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Characterization of Au/Au, Au/Ru and Ru/Ru ohmic contacts in MEMS switches improved by a novel methodology

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ABSTRACT
Comparisons between several pairs of contact materials have been done with a new methodology using a commercial nanoindenter coupled with electrical measurements on test vehicles specially designed to investigate the micro-scale contact physics. Experimental measurements are obtained to characterize the response of a 5 \(\mu\)m square contact bump under electromechanical stress with increased applied current. The data provide a better understanding of micro-contact behaviour related to the impact of current at low- to medium-power levels. Contact temperature rise is observed, leading to shifts of the mechanical properties of contact materials and modifications of the contact surface. The stability of the contact resistance, when the contact force increases, is studied for contact pairs of soft (Au/Au contact), harder (Ru/Ru contact) and mixed material configuration (Au/Ru contact). An enhanced stability of the bimetallic contact Au/Ru is demonstrated considering sensitivity to power increase, related to creep effects and topological modifications of the contact surfaces. These results are compared to previous ones and discussed in a theoretical way by considering the temperature distribution around the hottest area at the contact interface.

Keywords: MEMS, Switch, Micro contact, Contact material, Contact temperature, Creep, Gold, Ruthenium

1. INTRODUCTION
From DC up to tens of giga hertz, the bandwidth of RF MEMS switch specifications is spreading more and more. But is the challenge reachable for ohmic contact switches? Even if some successful companies now propose some of these breakthrough-told technologies, the gap is still large between the low TRL (Technology Readiness Level) academic demonstrations and the commercially available switches. Several reasons can be exposed to account for this, and one of them is probably the lack of adapted tools to investigate the physics occurring at such a tiny scale. Research on contact characterization for microelectromechanical system (MEMS) switches has been driven by the necessity to reach a high-reliability level for micro-switch applications. One of the main failure observed during aging of the devices is the increase of the electrical contact resistance. The key issue is the electromechanical behaviour of the materials used at the contact interface when the current is flowing through the contact asperities. In addition, design of MEMS switches is widely diversified and few comprehensive investigations of micro-contact physics have been reported on a single type of contact. For ohmic contact switches, the main issue is to control precisely the contact pressure and the corresponding temperature. These two parameters are interdependent as soon as some DC or RF signal is flowing through the contact.

As a consequence, new characterization techniques have been recently developed in order to measure the contact resistance functions of the load applied and the contact deformation. Moreover, a precise monitoring of the current intensity and of the associated potential drop is useful to get rid of any parasitic resistance using cross rod configuration.

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Works on FEM structural-contact analysis were reported at Memswave 2009 where temperature was studied at the contact interface 0[1]. Fortini presented a nanoscale comparison of Au and Ru contact behaviour using molecular dynamics simulation [2]. Experimental characterizations using nanoindenter [3], piezo actuator [4] and atomic force microscope (AFM) [5] were published by several researchers, to investigate the behaviour of Au and Au-Ni alloys as contact materials by mechanical actuation with increasing force. To our knowledge, the micro-scale heating effects on the contact material are not fully understood. However these phenomena are key issues for the designers to understand the capability of the contact to sustain the current before reaching critical temperature values for the used material [1]. In this paper, the analysis of the impact at low to medium power level on Au, Ru and Au/Ru based RF MEMS contacts differs from previous works as the study of the power sensitivity is based on experimental results. In addition, these tests are performed using the dynamic control module of the nanoindenter (DCM), which allows a higher resolution testing. This apparatus is coupled with an environmental chamber to control both temperature and humidity. In addition, the tests are simplify by measuring the DC contact resistance instead of the S-parameters.

This work proposes a method to compare the behavior of three couples of contact materials. This paper is organized as follows. The section 2 is an overview of the micro-contact theory. In section 3, the experimental apparatus and device are described. Experimental results are presented in section 4. These results are compared to previous ones and discussed in a theoretical way by considering the temperature distribution around the hottest area at the contact interface in section 5, followed by the summary and conclusions in section 6.

2. MICRO CONTACT THEORY

The electromechanical behaviour of the contacting surfaces during switching depends on several mechanical and operating parameters such as thermal ($\sigma_t$) and electrical ($\sigma_e$) conductivity of contact materials, mechanical properties of the contact spots ($E$ and $H$), roughness ($R_n$), adhesion energy of contact surfaces ($\gamma$, mechanical stiffness ($k$), humidity ($Rh$), contact voltage ($V_C$), contact force ($F_C$) (cf. Figure 1). The key issue is the stability of the contact surfaces through aging: the electrical, topological and mechanical properties have to be maintained the longest time.

![Complexity of micro-contact failure mechanisms](image)

Figure 1 : Complexity of micro-contact failure mechanisms

2.1 Analytic models

First, the micro-contact physically differs from the macro-contact due to the roughness of the contact surface. Only high points on each surface come in contact. Thereby the effective contact area, named asperities or a-spots, is largely smaller than the apparent one.

To simplify the model, the contact resistance ($R_C$) is assumed to be mainly governed by the constriction resistance as the current flow is constricted through the small asperities leading to the electrical contact ie. the electrical contact is considered strictly ohmic. This electrical resistance is directly linked to the constriction of current lines between both contacts. This constriction of current causes a local increase of the current density and tends to increase the electrical potential drop between the two sides of the contact.
Thus the expression of this constriction resistance, so called Maxwell resistance $R_{\text{Maxwell}}$, can be written for a single circular contact spot of radius $a$ as [6]:

$$R_{\text{Maxwell}} = \frac{\rho}{2a} \tag{1}$$

Where $\rho$ is the resistivity of the contact material. During the first contact establishment between the two surfaces, the load applied is generally higher than the yield stress of the contact material. Thus the deformation of the contact asperities is considered to be predominantly plastic. The contact area and the contact load can be linked to the radius of the contact spot $a$ using Abbott and Firestone’s plastic contact model [6].

$$a = \sqrt{\frac{A_c}{\pi}} = \sqrt{\frac{F_C}{H\pi}} \tag{2}$$

Where $A_c$ is the contact area, $F_C$ the contact force and $H$ the Meyer hardness of the softer material. Secondly, it’s necessary to keep in mind that, generally, the current flows by multiple asperities due to the surface roughness. The easier approach consists in considering that the whole conductance is the sum of conductances of the multiple contact spots with varying sizes. The effective contact resistance can be obtained as a first approximation by summing all the contact radii of all the individual asperities as the effective contact radius $a_{\text{eff}}$ of the effective contact area $A_{\text{eff}}$. This effective contact radius can be substituted into (1) and (2) to take into account the multiple asperity feature of the micro contact.

Because of tiny contact spot size due to small contact force available in micro-switches, mechanical, electrical and thermal properties shift from bulk to a specific physics of thin films lead by the geometry of the asperities. For example, heating of the contact spots is extremely localized when the current flows through the contact, whereas the device level remains at room temperature [7]. The highest contact spot temperature $T_c$ called supertemperature $T_{\theta}$ has already been expressed by Holm as a function of the contact voltage $V_c$ [6]:

$$T_{\theta} = \sqrt{\frac{\Gamma(K)R_{\text{Maxwell}}V_c^2}{R_c4L}} + T_0 \tag{3}$$

Where $L=2.45 \times 10^{-8} \text{W} \cdot \text{\Omega} \cdot \text{K}^2$ is the Lorentz constant, $\Gamma(K)$ a slowly varying Gamma function and $T_0$ the ambient temperature. By considering that electrical contact is strictly ohmic, this equation becomes:

$$T_{\theta} = \sqrt{\frac{V_c^2}{4L}} + T_0 \tag{4}$$

Eq. (3) is obtained from the Wiedemann-Franz law for a conductor heated by the current produced by the voltage $V_C$ between two arbitrary isotherms with the temperature $T_{\theta}$; the assumption is made that the thermal and electric currents obey similar laws thus with symmetric contacts the generated heat flows in the same path as the electric current.

### 2.2 Reliability of micro-contact

The second important point to be addressed is the reproducibility of the contact resistance over the actuations, and thus along its whole lifetime (cf. Figure 2). The reliability of the switch depends on its ability to withstand some degradations occurring at the contact interface. Three types of phenomena can be studied: the mechanical (cold welding, strain hardening), electrical (arcing, hot welding, annealing) and chemical ones (formation of insulating films at the extreme surface), all inducing modifications of the topological, mechanical and electrical properties of the contact.
As the experimental set-up developed here can be used for studying the fundamental physics of the micro-contact, the understanding of potential failure mechanisms can be made easier if the effort is put upon the identification of the key parameters for each mechanism. In addition, the contact temperature (Tc) of contact a-spots could be linked with the contact force depending on the contact voltage [4]. Indeed the plastic deformation of the asperities during the contact formation proceeds more rapidly when the softening temperature is reached [2]. Thus the effective contact area increases inducing a drop of the contact resistance. However the softening of the metal at the asperities of contact reduces the strain hardening of the a-spots and could accelerate the aging of the contact by the activation of thermal failure mechanisms. Therefore the softening temperature of contact metals is probably one of the most important critical parameter and it has to be study for reaching a stable and low contact resistance.

2.3 Contact material focus

The performance of the electrical contact in a MEMS switches is strongly linked with the materials used to perform the contact. The mechanical and electrical properties of the contact material will govern the behaviour of the contact resistance versus the contact load applied. It is important to find the best compromise between mechanical and electrical performances to reach the best reliable operations. The contact must have excellent electrical conductivity for low loss, high melting point to handle power, appropriate hardness to avoid stiction phenomenon and chemical inertness to avoid oxidation [8]. Furthermore, the used contact material influences the contact reliability because of material-dependent contact degradation through mechanical wear, fretting, creep, localized hardening, arcing, etc.

3. EXPERIMENTAL TECHNIQUES

3.1 Experimental set-up description

The test described in this paper is based on methods combining a nanoindenter and a high resolution source meter for the determination of the electrical contact resistance versus the contact force and the displacement of the free standing part of the contact. Nanoindentation is firstly designed for material characterization, which is based on the “displacement vs. load” curve obtained by driving a tip into the test material and by monitoring the applied load and the resulting displacement [9]. In our test bench, the nanoindenter’s spherical tip is used as a mechanical actuator. The actuation load is thus reproducible and known with a good accuracy, two points which remain hard to achieve with the actual actuations (for example, electrostatic actuation forces often drift due to dielectric charging [10]).
Once the bridge is correctly located under the column of nanoindentation, the tip is brought in contact with the surface of the mobile electrode, and then lowered until mechanical and electrical contact between the bridge and the line. The stiffness of the membrane can be measured in this step [11]. From this moment on, the applied load corresponds to the actuation load of the contact. A schematic view of the set-up can be seen on Figure 3.

The main difficulty for applying a load on test vehicles remains the accurate location of the tip above the contact. This issue has been here overcome by the use of a nanopositionning table achieving a 20nm X-Y resolution in a defined 100μm-side square. Thus, a pseudo-AFM scan is performed by probing the surface with a fixed stiffness of contact between the tip and the surface. An accurate location can then be reached on the scan (cf. Figure 4).

The electrical measurements are performed by the use of a four-wires probe: a current flow $I_C$ is applied and the potential drop $V_C$ is independently measured. This method for measuring contact resistance is the same method used in crossed rod contact resistance measurements of Holm [6]: only the contact resistance is measured independently of the voltage drop in the supply wires and access paths. Moreover the environmental chamber allows the control of the temperature of the DUT and of the relative humidity $Rh$.

### 3.2 Test vehicle description

Specific test vehicles have been designed for contact analysis, enhancing the extraction of characteristic curves and making possible the comparison between different contact shapes or materials. As illustrated in Figure 5, the tested device is composed of a bridge suspended over a contact line.
Figure 5: Scheme of the test structures

A 5\textmu mx5\textmu m square bump is processed underneath. Three kinds of test structures, described in table 1, are used.

<table>
<thead>
<tr>
<th>Nomenclature</th>
<th>Contact materials</th>
<th>Measured bridge stiffness</th>
</tr>
</thead>
<tbody>
<tr>
<td>P10</td>
<td>Au-to-Au</td>
<td>480 N/m</td>
</tr>
<tr>
<td>P09</td>
<td>Au-to-Ru</td>
<td>515 N/m</td>
</tr>
<tr>
<td>P05</td>
<td>Ru-to-Ru</td>
<td>225 N/m</td>
</tr>
</tbody>
</table>

Two contact materials are tested (cf. table 2): Au, which is the most popular material for electrical contact because of its high bulk conductivity, low contact resistance even at small contact forces, its high oxidation resistance, its low propensity to form alien surface films [11] and its fabrication compatibility with MEMS fabrication methods.

<table>
<thead>
<tr>
<th>Contact materials</th>
<th>Electrical conductivity (m\textOmega\cdot 1/mm\textsuperscript{2}) [9]</th>
<th>Softening temperature (\degree C)</th>
<th>Melting temperature (\degree C) [9]</th>
<th>Boiling temperature (\degree C) [9]</th>
<th>Estimated hardness (GPa) [9]</th>
</tr>
</thead>
<tbody>
<tr>
<td>Au</td>
<td>45.7</td>
<td>\sim100 [6]</td>
<td>1063</td>
<td>2966</td>
<td>\sim1.6</td>
</tr>
<tr>
<td>Ru</td>
<td>14.9</td>
<td>\sim430 [12]</td>
<td>2450</td>
<td>4900</td>
<td>\sim10.1</td>
</tr>
</tbody>
</table>

However, Au is a soft material, thus the probability of getting large modifications of the contact surface during aging is really high. Furthermore, Au is susceptible to contact wear, which affect the contact performance. This is the reason why new contact metals have to be introduced, as Ru. It is much harder than pure Au or any Au alloy. Ru/Ru contact and Ru/Au contact may provide better resistance stability in spite of a higher bulk resistivity [14]. The test structures were stored in clean \textsubscript{2}N\textsubscript{2} what could slow down environmental contamination of the contact surfaces but gradual contamination accumulation still occurred.

3.3 Test description

A campaign of tests has been made on the test vehicles described in the previous paragraph. These studies are focused on the self heating of the contact with increased levels of current flowing through the contact at low- to medium-power levels. Contact resistance versus contact force curves are obtained, up to 145 \textmu N for Au/Au and Ru/Au, and up to 200\textmu N for Ru/Ru. The contact super-temperature is then extracted. The measurements are performed by sweeping the contact current from 1 mA to 100 mA and measuring the contact voltage in a cold-switching way. The load rate is about 1\textmu N/s to allow the contact deformation to reach a steady-state. The following results are based on the same 11 successive tests.
for the three pairs of contact materials (Au/Au, Ru/Au and Ru/Ru). In addition dedicated validation tests and FEM modelling have already assessed the introduced methodology [15].

4. EXPERIMENTAL RESULTS

This part of the paper compares the heating of micro-contacts with various pairs of contact materials by increasing the contact current. To further examine the heating effect on contact resistance, the study is focused on the contact temperature when the maximum contact load is reached. The methodology consists on measuring the contact voltage while imposing the current flowing through the contact. The series of tests described previously will be used to compare the impact of increased contact current on the three test vehicles.

![Graph showing contact temperature versus contact current for Au/Au contact.]

Figure 6 - Contact temperature versus the contact current for Au/Au contact

The published softening temperature for gold contact is \( \sim100°C \), corresponding to a contact voltage of 70-80 mV for contact near room temperature [6]. When the contact current is increased after reaching the softening temperature from 40 mA, the contact resistance continues to decrease keeping the contact temperature roughly constant. The contact temperature increases with a constant slope from 1 mA to 40 mA. At this point the contact temperature stops to increase and fluctuates between 80°C and 120°C (65 mV to 75 mV). This temperature seems to be stabilize while the power within the contact rises.

The same behaviour is partially observed for the other symmetrical contact Ru/Ru. Indeed the published temperature for ruthenium contact is \( \sim430°C \), corresponding to a contact voltage of 200 mV for contact near from room temperature [12]. When the softening temperature is reached, the contact temperature does not depend strongly on contact current at high current levels. Beyond this value, the contact temperature is unstable even though it seems to oscillate around the softening temperature.
The behaviour for the asymmetrical contact Au/Ru defers from the two others because of the different materials in the both part of the contact.

As shown in Figure 8, the contact temperature increases with the current level without reaching a maximum. The temperature distribution within the contact constriction has change in comparison with a contact with the same contact metal. In the case of Au/Ru contact, the conductivities of the both material are different yielding to the apparition of thermoelectric effects.
5. DISCUSSION

The previous section pointed out the dissimilar behaviours of symmetrical and asymmetrical contact materials. The Figure 9.a is a scheme which represents both contact members of the same material M1. The assumption that the electrical and thermal currents flow in the same paths is always supported as described by the Wiedeman-Franz law. \( T_{b} \) is the bulk temperature in the both members of the contact. \( I_{c} \) corresponds to the section of the extreme contact interface across which no heat flows. The highest temperature \( T_{b} \) is localized at the contact interface where the contact voltage \( V_{c} \) is equal to 0V (Holm’s convention [6]). The distribution is symmetric around the hottest contact spot precisely localized at the intersection between the both part of the contact.

![Image of temperature distribution](https://example.com/figure9.png)

Figure 9 – Temperature distribution a) in a symmetric constriction b) in the constriction of a contact between two metals M1 and M2 with different conductivities

The Figure 9.b illustrates the temperature distribution at the contact constriction between two metals M1 and M2 with different conductivities. The distribution round the hottest area at the contact interface has changed because of the different nature of the both contact parts. Thus there is a change in the average distance of the thermoelectric heat flow. The maximum temperature \( T_{b} \) is located within the poorer conductor so the positions of interface around the contact between the metallic members M1 and M2 [6].

Thus the phenomenon observed in the Au/Ru contact can be explain as the maximum contact temperature calculated in the Figure 8 is thus located in the ruthenium member of the contact. The maximum contact temperature obtained is about 274°C, that is to say less than the softening temperature. The conclusion is that the contact temperature of the asymmetric contact is more stable because the softening temperature is theoretically not reached for the same contact current.
Figure 10 - Contact resistance versus contact force as a function of the current flowing through the contact for a) Au/Au, b) Ru/Ru and c) Au/Ru contacts
Measured contact force versus contact resistance curves for Au/Au, Ru/Ru and Au/Ru contacts are shown in Figure 10. When contact is made between two metallic electrodes, the contact area is very low at first and the resistance is high and unstable. At a certain minimum force, depending on the material under investigation, a significant reduction of the resistance occurs. At higher forces, the contact is stable and the resistance decreases slightly with further increasing force until a saturated regime. The contact resistance reduction with increased current is thermally induced.

The stability of the contacts is investigated in terms of dispersion of $R_C$ versus $F_C$ curves. The conclusion is that Au/Ru is clearly the more stable contact at the maximum contact load. Indeed, the contact temperature increases with the current until reaching the softening temperature for Au/Au contact and for Ru/Ru contact. The contact hardness is subsequently reduced because of the annealing of the contact. Then the roughness of the contact is lowered after softening of the surfaces. Thereby, the deformation of the contact asperities is higher with increased current making easier the breakdown of residual insulating films, and enlarging the total contact area. However, the variation of the contact resistance is determined by a competition between asperity flattening, which lowers the contact resistance and heating, which can affect the topography of the contact surface by the activation of thermal failure mechanisms. The stability of Au/Ru contact can then be helped by the lowering of material transfer because the contact materials are not softened. The low contact resistance of the Au/Ru contact could also be explain because this asymmetric contact includes a soft material and a harder one. The asperities of Au on the contact surface are more easily deformed due to the lower hardness compared to Ru. One assumption could be that the asperities of the softer material are much deformed by a hard material surface than a soft material one.

6. CONCLUSIONS

This new test facility enables new characterization tests of MEMS ohmic contacts under realistic conditions. The contact behaviour is typically multiphysic including mechanical, electrical and thermal interactions at the micro-contact interface. An emphasis was placed on the role of the low- to medium-power increase leading to contact heating. Various phenomena were observed as heating effects on contact resistance versus contact force dependence and the raw stabilization of the contact temperature from the softening. These mechanisms are commonly reported as failure vectors in the contact switch lifetime. Problems related to the softening of Au/Au and Ru/Ru electrodes are intimately related to the power handling capacity of the contact, which is limited by current density considerations. The main points revealed by these tests can be summarized as follow:

- The temperature rise induces contact softening. When the softening temperature is reached, the power dissipated in the contact contributes to anneal away the dislocations formed during plastic deformation.
- The softening of the contact surface is reached at ~3 mW for Au/Au contact and ~4.25 mW for Ru/Ru contact whereas the temperature of the Au/Ru contact is under the softening temperature at ~14 mW.
- The contact resistance becomes stable after reaching saturated regime with increasing force. This threshold is decreasing with increasing current level
- Au/Ru contact improved contact performance and enhanced power handling capability of the test structure thanks to electromechanical properties of the bimetallic contact configuration.

And finally, these tests bring to light the usability of this methodology to characterize the contact in term of reliability and performance under several test conditions. The knowledge gained is a better understanding of the contact physics in order to bring recommendations for further improvements of the contact reliability. The test vehicles used in this study could be also manufactured with various contact materials and several bump designs. The investigation of these two conception parameters is underway.

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