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▶ To cite this version:

Tahar Hammouda, Bertrand N. Moine, J.L. Devidal, C. Vincent. Trace element partitioning during partial melting of carbonated eclogites. Physics of the Earth and Planetary Interiors, 2009, 174 (1-4), pp.60. 10.1016/j.pepi.2008.06.009. hal-00533031

HAL Id: hal-00533031

https://hal.science/hal-00533031

Submitted on 5 Nov 2010

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Accepted Manuscript

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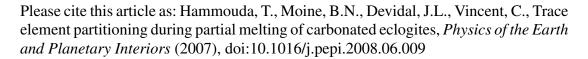
PII: S0031-9201(08)00128-3

DOI: doi:10.1016/j.pepi.2008.06.009

Reference: PEPI 4971

To appear in: Physics of the Earth and Planetary Interiors

Received date: 8-11-2007 Revised date: 14-4-2008 Accepted date: 13-6-2008



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Abstract:

Crystal / melt partition coefficients for Sr, U, and three Rare Earth Elements (La, Gd, Yb) have been measured during melting of carbonated eclogitic composition in the 4 - 7 GPa range. The experiments were conducted in a multi-anvil apparatus. Trace element concentrations were determined using electron microprobe and laser ablation ICP-MS in some cases. The obtained clinopyroxene / melt partition coefficients are very low whether the melt is silicate or carbonate in nature. This feature is due to Na presence and high vacancy concentration in the clinopyroxene. On the other hand, garnet / melt partition coefficients fall at the lower limit of or are slighly lower than litterature data dealing with carbonate melt-bearing systems.

We note that clinopyroxene / carbonated melt partition coefficient values for the Rare Earth Elements cannot be modelled using the lattice strain model. The reasons for the discrepancy lie on assumptions about melt structure as well as clinopyroxene solution model. In the latter case, high Na combined with high vacancy concentration precludes ideal mixing on octahedral sites to be a valid assumption. Taking charge balance on clinopyroxene crystallographic sites does not solve the discrepancy, suggesting that melt structure can be important. Garnet / melt partition coefficients can be predicted but we note a possible influence of the melt structure. The newly obtained partition coefficients have been used to test some simple models of mantle metasomatism.

1 Trace element partitioning during partial melting of carbonated eclogites.

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1. Introduction

Although they occur rarely on the Earth's surface, carbonatitic magmas are important
for our understanding of deep-seated processes such as the geodynamic cycle of volatiles and
trace element redistribution. The nature of the carbon source as well as the petrological
environment where carbonatitic magma genesis takes place has been the subject of numerous
experimental studies focusing either on peridotitic (Wyllie and Huang, 1976; Eggler, 1978;
Wallace and Green, 1988; Falloon and Green, 1989; White and Wyllie, 1992; Dalton and
Presnall, 1998) or on eclogitic sources (Hammouda, 2003; Dasgupta et al., 2004, Yaxley and
Brey, 2004). All authors showed that carbonatitic magmas could be generated in the mantle at
P and T conditions that can vary from those of the upper mantle to those of the transition
zone.
Erupted carbonatites are often enriched in incompatible elements (e.g. Nelson et al.,
1988; Wooley and Kempe, 1989). Carbon-rich liquids are also highly mobile in the mantle
(Watson et al.; 1990; Hammouda and Laporte, 2000), which gives them a high potential for
modifying lithologies they might encounter during their transport in the Earth's mantle
(Green and Wallace, 1988). The combination of high mobility and incompatible enrichment
prompted experimental studies aiming at determining trace element partition coefficients
between carbonatitic melts and silicate minerals (Brenan and Watson, 1991; Sweeney et al.,
1992, 1995; Klemme et al., 1995; Blundy and Dalton, 2000) in order to model trace element
incorporation in the melt at its source as well as element deposition during melt-rock
interaction. Most of the previous studies focused on minerals present in peridotites. However,
recent studies considered eclogitic environment as a potential source for carbonated liquids
(Hammouda, 2003; Dasgupta et al., 2004, Yaxley and Brey, 2004). In this perspective, it
appears that trace element partitioning data considering eclogitic mineralogy (omphacite,
calcic garnet, rutile) are lacking. The aim of the present study is to present experimental data

relevant to carbonatitic melt genesis in eclogitic environment. For this purpose, we have
performed partial melting experiments on trace element doped carbonated eclogitic material
at P and T conditions where earlier investigations showed that carbonate and / or silicate melt
can be produced. Run products have been analyzed by electron microprobe and laser ablation
ICP-MS for trace elements and partition coefficients have been determined.

2. Experiments and analyses

2.1. Experimental protocol

All experiments have been performed in multi-anvil apparatus. Starting material was a trace element doped mixture of basaltic glass and natural calcium carbonate (9:1 in weight). In order to introduce the trace elements, we mixed a synthetic basaltic glass (OTB issued from the same batch as material used in Hammouda, 2003) with carbonate containing the trace element mixture. Introduction of the latter was achieved by dry mixing of trace element oxides to the carbonate using successive dilution steps. This procedure which was at first intended for electron microprobe analysis yielded two batches of starting mixes that were doped with approximately 200 and 2000 ppm La, Gd, Yb, U, Pb, Sr expressed as oxides. Samples were loaded in $Au_{80}Pd_{20}$ capsules which were shut by arc welding.

High-pressure experiments were performed in the octahedral multi-anvil apparatus of the Laboratoire Magmas et Volcans in Clermont-Ferrand using a Walker (1990) module and following the procedure described in Hammouda (2003). High-pressure assemblies consisted of Cr-doped MgO octahedra and used pyrophyllite gaskets. All experiments were performed with an 18/11 assembly (octahedra side-length = 18 mm; WC cube truncation edge-length = 11 mm) calibrated against the α - γ transition in the Fe₂SiO₄ system (Morishima et al., 1994) and the coesite-stishovite transition (Akaogi et al., 1995). The recorded pressures of the experiments are believed to be accurate within 0.5 GPa. Heating was achieved using stepped

LaCrO₃ tubular furnaces. The samples were placed at the center of the furnace, inside MgO sleeves to prevent contact between the furnace and the gold capsules. Zirconia sleeves were placed between furnaces and octahedra to minimize heat-loss. The temperature was read with a W₅Re₉₅-W₂₆Re₇₄ thermocouple encased in a mullite sleeve and positioned axially in the assembly, in contact with the top of the sample container. The temperature was monitored using a Eurotherm 900 controller and remained constant within 1°C during the course of the experiments. In some cases where the thermocouple failed, temperature was monitored using the value of the electrical output power. No correction for the pressure effect on the thermocouple emf was applied. Each experiment consisted of first raising the pressure to the desired value before heating. Run duration was 20 hours at high temperature (Table 1). This duration was considered sufficient to achieve chemical equilibrium in a system containing melt and crystals at about 1300°C. For example, Blundy and Dalton (2000) used time series experiments and demonstrated chemical equilibrium at 3 GPa and 1375°C in less than 20 hours. The runs were terminated by switching off the power of the furnace, resulting in a temperature drop below 200°C in less than 2 seconds. Samples were then slowly decompressed during approximately 12 hours.

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2.2. Electron probe analyses (major and trace elements)

After recovery, experimental charges were mounted in resin and polished to $1/4~\mu m$ grit using diamond paste. The nature of the phase present and their textural relationships were determined using a scanning electron microscope at Laboratoire Magmas et Volcans. Operating conditions were 15 kV accelerating voltage. Major element compositions were determined with a Cameca SX100 electron microprobe at Laboratoire Magmas et Volcans. Operating conditions were 15 kV accelerating voltage and beam current of 15 nA, using a collapsed beam for crystalline phase and defocused beam (10 μ m) for melt pools. Standards

used were albite (Na, Si), Al₂O₃ (Al), Fe₂O₃ (Fe), olivine (Mg, Ni), wollastonite (Ca), synthetic MnTiO₃ (Mn, Ti), Cr₂O₃ (Cr) and orthoclase (K). Data reduction was performed using the ZAF procedure.

For trace element analyses (Sr, U, Pb, La, Gd and Yb) a 25 kV accelerating voltage was used with a 100 nA beam current. The three REE were chosen so as to avoid interferences on the L α fluorescence peaks. Measurements were performed on an LLiF analyzing crystal, using counting times of 100 s on peak. The three other trace elements were measured on a PET crystal using the L α line of Sr, with counting times of 300 s and the M β line of U and Pb, with counting times of 150 s. Standards were SrSO₄ for Sr, UO₂ for U, home made synthetic glass for Pb, and NIST standard glasses for REE. Detection limits were calculated following Ancey et al. (1987) and were less than 100 ppm for Sr and REE, 225 ppm for U and 350 ppm for Pb.

2.3. Laser ablation ICP-MS analyses

Trace elements in minerals and glass from experiment #273 were determined in polished section using a Perkin Elmer Elan 6100 ICP-MS coupled with a femtosecond laser (Pulsar 50 from Amplitude Technologies) at Laboratoire de Modélisation des Transferts Géochimiques (Observatoire Midi Pyrénées, Toulouse-France, Poitrasson et al., 2005). The laser was operated at 800 nm wavelength with 100μJ/pulse energy, 3 Hz frequency and 50 fs pulse duration. Crater size was 10-15 μm (Fig. 1) and ablation duration was 50-100 s. Ablation was done in pure He atmosphere and the analyte was carried to the ICP torch by a mixture of He +Ar. Detection limits range from 100 ppb for U to 10 ppm for Ni. NIST 610 glass standard (values from Norman et al. 1996) was used for calibration of relative sensitivities and each analysis was normalized using CaO content for clinopyroxene, garnet and melts, determined by electron microprobe. Signal intensity for major and minor elements

was monitored during analysis to make sure that the laser beam kept running within selected mineral grain. Trace-element reductions were done with the GLITTER software (Van Achterberg et al., 2001).

3. Results

3.1 Phase relationships and major element compositions

Experimental assemblages are given in Table 1 together with run conditions. Run products consist of garnet, clinopyroxene, coesite (at 6 and 7 GPa) and quenched melt (Figure 1). Electron microprobe analyses (major elements) of run products are given in Table 2. Garnets are calcium-rich whereas clinopyroxenes are sodium rich (omphacite). Additionally, we found that clinopyroxenes contain a large amount of the Ca-eskola compound, a feature that has been pointed out by Pertermann and Hirschmann (2002) in their study on eclogite melting. Mass balance has been performed on the basis of major element compositions determined by electron microprobe. Only experiments #273 and #274 could be treated. The others showed strong zoning across the capsule precluding meaningful mass balance.

Two types of melts were found. Up to 6 GPa, we found a silicate melt with about 51 wt% SiO₂ at 4.5 GPa to 40 wt% SiO₂ at 6 GPa. The electron probe total ranges from about 92 wt% at 4.5 GPa to 82 wt% at 6 GPa. Although low total can be anticipated in glass analyses using electron microprobe, the very low total found is partly attributed to dissolved CO₂ in the melts. At 7 GPa, the melt is purely carbonatitic, with no silicate fraction dissolved. This melt is calcium rich (Ca/ (Ca+Fe+Mg) ca. 0.77). These features are consistent with earlier work in the same system (Hammouda, 2003). Melt fraction is lower at 7 GPa (ca. 10 wt%) because only carbonate contributes to liquid whereas lower pressure melts (ca. 28 wt% at 6 GPa) are composed of a large silicate fraction.

3.2	Trace	element	concentrations	and	comparison	between	obtained	partition
coef	ficients	and litera	ture data					

Table 3 presents trace element concentrations and calculated partition coefficients. Lead could never be detected in the experimental run products. This absence is likely due to dissolution in noble metal capsules. When available (run 273) trace element concentrations obtained by the two techniques (EPMA and LA-ICPMS) are in agreement within 30% (Fig 2a). However, we note a systematic offset, where trace element concentrations obtained by EPMA are always lower than those obtained by LA-ICPMS. In addition, it seems that heavier elements display the largest offset. Important to note, however, is that calculated partition coefficients determined with both techniques agree very well (Fig. 2b).

Calculated mineral/melt partition coefficients and comparison with literature data on mineral / carbonated melt partition coefficients are presented in Figure 2a (clinopyroxene) and 2b (garnet). Literature data on trace element partitioning involving carbonatitic melts are available on clinopyroxene (Adam and Green, 2001; Blundy and Dalton, 2000; Brenan and Watson, 1991; Klemme et al., 1995; Sweeney et al., 1995, Keshav et al., 2005) and garnet (Adam and Green, 2001; Green et al., 1992; Sweeney et al., 1992, 1995). Previous experiments were performed between 1.5 and 6 GPa. Analytical techniques included electron microprobe, ion microprobe (SIMS, Blundy and Dalton, 2000), proton probe (PIXE, Green et al., 1992; Sweeney et al., 1992; Sweeney et al., 1995; Keshav et al., 2005).

For clinopyroxenes, we note that our values are lower by about one order of magnitude when compared with literature data, except for Sr. As will be discussed later, this is likely due to clinopyroxene crystal chemistry. On the other hand, our data for garnet are in

agreement with literature data. This general agreement for garnets is observed whatever the technique used by the previous authors.

When possible, we also calculated clinopyroxene / garnet partition coefficients (Table 3). The $D_{REE}^{cpx/gt}$ have rather low values, from about 1.5 for La to 0.024 for Yb. We note that the obtained values bear some resemblance to literature data on some Na- and vacancy-rich clinopyroxene and Ca-rich garnet found in mantle eclogites from Roberts Victor kimberlites (Caporuscio and Smyth, 1990; Harte and Kirkley, 1997). In particular, the higher the Na in clinopyroxene the lower the REE content while Ca-rich garnet are rather rich in REE.

4. Discussion

4.1. Mass balance and chemical equilibrium

Run durations (20 hours) have been chosen to favor chemical equilibrium in the samples. In order to check for the validity of our results, mass balance calculations have been performed on two samples where enough data were available. In addition, those calculations were used to discriminate which technique gave more consistent results when discrepancy was noted. The criterion was that mass balance should reproduce the composition of the starting material within 20%. Calculations performed on Sr, Gd, Yb, La and U show an overall agreement between computed mass balance based on analyses of experimental samples and initial concentrations in starting material (Table 4). We therefore confirm that experimental samples reached equilibrium. In detail, we note that in the case of sample #273, ICP-MS data are clearly overestimating true concentration for Gd and Yb, while EPMA data allow to better satisfy mass balance. Strontium mass balance is satisfactory whatever the analytical technique. Concerning La and U, only ICP-MS data are available. Mass balance for both U and La is satisfactory.

4.2. Lattice strain modeling and crystal chemical effects

Generally, experimentally obtained trace element partition coefficients are fitted using the lattice strain model such as that developed by Wood and Blundy (1997) for REE partitioning between clinopyroxene and silicate melts. The purpose of the lattice strain model is to predict values of unknown trace element partition coefficients using major element compositions of crystals and melts equilibrated at given P and T conditions. Here we used the theoretical expressions and compared the predicted values to our experimental results. We preferred this approach, instead of the usual fitting of experimental data found in the literature, because of the small number of elements analyzed in the present investigation. Using Wood and Blundy (1997), one can predict the value of the partition coefficient D_i of an element of ionic radius r_i as a function of r_0 , the radius of the crystallographic site, E, the Young's modulus of the site, and D_0 , the strain-free partition coefficient of a cation having the radius r_0 , with

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$$D_i = D_0 \exp\left(-4\pi E N_A \left[\frac{r_0}{2}(r_i - r_0)^2 + \frac{1}{3}(r_i - r_0)^3\right]/RT\right),$$

where N_A is Avogadro's number, R, the gas constant and T, the temperature in K. In the following, we discuss our results on clinopyroxene using Wood and Blundy (1997) and on garnet, using van Westrenen and Draper (2007) and Draper and van Westrenen (2007).

Pyroxenes

Following Wood and Blundy (1997) we focused on trivalent Rare Earth Elements. The predicted D_i vs. r_i relationship is shown on Figure 3a. It can be seen that the predicted partition coefficients are more than one order of magnitude larger than the actual data. This offset is larger than what is expected from Wood and Blundy (1997) model (92% probability that D_i is between 0.63 and 1.59 of the actual value). In order to gain further insight into the

model vs. data comparison, we have attempted to fit our REE data to the Wood and Blundy expression, keeping in mind that the fit is poorly constrained because we have only three elements. Nevertheless, La, Gd, and Yb covering almost the whole range of REE radii, we found the attempt worthwhile. The resulting fit is also shown on Figure 3. We note that the difference between the modeled and the fitted lattice strain parameters is essentially on the D_0 value. The value of D_0 is defined by Wood and Blundy (1997) in the following manner

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$$D_0 = \exp\left[\frac{88750 - 65.644T + 7050P - 770P^2}{RT}\right] \cdot \frac{Mg_{melt}^{\#}}{X_{Mg}^{M1}}$$

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This expression has been regressed and its coefficients ($\Delta H_T^0, \Delta S_T^0, \Delta V$, and $\frac{\partial \Delta V}{\partial P}$) have been adjusted for volatile-free silicate melts and vacancy-free clinopyroxenes. That our D_0 is poorly predicted by the model can be due either to crystal or melt properties. To address the first possibility we used the data of Pertermann and Hirschmann (2002) on vacancy-rich clinopyroxene at 3 GPa. The effect of melt properties have been tested using the high pressure data (6 GPa) of Keshav et al. (2005) on CaO-MgO-Al₂O₃-SiO₂-CO₂ (kimberlite) system and lower pressure (2.5 GPa) data on carbonatite system by Adam and Green (2001). For all studies, we used the lattice strain model to calculate D_0 , r_0 , and E. The results are added to Figure 3. Again, we note that r_0 and E are correctly predicted while calculated D_0 is larger by a factor 1.7 to 2.1 for vacancy-rich clinopyroxene, by a factor 2.8 to 3.3 in the kimberlite melt case, and by a factor 7.6 for the carbonatite system. As in our samples, the calculated D_i values fall outside the range [0.63 - 1.59] times the true values that has been judged acceptable by Wood and Blundy (1997). Since our samples combine both characteristics (i.e. vacancy-rich pyroxenes and carbonated melts), it is easy to understand that the lattice strain model fails to predict our experimental data.

One reason is that the model of Wood and Blundy assumes a quasicrystalline model wherein the melt consists of pyroxene-like structural units. Rare Earth Element partitioning is

then described as an equilibrium between a melt and a crystal having the same stoichiometry,
i.e., REEMgAlSiO ₆ (Eq. 30 in Blundy and Wood, 1997). However, silica undersaturated
carbonated melts such as those produced in the present investigation cannot be described in
terms of pyroxene-like entities. More likely, REE solubility should be described in terms of
carbonated complex because high-pressure studies on carbonate - silicate two-liquid
partitioning demonstrated that REE are preferentially partitioned in carbonate melt (Hamilton
et al., 1989; Wendtland and Harrison, 1979). Consequently, the heats, entropies, and volumes
of fusion of pyroxene-like units used in fitting Wood and Blundy's equations are not
applicable in our case.
The second reason for discrepancy between the model prediction and our data lies in
the presence of vacancies in the pyroxene. The composition-activity model of Wood and
Blundy (1997) uses $a_{REEMgAlSiO_6}^{cpx} = X_{REE}^{M2} \cdot X_{Mg}^{M1}$ on the assumption of a REE-bearing
clinopyroxene having the stoichiometry of REEMgAlSiO ₆ and ideal mixing on both M sites.
However, in our case, clinopyroxene M2 site has a low mean electrostatic charge (< +1.5)
because of the presence of a monovalent cation (Na) and of high vacancy concentration, while
the M1 site has a high electrostatic charge (> +2.5) due to the presence of Al ³⁺ associated with
the vacancies (Ca-Eskola molecule, $Ca_{0.5}\square_{0.5}AlSi_20_6$) and not compensated by lower charge
on the tetrahedra. This site filling is likely to be unfavorable for REE ³⁺ incorporation in the
M2 site and non-ideal mixing associated with short-range order in the clinopyroxene M sites
should be anticipated. Wood and Blundy (2001) addressed the effect of electrostatic charge
due to cation valence (non ideality) and they proposed a modification for the D_{o} expression
that takes into account the probability of having a favorable crystallographic site, in terms of
charge balance. Although they were mainly interested in CaTs molecule (i.e., tetrahedral Al
coupled to octahedral Al in M1 site) we attempted to model our vacancy-rich clinopyroxenes
using their formulation. We found that the predicted Do were slightly lowered as expected for

clinopyroxene with little Al^{IV} but never reached our experimental values. We note from Wood and Blundy (2001) that taking into account cation charge changes the D_o values for trivalent cations by no more than 25% relative, either increase or decrease, depending on tetrahedral Al content. Therefore, taking into account cation charge does not reconcile experimental data on carbonated melt-bearing systems with predictions using lattice strain modeling. When applied to Pertermann and Hirschmann's (2002) compositions, Wood and Blundy's (2001) expression increases D_o values because of high Al^{IV} contents, resulting in a poorer agreement with experimental values.

Garnet

Draper and van Westrenen (2007) and van Westrenen and Draper (2007) revised the earlier model of van Westrenen et al. (2001) and developed two predictive models for garnet / melt trivalent trace element partitioning (REE, Y, Sc). Both approaches are compared to our results on REE, on Figure 3b. Overall, it seems that the models satisfactorily predict the measured partition coefficients because the experimental points fall within the range defined by the statistical and thermodynamic models. We note, however, that the predicted E is too large resulting in impossibility of simultaneously predicting all REE D's with a single model. The D_o predicted with the statistical model agrees well with the experimentally determined value. Therefore, MREE and HREE can be reproduced but not LREE. On the other hand, the D_o predicted with the thermodynamic model is higher by a factor of about 50 yielding in poor prediction of MREE and HREE, while measured LREE falls on the predicted value. That the thermodynamic model fails to predict D_o accurately suggest that either garnet or melt properties are not well reproduced by the model. Concerning garnet, the critical point is the presence of quadruvalent cations in the octahedral site. Si and Ti were incorporated in the model for E (Young's modulus) by van Westrenen and Draper (2007) but neither the

thermodynamic nor the statistical model explicitly take into account the majoritic component in D_o modeling because its effect seemed to be minor. Nevertheless, our garnets are not majoritic and therefore, no compositional effect emanating from garnet is expected. On the other hand, melt composition or structure can be important. van Westrenen and Draper (2007) considered that melt effect is difficult to demonstrate given the available garnet / melt partitioning data. On the other, Draper and van Westrenen (2007) anticipate a potential effect of melt composition and structure in the presence of water. Here, we suggest that carbon presence in the melt could also have an impact on partition coefficients.

4.3. Effect of melt composition on partition coefficients and liquid/liquid partition coefficients

Previous studies on carbonate – silicate liquid/liquid partition coefficients have shown that REE partitioning strongly depends upon pressure. At low pressure (P < 0.1 GPa), Veksler et al. (1998) showed that all REE but La partitioned into the silicate melt. Liquid / liquid partition coefficients, $D_{REE}^{carb/sil}$, ranged between 1.33 (La) to 0.30 (Tm). Hamilton et al. (1989) studied the effect of pressure, temperature, and melt structure on $D_{REE}^{carb/sil}$. They observed that $D_{LREE}^{carb/sil} > D_{HREE}^{carb/sil}$ at all conditions. At low pressure, REE are partitioned in the silicate melt and the trend is reversed at P > 0.2 GPa for LREE and P > 0.4 GPa for HREE. According to Hamilton et al. (1989) increasing temperature favors REE partitioning into Na-rich silicate melts but temperature has limited effect in Ca-rich system. Wendlandt and Harrison (1979) showed that REE are preferentially partitioned in carbonate melt at 1300°C and 0.5 and 2 GPa.

The present data can be used to discuss the effect of melt composition on REE partitioning. Figure 4 shows that $D_{Gd}^{gr/liq}$ decreases smoothly with pressure for Gd in the 4.5 – 7 GPa range, while $D_{Yb}^{gr/liq}$ is constant within error in the 4.5 – 6 GPa range and decreases

abruptly between 6 and 7 GPa. Decreasing trends for MREE and HREE have been observed
by Draper et al. (2003) for garnet / melt in chondritic system. These authors attributed their
observations to the onset of majoritic substitution in garnet. In the present study, we have no
evidence for majoritic substitution in garnet. More likely, D_{Gd} and D_{Yb} evolution between 4.5
and 6 GPa is solely due to pressure effect on garnet bulk modulus and consequently on its
Young's modulus (E). Given the form of the pseudo-parabolae that describe trivalent cation
partition coefficients in the octahedral site of garnets (e.g. van Westrenen et al., 1999), we can
anticipate that HREE will be only slightly affected by E increase because they sit on the
maximum of the curve (i.e. HREE effective radii are close to the ideal value corresponding to
strain-free substitution). On the other hand, as MREE sit on a descending branch, lower D_{REE}
can be expected with pressure increase accompanied by higher E values.
The sharp decrease observed for Yb between 6 and 7 GPa (Figure 4) cannot be related
solely to pressure nor to garnet composition because garnets are almost identical in the 6 and
7 GPa experiments. We must therefore conclude that the change in $D_{Yb}^{gt/melt}$ is mainly due to
change in liquid composition from silicate to carbonate. Wendlandt and Harrison showed that
not only are the REE preferentially partitioned in carbonate melt at high pressure but also that
$D_{HREE}^{carbonate / silicate} > D_{MREE}^{carbonate / silicate} > D_{LREE}^{carbonate / silicate}$. The consequence is that Yb is more retained in
the melt when it becomes carbonatitic, a feature that explains the sharp decrease of $D_{Yb}^{gt/melt}$
between 6 and 7 GPa.
Although the present study was not designed to investigate trace element liquid/liquid
partition coefficients, our data allow to discuss the effect of melt nature by considering the
composition of silicate melt and coexisting carbonate globules in the lower pressure run (6
GPa; 1300°C). If we assume that it is an equilibrium texture, we can calculate the following
partition coefficients using electron probe analyses: $D_{Sr}^{carb/sil} = 1.58$; $D_{U}^{carb/sil} = 1.38$;
$D_{La}^{carb/sil} = 1.96$; $D_{Gd}^{carb/sil} = 1.83$; $D_{Yb}^{carb/sil} = 1.69$. We find that Sr is preferentially partitioned into

carbonate melt in agreement with previous investigators (Veksler et al., 1998). In the case of REE, we agree with Wendlandt and Harrison (1979) and Hamilton et al. (1989) in that REE are preferentially partitioned in carbonate melt at high pressure. Therefore, this trend appears to be robust at mantle pressure, whatever the melt composition. Almost all previous studies used alkali-rich melt compositions, mainly because the authors were interested in direct comparison with natural melt compositions. In our case, melts are poorer in alkalis, mainly because of the presence of Na-rich clinopyroxene in the experimental assemblages. Nevertheless, high-pressures favor REE partitioning into the carbonate melt.

5. Implications for mantle metasomatism

Melts produced by carbonated eclogite melting might be liberated from the slab and interact with surrounding mantle. The processes involved are illustrated in Figure 5 on the basis of REE, Sr and U only. (We are presently lacking information on other elements.) In our models, we considered a starting composition of 90 percent basalt + 10 percent carbonated pelagic sediments (in weight). Initial REE, Sr and U concentrations were taken from McDonough and Sun (1989) and Planck and Langmuir (1998) for basalt and sediments, respectively. The trace element spectrum of the starting mixture is shown on Figure 5a. All simulations considered batch-melting hypothesis.

The spectra of melts produced at 6 and 7 GPa, are almost identical (Fig. 5a). The contrast in major element composition between melts produced at 6 GPa (silicate) and 7 GPa (carbonate) appears to have limited impact. Main features are enrichment of 80 – 100 for the most incompatible elements, and enrichment of slightly less than 10 for HREE, relative to primitive mantle (McDonough and Sun, 1995). HREE fractionation relative to LREE is due to garnet presence in the residue. Positive U and Sr anomalies reflect the simulated source characteristics.

Next we consider two end-member models for a carbonatite melt interacting with
peridotite mantle (Figure 5b). In the first case, a low melt fraction (1 wt%) is simply added to
the peridotite. The resulting REE spectra exhibit slight LREE enrichment relative to HREE
([La/Yb] $_{ m N}$ < 2). In the second type of model, we considered large volumes of carbonatitic
melt percolating through the mantle and subsequent melt-mineral equilibrium. In this model
we used the partition coefficient fitted from the data of Adam and Green (2001) for garnet,
and from Blundy and Dalton (2000) for clinopyroxene. The percolated mantle contains 15
wt% clinopyroxene and 8 wt% garnet. The calculated REE spectrum exhibits HREE
enrichment relative to LREE ([La/Yb] $_N$ < 0.01) with [LREE] $_N$ < 1. Very low [LREE] $_N$ is due
to low values of the clinopyroxene / melt and garnet / melt partition coefficients for the LREE
and the assumption of very large volume of percolating melt that leached the mantle. A more
plausible scenario is that of infinite mantle reservoir. In that case, the LREE part of the
spectrum levels off with values slightly less than 1.
The third step of the model consists in remelting the mantle region metasomatized by

The third step of the model consists in remelting the mantle region metasomatized by the carbonatitic melt (Figure 5c). One case considers mantle wherein melt has been simply added (melting degree 0.5 and 1.5 wt%), whereas the other case (1.5 wt% melting degree) is for the mantle region equilibrated with the percolating melt. The first model yields a steep pattern with high LREE enrichment relative to HREE ([La/Yb]_N about 40 to > 100). On the opposite, the second model yields pattern with [REE]_N in the range 5 - 10 exhibiting slight MREE enrichment. This type of spectrum is unknown for natural melts but some alkaline cumulates have been found with similar patterns (Grégoire et al., 1998).

The REE pattern obtained by melting of mantle with simple addition of 1 wt% carbonated melt has about the same level of HREE enrichment and twice LREE levels when compared to melting of fertile pyrolitic mantle (with 1.5 wt% melting, Fig. 5c). Our modeled melts bears strong resemblance with those of kimberlite particularly when very small melting

degrees are considered. Modeling of trace element patterns of kimberlites has been attempted by Keshav et al. (2005). Those authors modeled melting of a refractory mantle portion mesomatized by silicate melt by 1 wt% melting and found that similarities with kimberlites could be obtained only after repeated steps of melting and metasomatism. Here we find that a single metasomatic stage is sufficient in order to obtain kimberlite REE signatures.

6. Conclusions

We have obtained new partition coefficients values relevant to melting of subducted carbonate-bearing eclogitic composition in the 4-7 GPa range. The nature of Na- and vacancy-rich clinopyroxenes yields very low values for crystal / melt partition coefficients of REE, U, and Sr to a lesser extent, compared to literature data. Garnet / melt partition coefficients fall at the lower limit of literature data dealing with carbonate melt-bearing systems.

Clinopyroxene / carbonated melt partition coefficient values for the Rare Earth Element cannot be modeled using the lattice stain model (Wood and Blundy, 1997) in its current state, mainly because the pyroxene and melt respective solution models differ strongly in the present case. In addition, ideal solution on each octahedral (M) site in Na- and vacancy-rich clinopyroxenes might no longer be a valid assumption. Taking into account cation charge budget on octahedral sites (Wood and Blundy, 2001) does not allow reducing the discrepancy significantly. We also find that garnet partition coefficients are not predicted in a satisfactory manner, suggesting that carbonated melts have an effect that is not included in the models in their present state as suspected in the case of water by Draper and van Westrenen (2007).

Simple modeling of single stage reaction of melts liberated by carbonated eclogites in the mantle shows that it possible to create sources whose subsequent melting might yield liquid bearing strong resemblance with kimberlites, on the basis of REE pattern, mainly.

4	1	4	
4	1	5	

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Acknowledgements: We thank M. Veschambre for assistance with the electron microprobe and J.-M. Hénot and F. Faure for assistance on the scanning electron microscope. We much appreciated the help from F. Poitrasson and R. Freydier on the laser ablation ICP-MS. Comments by J. Longhi and an anonymous reviewer helped improve the manuscript. Financial support from CNRS-INSU (DyETI program) is acknowledged. The multi-anvil apparatus of Laboratoire Magmas et Volcans is financially supported by the Centre National de la Recherche Scientifique (Instrument National de l'INSU).

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Table 1
Details on experimental run conditions and resulting phase assemblages and phase proportions.

Run#	P (GPa)	T (°C)	t (h)	Result	gt	срх	coes	melt	Σ sq. res.
258	4.5	1300	20	gt-cpx-sil melt					
259	5.0	1300	20	gt-cpx-sil melt					
273	6.0	1300	20	gt-cpx-coes-sil melt	23.76	44.58	2.23	27.78	1.415
274	7.0	1300	20	gt-cpx-coes-carb melt	35.06	47.13	6.62	10.25	0.306

Table 2 Major element composition of selected phases in experimental assemblages. (n.a., not analyzed.)

Oxide (wt %)	SiO ₂	Al_2O_3	FeO	MgO	CaO	Na ₂ O	K ₂ O	Total
ME258 (4.5 Gpa - 1300°C)								
CPX	51.1081	16.1666	8.882	6.5218	14.1514	3.4629	n.a.	100.3045
GARNET	39.822	22.2105	18.1422	8.1454	12.579	0.0383	0.0151	100.9525
MELT	51.2436	15.0045	7.6117	3.1035	13.0825	2.2664	0.1731	92.4854
ME259 (5 Gpa- 1300°C)								
CPX	50.8388	15.333	7.732	6.9036	15.4316	3.5306	n.a.	99.7696
GARNET	39.317	22.0813	18.2679	8.161	12.0877	0.0688	0.0242	100.0079
MELT	47.9129	13.8499	8.9488	3.0546	14.5676	2.0411	0.1983	90.5732
ME273 (6 Gpa - 1300°C)								
CPX	52.716	17.598	5.565	6.5	13.27	4.196	n.a.	99.857
GARNET	39.168	22.195	15.111	8.731	13.62	0.109	n.a.	98.938
MELT	39.78	10.193	9.324	3.127	18.247	1.658	0.118	82.45
ME274 (7 Gpa - 1300°C)								
CPX	54.494	15.607	5.981	6.702	12.651	5.49	n.a.	100.929
GARNET	40.065	21.81	15.282	9.017	13.709	0.125	0.035	100.046
MELT	0.086	0.161	7.236	4.555	40.67	0.472	0.066	53.246

Table 3 Structural formulae and site filling for clinopyroxenes and garnets from Table 2. All iron was assumed to be in the form Fe^{2+} . For clinopyroxene, satured M1 octahedral site was considered and all vacancies were assigned to M2 site.

CPX	ME258	ME259	ME273	ME274	GARNET	ME258	ME259	ME273	ME274
P (GPa)	4.5	5	6	7	P (GPa)	4.5	5	6	7
T (°C)	1300	1300	1300	1300	T (°C)	1300	1300	1300	1300
Si	1.842	1.843	1.870	1.918	Si	2.991	2.983	2.976	3.010
AIIV	0.158	0.157	0.130	0.082	Al	1.967	1.975	1.988	1.931
					Fe	1.140	1.159	0.960	0.960
M1					Mg	0.912	0.923	0.989	1.009
AIVI	0.529	0.498	0.605	0.566	Ca	1.012	0.983	1.109	1.103
Mg	0.350	0.373	0.344	0.352	Na	0.006	0.010	0.016	0.018
Fe	0.121	0.129	0.051	0.083	K	0.001	0.002	0.000	0.003
					sum cations	8.029	8.036	8.038	8.035
M2									
Ca	0.547	0.599	0.504	0.477	Ca%	33.04	32.07	36.26	35.91
Na	0.242	0.248	0.289	0.375	Mg%	29.76	30.11	32.33	32.85
Fe	0.147	0.106	0.114	0.093	Fe%	37.20	37.82	31.40	31.24
Mg	0.000	0.000	0.000	0.000	Mg#	44.45	44.32	50.73	51.25
sum cations	3.936	3.953	3.907	3.945					
vac	0.064	0.047	0.093	0.055					
Mg#	56.68	61.41	67.55	66.63					

Table 4
Trace element concentrations (in ppm) and resulting partition coefficients. Number in brackets indicate the uncertainty (1 sigma) on the last reported digits; n is the number of analyses in each phase.

Gd

Yb

189(1)

205(12)

6361(180)

2459(43)

2033(55)

8355(134)

or analyses in	each phase.							
By electron r	microprobe							
Concentratio	ons							
run# P (GPa) T (°C) phase n	silicate melt	ME258 4.5 1300 Garnet 7	D gt/silicate melt		silicate melt	ME259 5 1300 Garnet 4	D gt/silicate melt	
Sr U La Gd Yb	728(47) 1023(83) 690(47) 330(33) 190(83)	234(64) 679(219)	0.71(27) 3.57(271)		700(61) 650(63) 778(29) 300(22) 125(35)	170(51) 458(45)	0.57(21) 3.66(1.39)	
run# P (GPa)					273 6			
T (°C) phase n	silicate melt 11	carbo. melt 2	Pyroxene 5	13 Garnet 6	300 D cpx/silicate melt	D cpx/carbo. melt	D gt/silicate melt	D gt/carbo. melt
Sr U La Gd Yb	5911(44) 3680(76) 4738(30) 5012(53) 1702(27)	9385(2595) 5080(2036) 9265(2680) 9950(3041) 2985(587)	224(21) 166(8) 148(24)	1568(145) 6135(121)	0.038(4) 0.033(2) 0.087(15)	0.024(9) 0.017(6) 0.050(18)	0.31(3) 3.60(13)	0.16(6) 2.06(44)
run# P (GPa) T (°C) phase n	carbo. melt	Pyroxene 5	ME274 7 1300 Garnet 2	cpx/carbo. melt	D gt/carbo melt			
Sr U La Gd Yb	18845(2530) 12713(2813) 16408(2682) 11643(1790) 1855(253)	422(50) 142(47) 106(41)	1670(293) 3877(450)	0.022(6) 0.012(6) 0.057(30)	0.14(5) 2.09(53)			
By laser abla	ntion ICP-MS -	sample ME27	73					
Concentratio	ons				Partition coe	fficients - san	nple ME273	
n	silicate melt 3	pyroxene 3	garnet 2		D cpx/melt	D gt/melt	D cpx/gt	
Sr U La	6320(152) 5475(252) 5221(44)	137(10) 1.7 9.1(1)	15.3(27) 18.7(2) 6(1)		0.022(2) 0.00031(2) 0.00175(3)	0.0024(5) 0.0034(2) 0.0012(3)	9.0(2.2) 0.091(1) 1.51(33)	

0.030(1)

0.083(6)

0.32(2)

3.4(1)

0.093(3)

0.024(2)

Table 5 Mass balance calculations on experimental samples.

sample #273 : 6 G	SPa - 1300°C								
•		EPMA	EPMA	EPMA	ICPMS	ICPMS	ICPMS	ICPMS	ICPMS
	fractions	Sr	Gd	Yb	Sr	La	Gd	Yb	U
Coesite	2.23	0	0	0	0	0	0	0	0
Clinopyroxene	44.58	224	166	148	137	9	189	204	1.7
Garnet	23.76	0	1568	6135	15	6	2033	8355	19
Melt	27.78	5911	5012	1702	6320	5221	6361	2459	5475
resulting balance		1742	1839	1996	1820	1456	2334	2759	1526
starting concentrat	tion	1757	1768	1740	1757	1739	1768	1740	1765
offset (rel%)		-0.86	4.01	14.74	3.61	-16.29	32.02	58.58	-13.53
sample #274 : 7 G	GPa - 1300°C								
·		EPMA	EPMA	EPMA					
	fractions	Sr	Gd	Yb					
Coesite	6.62	0	0	0					
Clinopyroxene	47.13	422	142	106					
Garnet	35.06	0	1670	3877					
MeIt	10.25	18845	11643	1855					
resulting balance		2131	1846	1599					
starting concentrate	tion	1757	1768	1740					
offset (rel%)		21.26	4.40	-8.08					

Table 6 Parameters obtained using the lattice strain model for trivalent cations in clinopyroxene and garnet. Top: fit on the experimental data on La, Gd, Yb. Bottom: calculated parameters using predictive model of Wood and Blundy (1997) for cpx and van Westrenen and Draper (2007, D_o thermodynamical) and Draper and van Westrenen (2007, D_o statistical) for garnet.

			CPX			GARNET			
Exp #	P (GPa)	T (°C)	r _o (Å)	D _o	E ³⁺ (GPa)	r _o (Å)	D _o exp.		E³+ (GPa)
Experimental									
273	6	1300	0.945	0.0871	251	0.945	3.46		572
Model							D₀ thermo.	D₀ stat.	
273	6	1300	0.974	1.46	303	0.922	175	7.69	664
274	7	1300	0.977	1.47	310	0.917	147	6.83	664

Figure captions

Figure 1: (a) Backscattered electron micrograph of run products showing coesite needles

(darkest gray), Na-rich clinopyroxene (dark gray), garnet (lightest gray) in glassy matrix. (b)

Close-up view of the region shown in (a) illustrating the size of the craters due to laser

ablation.

Figure 2: Obtained partition coefficients and comparison with literature data (references used

are given in the text).

Figure 3: Comparison between experimentally obtained partition coefficient values and

values predicted using lattice strain modeling. Ionic radii are those of Shannon (1976). (3a)

Clinopyroxenes: experimental data (exp.) are compared to predictions (pred.) using Wood

and Blundy (1997). In addition to data obtained in the present investigation, references used

are AG, Adam and Green (2001); K, Keshav et al. (2005); PH, Pertermann and Hirschmann

(2002). (3b) Garnets: experimental data (exp.) are compared to predictions using the

statistical approach (pred. stat) of Draper and van Westrenen (2007) and the thermodynamic

approach (pred. thermo) of van Westrenen and Draper (2007). See text for further discussion.

Figure 4: Partition coefficient values for Gd and Yb in garnet as a function of pressure.

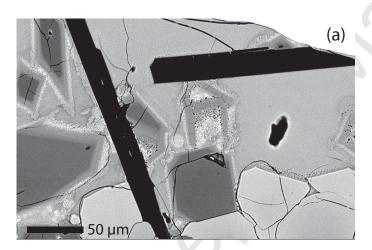
Figure 5: Modeling of REE spectra associated to mantle metasomatism due to melting of

MORB + carbonate mixtures. All spectra are normalized relative to primitive mantle

(McDonough and Sun, 1995). (5a) Modeled REE spectra of a source composed of 0.9 MORB

+ 0.1 carbonated sediments as used in the experiments and liquids produced by 10 wt%

melting degree (batch melting) at 6 and 7 GPa. (5b) Modeled REE spectra of a lherzolitic mantle after interaction with carbonatitic melts produced at 7 GPa. Two case have been considered: simple adjunction of 1 wt% melt (10 wt% melting degree) to mantle; melt mantle equilibration wherein the mantle portion is impregnated by an interconnected melt network (also produced by 10 wt% melting degree). In the latter case, partition coefficients used where those of Blundy and Dalton (2000) for clinopyroxene and Adam and Green (2001) for garnet. Both melt type are used in the models detailed in the next panel. (5c) Modeled REE spectra for melting of a metasomatized mantle portion and comparison with low degree melting of fertile pyrolite. The following cases are illustrated: (■) 1.5 wt% melting of pyrolite with 1 wt% carbonatite melt added; (△) 0.5 wt% melting of pyrolite equilibrated with interconnected carbonatitic melt network treated as an infinite reservoir; (△) 1.5 wt% melting of pyrolite. All models assume batch melting.



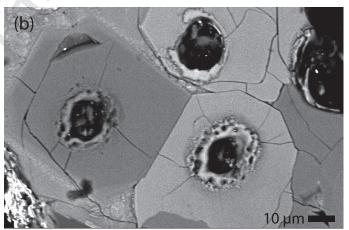


FIGURE 1

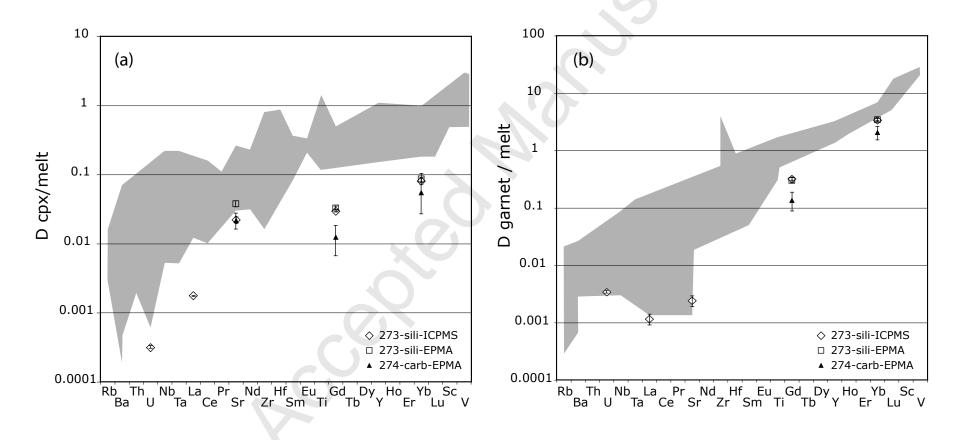
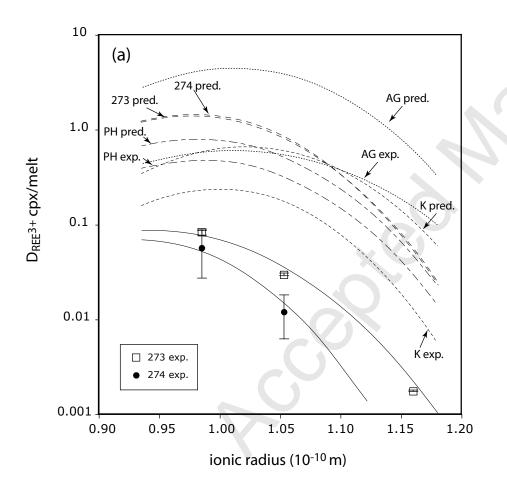


FIGURE 2



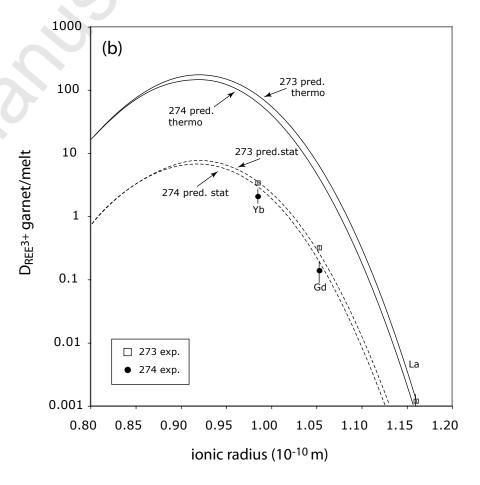


FIGURE 3

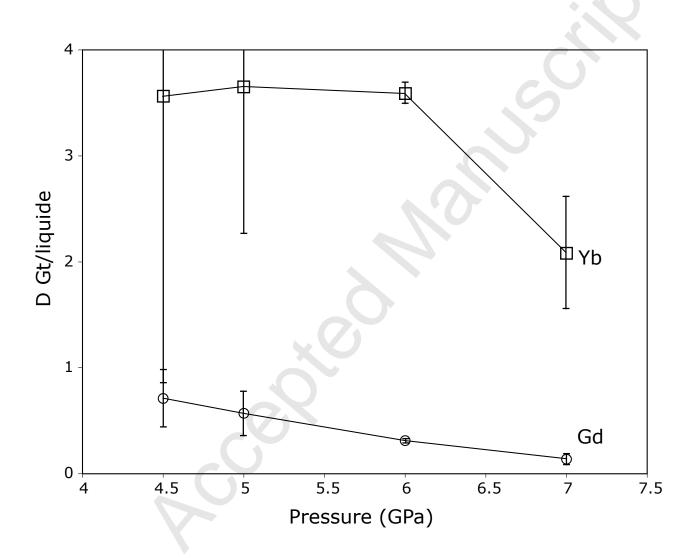


FIGURE 4

