Cool gas and dust in M33: Results from the Herschel M33 extended survey (HERM33ES)

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Cool gas and dust in M33: Results from the Herschel M33 extended survey (HERM33ES). *


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ABSTRACT

We present an analysis of the first space-based far-IR-submm observations of M 33, which measure the emission from the cool dust and resolve the giant molecular cloud complexes. With roughly half-solar abundances, M33 is a first step towards young low-metallicity galaxies where the submm may be able to provide an alternative to CO mapping to measure their H\(_2\) i abundance and with important participation from NASA.


1. Introduction, data, and dust temperature

Understanding star formation requires studying the interplay between the phases of the interstellar medium (ISM). Dust processes most of the energy transiting the ISM, but the cool dust component, although representing the vast majority of the dust mass, is difficult to observe from the ground. Herschel SPIRE observations (Pilbratt 2010; Griffin 2010) are the first space-based 250-500 \(\mu\)m data and as such provide a unique occasion to put together a global picture of the cool gas and dust in M 33. In particular, we compare the morphology of the Far-IR emission and that of the gas as determined from CO and H\(_i\) measurements and attempt to measure how the dust cross-section varies in M 33. A longer term goal is to be able to use the dust emission to constrain the variation of the \(N(H_2)/I_{CO}\) factor within M 33 and elsewhere.

This Letter is one of a series on the HERM33ES project on the ISM of the Local Group galaxy M 33, an overview of which is given in Kramer et al. (2010), hereafter K10. For consistency with the other M 33 papers in this volume, we adopt a distance of \(D = 840\) kpc for M 33 (i.e., 25 arcsec = 100 pc) and orientation parameters of \(PA = 22.5\) degrees and \(i = 56^\circ\). We use the recent HERM33ES SPIRE observations at 250, 350, and 500 \(\mu\)m combined with CO(2-1) observations from the M33CO@IRAM project, as a tracer of the molecular component, and a high-resolution mosaic of VLA HI data (both from Gratier et al. 2010). The SPIRE data were first processed as described in K10 and then convolved from Jy/beam to brightness units (MJy/sr). To estimate the dust temperature, the 250\(\mu\)m data were convolved to the 350\(\mu\)m beamsize (\(\sim 25^\circ\)) and, assuming a single temperature grey body with an emissivity \(\propto \nu^\beta\) with \(\beta = 2\), a temperature was derived from the flux ratio. At the temperatures of the cool component seen in M 33, Dupac et al. (2003) find \(\beta = 2\). To minimize the effect of the uncertainty in \(\beta\), we chose to use adjacent bands to estimate temperature. The 250/350\(\mu\)m ratio provides more accurate temperatures than the 350/500 micron ratio even for temperatures below 10 K. A 15% variation (or uncertainty) in the 250/350\(\mu\)m ratios corresponds to a temperature change of 0.9, 2.1, and 4.1 K at temperatures of 10, 15, and 20 K but the same uncertainty in the 350/500\(\mu\)m ratios yields 1.3, 3.2, and 6.3 K errors for the same dust temperatures. A further advantage is that we obtain the dust temperature at a resolution typical of giant molecular clouds (GMC), \(\sim 100\) pc.

Ideally, a multi-component fit would be used but this requires high S/N data at many wavelengths over the whole disk. By assuming that the cool dust component dominates the emission beyond 250\(\mu\)m, we could calculate the dust temperature out to the optical radius of M 33. The resulting temperature map is shown in Fig. 1. Figure 2 shows a comparison with the temperature of the cool component of the preferred two-component (warm plus cool

* Herschel is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA.
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The two-component model uses β frequencies, the CO lines contribute very little to the contamination was subtracted from the SPIRE data. At these cool dust at these wavelengths. No correction for line contamination was subtracted from the SPIRE data. At these frequencies, the CO lines contribute very little to the continuum flux, unlike at 1.3mm or 850µm (e.g. Braine et al. 1997).

Figure 3 shows the SPIRE 250µm emission with the CO(2–1) emission as contours at 25′′ resolution to show how closely the CO emission follows the dust emission peaks. Gratier et al. (2010) show a detailed comparison of the CO and HI emission with star formation tracers such as Hα and Spitzer 8 and 24µm maps. The 250µm-bright regions are detected in CO and the general morphology of the cool dust emission is seen in the HI image.

2. Dust cross-section

In order to estimate the total gas mass of a galaxy from its dust emission, it is first necessary to measure the dust emission cross-section per H-atom σ. We do this by selecting regions with HI emission and well-constrained dust temperatures but with little or no CO emission (in order to avoid the assumption of a N(H$_2$)/I$_{CO}$ value and CO line ratio). Thereby, we can relate the observed neutral hydrogen column density, N$_{HI}$, to the dust optical depth, τ = σN$_{HI}$ = $S_\nu/B_{\nu,T}$, and infer the cross section σ. In practice, two methods were used: (i) we selected large contiguous regions as shown on the map in Fig. 3 and (ii) all pixels with HI emission, a defined dust temperature, and no detected CO emission (I$_{CO}$ < 0.05 K km s$^{-1}$) were taken, using maps at the resolution of the 350µm SPIRE data. In the optically thin limit, $N_{tot} = N_{HI} + N_{H_2} = S_\nu/(\sigma B_{\nu,T})$ (1)

where $S_\nu$ is the flux density and $N_{tot}$ is the H column density. Where CO is not detected (i.e. $N_{tot} \approx N_{HI}$), we can calculate σ from the dust temperature and the observed HI column density and SPIRE flux by σ $\approx S_\nu/(N_{HI} B_{\nu,T})$.

From the regions indicated in Fig. 3, we estimate σ (using Eq. 1) for different radii and the results are shown in the first line of Table 1. We assume that the cross section within the inner 4 kpc is the same as at 4kpc for two reasons: (i) there are few truly CO-free regions in the central part of M 33 and (ii) Garanz et al. (2007) showed (their Fig. 13) that the H-I-H$_2$ relation was similar within 4kpc but beyond 4 kpc the H$_2$ fraction decreased sharply even at constant total Hydrogen column density, so we consider $R \sim 4kpc$ as a transition in ISM properties. It is very likely
that some molecular gas is present in these regions, particularly those closer to the center, so $\sigma$ is likely overestimated. In the second approach, in order to exclude averaging in undetected molecular clouds, we take the peak of the histogram of $\sigma$ values for the areas without detected CO emission but with HI emission and a constrained dust temperature, yielding the lower values for $\sigma$, as shown in Table 1. Using this method, we find a north-south difference in M33 at equivalent radii. The uncertainty in the histogram method is about ±0.1; although the distribution is sometimes broad, the peak value is well defined. The histograms of $\sigma$ values are clearly different in the north and south. The polygon method was difficult to use in the south due to the absence of large CO-free regions with HI and dust emission. An advantage of the histogram method is that it excludes the tail of high $\sigma$ which may be due to H$_2$ without detected CO emission.

The intrinsic expectation is that because the oxygen abundance is about half solar, the dust cross-section should be as well (Draine et al. 2007), and that given the shallow abundance gradient in M33 (Rosolowsky et al. 2007; Magrini et al. 2009), this should hold for the entire galaxy. However, it quickly became apparent that the dust cross-section per H-atom varied over the galaxy, decreasing with radius by close to a factor 2, with a higher abundance in the south – similar to the variation found by Magrini et al. (2010) from optical HII region and PNe abundance measurements. In the solar neighborhood, $\sigma$ is about $\sigma \approx 1.1 \times 10^{-22}$ cm$^2$ per H-atom (Draine & Lee 1984; Draine & Li 2007) at 250$\mu$m and varying as $\nu^2$. If we apply a $\sigma$ of half this value to M33 using the dust temperature map in Fig. 1, we obtain unrealistically high gas masses in the inner disk: $\sim 40$ $M_\odot$ pc$^{-2}$ on average, whereas the H I is about 11 $M_\odot$ pc$^{-2}$ and the H$_2$ considerably less (Gratier et al. 2010). For a half-solar $\sigma$, the (cool) dust temperature would have to be about 30K, beyond any experimental uncertainties, so we conclude that $\sigma$ must be higher in the inner disk. In the outer disk, however, the H I column density (where no CO emission is present) is roughly equal to what is derived from the dust with $\sigma = 0.5\sigma_0$, suggesting that this value is appropriate for the outer disk.

Table 1. Dust cross-section $\sigma$ at 250$\mu$m as a function of radius in M33, expressed in units of $1.1 \times 10^{-22}$ cm$^2$ per H-atom. The polygons are shown in Fig. 3. The histogram method was applied separately in the north and south (lines 2 and 3), leading to the "Model" values used to estimate the total H column density to make Fig. 4.

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<th>5.5</th>
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</table>

3. Total and molecular gas mass

To really obtain an accurate gas mass from the cold dust emission data, the calibration uncertainties in the SPIRE bands must be reduced because the 15% uncertainty is not sufficient to distinguish between different solutions of $\beta$ and temperature. A second requirement is that the whole disk of M33 must be observed in CO so that a well-determined average of $\sigma$, based on the second technique, can be constructed. The H I data provide a reliable picture of the H I column density, so subtracting this from the total gas column density derived from the dust emission yields an estimate of the molecular gas mass which does not rely on CO emission. A further caveat is that $\sigma$ must not change from the atomic to molecular medium. Once the [C I] and [C II] lines have been measured, these uncertainties can be reduced by the additional constraint that the C/H ratio (from the sum of the C reservoirs) should be equal to the C abundance derived by other means.

Using the information currently available, we have used Table 1 (second method) to build an image of the dust cross-section $\sigma$, erring on the low-$\sigma$ side, as indicated in the last line of Table 1. We then used this $\sigma$, scaled to 500$\mu$m with $\beta = 2$, and the dust temperature (Fig. 1) to estimate the total gas mass from the 500$\mu$m data and then subtracted the H I column density to obtain the H$_2$ column density using the interferometric H I maps from Gratier et al. (2010), which recover more than 90% of the flux found by Putman et al. (2009) using Arecibo. The missing 21cm flux is expected to be located in the mid-to-out disk where the
rotation curve is fairly flat. A morphological comparison of the dust-derived H$_2$ column and the CO contours in Fig. 4 shows that while major regions were not missed by the CO observations, the dust-derived H$_2$ map does not correlate perfectly with the CO. For example, the strong CO peak at (01°34'09".4, +30°49'06") corresponds to a red (not white) region in Fig. 4; this may of course reflect a real variation in the $N$(H$_2$)/$I_{CO}$ ratio.

The data used in the scatter plot in Fig. 5 (see also Fig. 7 in Leroy et al. 2009) yield a formal fit with a $N$(H$_2$)/$I_{CO(2-1)}$ = 3 × 10$^{20}$ H$_2$ per K km s$^{-1}$ value, corresponding to a roughly galactic $N$(H$_2$)/$I_{CO(1-0)}$ = 2.1 × 10$^{20}$ H$_2$ per K km s$^{-1}$ for a line ratio CO(2-1)/1$^{-1}$ = 0.7. The inner disk $N$(H$_2$)/$I_{CO(2-1)}$ is lower and rises to about 4 × 10$^{20}$ H$_2$ per K km s$^{-1}$ beyond 2 kpc. The offset, in principle, implies that some low column density H$_2$ is not seen in CO. Both values depend strongly on the details of the dust cross-section, illustrating the need for a complete map of M 33 in CO.

The total hydrogen mass derived from the dust emission within 8 kpc is estimated to be 1.63 × 10$^9$ M$_\odot$, very similar to the more directly measured H I+H$_2$ masses in Gratier et al. (2010). The dust cross-sections we derive suggest dust-to-gas mass ratios (including He) ranging from ∼125 in the inner 4 kpc to 200 near R$_{25}$ for Milky Way like dust (cf. K10). Comparing the total dust-derived H$_2$ mass with the CO luminosity (Gratier et al. 2010) yields a $N$(H$_2$)/$I_{CO}$ of about 1.5 times the Galactic $N$(H$_2$)/$I_{CO}$. Because the H I dominates, however, a small error in the total dust-derived gas mass translates into a large uncertainty in the dust-derived H$_2$ mass. It is clearly necessary to obtain whole-galaxy measurements of $\sigma$ in CO-free zones to refine these estimates.

**Fig. 4.** H$_2$ column density map estimated from the 500\(\mu\)m SPIRE map, $\beta = 2$, the 250/350\(\mu\)m dust temperature as shown in Fig. 1, the H I data from Gratier et al. (2010), and the cross-sections given in Table 1. The column density scale is shown at the top in units of H$_2$ molecules per cm$^2$. CO contours at 0.4 and 1.6 K km/s (black) and 5 K km/s (blue) are superposed.

**Fig. 5.** Scatter plot of H$_2$ column density as derived from the dust emission and the CO(2-1) intensity. The solid red line is a fit to the entire dataset and corresponds to a $N$(H$_2$)/$I_{CO(2-1)}$ value of 3 × 10$^{20}$ H$_2$ per K km s$^{-1}$ with an offset as described in the text. A radial variation is present: fitting within 2 kpc yields a $N$(H$_2$)/$I_{CO(2-1)}$ value of 2.2 × 10$^{20}$ H$_2$ per K km s$^{-1}$ and 4.1 × 10$^{20}$ H$_2$ per K km s$^{-1}$ beyond 2 kpc (dotted and dashed lines, respectively).

**References**


