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Kinetic study of forest fuels by TGA: Model-free kinetic approach for the prediction of phenomena

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1. Introduction

Thermal decomposition kinetics of biomass is an important key in thermochemical conversion processes aimed at the production of energy and chemical products \cite{1–3}. Biomass is also implicated in wildland fire. Indeed, the rate of mass loss due to thermal decomposition determines the available volatile fuel in the flaming zone. To a lesser extent, the mass loss rate also determines the heat release rate (product of the heat of combustion and the mass of fuel burned). Therefore, the analysis of the thermal degradation of plant fuels is decisive for wildland fire modelling and fuel hazard studies \cite{4–10}. Numerous studies led to different degradation schemes in inert \cite{11–13} environment but only a few were monitored in air atmosphere \cite{14–16}. We propose here to study the thermal degradation of a forest fuel under air with thermogravimetry and kinetics analysis on the data. Degradation of lignocellulosic biomass is a very complex process of interdependent reactions; nevertheless it can be reduced to the reaction illustrated in Fig. 1.

Between 373 and 553 K only non-combustible gases are produced, primarily water vapour with some carbon monoxide and traces of formic and acetic acids \cite{17}. From 553 to 773 K active pyrolysis takes place. Pyrolysis breaks down the substance molecules into low molecular mass gases (volatile), highly flammable tars and carbonaceous char. The whole process is complex and lead to solid degradation and gaseous reactions. Thermogravimetric analyses were focused on solid phase degradation. In solid-state, a variation in apparent activation energy could be observed for an elementary reaction due to the heterogeneous nature of the solid or due to a complex reaction mechanism. This variation can be detected by isoconversional or model-free methods \cite{18}. The isoconversional analysis provides a fortunate compromise between the oversimplified but widely used single-step Arrhenius kinetic treatment and the prevalent occurrence of processes whose kinetics are multi-step and/or non-Arrhenius \cite{19}. These methods allow estimates of the apparent activation energy at progressive degrees of conversion for an independent model. These data are obtained by conducting multiple experiments at different heating rates. Application of model-free methods was highly recommended in order to obtain a reliable kinetic description of the investigated process. In a previous work we demonstrate the utility of thermal analysis in forest fuel hazard study \cite{20}, here we propose a kinetic study based on an isoconversional method and a simulation of the solid temperature in conditions outside the experimental range of thermal analysis.

2. Experimental

2.1. Sampling

We sampled the foliage and aerial parts of \textit{Arbutus Unedo} (Strawberry tree). This is an abundant species in the Corsican vegetation concerned by wildland fires. Plant materials were collected from a...
experiments were performed three times. A systematic deviation from preset linear temperature programs. All over the crucible base in all the experiments. The sample conformity of the sample was maintained by spreading it uniformly.

Reactions such as thermal cracking, repolymerisation and recrystallisation occur continuously in the reaction zone, thereby reducing the extent of secondary reactions. Metered flow rate 50 mL/min to sweep the evolved gases from the reaction zone and pre-exponential factor using Arrhenius equation. The model-fitting method involves two fits: the first establishes the model that best fits data while the second determines specific kinetic parameters such as activation energy and pre-exponential factor using Arrhenius equation. The model-fitting approach has the advantage that only one TGA measurement is needed. However, almost any fit(α) can satisfactorily fit the data by virtue of the Arrhenius parameters compensation effects and only a single pair of Arrhenius parameters results from the model-fitting procedure. Consequently, researchers give up this kind of method for the benefit of isokinetic methods, which can compute kinetic parameters without modelling assumptions [31,32]. The isoconversional method has the ability to reveal complexity of the process in the form of a functional dependence of the activation energy on the extent of conversion x. The basic assumption of these methods is that the reaction rate for a constant extent of conversion, α, depends only on the temperature [33,34]. To use these methods,

natural Mediterranean ecosystem situated away from urban areas in order to prevent any pollution on the samples. A bulk sample from six individual plants was collected in order to minimize interspecies differences. About 500g of plant were brought to the laboratory, washed with deionized water and oven-dried for 12 h at 333 K [21]. Only small particles (<6 mm) are considered in governing the dynamic of fire spread [22–24]. Dry samples were then ground and sieved to pass through a 600 μm mesh, then kept to the desiccator. The sieved powdery sample was stored in airtight plastic containers for future use. Thermal analyses were performed on these powders. The moisture content coming from the air was removed by using the dependence of the activation energy, in association with thermographical data [27–29]. The parameters of the reaction kinetics were determined using the following procedure: under non-isothermal conditions in which a sample is heated at a constant rate, the explicit temporal dependence in Eq. (2) is eliminated through the following trivial transformation.

\[
\frac{dx}{dt} = \frac{A}{\beta} e^{-\frac{E_a}{RT}} f(\alpha) \tag{3}
\]

with \( \alpha = \frac{m_0 - m}{m_0 - m_{\infty}} \) and \( p(x) = \int_{x_0}^{x} e^{-x/2} dx \).

Kinetics analyses are traditionally expected to produce an adequate kinetic description of the process in terms of the reaction model and the Arrhenius parameters. There are many methods for analysing solid-state kinetic data. These methods may be classified according to the experimental conditions selected and the mathematical analysis performed. Experimentally, either isothermal or non-isothermal methods are employed. The earliest kinetics studies were performed under isothermal conditions [29] but a major problem is that a typical solid-state process has its maximum reaction rate at the beginning of the transformation. However, the advantages of the non-thermal experimental technique are at least partially compensated by serious computational difficulties associated with the kinetic analysis [30].

Several mathematical methods can be used to calculate the kinetic of solid-state reactions: model-fitting and isoconversional (model-free) methods. A model-fitting method involves two fits: the first establishes the model that best fits data while the second determines specific kinetic parameters such as activation energy and pre-exponential factor using Arrhenius equation. The model-fitting approach has the advantage that only one TGA measurement is needed. However, almost any fit(α) can satisfactorily fit the data by virtue of the Arrhenius parameters compensation effects and only a single pair of Arrhenius parameters results from the model-fitting procedure. Consequently, researchers give up this kind of method for the benefit of isokinetic methods, which can compute kinetic parameters without modelling assumptions [31,32]. The isoconversional method has the ability to reveal complexity of the process in the form of a functional dependence of the activation energy on the extent of conversion x. The basic assumption of these methods is that the reaction rate for a constant extent of conversion, α, depends only on the temperature [33,34]. To use these methods,

### Nomenclature

<table>
<thead>
<tr>
<th>Symbol</th>
<th>Description</th>
</tr>
</thead>
<tbody>
<tr>
<td>α</td>
<td>conversion degree</td>
</tr>
<tr>
<td>m</td>
<td>mass of the sample (mg)</td>
</tr>
<tr>
<td>m₀</td>
<td>initial sample mass (mg)</td>
</tr>
<tr>
<td>mₘ</td>
<td>final sample mass (mg)</td>
</tr>
<tr>
<td>T</td>
<td>temperature (K)</td>
</tr>
<tr>
<td>t</td>
<td>time (min)</td>
</tr>
<tr>
<td>f(α)</td>
<td>kinetic model reaction</td>
</tr>
<tr>
<td>g(α)</td>
<td>integral form of the kinetic model reaction</td>
</tr>
<tr>
<td>p(x)</td>
<td>exponential integral</td>
</tr>
<tr>
<td>A</td>
<td>pre-exponential factor (1/s)</td>
</tr>
<tr>
<td>Eₐ</td>
<td>activation energy (kJ/mol)</td>
</tr>
<tr>
<td>R</td>
<td>gas constant = 8.314J/mol/K</td>
</tr>
<tr>
<td>β</td>
<td>heating rate (K/min)</td>
</tr>
<tr>
<td>W</td>
<td>product log function</td>
</tr>
</tbody>
</table>

### 2.2 Thermogravimetric experiments

TGA dynamic experiments were performed using a PerkinElmer TGA Pyris 1. Temperature, furnace and weight of the TGA were measured according to the manufacturer’s recommendation by measuring Curie point of magnetic standards: perkalloy® and alunel. The pulverized samples were conditioned in 33 μL open platinum crucibles with 10,000 ± 0.005 mg. The furnace was continuously purged with a flowing atmosphere of air at a fixed metered flow rate 50 mL/min to sweep the evolved gases from the reaction zone, thereby reducing the extent of secondary reactions such as thermal cracking, repolymerisation and recrystallisation. Experiments were performed at different heating rates (β = 10–40 K/min) from room temperature to 900 K and the uniformity of the sample was maintained by spreading it uniformly over the crucible base in all the experiments. The sample temperature, which is controlled by a thermocouple, did not exhibit any systematic deviation from preset linear temperature programs. All experiments were performed three times.

### 3. Kinetic procedure

#### 3.1 Model-free kinetic MFK algorithm

As we used TGA records the conversion degree is defined as:

\[
\alpha = \frac{m_0 - m}{m_0 - m_{\infty}} \tag{1}
\]

The rate of heterogeneous solid-state reactions can generally be described by:

\[
\frac{dx}{dt} = \frac{A}{\beta} e^{-\frac{E_a}{RT}} f(\alpha) \tag{2}
\]

The temperature dependence of the rate constant is described by the Arrhenius equation. Galwey and Brown [25] presented a theoretical justification for the application of the Arrhenius equation to the kinetics of solid-state reactions, and it has now been recognized that this empirical equation can represent the experimental rate data as function of temperature accurately for both homogeneous and heterogeneous reactions [26]. Some papers have demonstrated how a complete isoconversional kinetic analysis can be performed using the dependence of the activation energy, in association with thermo-analytical data [27–29].

The parameters of the reaction kinetics were determined using the following procedure: under non-isothermal conditions in which a sample is heated at a constant rate, the explicit temporal dependence in Eq. (2) is eliminated through the following trivial transformation.

\[
\frac{dx}{dt} = \frac{A}{\beta} e^{-\frac{E_a}{RT}} f(\alpha) \tag{3}
\]

The rate dependence in Eq.(2) is eliminated through the following trivial transformation.

\[
\int_{x_0}^{x} e^{-x/2} dx.
\]

Kinetics analyses are traditionally expected to produce an adequate kinetic description of the process in terms of the reaction model and the Arrhenius parameters. There are many methods for analysing solid-state kinetic data. These methods may be classified according to the experimental conditions selected and the mathematical analysis performed. Experimentally, either isothermal or non-isothermal methods are employed. The earliest kinetics studies were performed under isothermal conditions [29] but a major problem is that a typical solid-state process has its maximum reaction rate at the beginning of the transformation. However, the advantages of the non-thermal experimental technique are at least partially compensated by serious computational difficulties associated with the kinetic analysis [30].

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a series of experiments has to be performed at different heating rates [34–35]. The knowledge of $E_a$ vs. $\alpha$ allows detecting multi-step processes and predicting the reaction “kinetics scheme” [27] over a wide temperature range. Several methods do exist [35–37], for this work we chose the method of Kissinger–Akahira–Sunose (KAS) [37] applied without any assumption concerning the kinetic model. The Kissinger–Akahira–Sunose method uses the approximation presented in Eq. (5):

$$p(x) \approx \frac{e^{-x}}{x^2}$$ (5)

Taking into account the approximation, the logarithm of Eq. (4) gives:

$$\ln \left( \frac{\beta_j}{T_{jk}^2} \right) = \ln \left( \frac{A_\alpha R}{E_a(\alpha)} \right) - \ln(g(\alpha)) - \frac{E_a(\alpha)}{RT_{jk}}$$ (6)

where $E_a$ and $A_\alpha$ are respectively the apparent activation energy and the pre-exponential factor at a given conversion degree $\alpha$, and the temperatures $T_{jk}$ are those which the conversion $\alpha$ is reached at a heating rate $\beta_j$. During a series of measurements the heating rate are $\beta = \beta_1 \ldots \beta_j \ldots$

The apparent activation energy was obtained from the slope of the linear plot of $\ln(\beta_j/T_{jk}^2)$ vs. $1/T_{jk}$ performed thanks to a Microsoft® Excel® spreadsheet developed for this purpose.

To apply the method, it is necessary to obtain at least three different heating rates, we work with $\beta = 10, 20, 30$ K/min. Linear plots were computed for $\alpha \in [0.05, 0.95]$ with a 0.05 step. In the following, Table 1 presents the KAS plots and Fig. 3 shows the results of activation energy vs. conversion degree.

### 3.2. Numerical method predicting algorithm

The prediction of theoretical temperature vs. reaction degree for high heating rates from dynamic scans has a scientific interest for fuels degradation. Indeed, wildfires heating rates can be higher than the classical heating rates of thermal analyser apparatus.

This numerical method allows the calculation of the solid temperature at different conversion degree and outside the experimental range of heating rate authorized by the thermal analyser.

**Table 1**

<table>
<thead>
<tr>
<th>$\alpha$</th>
<th>$y = \alpha x + b^*$</th>
<th>$r^2$</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.05</td>
<td>$y = -7587.6 x + 10.119$</td>
<td>0.9986</td>
</tr>
<tr>
<td>0.10</td>
<td>$y = -7700.8 x + 10.236$</td>
<td>0.9999</td>
</tr>
<tr>
<td>0.15</td>
<td>$y = -8788.9 x + 11.442$</td>
<td>0.9998</td>
</tr>
<tr>
<td>0.20</td>
<td>$y = -9484.8 x + 11.909$</td>
<td>0.9986</td>
</tr>
<tr>
<td>0.25</td>
<td>$y = -10169.0 x + 12.774$</td>
<td>0.9996</td>
</tr>
<tr>
<td>0.30</td>
<td>$y = -10709.6 x + 13.254$</td>
<td>0.9995</td>
</tr>
<tr>
<td>0.35</td>
<td>$y = -10577.6 x + 12.553$</td>
<td>0.9948</td>
</tr>
<tr>
<td>0.40</td>
<td>$y = -10709.6 x + 12.349$</td>
<td>0.9939</td>
</tr>
<tr>
<td>0.45</td>
<td>$y = -10199.0 x + 11.073$</td>
<td>0.9915</td>
</tr>
<tr>
<td>0.50</td>
<td>$y = -9781.2 x + 9.597$</td>
<td>0.9858</td>
</tr>
<tr>
<td>0.55</td>
<td>$y = -8782.7 x + 7.8668$</td>
<td>0.9801</td>
</tr>
<tr>
<td>0.60</td>
<td>$y = -7824.7 x + 5.8464$</td>
<td>0.9806</td>
</tr>
<tr>
<td>0.65</td>
<td>$y = -6741.2 x + 3.7126$</td>
<td>0.9789</td>
</tr>
<tr>
<td>0.70</td>
<td>$y = -5940.2 x + 2.7153$</td>
<td>0.9935</td>
</tr>
<tr>
<td>0.75</td>
<td>$y = -5132.2 x + 0.7962$</td>
<td>0.9975</td>
</tr>
<tr>
<td>0.80</td>
<td>$y = -4368.2 x - 0.4271$</td>
<td>0.9990</td>
</tr>
<tr>
<td>0.85</td>
<td>$y = -3777.5 x - 1.3623$</td>
<td>0.9994</td>
</tr>
<tr>
<td>0.90</td>
<td>$y = -3265.5 x - 2.1571$</td>
<td>0.9999</td>
</tr>
<tr>
<td>0.95</td>
<td>$y = -3194.1 x - 2.0836$</td>
<td>0.9999</td>
</tr>
</tbody>
</table>

* $\alpha = E_a/R$, $b = [\ln(A_\alpha R/E_a(\alpha))] - \ln(g(\alpha))]$ and $x = 1/T_{jk}$.

Eq. (6) was rewritten with $b = [\ln(A_\alpha R/E_a(\alpha))] - \ln(g(\alpha))]$:

$$\ln \left( \frac{\beta_j}{T_{jk}^2} \right) = - \frac{E_a(\alpha)}{RT_{jk}} + b$$ (7)

In order to give the explicit formulas for the solutions to the nonlinear Eq. (7) we use a Mathematica® Package (Wolfram, 2003): “Solve”. This function gives the complete set of possible solutions (see Eqs. (8a) and (8b)):

$$T^{(1)}(\alpha_k) = - \frac{E_a(\alpha_k)/R}{2W(1/2 \times 10^k(E_a(\alpha_k)^2/R) / \beta_j)}$$ (8a)

$$T^{(2)}(\alpha_k) = - \frac{E_a(\alpha_k)/R}{2W(-1/2 \times 10^k(E_a(\alpha_k)^2/R) / \beta_j)}$$ (8b)

$W$ represents the “product log” function giving the solution for $w$ in $z = w^w$. The function can be viewed as a generalisation of a logarithm. To calculate the temperature for specific cases we select $T^{(1)}(\alpha_k)$. Indeed, $T^{(1)}(\alpha_k)$ gives a real solution and $T^{(2)}(\alpha_k)$ a complex one.

The values of $E_a$, and $b$ at different conversion degree are those previously calculated performing KAS method. $T^{(1)}(\alpha_k)$ is then the temperature reached by the solid at a specific conversion degree i.e. specific $E_a(\alpha_k)$ and $b$ for the heating rate $\beta_j$. In the following section we present results of temperature vs. conversion degree computed from Eq. (8a) at different values of $\beta_j$ (i.e. 5, 10, 20, 30, 40, 60, 80 and 100 K/min). Our approach is similar to the one proposed by Vyazovkin [27].

The scope of this study is to provide a global simplified kinetics of the degradation of different wood species with the final aim to be incorporate in a physical fire spread model, that is the reason why we chose to work on dynamic runs and with these high heating rates.

### 4. Results

#### 4.1. Features of the thermal processes

For each sample 3 runs were performed with the same experimental conditions, and the reproducibility was verified by good agreement between the 3 obtained mass loss curves (accuracy ±2%). Fig. 2 presents the experimental results on thermal degradation of Arbutus Unedo fuel in the range of 400–900 K under air sweeping.

TGA curves of Arbutus Unedo show two steps of weight loss confirmed in DTG by 2 peaks. During the first process the plant is pyrolysed in the temperature range [400–700 K]. Gases emission and oxidation are visualized with a mass loss around 60%.

---

Fig. 2. TGA and DTG curves of Arbutus Unedo obtained with a linear heating rate of 20 K/min under air atmosphere.
The second process can be considered like the surface combustion of charred residue. The char forms ashes in the temperature range [700–900 K] with a mass loss around 40% for this step. The same phenomena were observed and recorded by other authors [5,39–41].

4.2. MFK algorithm

For the whole degradation $E_a(\alpha)$ values were calculated for $\alpha \in [0.05, 0.95]$ with a 0.05 step. Table 1 resumes the KAS plots obtained on experiments performed at 10, 20 and 30 K/min. Fig. 3 shows the transformation-degree dependence of the effective activation energy.

To explain the complexity of phenomena we used the approach of Vyazovkin and Lesnikovich [38] based on the inter-relationship between the shape of the curve and the type of process.

In Table 1, the values of $a$ (indirectly $E_a$) have a tolerable change, but the vast variation of $b$ due to its large instability is hardly usable. For this reason model-free kinetic approach have the inconvenient to treating the experimental value of $A$ as a dependent and minor parameter [42]. Nevertheless, we chose to use the KAS method in order to get $E_a$ for different $\alpha$ and the method proposed herein is dedicated to the simulation of the sample temperature at high heating rates as encountered in wildland fire. We think these results will be useful for fire spread modelling.

As shown in Fig. 3, the dependence has a pronounced increasing character for $0.1 < \alpha < 0.3$ characteristic of processes involving competing reactions [27]. This fact is in agreement with the phenomena known during the degradation of plants. Indeed, we have seen that two processes occurred in the temperature range from 400 to 700 K: pyrolysis gases emission and their oxidation. These reactions are competing ones because as gases are produced, they are immediately oxidized by the ambient air. For $0.3 < \alpha < 0.65$, reactions of pyrolysis and smoldering are in competition. This state involves strong disturbances for the determination of the activation energy. We can see in Table 1 a deviation from the unity for the correlation coefficient in this range of conversion degree. For $0.65 < \alpha < 0.9$, we noticed an important decreasing dependence of $E_a$ in Fig. 3. Vyazovkin and Lesnikovich [38] and Vyazovkin [27] have observed decreasing dependences when studying the thermal degradation of wood and interpreted this fact as the transition of the process to the diffusion regime.

4.3. Predictions algorithm

MFK algorithm allowed the calculation of $E_{a\alpha}$ and $b$, for the whole phenomenon at 10, 20 and 30 K/min. Plots of experiments vs.

![Fig. 3. Transformation-degree dependence of the effective activation energy.](image)

numerical results (Eq. (6)) are shown in Fig. 4 and a good agreement is observed.

Values of $E_{a\alpha}$ and $b$ (obtained with $\beta = 10, 20, 30 K/min$) can be used to fairly describe the dependence of conversion degree on temperature for different heating rates.

The decomposition of biopolymers in thermogravimetric experiments is a heterogeneous process. Reactions occur at a gas-solid interface, involved certainly a multi-step mechanism for the degradation of this complex sample (several kinetic steps). The derivative of mass loss (Fig. 2) shown two peaks. The first peak has a complex structure, at least two different processes are imposed inside this peak.

The second peak is an individual process confirmed by a better stability of the $E_a$ values. However, we do the choice of a global reaction step because our scope is to provide simplified data for a multi-phasic fire spreading model [22]. Fig. 5 is the proof of the correctness of our considerations.

In wildland fire, the solid sample temperature is a crucial data which is difficult to predict since the heating rate of a particle can vary on a large domain. Fig. 6 is a plot of the sample temperature vs. the conversion degree at different heating rates. Values presented here are numerical ones and come from the calculation of the sam-
As the heating rate is increasing, the shape of the temperature vs. conversion degree superior to 0.75. For example at: 

\[
\begin{align*}
\alpha &= 0.75, \ T = 680 \text{ K for } \beta = 5 \\
\alpha &= 0.95, \ T = 715 \text{ K for } \beta = 5 \\
\alpha &= 0.95, \ T = 963 \text{ K for } \beta = 100
\end{align*}
\]

For values of the conversion degree varying from 0.4 to 0.75, the heating rate has a moderate influence on the reaction temperature. Thus, high heating rates do have a significant influence on the temperature of the thermal degradation during the smouldering phase. As the heating rate is increasing, the shape of the temperature vs. conversion degree curve is a combination of two linear plots. For the simulation at \( \beta = 100 \text{ K/min} \) we found the linear equations:

- \( 0.1 < \alpha < 0.45, \ T = 476 + 329\alpha \)
- \( \alpha > 0.45, \ T = 259 + 732\alpha \)

The present work demonstrates that at high heating rates, two linear regimes govern the thermal degradation of thin particles. The first step refers to the gases emission and oxidation and the second step concerns the char combustion (smouldering combustion). The transition between those two regimes is easily visualized in Fig. 7 with a strong slope variation of temperature at \( \alpha = 0.45 \) for \( \beta = 100 \text{ K/min} \). We present here a way to get a prediction of the solid temperature vs. the conversion degree at high heating rates. This knowledge is important in fire risk quantification since fire spreads are caused by the transition from smouldering to flaming combustion.

5. Conclusion

We used TGA in order to study the thermal degradation of a Mediterranean forest fuel. Two stages were visualized under air. The KAS method allowed calculating the activation energy vs. the conversion degree of the thermal degradation under air in the range [400–900 K]. In a first step, the activation energy was calculated for heating rates of 10, 20 and 30 K/min. We found that the activation energy increases from 90 to 125 kJ/mol when the conversion degree is varying from 0.1 to 0.4 and the activation energy decreases from 125 to 40 kJ/mol when the conversion degree is varying from 0.4 to 0.9. In a second step, values of activation energy were used to simulate the temperature of the sample for heating rates up to 30 K/min. As the heating rate increases the results show that the temperature behaviour vs. the conversion degree is affected. At low heating rates, the plot of the temperature exhibits a (sigmoid) shape whereas at high heating rates two linear plots represent the temperature curve. In the field of wildland fire, with high heating rates (around 100 K/min), this work show that the sample temperature is linearly correlated to the conversion degree. With a heating rate of 100 K/min, the linear plots are: \( T = 476 + 329\alpha \) for \( 0.1 < \alpha < 0.45 \) and \( T = 259 + 732\alpha \) for \( \alpha > 0.45 \).

References


