Unconditional security proof of long-distance continuous-variable quantum key distribution with discrete modulation

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We present a continuous-variable quantum key distribution protocol combining a discrete modulation and reverse reconciliation. This protocol is proven unconditionally secure and allows the distribution of secret keys over long distances, thanks to a reverse reconciliation scheme efficient at very low signal-to-noise ratio.

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The first practical application of quantum information theory is certainly quantum key distribution (QKD) [1], which allows two distant parties to communicate with absolute privacy, even in the presence of an eavesdropper. Two families of QKD protocols coexist today, relying either on photon counting techniques or homodyne detection, which correspond to discrete and continuous-variable protocols, respectively. The better efficiency of homodyne detection over single photon counting at the telecom wavelength has stimulated the study of continuous-variable protocols in the last few years [2,3]. However, despite its technological advantages, continuous-variable QKD (CVQKD) is still not considered as a true alternative to discrete QKD, mostly because it seems restricted only to short distances. The main reason for that lies in the classical postprocessing of the data shared by Alice and Bob who need to construct a key from continuous random values, which is a task far more complicated than its discrete counterpart.

In this Letter, we introduce a specific CVQKD scheme, which exhibits two specific related advantages: first, it allows us to simplify significantly both the modulation scheme and the key extraction task, and second, it makes it possible to distill secret keys over much longer distances. Continuous-variable protocols have recently been shown to be unconditionally secure, that is, secure against arbitrary attacks [4]. In particular, collective attacks are asymptotically optimal, meaning that the theoretical secret key rate $K$ obtained using one-way (reverse) reconciliation is bounded below by:

$$K \geq I(x;y) - S(y;E) = K_{th},$$

where $x$, $y$ represent the classical data of Alice and Bob, and $E$ is Eve’s quantum state. Here $I(x;y)$ refers to the Shannon mutual information [5] between classical random values $x$ and $y$, and $S(y;E)$ is the quantum mutual information [6] between $y$ and the quantum state $E$. The reason for using two different measures of information is that Eve has no restriction on her capabilities (other than the ones imposed by quantum mechanics), while Alice and Bob must be able to extract a key with current technology. This secret key rate is valid for reverse reconciliation [3]: the final key is extracted from Bob’s data and Bob sends some side information to Alice on the authenticated classical channel to help her correct her errors. In addition, one should note that $K_{th}$ corresponds to a scenario where Alice and Bob could perform perfect error correction, which is never the case in practice. For this reason, the key rate must be modified in the following way [7,8]:

$$K_{real} = \beta I(x;y) - S(y;E),$$

where $\beta$ is the so-called reconciliation efficiency. The term $\beta I(x;y)$ simply corresponds to the amount of information Alice and Bob have been able to extract through reconciliation. The second term, $S(y;E)$, is bounded from the correlation between Alice and Bob’s data, using a Heisenberg-type inequality.

Whereas the reconciliation efficiency is not usually taken into account to estimate asymptotic bounds, we must include it in our analysis because it is currently the limiting factor for the range of CVQKD with Gaussian modulation. In [8], it was argued that working at low signal-to-noise ratio (SNR) increases the range of the protocol. Unfortunately, maintaining a good reconciliation efficiency at very low SNR is even more difficult to achieve. This point is exactly the limitation that the protocol presented in this Letter manages to overcome, hence allowing QKD over longer distances.

This Letter is organized as follows: after detailing the limitations of the Gaussian modulation, we present our new four-state protocol as well as its unconditional security proof. Then we describe the reconciliation step and show that its efficiency remains remarkably high, even at very low SNR. Finally, we show the expected performances of the protocol and discuss some perspectives.
Gaussian vs discrete modulation.—Most CVQKD protocols use a Gaussian modulation since it is the one maximizing the mutual information between Alice and Bob over a Gaussian channel. In such a protocol, Alice draws two random values \(q_A, p_A\) with a Gaussian distribution \(\mathcal{N}(0, V_A)\) and sends a coherent state \(|q_A + ip_A\rangle\) to Bob. The main problem of this modulation arises when one wants to perform QKD over long distances. In this case, there are two possibilities to fight the noise induced by the losses in the channel: either increase the variance of the modulation so that the SNR remains reasonably high, or work at low SNR. Unfortunately both approaches tend to fail over a few tens of kilometers.

Working at high SNR requires us to achieve a very good reconciliation efficiency, otherwise the secret key rate goes to zero \([7,8]\). Capacity-achieving error correcting codes are therefore required for this task. Unfortunately, even with the best codes presently available, such as low-density parity-check (LDPC) codes \([9]\) or turbo codes \([10]\), one cannot expect to extend the range of the protocol well over 30 kilometers \([7]\).

Working at low SNR relieves a little bit the need for capacity-achieving codes, but reasonably good low-rate codes are still hard to combine with the Gaussian modulation. Some interesting algebraic properties of \(\mathbb{R}^d\) can be useful in this situation, and help with increasing the achievable distance to over 50 kilometers \([8]\).

At the present time both of these approaches seem to have been pushed at their maximum using the state-of-the-art channel coding techniques, and breaking this 50 kilometers limit seems unlikely with a Gaussian modulation.

One should emphasize the following point: while a Gaussian modulation performs much better than a binary modulation at high SNR (simply because a binary modulation cannot send more than one bit of information per signal), it is not the case for low SNR. Adding to this fact that a binary modulation allows for a much better reconciliation efficiency at low SNR, we infer that the modulation required to achieve long distances is not Gaussian. Examples of binary (or quaternary depending on the number of quadratures considered) modulation have been proposed in the past \([11,12]\), but were often combined with a postselection procedure \([13]\), and are not known to be unconditionally secure.

The four-state protocol.—The protocol we propose runs as follows. Alice sends randomly one of the four coherent states: \(|ae^{i(2k+1)\pi/4}\rangle\) with \(k \in \{0, 1, 2, 3\}\). The amplitude \(a\) (taken as a real number) is chosen so as to maximize the secret key rate one can expect from the expected experimental parameters (transmission of the line and excess noise). Bob measures randomly one of the quadratures in the case of the homodyne protocol \([14]\) and gets the result \(y\). The sign of \(y\) encodes the bit of the raw key while Bob reveals the absolute value \(|y|\) to Alice through the classical authenticated (but not secure) channel. At this point, Alice and Bob share correlated strings of bits. In order to help Alice correct her data, Bob sends some side information over the classical channel, typically the syndrome of his string relative to a binary code they agreed on beforehand. From a classical communication perspective, the error correction (reconciliation) is then a problem of channel coding for the so-called BIAWGN channel, where a binary modulation is sent over an Additive White Gaussian Noise channel, and for which there exist very good codes, even for extremely low SNR.

The present protocol can thus be seen as a hybrid between the Gaussian modulation protocol, with which it shares the physical implementation as well as the security proofs based on the optimality of Gaussian states, and protocols combining a discrete modulation with postselection, for which the error correction is substantially easier to perform, but whose unconditional security has not yet been established.

Let us now prove that the four-state protocol is unconditionally secure. First, it is enough to prove the security against collective attacks as they are the most powerful attacks in the asymptotic limit \([4]\). Then, as usual, the security is established by considering the equivalent entanglement-based version of the protocol. The state sent to Bob in the prepare and measure scheme is a mixture of four coherent states: \(\rho = \frac{1}{4} \sum_{k=0}^{3} |\alpha_k\rangle \langle \alpha_k|\) with \(\alpha_k = \alpha \exp(i(2k+1)\pi/4)\). The entanglement-based version uses a purification \(|\Phi\rangle\) of this state such that: \(\rho = \tr_s(\langle \Phi |\Phi\rangle)\). This state \(\rho\) can be diagonalized as

\[
\rho = \lambda_0 |\Phi_0\rangle \langle \Phi_0| + \lambda_1 |\Phi_1\rangle \langle \Phi_1| + \lambda_2 |\Phi_2\rangle \langle \Phi_2| + \lambda_3 |\Phi_3\rangle \langle \Phi_3|
\]

where \(\lambda_{0,2} = \frac{1}{4} e^{-\alpha^2} [\cosh(\alpha^2) \pm \cos(\alpha^2)]\), \(\lambda_{1,3} = \frac{1}{4} e^{-\alpha^2} [\sinh(\alpha^2) \pm \sin(\alpha^2)]\) and

\[
|\phi_k\rangle = \frac{e^{-\alpha^2/2}}{\sqrt{\lambda_k}} \sum_{n=0}^{\infty} \frac{\alpha^{n+k}}{(4n+k)!} (-1)^n |4n+k\rangle
\]

for \(k \in \{0, 1, 2, 3\}\). Therefore, a particular purification of \(\rho\) obtained by the Schmidt decomposition is \(|\Phi\rangle = \frac{1}{2} \sum_{k=0}^{3} \sqrt{\lambda_k} |\phi_k\rangle \langle \phi_k|\) which can be rewritten as \(|\Phi\rangle = \frac{1}{2} \sum_{k=0}^{3} \sum_{m=0}^{\infty} |\psi_k\rangle \langle \psi_k|\) where the states

\[
|\psi_k\rangle = \frac{1}{2} \sum_{m=0}^{\infty} e^{-i(1+2k)m(\pi/4)} |\phi_m\rangle
\]

are orthogonal non-Gaussian states.

The entanglement-based version of the four-state protocol can be described as follows. Alice prepares the entangled state \(|\Phi\rangle\) and performs the projective measurement \(|\psi_0\rangle \langle \psi_0|, |\psi_1\rangle \langle \psi_1|, |\psi_2\rangle \langle \psi_2|, |\psi_3\rangle \langle \psi_3|\) on her half thus preparing the coherent state \(|\alpha_k\rangle\) when her measurement gives the result \(k\). This state is sent through the quantum channel to Bob who measures either one of the quadratures with a homodyne detection.

In order to prove the security of the protocol, we use the fact that \(S(y;E)\), the Holevo information between Eve and Bob’s classical variable, is maximized when then the state \(\rho_{AB}\) shared by Alice and Bob is Gaussian \([15]\). Therefore, \(S(y;E)\) can be bounded from above by a function of the
covariance matrix $\Gamma$ of $\rho_{AB}$ [7]. For a quantum channel characterized by its transmission $T$ and excess noise $\xi$, $\Gamma$ is given by:

$$
\Gamma = \begin{pmatrix}
(V_A + 1)I_2 & \sqrt{T}Z\sigma_z \\
\sqrt{T}Z\sigma_z & (TV_A + 1 + T\xi)I_2
\end{pmatrix}
$$

where $V_A$ is the variance of Alice’s modulation in the prepare and measure scheme. This covariance matrix has the same form as in the Gaussian modulation scheme where $Z$ would be replaced by the correlation of a two-mode squeezed vacuum $Z_{EPR} = \sqrt{V_A + 2V_A}$. The correlation $Z$ for the state $|\Phi\rangle$ does not take such a simple mathematical form but turns out to be almost equal to $Z_{EPR}$ for small variances (see Fig. 1). Hence, for a sufficiently low modulation variance, the bound on $S(y:E)$ is almost identical to the one obtained for a Gaussian modulation. However, the efficient reconciliation $\beta$ available at low SNR (see Fig. 1) allows us to extract more information, and in fine, to distill a key in conditions where the Gaussian modulation protocol is ineffective.

**Realistic reconciliation.**—The main advantage of a binary modulation compared to a Gaussian modulation is that one can find binary codes allowing high reconciliation efficiency, e.g., 80%, even with a SNR close to 0. This is quite remarkable since all practical reconciliation schemes for a Gaussian modulation [8,18] see their efficiency drop to zero as the SNR becomes too low (see Fig. 1). In order to achieve an efficient reconciliation at low SNR, one needs good low-rate codes. These can be constructed rather easily with a concatenation of a capacity-achieving code and a repetition code that we describe now.

At the end of the quantum exchange, Alice and Bob share two correlated vectors $x = (x_1, \ldots, x_N)$ (with $x_i = \pm x/\sqrt{2}$) and $y = (y_1, \ldots, y_N)$. We will use the concatenation of a capacity-achieving code $C$ of length $m$ and a repetition code of length $k$, assuming that $N = mk$. Bob starts by defining the vector $Y = (Y_1, \ldots, Y_m)$ where $Y_i = \text{sgn}(y_{k(i-1)+i})$ for $i \in \{1, \ldots, m\}$. The goal of the reconciliation is for Alice to be able to compute the vector $Y$. To do this, Bob sends some side information: the vector $\{y_1, \ldots, y_N\}$, the $m$ vectors $\{(1, \text{sgn}(y_{k(i-1)+1} \times y_{k(i-1)+2}), \ldots, \text{sgn}(y_{k(i-1)+1} \times y_{k(i-1)+k})\}$, and the syndrome of $Y$ for the code $C$. This scheme allows Alice and Bob to extract $m$ bits out of their $N = km$ data.

This repetition scheme is a simple way to build a good code of rate $R/k$ out of a code of rate $R$. This construction is not optimal compared to using a very good error correcting code at the considered signal-to-noise ratio but exhibits some interesting features. First, designing very good codes at low SNR is not easy, and has not been intensively studied so far, mainly because the telecom industry does not operate in this regime: this would not be economical since an important number of physical signals would be required to send one information bit. The problem is very different in QKD, where quantum noise is an advantage rather than a drawback. A second advantage of this repetition scheme lies in its simplicity. As we mentioned earlier, the main bottleneck of CVQKD is the reconciliation: it was limiting both the range and the rate of the protocol. In particular, the rate is limited by the complexity of decoding LDPC codes, which is roughly proportional to the size of the code considered (in fact $O(N\log N)$). If one uses a repetition scheme of parameter $k$, then the length of the LDPC code becomes $m = N/k$ allowing a speedup of a factor $k$. The speed of the reconciliation is not proportional to the number of signals exchanged by Alice and Bob anymore, but to the mutual information they share, which is a major improvement for noisy channels, i.e., long distance. Finally, the penalty in terms of reconciliation efficiency imposed by using this scheme instead of a dedicated low-rate error correcting code is actually quite small. Roughly speaking, a repetition code of length $k$ allows us to decode at a SNR $k$ times smaller. It is indeed easy to show that the efficiency $\beta_R(s/k)$ obtained at a SNR $s/k$ with such a repetition code is related to the efficiency $\beta_{LDPC}(s)$ available at SNR $s$ through $\beta_R(s/k) = \frac{\log(1+s/k)}{\log(1+s)} \beta_{LDPC}(s)$, that is, $\beta_R(s/k) = (1-\frac{s}{k})\beta_{LDPC}(s)$ when $s$ is small enough. For instance, there exist good LDPC codes of rate $1/10$ decoding at SNR 0.17 [19], meaning that $\beta_{LDPC}(0.17) \approx 88\%$ and $\forall k \geq 1$, $\beta_R(0.17/k) \geq 80\%$. By using different codes, one can have a reconciliation efficiency greater than 80% for all SNRs below 1.

The reconciliation scheme presented above performs indeed much better at low SNR (lower than 1) than reconciliation schemes used for a Gaussian modulation. This behavior is inverted for higher SNR as a binary modulation is unable to send more that one bit per channel use. As a consequence, the four-state protocol is particularly relevant in a long distance scenario, whereas the Gaussian modulation protocol is still better suited to distribute high key rate at short distances.

**FIG. 1** (color online). Left: correlation $Z_{EPR}$ of an EPR pair (solid line) and correlation $Z$ of state $|\Phi\rangle$ (dashed line) as a function of the modulation variance. Right: practical reconciliation efficiency for a binary modulation (dashed line) and for a Gaussian modulation (solid line) [18].
The quantum channel is characterized by its transmission. The performances of the four-state protocol are displayed in Fig. 2. The excess noise in excess compared to the shot noise. It should be noted that these performances are comparable with discrete-variable protocols, and are much better than previous CVQKD schemes.

Whereas Alice usually sends coherent states with a few photons per pulse (between 3 and 10) in the Gaussian modulation protocol, here, the optimal number of photons per pulse typically ranges from 0.2 to 1. Therefore, the similarities with discrete-variable QKD are important: the information is encoded onto low amplitude coherent states with generally less than 1 photon per pulse. The main difference is that homodyne detection replaces photon counting. In our protocol, however, the error rate is not upper bounded (and can be as close as 0.5 as the reconciliation efficiency allows). This sounds in disagreement with the vacuum noise is processed in two very different ways: whereas it creates “deletion errors” (which are ignored) in the photon counting scheme, it produces “real errors” (which have to be corrected) in the continuous-variable scheme. But in both cases, these errors due to vacuum noise cannot be exploited by anybody, neither by the legitimate parties, nor by Eve.

As a conclusion, we presented a new unconditionally secure continuous-variable QKD protocol based on a discrete modulation. The use of good error correcting codes at low SNR allows us to achieve long distances, which was impossible with a Gaussian modulation. Further work will include analysis of the finite-key effects [20], as well as the implementation of the present protocol.

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Results and perspectives.—The theoretical performances of the four-state protocol are displayed in Fig. 2. The quantum channel is characterized by its transmission $T = \eta 10^{-0.02d}$ where $\eta$ is the quantum efficiency of the homodyne detection and $d$ is the distance between Alice and Bob, and its excess noise $\xi$, that is the noise in excess compared to the shot noise. It should be noted that these performances are comparable with discrete-variable protocols, and are much better than previous CVQKD schemes.

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FIG. 2 (color online). Secret key rate as a function of the distance for different values of the excess noise: from top to bottom, $\xi = 0.002, 0.004, 0.006, 0.008, 0.01$. The quantum efficiency of Bob’s detection is $\eta = 0.6$. 

[14] Alternatively, Bob can measure simultaneously both quadratures with a heterodyne detection, and the security proof will proceed along the same lines.
[15] This can be proved by applying the reasoning presented in [16] to the quantity $S(y:E)$ instead of $K_{th} = I(x:y) - S(y:E)$, see also [17].