stiffness matrix $[K^k]$ such that the elementary nodal force vector acting on a tetrahedron is as follows:

\[
\{f^k\} = [K^k] \{u^k\}.
\]

(3)

This stiffness matrix is composed of a plurality of elementary submatrices, each connecting the elementary force acting on the node $i$ to the displacement of the node $j$

\[
[K^k]_{ij} = \frac{1}{36 V^k} \left( \lambda \{m_i\} \{m_j\}^T + \mu \{m_i\} \{m_j\}^T + \mu \{m_i\} \{m_i\}^T \{m_j\} \right)
\]

(4)

where $\{m_i\}$ are unit outward-pointing normals to triangular faces and $V^k$ is the volume of the tetrahedron $T^k$.

Taking into account the contribution of all adjacent tetrahedra, the global internal force acting on a node $l$ can be expressed as follows:

\[
\{F_{int}^l\} = \sum_{k \in V_l} \left( \sum_{j=1}^{4} [K^k]_{ij} \{u_j\} \right)
\]

(5)

where $V_l$ is the neighborhood of vertex $l$ (i.e., the tetrahedra containing node $l$).

The tensors $[K]_{ij}$ depending on the remaining geometry and Lame’s coefficients are constant. They can be precomputed in an offline phase. This is the essential advantage of the mass–tensor approach, which makes it useful for real-time application.

A. Cell Modeling Using FEM

In our paper, the mESC is meshed with 3D first-order tetrahedral elements, as shown in Fig. 2. The geometrical characteristics of the cells were obtained from experimental data using the AFM system. This includes shape, diameter, and height of the cells. The mesh that has been used for the following simulations is composed of 835 vertices and 2456 tetrahedra, where all vertices are free to move except the ones in contact with the petri dish, which are fixed.

B. Dynamic Model

The equation of motion of a vertex $l$ of the cell mesh can be written as follows:

\[
M^l \{\ddot{u}_l\} + \gamma^l \{u_l\} + \sum_{k \in V_l} \left( \sum_{j=1}^{4} [K^k]_{ij} \{u_j\} \right) = \{F_{ext}^l\}
\]

(6)

where $M^l$ and $\gamma^l$ are, respectively, the mass and damping coefficients of each vertex.
To solve the dynamic system, we tested different integration schemes (implicit and explicit) taking into account the tradeoff between real-time simulation and haptic stability requirements. We choose the explicit centered finite-difference scheme defined by

\[
\begin{align*}
\dot{u}(t+h) &= \dot{u}(t) + \frac{1}{2} \ddot{u}(t) h + \frac{1}{2} \dddot{u}(t) h^2 + O(h^3) \\
\dddot{u}(t+h) &= \dddot{u}(t) + \frac{1}{2} \ddot{u}(t) h + \frac{1}{2} \dot{u}(t+h) h + O(h^3)
\end{align*}
\]

(7)

where \( h \) is the temporal integration step size chosen to satisfy the Courant–Friedrich–Lewy condition.

The discrete equation of motion of a vertex \( l \) reads

\[
\begin{align*}
\left( M^l + \frac{h^2}{2} \gamma^l \right) \{ \dddot{u}_l \}_{(t+h)} &= \{ F^l_{\text{ext}} \}_{(t+h)} - \gamma^l \left( \{ \dddot{u}_l \}_{(t)} + \frac{h}{2} \{ \ddot{u}_l \}_{(t)} \right) \\
&\quad - \sum_{k \in \Omega} \left( \frac{4}{3} \sum_{j=1}^{4} \left[ K^l_{jk} \right] \{ u_j \}_{(t)} + h \{ \dot{u}_j \}_{(t)} + \frac{h^2}{2} \{ \ddot{u}_j \}_{(t)} \right)
\end{align*}
\]

(8)

from which we can compute the new vertices for accelerations, velocities, and positions.

III. VE SYSTEM FOR CELL INDENTATION

Fig. 3 shows the architecture of the real-time virtual-reality-based cell indentation simulator. This includes the computer generated models of stem cells, the AFM nanoindentor, the collision detection algorithm, physics-based models of deformable mESC, and the haptic interaction controller. The operator will be able to interact with the 3-D model of the cell using his sense of vision as well as actively manipulate using his sense of touch. Usually, real-time graphics translates to an update rate of 25 Hz, stable haptic interaction in VE’s requires much higher update rate of around 1 kHz. As soft biological tissues exhibit complex viscoelastic and nonlinear properties, the real-time interaction imposes severe restrictions on the FEM model.

A significant difficulty of using the finite-element technique for real-time simulation is that it is computationally costly. These mesh-based schemes also require an expensive numerical integration operation for the computation of the system stiffness matrices. For this reason, we adopted a computational architecture with two simulation stages: 1) an offline computation of stiffness matrices for each triangular element and 2) visual and haptic interaction for real-time indentation simulation.

A. Offline Computation

The proposed computational methodology contains an offline precalculation step. The most costly and time-consuming operations are realized during this step. The database contains various developed models of the different stem cells geometry (bone marrow, pancreatic, heart muscle, nerve cells, etc.). Using AFM image analysis (ImagePro software from AFM Asylum), we derived different geometric stem cell models based on a commercial computer-aided design (CAD) package (MARC-ADAMS).

The meshing of the 3-D internal structure of the mESC is then carried out through a dedicated 3-D meshing software (GID software) modeled in exact dimensions. Although they were displayed as 3-D texture-mapped objects to the user, they were modeled as connected line segments to reduce the number of collision computations during real-time interactions.

The adherent cell was approximated as an assembly of discrete tetrahedral elements interconnected to each other through a fixed number of nodes (see Fig. 2). The displacements of these nodal points for applied external indentation forces were the basic unknowns of our FEM analysis. The coordinates...
An important part of biological can-
248, sions, such as punction or cutting. Physical simulations suppose
250, in the velocities of the involved objects, and topology modifica-
252, sion between two objects may generate local and global defor-
254, mplex 3-D object, composed of numerous surfaces, edges, and
256, vertices, but this constitutes a challenging task. For the purpose
258, of detecting collisions, we used a point-based representation of
260, the cantilever tip as a simple spherical object and employed a
262, simple sphere-points intersection algorithm with local search
264, technique. It calculates the distance between the center of the
266, spherical tip and the nodes of the cell surface is calculated in real time.
267, of vertices, the tetrahedron indexing, and the connectivity of
269, vertices were derived from the geometric model. Then, the
270, mechanical (Young’s modulus and stiffness) and geometrical
271, (height and shape) properties of the cell structure are determined
273, through AFM indentation experiments. Finally, these properties
274, are passed to the FEM algorithm for calculating the physically
276, based behavior of biological soft tissues. The offline calcula-
278, tion of all tensors $K_{ij}$ allows precalculation stiffness values for
279, preselected cells such as, for example, live and fixed cells.

280, **B. Real-Time Indentation**

281, Since haptics and graphics have different update frequencies,
282, we implemented separate threads to update the loops. During
283, the indentation task, the contact between the tip and the cell
284, must occur at a special set of points (nodal points).

285, 1) **Collision Detection:** An important part of biological can-
286, tilever tip-soft cell modeling is the fast collision detection al-
287, gorithm. The collision detection is the first step to carry out
288, realistic interactions. However, locating the contact primitive
289, (e.g., facets) between two objects may be computationally expen-
290, sive, especially if the objects are composed of a large number
291, of polygons. In general, to determine a contact between two vir-
292, tual objects, we compute the distance between them. If this
293, distance is negative, then the objects are in contact. We may
294, continue by determining the set of geometric objects or entities
295, that collide. The cantilever tip in VE can be modeled as a com-
296, plex 3-D object, composed of numerous surfaces, edges, and
297, vertices, but this constitutes a challenging task. For the purpose
298, of detecting collisions, we used a point-based representation of
299, the cantilever tip as a simple spherical object and employed a
300, simple sphere-points intersection algorithm with local search
301, technique. It calculates the distance between the center of the
302, spherical tip and the nodes of the cell surface (see Fig. 4).

303, 2) **Penalty-Based Models:** In an interaction model, a colli-
304, sion between two objects may generate local and global defor-
305, mations of the colliding objects, new internal forces, changes
306, in the velocities of the involved objects, and topology modifica-
307, tions, such as punction or cutting. Physical simulations suppose
308, that these events are related to the mechanical properties of the
309, objects. In our simulations, we used the penalty-based models.

These models consider the reactive force that first reduces the
relative speed of the two objects, and finally, repulses them from
each other, as proportional to the deformation. Thus, the reactive
force is a function of the local interpenetration depth $\delta$. The
normal repulsive force $F_{\text{rep}}$ can be written as follows:

$$ F_{\text{rep}} = \lambda \delta \{n\} $$

(9)

where $\lambda$ is the penalty of the interpenetration and $\{n\}$ is the
outward-pointing normal.

These contact forces are evaluated at each time step $t$ (see
Fig. 5). They represent the new updated external forces acting
upon our stem cell.

3) **Virtual Coupling for Stability of the Haptic Rendering:**

The internal operating loop of the haptic interface requires an
update frequency around $1 \text{kHz}$. However, the FEM update fre-
quency is in the range of $25–30 \text{Hz}$. This frequency difference
threatens the coherence between both systems leading to insta-
bilites of the user haptic rendering. We adopted the solution to
use a virtual coupling model defined in [24].

This approach introduces a virtual passive link between the
simulation model and the haptic interface in order to ensure
the stability and the performance of the system [see Fig. 6(a)].

When we combine the impedance display implementation with
an appropriate virtual coupling network, we get the admittance

Fig. 4. Scheme illustrating the multiple collision detection algorithm using a
simple sphere-points intersection problem: the distance between the center of
the spherical tip and the nodes of the cell surface is calculated in real time.

Fig. 5. Main flowchart for multi collision detection: $M$ represents the position
of the center of the spherical tip and $R$ its radius, $i$ represents the current external
cell surface node, $d_i$, $\{F_{\text{rep}}\}$ and $\{n_i\}$ represent, respectively, the computed
distance between the center of the tip and the current node, the repulsive force,
and the outward-pointing contact normal at node $i$. **
matrix for the combined interface. The linear two-port is said to be absolutely stable if there exists no set of passive terminating one-port impedances for which the system is unstable. Llewellyn’s stability criteria [25] provides both necessary and sufficient conditions for absolute stability of linear two-ports. We get the conditions for absolute stability of the haptic interface

\[
\text{Re}(Z_{di}(z)) \geq 0, \quad \frac{1}{Z_{cvl}(z)} \geq 0 \quad (10)
\]

\[
\cos(\angle\text{ZOH}(z)) + \frac{2\text{Re}(Z_{di}(z))\text{Re}(1/Z_{cvl}(z))}{|\text{ZOH}(z)|} \geq 1 \quad (11)
\]

where \(Z_{cvl}(z)\) is the virtual coupling impedance \((k_c, b_c)\), \(\text{ZOH}(z)\) is a zero-order holder, and \(Z_{di}(z)\) is the PHANToM impedance. The inequality (11) can be rewritten to get an explicit expression of absolute stability of the haptic interface

\[
\text{Re}\left(\frac{1}{Z_{cvl}(z)}\right) \geq 1 - \frac{\cos(\angle\text{ZOH}(z))}{2\text{Re}(Z_{di}(z))} |\text{ZOH}(z)|. \quad (12)
\]

For the virtual coupling, the impedance display induces a limit on the maximum impedance which can be rendered. We use (12) to find the virtual coupling, which makes the haptic interface absolutely stable. Using the design bound of the best performing, absolutely stabilizing virtual coupling. These parameters are found as \(b_c = 0.008 \text{ N/(mm/s)}\) and \(k_c = 2.3 \text{ N/mm}\) (plotted on Fig. 6(a) as a thin red line). The left side of (12) with the resulting values is plotted on Fig. 6(b) as a bold blue line.

Fig. 6. Stability conditions using the virtual coupling model. (a) Model of virtual coupling. (b) Impedance display virtual coupling.

### IV. EXPERIMENTAL SETUP

#### A. Description

The experimental tests of the cell indentations were performed using an AFM (Model: MFP-3D-BIO, Asylum Research, Santa Barbara, CA). The AFM is integrated with a top view module and mounted on an active vibration isolation table manufactured by Herzan (Laguna Hills, CA), as shown in Fig. 7. The top view module enables viewing of cells and easy alignment of the laser beam on the AFM cantilever. The XY stage (manual) allows the user to position the cell beneath the cantilever tip of the AFM. The entire AFM setup is enclosed in an acoustic isolation chamber to prevent acoustic noise from interfering with the AFM measurements. The \(x\)- and \(y\)-axis ranges of the scan head are 90 \(\mu\text{m}\). The \(z\)-axis scan range is 40 \(\mu\text{m}\). The AFM system is used to obtain force and cell deformation data from biological samples.

The cantilever is moved by the piezoelectric scanner in \(z\)-direction toward the cell. The deflection of the cantilever is detected by a photodiode when the tip comes in contact with the cell. When the tip of the cantilever is in contact with the cell, the initial cantilever deflection \(d_0\), and initial cantilever movement in \(z\)-direction \(z_0\) are stored. As the cantilever moves further in \(z\)-direction and deforms the cell, the final cantilever deflection \(d_1\) and the cantilever movement \(z_1\) are obtained [21].

The stiffness depends not only on the Young’s modulus but also on the geometry of the tip-surface contact. Therefore, the geometry and spring constant of the cantilever are calibrated in the same way for live and fixed cells. In the following experiments, we employed two types of cantilevers attached with a spherical probe (5 \(\mu\text{m}\) in diameter).

1. Silicon nitride cantilever with a spring constant of 0.06 \(\text{N/m}\) (Novascan Technologies, Inc., Ames, IA) for live cells.
2. Silicon cantilever with a spring constant of 1.75 \(\text{N/m}\) (Novascan Technologies, Inc.) for fixed cells.
Fig. 8. Thermal graph measurement of a spherical tip cantilever ($K_c = 1.29$ N/m) used for fixed cells indentation. The left corner inset represents the contact slope from a force curve to determine the sensitivity of the cantilever and (deflection volt = 98.01 nm/v) and the right corner inset represents the zoom-in view of the resonant frequency plot ($F = 46.043$ kHz).

**B. AFM Calibration**

Before any experiment is done, it is necessary to perform a calibration of the cantilever to determine its resonant frequency. The resonant frequency measurement enables the user to determine if it lies within the specified range provided by the manufacturer in order to detect probe defects. The determination of spring constant $K_n$ by the thermal method involves both measurements: 1) sensitivity in nanometer per volt and 2) first-resonant frequency of the cantilever.

1) Determine the contact slope of the indentation curve on a hard surface to determine the sensitivity of the cantilever (see left corner inset of Fig. 8).

2) Measure the thermal power spectral density (PSD) to determine the resonant frequency of the cantilever and to confirm that the cantilever and light source are aligned. This confirms that the system is functioning properly. The results are plotted as the deflection amplitude data using Fourier transform in meter per square root hertz versus the frequency in hertz (see right corner inset of Fig. 8). The spring constant of the cantilever was determined experimentally for each tip used in our studies using the IGOR software interface supplied by Asylum Research.

Finally, the exact $K_c$ values after the calibration procedure are $K_c = 0.08745$ N/m for live cells and $K_c = 1.29$ N/m for fixed cells.

**V. mESC Characterization**

**A. Cell Culture Preparation**

It is recommended to use specially coated dishes or slides to facilitate adhesion. Such petri dishes or slides are available commercially. Cells should be attached to some rigid substrate, usually either a slide or the bottom of a petri dish. In our study, the mESC R1 (SCRC-1011, American Type Culture Collection (ATCC), Manassas, VA) were grown on 0.1% gelatin-coated plates in the absence of feeder cells. The ES medium consisted of 1000 U/mL leukemia inhibitory factor (LIF, ESGRO, Chemicon, Temecula, CA), 15% fetal bovine serum (FBS) (Invitrogen), 2 mM L-glutamine, 1x nonessential amino acids, and 0.1 mM mercaptoethanol. Differentiation was induced by removal of LIF from the medium. Prior to the experiments, cells were dispersed using trypsin to obtain single cells and were plated on 60 mm tissue culture petri dish. Fixed mESC were obtained by treating live cells with 4% formaldehyde for 10 min, washed and stored in phosphate buffered saline.

**B. Geometry of the Cell and Cantilever Interaction With the Surface**

The cells were determined to be 10 and 15 μm in diameter using the cantilever tip to measure the interaction force between the cell and the tip. The geometric modeling of the cell requires the determination of the height, adhesion surface, and cell’s contour [see Fig. 9(a)]. The mESC height was calculated from the force-displacement curve obtained on hard and mESC surfaces. First, the initial height $h_0$ is measured by bringing...
the AFM tip in contact with the hard surface near mESC [see Fig. 9(b)], and then, the second height \( h_1 \) by bringing the AFM tip on mESC surface [see Fig. 9(c)]. The mESC height is the difference \( h = h_0 - h_1 \). Second, to determine the mESC shape, AFM imaging experiments are carried out in order to determine the mESC’s contour and diameter using image processing techniques [19]. From Fig. 10, we selected mESC shapes with circular adhesion surfaces. As illustration, the left corner inset shows a typical circular mESC, where \( D_1 \) and \( D_2 \) stands for the horizontal cells diameters (\( D_1 \approx D_2 \) close to 12 \( \mu \)m).

**C. Mechanical Characterization of mESC Using Spherical Tip**

To estimate the mechanical properties of biological cells using the AFM, various analytical models can be used to identify the Young’s modulus of mESC in live as well as fixed cells.

1) **Hertz Contact:** The Hertz contact model has been used extensively by the AFM community to quantify the mechanical property of biological samples using AFM [8], [9]. The Hertz contact model describes the simple case of elastic deformation of two perfectly homogeneous smooth surfaces touching under load (see Fig. 11). The forces measured are dominated by the elastic properties. In our work, the geometry of the tip used for AFM imaging experiments are carried out in order to determine the mESC’s contour and diameter using image processing techniques [19]. From Fig. 10, we selected mESC shapes with circular adhesion surfaces. As illustration, the left corner inset shows a typical circular mESC, where \( D_1 \) and \( D_2 \) stands for the horizontal cells diameters (\( D_1 \approx D_2 \) close to 12 \( \mu \)m).

2) the normal contact of the two bodies is *adhesionless and frictionless*;
3) the contact geometry is assumed to be *axisymmetric, smooth, and continuous*.

The relationship between the indentation \( \delta \), and the loading force \( F \) is given by

\[
F = \frac{4}{3} E^* R^2 \delta^2
\]

where \( E^* \) and \( R \) are the combined modulus and the relative curvature of the tip and the cell, respectively,

\[
\frac{1}{E^*} = \frac{1 - \nu_1^2}{E_1} + \frac{1 - \nu_2^2}{E_2}
\]

\[
\frac{1}{R} = \frac{1}{R_1} + \frac{1}{R_2}
\]

where \( (E_1, \nu_1, R_1) \) and \( (E_2, \nu_2, R_2) \) represents the elastic modulus, Poisson’s ratio, and the radius of the cell and spherical indenter, respectively. The elastic modulus of the silicon nitride cantilever and silicon cantilever were 222.22 and 168.17 GPa used for the live and fixed cells, respectively. The elastic modulus of the cells is in range of kilopascals [11]. Hence, our assumption that the tip used for probing is infinitely stiff (\( E_2 \gg E_1 \)) compared to the cell and Hertz contact model is valid. Thus, (13) can be rewritten as follows:

\[
F = \frac{4E_1}{3(1 - \nu_1^2)} R^2 \delta^2
\]

where \( E_1 \) and \( \nu_1 \) represent the elastic modulus and Poisson’s ratio of the cell assuming \( \nu_1 = 0.5 \).
The force ($F$) and relative deformation ($\varepsilon$) relationship is given by

$$F = \frac{2\pi h R_0 [1 + R_0/(2R_s)]^2 E}{(1 - \nu)(1 + R_0/R_s)} \varepsilon^3 + \frac{\pi E h^2}{2 \sqrt{2}} \varepsilon$$

(17)

where $E$, $\nu$, and $h$ represent the elastic modulus, Poisson’s ratio, and the thickness of the capsule membrane, respectively. $R_0$ and $R_s$ are the radius of the capsule and the spherical indenter, respectively. $\varepsilon$ is the relative deformation of the capsule related to the indentation by $\varepsilon = \delta/2R_0$. The first and the second term in (17) represent stretching and bending of the cell membrane, respectively, however, in our experiments, the bending deformation term can be neglected for $\varepsilon > 0.15$ [22].

D. Comparison Between Experimental Data and Analytical Models

We evaluated and compared the different analytical models to determine whether they appropriately predict the force-indentation relationship of mESC. Figs. 13 and 16 present the force (loading and unloading) versus time of fixed and live undiff mESC. From all these experiments, we observed adhesion force in few of the cases. For live mESC probed by a spherical tip, the adhesion force existed only for one sample. However, their amplitudes were negligible (less than 0.05 times the maximum indentation) compared to the maximum indentation forces (more than 0.1 times the maximum indentation) observed by other researchers [28], [29]. In our experiments, we observed that the adhesion force exists for only one of the live undiff and does not exist for the fixed cells. This could be due to the wear of the spherical tip [26]. Since we did not observe any adhesion forces, we have not used the JKR or DMT models.

Figs. 14 and 15 show the experimental force versus cell indentation curves for live and fixed undiff mESC and the simulated ones using the capsule and the Hertz models, respectively, for small and rather larger indentation range. Based on the experimental investigation, a least-square fit on the dataset (for the whole cells) and the corresponding $r^2$ values are calculated. These results show that the capsule model is inappropriate for describing mESC indentation. This conclusion is confirmed by the calculation of $r^2$ values performed on each mESC (see Table I). Finally, from our experiments, we infer that Hertz model appropriately describes the mechanical behavior of the live and fixed undiff mESC (see Figs. 16–18). Equation (16) was used to determine the global elastic modulus of the mESC. The average elastic modulus were 17.87 and 0.217 kPa for the fixed and live undiff, respectively, (see Fig. 19). The standard deviations for the fixed and live undiff cells were 3.37 and 0.05 kPa.
It should be noticed that alterations in the mechanical properties of cells directly reflect changes in their cellular composition, internal structure (cytoskeleton), and external interactions (cell–cell and/or cell–surface) that occur during differentiation, aging, and other changes in physiological status that are not accounted in the cell indentation model.

VI. FEM SIMULATION OF CELL INDENTATION

Fig. 20 presents a photo of the visual and haptic user interface setup. To test the accuracy and reliability of the proposed user interface system with haptics-enabled simulation, a set of experiments is designed using the setup given in Section II. The geometrical dimensions of the cell are determined through AFM image processing and force curve deflection (diameter and height), and the mechanical properties through AFM indentation studies due to the use of an analytical Hertz contact model.

These properties (see Table II) are passed to the finite-element simulations.

In order to demonstrate the validity of our finite-element indentation model in real time, we compare the simulation indentation using the physics-based FEM model to the incremental data provided by AFM indentation tool and to the response simulated with a commercial FE software. Figs. 21 and 22 present the nonlinear relationship between the penetration distance (indentation) and the reaction force for live and fixed undiff cells. The finite-element simulations using the commercial software show good agreement with both Hertz contact model and the experimental indentation data. The developed real-time physics-based FEM simulator demonstrates its accuracy in predicting...
indentation forces for large deformations (>25%) of relative nominal strain (see Figs. 21 and 22). Fig. 23 illustrates a real-time cell indentation procedure in a VE for training. It shows a screenshot of the cell surface deformation. The interaction between tip and surface can be visualized to the user in 3-D (sphere and cone represent AFM tip and forces acting on tip, respectively). It shows the indentation (elastic deformation) on the cell surface and repulsive contact forces while pressing on the surface. In order to improve the force feedback rendering, we added a visual metaphor on the computer screen.

The rendered reflection force is depicted by a cone representing the force vector acting on the AFM tip. The visual metaphor has efficiently proved its effectiveness in AFM-based nanomanipulation simulation [23]. The vector amplitude reflects the force feedback amplitude and the vector orientation represents the vertical/horizontal force components. Fig. 24 shows in 3-D real-time the probe tip, zoomed cell indentation, and micro-forces. The indentation force rendered to the operator is composed of a vertical $F_z$ and horizontal $F_x$ component. For the specific case corresponding to the center indentation, the tangential component is very small $F_x \approx 0$. Only vertical reflection force is provided to the user. He cannot haptically feel it during simulated indentation tests (see Figs. 25(a) and 26). In the case, where the AFM probe is pushed away from the center of the cell, a tangential component exists ($F_x \neq 0$) and the reflection force rendered to the user contains a vertical and tangential component (see Figs. 25(b) and 27). We did not reproduce this specific case in our results since the main interest of the paper is to prove experimentally the real-time explicit simulator results during centered cell indentations (since only vertical components of the force experiments are measured through the AFM cantilever).
We have developed a computer-based training system to simulate real-time cell indentation procedures in VEs for training biologists. The simulator provides the user with visual and haptic feedback. The simulation of this procedure involves real-time rendering of computer generated graphical images of the AFM tip, physics-based modeling of soft biological tissue, and display of touch and force sensations to the trainee through the simulation of haptic interactions. The real-time haptics-enabled simulator is realistic, since it is mainly based on experimental AFM data (mechanical and geometrical property of mESC) to accurately replicate the indentation task and predict the cell deformation during indentation. We first investigated the challenging issues in the real-time modeling of the biomechanical properties of the cell indentation through FEMs. Compared to experimental AFM indentation results performed on mESC, we can see clearly that the proposed physically based FEM model is able to simulate the cell deformation through real-time simulation constraints. Currently, we are working on integrating nanoscale effects, such as friction, viscosity, puncture, and adhesion forces in the nonlinear FEM model. All modalities will be merged in an ergonomic tool and intelligent biological simulator for stem cell characterization.

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Q9. Author: Please provide the title of Engineer and Master’s degrees for H. Ladjal.
Q10. Author: Please provide the title of Undergraduate degree for A. Pillarisetti.
Q11. Author: Please provide the title (B.Sc., M.Sc., Ph.D., etc.) and university names from which C. Keefer received these degrees.
Abstract—Training simulators that provide realistic visual and haptic feedback during cell indentation tasks are currently investigated. Complex cell geometry inherent to biological cells and intricate mechanical properties drive the need for precise mechanical and numerical modeling to assure accurate cell deformation and force calculations. Advances in alternative finite-element formulation, such as the mass—tensor approach, have reached a state, where they are applicable to model soft-cell deformation in real time. The geometrical characteristics and the mechanical properties of different cells are determined with atomic force microscopy (AFM) indentation. A real-time, haptics-enabled simulator for cell centered indentation has been developed, which utilizes the AFM data (mechanical and geometrical properties of embryonic stem cells) to accurately replicate the indentation task and predict the cell deformation during indentation in real time. This tool can be used as a mechanical marker to characterize the biological state of the cell. The operator is able to feel the change in the stiffness during cell deformation between fixed and live cells in real time. A comparative study with finite-element simulations using a commercial software and the experimental data demonstrate the effectiveness of the proposed physically based model.

Index Terms—Atomic force microscope (AFM), finite-element modeling (FEM), haptics, modeling, real-time interaction, stem cell.

I. INTRODUCTION

Mechanical manipulation and characterization of biological cells is currently one of the most promising research areas in the field of biorobotics applied to cellular level interactions. The mechanical properties, such as elasticity, membrane tension, cell shape, and adhesion strength, may play an important role in cell development and differentiation [1]–[3]. In contrast to adult stem cells, embryonic stem cells differentiate into every cell type in the body, and thus, have a huge therapeutic potential. Stem-cell-based systems offer a very promising and innovative solution for obtaining large numbers of cells for early efficacy and toxicity screening. Stem cell technology provides a new tool for drug development and a better insight to understand mammalian gene function (see Fig. 1). The mechanical properties of biological cells have been studied with different techniques [4], the most popular being optical tweezers [5], magnetic beads [6], and micropipette aspiration [7]. Different microrobotic systems with various end effectors have been proposed [33], [34]. However, these methods cannot compete with the precision that can be obtained with the atomic force microscopy (AFM) method [8], [32]. AFMs have been widely used in the study of micro- and nanostructures including living cells. Modern AFM techniques allow solving several problems of cell biomechanics by simultaneous evaluation of the local mechanical properties and the topography of living cells, at a high spatial resolution and force sensitivity [9]. In these experiments, an AFM cantilever serves as a microindenter to probe the cell, and further analysis of force-indentation data yields the local Young’s modulus. In addition, the AFM indentation technique can be used to characterize the viscoelastic behavior of the cell cytoskeleton [10], including viscosity [11], loss and storage moduli [12], and stress relaxation times. Mechanical phenotyping prove actually to be a valuable tool in the development of improved methods of targeted cellular differentiation of embryonic, induced pluripotent (iPS), and adult stem cells for therapeutic and diagnostic purposes [30], [31]. In this study, we propose to characterize the mechanical behavior of an individual mouse embryonic stem cells (mESC) in undifferentiated state.
using AFM indentation [13]. As we hypothesize that the cell stiffness of mESC may reflect the nuclear physical changes of the cell, it should be possible to detect modifications in stiffness by probing the cell directly. However, different limitations are inherent in AFM-based indentation techniques. First, deformations induced by the AFM probe and cell positions cannot be seen by the user. Second, the force sensing of the AFM cannot resolve the force interaction between the AFM tip and the cell’s surface independently in all three dimensions to avoid cytoskeleton-disrupted experiments. Therefore, the user performs extremely challenging manipulations requiring advanced manual skills. Modeling the physics of the interaction between tip and cell is necessary to realize complicated indentation tasks like cancer cell probing [14].

We developed finite-element models (FEMs) that include the topological information of adherent stem cells (shape and dimensions), AFM tips (conical and spherical), and biological structure (cell models with cytoplasm layers, cytoskeletons, and nucleus) [15]. To validate the viscoelastic and nonlinear dynamics of the cell in real time using finite-element procedures, we report herein the mechanical characterization of single mESC indentation-relaxation for different colonies of fixed and live undifferentiated mESCs (termed, undiff mESC). The proposed indentation tests use microsphere-modified AFM probes in order to estimate the global elastic modulus of the cell to reflect the true global response of a mESC [14], [15]. This paper is a subsequent contribution to [27], where various analytical models are tested and compared to estimate the mechanical properties of the biological cells. Based on the proposed physical model, we have developed a computer-based training system using force feedback to simulate cell indentation procedures in virtual environments (VEs) for biologist training. The paper consists of five sections. In Section II, we present the biomechanical finite-element approach dedicated to real-time indentation. In Section III, we present the methodology of the VE system for cell indentation. In Section IV, we present the experimental setup using the force microscopy. In Section V, we make a comparative study between the physics-based model and experimental data. In Section VI, we present our real-time cell indentation simulator using a spherical tip. Finally, we give some concluding remarks and the directions for future work.

II. PHYSICS-BASED DEFORMABLE MODELS FOR CELL NANOINDENTATION SIMULATION

In this section, we present further developments to the linear elastic mass–tensor model introduced by [16], [17] and extended by Schwartz et al. [18]. In our paper, the cell is modeled by a volume object discretized into a conformal tetrahedral mesh as defined by finite-element theory. Using classical notations, inside each tetrahedron \( T^k \), the displacement field, defined by a linear interpolation \([N^k]\) of the nodal displacement vector \( \{u^k\} \) of the four vertices of tetrahedron, is written as follows:

\[
\{U(x)^k\} = [N^k(x)] \{u^k\},
\]

(1)

For a linear elastic material, the relation between the Cauchy stress tensor and the linearized strain tensor is written with Lame’s coefficient in condensed vector notation as follows:

\[
\{\sigma\} = \lambda \left( \{\varepsilon\}_1 + \{\varepsilon\}_2 + \{\varepsilon\}_3 \right) \{I\} + 2\mu \{\varepsilon\}
\]

(2)

where \( \{I\} \) is the identity matrix.

The principle of virtual work applied to a single tetrahedron \( T^k \) leads to the elementary stiffness matrix \([K^k]\) such that the elementary nodal force vector acting on a tetrahedron is as follows:

\[
\{f^k\} = [K^k] \{u^k\}.
\]

(3)

This stiffness matrix is composed of a plurality of elementary submatrices, each connecting the elementary force acting on the node \( i \) to the displacement of the node \( j \)

\[
[K^k_i] = \frac{1}{36V^k} \left( \lambda \{m_j\} \{m_j\}^T + \mu \{m_j\} \{m_j\}^T + \mu \{m_i\} \{m_i\}^T \{m_j\} \{I\} \right)
\]

(4)

where \( \{m_i\} \) are unit outward-pointing normals to triangular faces and \( V^k \) is the volume of the tetrahedron \( T^k \).

Taking into account the contribution of all adjacent tetrahedra, the global internal force acting on a node \( i \) can be expressed as follows:

\[
\{F^i_{int}\} = \sum_{k \in V_i} \left( \sum_{j=1}^4 [K^k_i] \{u_j\} \right)
\]

(5)

where \( V_i \) is the neighborhood of vertex \( i \) (i.e., the tetrahedra containing node \( i \)).

The tensors \([K^k_{ij}]\) depending on the remaining geometry and Lamé’s coefficients are constant. They can be precomputed in an offline phase. This is the essential advantage of the mass–tensor approach, which makes it useful for real-time application.

A. Cell Modeling Using FEM

In our paper, the mESC is meshed with 3-D first-order tetrahedral elements, as shown in Fig. 2. The geometrical characteristics of the cells were obtained from experimental data using the AFM system. This includes shape, diameter, and height of the cells. The mesh that has been used for the following simulations is composed of 835 vertices and 2456 tetrahedra, where all vertices are free to move except the ones in contact with the petri dish, which are fixed.

B. Dynamic Model

The equation of motion of a vertex \( l \) of the cell mesh can be written as follows:

\[
M^l \{\ddot{u}_l\} + \gamma^l \{\dot{u}_l\} + \sum_{k \in V_l} \left( \sum_{j=1}^4 [K^k_{lj}] \{u_j\} \right) = \{F^l_{ext}\}
\]

(6)

where \( M^l \) and \( \gamma^l \) are, respectively, the mass and damping coefficients of each vertex.
To solve the dynamic system, we tested different integration schemes (implicit and explicit) taking into account the tradeoff between real-time simulation and haptic stability requirements. We choose the explicit centered finite-difference scheme defined by

$$\begin{align*}
\mathbf{u}(t + h) &= \mathbf{u}(t) + \frac{h}{2} \mathbf{\dot{u}}(t) h + \frac{h^2}{2} \mathbf{\ddot{u}}(t) h + O(h^3) \\
\mathbf{\dot{u}}(t + h) &= \mathbf{\dot{u}}(t) + \frac{h}{2} \mathbf{\ddot{u}}(t) h + \frac{h^2}{2} \mathbf{\dddot{u}}(t) h + O(h^3)
\end{align*}$$

(7)

where $h$ is the temporal integration step size chosen to satisfy the Courant–Friedrich–Lewy condition.

The discrete equation of motion of a vertex $l$ reads

$$\begin{align*}
\left( M^l + \frac{h}{2} \mathbf{\gamma}^l \right) &\left\{ \mathbf{\ddot{u}}_l \right\}_{(t+h)} \\
= &\left\{ \mathbf{F}^l_{\text{ext}} \right\}_{(t+h)} - \gamma^l \left( \left\{ \mathbf{u}_l \right\}_{(t)} + \frac{h}{2} \left\{ \mathbf{\dot{u}}_l \right\}_{(t)} \right) \\
- &\sum_{k \in V} \left( \frac{4}{\mathbf{K}^l_{jk}} \right) \left( \left\{ \mathbf{u}_j \right\}_{(t)} + h \left\{ \mathbf{\dot{u}}_j \right\}_{(t)} + \frac{h^2}{2} \left\{ \mathbf{\ddot{u}}_j \right\}_{(t)} \right)
\end{align*}$$

(8)

from which we can compute the new vertices for accelerations, velocities, and positions.

### III. VE SYSTEM FOR CELL INDENTATION

Fig. 3 shows the architecture of the real-time virtual-reality-based cell indentation simulator. This includes the computer generated models of stem cells, the AFM nanoindentor, the collision detection algorithm, physics-based models of deformable mESC, and the haptic interaction controller. The operator will be able to interact with the 3-D model of the cell using his sense of vision as well as actively manipulate using his sense of touch. Usually, real-time graphics translates to an update rate of 25 Hz, stable haptic interaction in VE’s requires much higher update rate of around 1 kHz. As soft biological tissues exhibit complex viscoelastic and nonlinear properties, the real-time interaction imposes severe restrictions on the FEM model.

A significant difficulty of using the finite-element technique for real-time simulation is that it is computationally costly. These mesh-based schemes also require an expensive numerical integration operation for the computation of the system stiffness matrices. For this reason, we adopted a computational architecture with two simulation stages: 1) an offline computation of stiffness matrices for each triangular element and 2) visual and haptic interaction for real-time indentation simulation.

#### A. Offline Computation

The proposed computational methodology contains an offline precalculation step. The most costly and time-consuming operations are realized during this step. The database contains various developed models of the different stem cells geometry (bone marrow, pancreatic, heart muscle, nerve cells, etc.). Using AFM image analysis (ImagePro software from AFM Asylum), we derived different geometric stem cell models based on a commercial computer-aided design (CAD) package (MARC-ADAMS). The meshing of the 3-D internal structure of the mESC is then carried out through a dedicated 3-D meshing software (GID software) modeled in exact dimensions. Although they were displayed as 3-D texture-mapped objects to the user, they were modeled as connected line segments to reduce the number of collision computations during real-time interactions.

The adherent cell was approximated as an assembly of discrete tetrahedral elements interconnected to each other through a fixed number of nodes (see Fig. 2). The displacements of these nodal points for applied external indentation forces were the basic unknowns of our FEM analysis. The coordinates...
of vertices, the tetrahedron indexing, and the connectivity of vertices were derived from the geometric model. Then, the mechanical (Young’s modulus and stiffness) and geometrical (height and shape) properties of the cell structure are determined through AFM indentation experiments. Finally, these properties are passed to the FEM algorithm for calculating the physically based behavior of biological soft tissues. The offline calculation of all tensors $K_{ij}$ allows precalculation stiffness values for preselected cells such as, for example, live and fixed cells.

B. Real-Time Indentation

Since haptics and graphics have different update frequencies, we implemented separate threads to update the loops. During the indentation task, the contact between the tip and the cell must occur at a special set of points (nodal points).

1) Collision Detection: An important part of biological cantilever tip-soft cell modeling is the fast collision detection algorithm. The collision detection is the first step to carry out realistic interactions. However, locating the contact primitive (e.g., facets) between two objects may be computationally expensive, especially if the objects are composed of a large number of polygons. In general, to determine a contact between two virtual objects, we compute the distance between them. If this distance is negative, then the objects are in contact. We may continue by determining the set of geometric objects or entities that collide. The cantilever tip in VE can be modeled as a complex 3-D object, composed of numerous surfaces, edges, and vertices, but this constitutes a challenging task. For the purpose of detecting collisions, we used a point-based representation of the cantilever tip as a simple spherical object and employed a simple sphere-points intersection algorithm with local search technique. It calculates the distance between the center of the spherical tip and the nodes of the cell surface (see Fig. 4).

2) Penalty-Based Models: In an interaction model, a collision between two objects may generate local and global deformations of the colliding objects, new internal forces, changes in the velocities of the involved objects, and topology modifications, such as punction or cutting. Physical simulations suppose that these events are related to the mechanical properties of the objects. In our simulations, we used the penalty-based models. These models consider the reactive force that first reduces the relative speed of the two objects, and finally, repulses them from each other, as proportional to the deformation. Thus, the reactive force is a function of the local interpenetration depth $\delta$. The normal repulsive force $F_{\text{rep}}$ can be written as follows:

$$F_{\text{rep}} = \lambda \delta \{n\}$$

where $\lambda$ is the penalty of the interpenetration and $\{n\}$ is the outward-pointing normal.

These contact forces are evaluated at each time step $t$ (see Fig. 5). They represent the new updated external forces acting upon our stem cell.

3) Virtual Coupling for Stability of the Haptic Rendering:

The internal operating loop of the haptic interface requires an update frequency around 1 kHz. However, the FEM update frequency is in the range of 25–30 Hz. This frequency difference threatens the coherence between both systems leading to instabilities of the user haptic rendering. We adopted the solution to use a virtual coupling model defined in [24].

This approach introduces a virtual passive link between the simulation model and the haptic interface in order to ensure the stability and the performance of the system [see Fig. 6(a)]. When we combine the impedance display implementation with an appropriate virtual coupling network, we get the admittance...
For the virtual coupling, the impedance display induces a limit where
the values is plotted on Fig. 6(b) as a bold blue line. The left side of (12) with the resulting
values is plotted on Fig. 6(b) as a bold blue line.

matrix for the combined interface. The linear two-port is said
to be absolutely stable if there exists no set of passive ter-
minating one-port impedances for which the system is unsta-
ble. Llewellyn’s stability criteria [25] provides both necessary
and sufficient conditions for absolute stability of linear two-
ports. We get the conditions for absolute stability of the haptic
interface

\[ \text{Re}(Z_{di}(z)) \geq 0, \quad \frac{1}{Z_{cvi}(z)} \geq 0 \]  

(10)

\[ \cos(\angle \text{ZOH}(z)) + \frac{2\text{Re}(Z_{di}(z))\text{Re}(1/Z_{cvi}(z))}{|\text{ZOH}(z)|} \geq 1 \]  

(11)

where \( Z_{cvi}(z) \) is the virtual coupling impedance \( (k_c, b_c) \),\n\( \text{ZOH}(z) \) is a zero-order holder, and \( Z_{di}(z) \) is the PHANToM
impedance. The inequality (11) can be rewritten to get an ex-
licit expression of absolute stability of the haptic interface

\[ \text{Re} \left( \frac{1}{Z_{cvi}(z)} \right) \geq 1 - \cos(\angle \text{ZOH}(z)) \frac{2\text{Re}(Z_{di}(z))}{|\text{ZOH}(z)|} \]. \]  

(12)

For the virtual coupling, the impedance display induces a limit
on the maximum impedance which can be rendered. We use (12)
to find the virtual coupling, which makes the haptic interface
absolutely stable. Using the design bound of the best performing,
absolutely stabilizing virtual coupling. These parameters are
found as \( b_c = 0.008 \text{ N/(mm/s)} \) and \( k_c = 2.3 \text{ N/mm} \) (plotted on
Fig. 6(a) as a thin red line). The left side of (12) with the resulting
values is plotted on Fig. 6(b) as a bold blue line.

Fig. 6. Stability conditions using the virtual coupling model. (a) Model of
virtual coupling. (b) Impedance display virtual coupling.

Fig. 7. The top view module enables viewing of cells and easy
alignment of the laser beam on the AFM cantilever. The XY
stage (manual) allows the user to position the cell beneath the
cantilever tip of the AFM. The entire AFM setup is enclosed in
an acoustic isolation chamber to prevent acoustic noise from inter-
fering with the AFM measurements. The \( x \)- and \( y \)-axis ranges
of the scan head are 90 \( \mu \text{m} \). The \( z \)-axis scan range is 40 \( \mu \text{m} \).
The AFM system is used to obtain force and cell deformation data from biological samples.

The cantilever is moved by the piezoelectric scanner in \( z \)-direction toward the cell. The deflection of the cantilever is
detected by a photodiode when the tip comes in contact with the
cell. When the tip of the cantilever is in contact with the cell, the
initial cantilever deflection \( (d_0) \), and initial cantilever movement
in \( (z) \) direction \( (z_0) \) are stored. As the cantilever moves further
in \( z \)-direction and deforms the cell, the final cantilever deflection
\( (d_1) \) and the cantilever movement \( (z_1) \) are obtained [21].

The stiffness depends not only on the Young’s modulus but also on the geometry of the tip-surface contact. Therefore, the
geometry and spring constant of the cantilever are calibrated in
the same way for live and fixed cells. In the following experi-
ments, we employed two types of cantilevers attached with a
spherical probe (5 \( \mu \text{m} \) in diameter).

1) Silicon nitride cantilever with a spring constant of 0.06
\( \text{N/m} \) (Novascan Technologies, Inc., Ames, IA) for live
cells.

2) Silicon cantilever with a spring constant of 1.75 \( \text{N/m} \) (No-
vascan Technologies, Inc.) for fixed cells.
B. AFM Calibration

Before any experiment is done, it is necessary to perform a calibration of the cantilever to determine its resonant frequency. The resonant frequency measurement enables the user to determine if it lies within the specified range provided by the manufacturer in order to detect probe defects. The determination of spring constant $K_n$ by the thermal method involves both measurements: 1) sensitivity in nanometer per volt and 2) first-resonant frequency of the cantilever.

1) Determine the contact slope of the indentation curve on a hard surface to determine the sensitivity of the cantilever (see left corner inset of Fig. 8).

2) Measure the thermal power spectral density (PSD) to determine the resonant frequency of the cantilever and to confirm that the cantilever and light source are aligned. This confirms that the system is functioning properly. The results are plotted as the deflection amplitude data using Fourier transform in meter per square root hertz versus the frequency in hertz (see right corner inset of Fig. 8). The spring constant of the cantilever was determined experimentally for each tip used in our studies using the IGOR software interface supplied by Asylum Research.

Finally, the exact $K_c$ values after the calibration procedure are $K_c = 0.08745 \text{ N/m}$ for live cells and $K_c = 1.29 \text{ N/m}$ for fixed cells.

V. mESC Characterization

A. Cell Culture Preparation

It is recommended to use specially coated dishes or slides to facilitate adhesion. Such petri dishes or slides are available commercially. Cells should be attached to some rigid substrate, usually either a slide or the bottom of a petri dish. In our study, the mESC R1 (SCRC-1011, American Type Culture Collection (ATCC), Manassas, VA) were grown on 0.1% gelatin-coated plates in the absence of feeder cells. The ES medium consisted of 1000 U/mL leukemia inhibitory factor (LIF, ESGRO, Chemicon, Temecula, CA), 15% fetal bovine serum (FBS) (Invitrogen), and basic medium that included Knockout Dulbecco’s modified Eagle’s medium (Invitrogen), 2 mM L-glutamine, 1x nonessential amino acids, and 0.1 mM mercaptoethanol. Differentiation was induced by removal of LIF from the medium. Prior to the experiments, cells were dispersed using trypsin to obtain single cells and were plated on 60 mm tissue culture petri dish. Fixed mESC were obtained by treating live cells with 4% formaldehyde for 10 min, washed and stored in phosphate buffered saline.

B. Geometry of the Cell and Cantilever Interaction With the Surface

The cells were determined to be 10 and 15 μm in diameter using the cantilever tip to measure the interaction force between the cell and the tip. The geometric modeling of the cell requires the determination of the height, adhesion surface, and cell’s contour [see Fig. 9(a)]. The mESC height was calculated from the force–displacement curve obtained on hard and mESC surfaces. First, the initial height $h_0$ is measured by bringing...
Fig. 10. How to determine the diameter of the stem cell through image analysis (ImagePro software from Asylum).

Fig. 11. Hertzian contact with a spherical tip: $F$ is the loading force, and $\delta$ and $a$ are the indentation and contact radius, respectively.

the AFM tip in contact with the hard surface near mESC [see Fig. 9(b)], and then, the second height $h_1$ by bringing the AFM tip on mESC surface [see Fig. 9(c)]. The mESC height is the difference $h = h_0 - h_1$. Second, to determine the mESC shape, AFM imaging experiments are carried out in order to determine the mESC’s contour and diameter using image processing techniques [19]. From Fig. 10, we selected mESC shapes with circular adhesion surfaces. As illustration, the left corner inset shows a typical circular mESC, where $D_1$ and $D_2$ stands for the horizontal cells diameters ($D_1 \approx D_2$ close to 12 $\mu$m).

C. Mechanical Characterization of mESC Using Spherical Tip

To estimate the mechanical properties of biological cells using the AFM, various analytical models can be used to identify the Young’s modulus of mESC in live as well as fixed cells.

1) Hertz Contact: The Hertz contact model has been used extensively by the AFM community to quantify the mechanical property of biological samples using AFM [8], [9]. The Hertz contact model describes the simple case of elastic deformation of two perfectly homogeneous smooth surfaces touching under load (see Fig. 11). The forces measured are dominated by the elastic properties. In our work, the geometry of the tip used for all experiments is spherical. The mechanical interaction between the spherical tip and mESC can be described by the Hertz contact model of two elastic bodies [20]. The model assumes that:

1) the material properties of the tip and the cell are isotropic and homogeneous;

2) the normal contact of the two bodies is adhesionless and frictionless;

3) the contact geometry is assumed to be axisymmetric, smooth, and continuous.

The relationship between the indentation $\delta$, and the loading force $F$ is given by

$$F = \frac{4}{3} E^* R^2 \delta^2$$

where $E^*$ and $R$ are the combined modulus and the relative curvature of the tip and the cell, respectively,

$$\frac{1}{E^*} = \frac{1}{E_1} + \frac{1}{E_2}$$

$$\frac{1}{R} = \frac{1}{R_1} + \frac{1}{R_2}$$

where $(E_1, \nu_1, R_1)$ and $(E_2, \nu_2, R_2)$ represents the elastic modulus, Poisson’s ratio, and the radius of the cell and spherical indenter, respectively. The elastic modulus of the silicon nitride cantilever and silicon cantilever were 222.22 and 168.17 GPa used for the live and fixed cells, respectively. The elastic modulus of the cells is in range of kilopascals [11]. Hence, our assumption that the tip used for probing is infinitely stiff $(E_2 \gg E_1)$ compared to the cell and Hertz contact model is valid. Thus, (13) can be rewritten as follows:

$$F = \frac{4E_1}{3(1 - \nu_1^2)} R^2 \delta^2$$

where $E_1$ and $\nu_1$ represent the elastic modulus and Poisson’s ratio of the cell assuming $\nu_1 = 0.5$.

2) Capsule Contact: The second model considers the biological cell to be composed of a cell membrane and cytoplasm (see Fig. 12). It is the capsule model. The model assumes that the cell membrane is a thin film and that the inner cytoplasm provides a uniform hydrostatic pressure on the membrane [22]. The model assumes the following:

1) The cell membrane is linearly elastic.

2) The deformation of the cell membrane is caused by stretching and bending of the cell membrane.

3) The cell is free of initial membrane stress or residual stress.

4) The cell volume is constant.
The force \(F\) and relative deformation \(\varepsilon\) relationship is given by [22]

\[
F = 2\pi h R_0 \left[1 + \frac{R_0}{(2R_s)}\right]^2 \varepsilon^3 + \frac{\pi}{2\sqrt{2}} Eh^2 \sqrt{\varepsilon}
\]

where \(E\), \(\nu\), and \(h\) represent the elastic modulus, Poisson’s ratio, and the thickness of the capsule membrane, respectively. \(R_0\) and \(R_s\) are the radius of the capsule and the spherical indenter, respectively. \(\varepsilon\) is the relative deformation of the capsule related to the indentation by \(\varepsilon = \delta/2R_0\). The first and the second term in (17) represent stretching and bending of the cell membrane, respectively. However, in our experiments, the bending deformation term can be neglected for \(\varepsilon > 0.15\) [22].

D. Comparison Between Experimental Data and Analytical Models

We evaluated and compared the different analytical models to determine whether they appropriately predict the force-indentation relationship of mESC. Figs. 13 and 16 present the force (loading and unloading) versus time of fixed and live undiff mESC. From all these experiments, we observed adhesion force in few of the cases. For live mESC probed by a spherical tip, the adhesion force existed only for one sample. However, their amplitudes were negligible (less than 0.05 times the maximum indentation) compared to the maximum indentation forces (more than 0.1 times the maximum indentation) observed by other researchers [28], [29]. In our experiments, we observed that the adhesion force exists for only one of the live undiff and does not exist for the fixed cells. This could be due to the wear of the spherical tip [26]. Since we did not observe any adhesion forces, we have not used the JKR or DMT models.

Figs. 14 and 15 show the experimental force versus cell indentation curves for live and fixed undiff mESC and the simulated ones using the capsule and the Hertz models, respectively, for small and rather larger indentation range. Based on the experimental investigation, a least-square fit on the dataset (for the whole cells) and the corresponding \(r^2\) values are calculated. These results show that the capsule model is inappropriate for describing mESC indentation. This conclusion is confirmed by the calculation of \(r^2\) values performed on each mESC (see Table I). Finally, from our experiments, we infer that Hertz model appropriately describes the mechanical behavior of the live and fixed undiff mESC (see Figs. 16–18). Equation (16) was used to determine the global elastic modulus of the mESC. The average elastic modulus were 17.87 and 0.217 kPa for the fixed and live undiff, respectively, (see Fig. 19). The standard deviations for the fixed and live undiff cells were 3.37 and 0.05 kPa.
It should be noticed that alterations in the mechanical properties of cells directly reflect changes in their cellular composition, internal structure (cytoskeleton), and external interactions (cell–cell and/or cell–surface) that occur during differentiation, aging, and other changes in physiological status that are not accounted in the cell indentation model.

VI. FEM SIMULATION OF CELL INDENTATION

Fig. 20 presents a photo of the visual and haptic user interface setup. To test the accuracy and reliability of the proposed user interface system with haptics-enabled simulation, a set of experiments is designed using the setup given in Section II.

The geometrical dimensions of the cell are determined through AFM image processing and force curve deflection (diameter and height), and the mechanical properties through AFM indentation studies due to the use of an analytical Hertz contact model.

These properties (see Table II) are passed to the finite-element simulations.

In order to demonstrate the validity of our finite-element indentation model in real time, we compare the simulation indentation using the physics-based FEM model to the incremental data provided by AFM indentation tool and to the response simulated with a commercial FE software. Figs. 21 and 22 present the nonlinear relationship between the penetration distance (indentation) and the reaction force for live and fixed undiff cells. The finite-element simulations using the commercial software show good agreement with both Hertz contact model and the experimental indentation data. The developed real-time physics-based FEM simulator demonstrates its accuracy in predicting...
indentation forces for large deformations (>25%) of relative nominal strain (see Figs. 21 and 22). Fig. 23 illustrates a real-time cell indentation procedure in a VE for training. It shows a screenshot of the cell surface deformation. The interaction between tip and surface can be visualized to the user in 3-D (sphere and cone represent AFM tip and forces acting on tip, respectively). It shows the indentation (elastic deformation) on the cell surface and repulsive contact forces while pressing on the surface. In order to improve the force feedback rendering, we added a visual metaphor on the computer screen. The rendered reflection force is depicted by a cone representing the force vector acting on the AFM tip. The visual metaphor has efficiently proved its effectiveness in AFM-based nanomanipulation simulation [23]. The vector amplitude reflects the force feedback amplitude and the vector orientation represents the vertical/horizontal force components. Fig. 24 shows in 3-D real-time the probe tip, zoomed cell indentation, and microforces. The indentation force rendered to the operator is composed of a vertical $F_z$ and horizontal $F_x$ component. For the specific case corresponding to the center indentation, the tangential component is very small $F_x \simeq 0$. Only vertical reflection force is provided to the user. He cannot haptically feel it during simulated indentation tests (see Figs. 25(a) and 26). In the case, where the AFM probe is pushed away from the center of the cell, a tangential component exists ($F_x \neq 0$) and the reflection force rendered to the user contains a vertical and tangential component (see Figs. 25(b) and 27). We did not reproduce this specific case in our results since the main interest of the paper is to prove experimentally the real-time explicit simulator results during centered cell indentations (since only vertical components of the force experiments are measured through the AFM cantilever).
We have developed a computer-based training system to simulate real-time cell indentation procedures in VEs for training biologists. The simulator provides the user with visual and haptic feedback. The simulation of this procedure involves real-time rendering of computer generated graphical images of the AFM tip, physics-based modeling of soft biological tissue, and display of touch and force sensations to the trainee through the simulation of haptic interactions. The real-time haptics-enabled simulator is realistic, since it is mainly based on experimental AFM data (mechanical and geometrical property of mESC) to accurately replicate the indentation task and predict the cell deformation during indentation. We first investigated the challenging issues in the real-time modeling of the biomechanical properties of the cell indentation through FEMs. Compared to experimental AFM indentation results performed on mESC, we can see clearly that the proposed physically based FEM model is able to simulate the cell deformation through real-time simulation constraints. Currently, we are working on integrating nanoscale effects, such as friction, viscosity, puncture, and adhesion forces in the nonlinear FEM model. All modalities will be merged in an ergonomic tool and intelligent biological simulator for stem cell characterization.

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