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Synchronization of chaotic fractional-order systems via linear control

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Abstract The chaotic dynamics of fractional-order systems has attracted much attention recently. Chaotic synchronization of fractional-order systems is further studied in this paper. We investigate the chaos synchronization of two identical systems via a suitable linear controller applied to the response system. Based on the stability results of linear fractional-order systems, sufficient conditions for chaos synchronization of these systems are given. Control laws are derived analytically to achieve synchronization of the chaotic fractional-order Chen, Rössler and modified Chua systems. Numerical simulations are provided to verify the theoretical analysis.

Keywords: Chaos synchronization; Fractional-order system; Caputo fractional derivative; Mittag-Leffler function; Stability; Linear control.

1 Introduction

Chaos, as an important topic in nonlinear science, has been investigated and studied in mathematical and physical communities in the last few decades and the research efforts have devoted to chaos control and chaos synchronization problems in many dynamical systems. During the last two decades, synchronization in chaotic dynamical systems has attracted a great deal of attention [Yamada & Fujisaka, 1983; Pecora & Carrol, 1990; Ott *et al.*, 1990; Chen & Dong, 1998; Derivière & Aziz-Alaoui, 2003; Aziz-Alaoui, 2006] since the seminal papers by Yamada and Fujisaka [1983] and Pecora and Carrol [1990]. A way to study synchronization is to use a controller to make the output of the slave (reponse) system copy in some manner the master (drive) system one.

On the other hand, the development of models based on fractional-order differential systems has recently gained popularity in the investigation of dynamical systems. Fractional derivatives provide an excellent instrument to describe memory and hereditary properties of various materials and processes. The main reason for using the integer-order models was the absence of solution methods for fractional differential equations. The advantages or the real objects of the fractional-order systems [Petráš, 2008] are that we have more degrees of freedom in the model and that a “memory” is included in the model (fractional-order systems have an unlimited memory).

Recently, studying fractional-order systems has become an active research area. The chaotic dynamics of fractional-order systems began to attract much attention in recent years [Podlubny, 1999; Hil-

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fer, 2000]. It has been shown that fractional-order systems, as generalizations of many well-known systems, can also behave chaotically, such as the fractional Duffing system [Ge & Ou, 2007], the fractional Chua system [Petráš, 2008; Hartley *et al.*, 1995], the fractional Rössler system [Li & Chen, 2004a], the fractional Chen system [Li & Peng, 2004; Li & Chen, 2004b; Lu & Chen, 2006], the fractional Lorenz system [Grigorenko & Grigorenko, 2003], the fractional Arneodo system [Arneodo *et al.*, 1985] and the fractional Lü system [Deng & Li, 2005]. In [Hartley *et al.*, 1995; Li & Chen, 2004a; Li & Peng, 2004; Li & Chen, 2004b] it has been shown that some fractional-order systems can produce chaotic attractors with order less than 3.

Recent studies show that chaotic fractional-order systems can also be synchronized [Lu & Chen, 2006; Deng & Li, 2005; Li *et al.*, 2003; Li & Zhou, 2005; Zhou *et al.*, 2008; Peng, 2007]. In many literatures, synchronization among fractional-order systems is only investigated through numerical simulations. A simple method to synchronize chaotic fractional-order systems based on the stability criteria of linear differential systems, see [Matignon, 1996; Deng *et al.*, 2007], is presented in [Sheu *et al.*, 2007; Yan & Li, 2007; Li & Yan, 2007]. Numerical algorithms for chaos synchronization of fractional-order systems based on Laplace transform theory are presented in [Wang *et al.*, 2006; Li *et al.*, 2006; Zhu *et al.*, in press; Wu *et al.*, 2008; Yu & Li, 2008].

In this paper, we recall and present some stability results for linear fractional-order systems. Then, based on these results, we study the chaos synchronization of fractional-order systems. Using the drive-response concept, a linear feedback control law is derived to achieve synchronization of the two chaotic fractional-order systems. The synchronization controllers are investigated theoretically and numerical simulations are then presented in the case of fractional-order Chen, Rössler and modified Chua systems.

2 Preliminaries and definitions

There are several definitions of a fractional derivative of order $\alpha > 0$ [Podlubny, 1999; Caputo, 1967].

The two most commonly used are the Riemann-Liouville and Caputo definitions. Each definition uses Riemann-Liouville fractional integration and derivatives of whole order. The difference between the two definitions appears in the order of evaluation. The Riemann-Liouville fractional integral operator of order $\alpha \geq 0$ of the function $f(t)$ is defined as,

$$J^\alpha f(t) = \frac{1}{\Gamma(\alpha)} \int_0^t (t - \tau)^{\alpha-1} f(\tau) d\tau, \quad t > 0. \quad (1)$$

Some properties of the operator J^α can be found, for example, in [Podlubny, 1999; Gorenflo & Mainardi, 1997; Oldham & Spanier, 1974]. We recall only the two following properties. For $\mu \geq -1$, $\alpha, \beta \geq 0$ and $\gamma > -1$, we have,

$$J^\alpha J^\beta f(t) = J^{\alpha+\beta} f(t),$$

$$J^\alpha t^\gamma = \frac{\Gamma(\gamma+1)}{\Gamma(\alpha+\gamma+1)} t^{\alpha+\gamma}.$$

In this study, Caputo definition is used and the fractional derivative of $f(t)$ is defined as, $J^{m-\alpha} D^m f(t)$,

$$D^\alpha f(t) = \frac{1}{\Gamma(m-\alpha)} \int_0^t \frac{f^{(m)}(\tau)}{(t-\tau)^{\alpha-m+1}} d\tau, \quad (2)$$

for $m-1 < \alpha \leq m$, $m \in \mathbb{N}$, $t > 0$.

Caputo definition has the advantage of dealing properly with initial value problems in which the initial conditions are given in terms of the field variables with their integer order which is the case in most physical processes. Fortunately, the Laplace transform of the Caputo fractional derivative satisfies,

$$L\{D^\alpha f(t)\} = s^\alpha L\{f(t)\} - \sum_{k=0}^{m-1} f^{(k)}(0^+) s^{\alpha-1-k}, \quad (3)$$

where $m-1 < \alpha \leq m$.

The Laplace transform of Caputo fractional derivative requires the knowledge of the (bounded) initial values of the function and of its integer derivatives of order $k = 1, 2, \dots, m-1$. It has been found that the solution of many fractional-order differential equations can be written in terms of Mittag-Leffler functions, which are relevant for their connection with fractional calculus. The one-parameter Mittag-Leffler function $E_\alpha(z)$ is defined by [Podlubny, 1999],

$$E_\alpha(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + 1)}, \quad \alpha > 0, \quad z \in \mathbb{C}, \quad (4)$$

which is an entire function. The Laplace transform for $E_\alpha(z)$ is given by [Podlubny, 1999],

$$L\{E_\alpha(-\lambda t^\alpha)\} = \frac{s^{\alpha-1}}{s^\alpha + \lambda}. \quad (5)$$

The Mittag-Leffler function provides a simple generalization of the exponential function. For studying the stability of fractional order systems, we need the final value theorem.

Theorem 1 *Let $F(s)$ be the Laplace transform of function $f(t)$. If all poles of $sF(s)$ are in the open left half plane, then,*

$$\lim_{t \rightarrow \infty} f(t) = \lim_{s \rightarrow 0} sF(s). \quad (6)$$

The final value theorem gives information about the asymptotic behavior of $f(t)$ (how $f(t)$ behaves as $t \rightarrow +\infty$) directly from $F(s)$. This result plays an important role in the derivation of the linear feedback control laws to achieve synchronization of the chaotic fractional-order systems.

3 Stability results of linear fractional systems

Stability of linear fractional systems, which is of main interest in control theory, has been thoroughly investigated where necessary and sufficient conditions have been derived, see [Matignon, 1996; Deng *et al.*, 2007] or references therein. In [Matignon, 1996], Matignon introduced the stability properties for some linear fractional-order systems. In [Deng *et al.*, 2007], Deng *et al.* studied the stability of n -dimensional linear fractional differential equations with time delays. In this section, we recall the main stability properties and then, we present a new property for a general class of fractional-order systems. For this object, we consider the following linear system of fractional differential equations,

$$\frac{d^\alpha}{dt^\alpha} \mathbf{x}(t) = A\mathbf{x}(t), \quad \mathbf{x}(0) = \mathbf{x}_0, \quad (7)$$

where $\mathbf{x} \in \mathbb{R}^n$, the matrix $A \in \mathbb{R}^n \times \mathbb{R}^n$, $\alpha = [\alpha_1, \alpha_2, \dots, \alpha_n]$ indicates the fractional orders, $\frac{d^\alpha}{dt^\alpha} = [\frac{d^{\alpha_1}}{dt^{\alpha_1}}, \frac{d^{\alpha_2}}{dt^{\alpha_2}}, \dots, \frac{d^{\alpha_n}}{dt^{\alpha_n}}]$ and $\frac{d^{\alpha_i}}{dt^{\alpha_i}}$ is the Caputo fractional derivative of order α_i , where $0 < \alpha_i \leq 1$ for $i = 1, 2, \dots, n$. If $\alpha = \alpha_1 = \alpha_2 = \dots = \alpha_n$, then the stability of the fractional-order system (7) has been studied in [Matignon, 1996], where necessary and sufficient conditions have been introduced. The following stability result was derived in [Matignon, 1996].

Theorem 2 *The autonomous fractional-order system (7), when $\alpha = \alpha_1 = \alpha_2 = \dots = \alpha_n$, is asymptotically stable iff $|\arg(\text{spec}(A))| > \alpha\pi/2$. In this case the components of the state decay towards 0 like $t^{-\alpha}$.*

In case of $\alpha = 1$, the stability occurs if no poles of the linear system (7) lie in the closed right half plane of the Laplace plane which agrees with well-known results for ordinary linear differential systems.

If $\alpha_1, \alpha_2, \dots, \alpha_n$ are rational positive numbers, then we have the following result, which is introduced in [Deng *et al.*, 2007].

Theorem 3 *Suppose that α_i 's are rational numbers between 0 and 1, for $i = 1, 2, \dots, n$. Let $\gamma = 1/m$ where m is the least common multiple of the denominators m_i of α_i 's, where $\alpha_i = k_i/m_i$, $k_i, m_i \in \mathbb{N}$, $i = 1, 2, \dots, n$. Then system (7) is asymptotically stable if all roots λ of the equation $\det(\text{diag}(\lambda^{m\alpha_1}, \lambda^{m\alpha_2}, \dots, \lambda^{m\alpha_n}) - A) = 0$ satisfy $|\arg(\lambda)| > \gamma\pi/2$.*

Now, if $\alpha_1, \alpha_2, \dots, \alpha_n$ are any positive numbers between 0 and 1, not necessary equals or rational numbers, then we have the following new result.

Theorem 4 *The autonomous fractional-order system (7) is asymptotically stable if $|\arg(\text{spec}(A))| > \pi/2$, that is $a_{11}, a_{22}, \dots, a_{nn} < 0$, when A is an upper or lower triangular matrix.*

Proof. For simplicity, we give the proof when A is an 3×3 lower triangular matrix. A similar way can be followed to prove the general theorem. Consider the fractional-order system,

$$\begin{pmatrix} \frac{d^{\alpha_1} x}{dt^{\alpha_1}} \\ \frac{d^{\alpha_2} y}{dt^{\alpha_2}} \\ \frac{d^{\alpha_3} z}{dt^{\alpha_3}} \end{pmatrix} = \begin{pmatrix} a_{11} & 0 & 0 \\ a_{21} & a_{22} & 0 \\ a_{31} & a_{32} & a_{33} \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix} \quad (8)$$

Applying Laplace transform to system (8), letting $X(s) = L(x(t))$, $Y(s) = L(y(t))$, $Z(s) = L(z(t))$ and using formula (3), we obtain,

$$\begin{cases} s^{\alpha_1} X(s) - s^{\alpha_1-1} x(0) = a_{11} X(s) \\ s^{\alpha_2} Y(s) - s^{\alpha_2-1} y(0) = a_{21} X(s) + a_{22} Y(s) \\ s^{\alpha_3} Z(s) - s^{\alpha_3-1} z(0) = a_{31} X(s) + a_{32} Y(s) \\ \quad + a_{33} Z(s) \end{cases} \quad (9)$$

It follows from Eqs. (9) that,

$$\begin{cases} X(s) = \frac{s^{\alpha_1-1} x(0)}{s^{\alpha_1} - a_{11}} \\ Y(s) = \frac{s^{\alpha_2-1} y(0)}{s^{\alpha_2} - a_{22}} + \frac{a_{21} s^{\alpha_1-1} x(0)}{(s^{\alpha_1} - a_{11})(s^{\alpha_2} - a_{22})} \\ Z(s) = \frac{s^{\alpha_3-1} z(0)}{s^{\alpha_3} - a_{33}} + \frac{a_{31} s^{\alpha_1-1} x(0)}{(s^{\alpha_1} - a_{11})(s^{\alpha_3} - a_{33})} \\ \quad + \frac{a_{32} s^{\alpha_2-1} y(0)}{(s^{\alpha_2} - a_{22})(s^{\alpha_3} - a_{33})} \\ \quad + \frac{a_{21} a_{32} s^{\alpha_1-1} x(0)}{(s^{\alpha_1} - a_{11})(s^{\alpha_2} - a_{22})(s^{\alpha_3} - a_{33})} \end{cases} \quad (10)$$

In particular, using (5), we have $x(t) = x(0) E_{\alpha_1}(a_{11} t^{\alpha_1})$, where $E_{\alpha}(t)$ is the Mittag-Leffler function. From the properties of this function, see [Gorenflo & Mainardi, 1997], we can conclude that $x(t) \rightarrow 0$ as $t \rightarrow +\infty$ if $a_{11} < 0$. In view of (10), all poles of $sX(s)$, $sY(s)$ and $sZ(s)$ lie in the open left half plane if $a_{11}, a_{22}, a_{33} < 0$. Hence, using the final value theorem we get,

$$\begin{cases} \lim_{t \rightarrow \infty} x(t) = \lim_{s \rightarrow 0} sX(s) = 0, \\ \lim_{t \rightarrow \infty} y(t) = \lim_{s \rightarrow 0} sY(s) = 0, \\ \lim_{t \rightarrow \infty} z(t) = \lim_{s \rightarrow 0} sZ(s) = 0. \end{cases} \quad (11)$$

As a result, the states $x(t), y(t), z(t)$ decay towards 0, as $t \rightarrow +\infty$, if $a_{11}, a_{22}, a_{33} < 0$. Therefore, in this case, system (8) is asymptotically stable.

4 Synchronization of fractional-order systems

In this section, based on the stability results of fractional differential equations, we briefly discuss the is-

sue of controlling fractional-order chaotic systems to realize synchronization with linear error feedback control. In order to observe synchronization behavior, we construct the master (drive) system and the slave (response) system as,

$$\text{Master} \begin{cases} \frac{d^{\alpha_1} x_m}{dt^{\alpha_1}} = b_{11} x_m + b_{12} y_m + b_{13} z_m \\ \quad + g_1(x_m, y_m, z_m, t) \\ \frac{d^{\alpha_2} y_m}{dt^{\alpha_2}} = b_{21} x_m + b_{22} y_m + b_{23} z_m \\ \quad + g_2(x_m, y_m, z_m, t) \\ \frac{d^{\alpha_3} z_m}{dt^{\alpha_3}} = b_{31} x_m + b_{32} y_m + b_{33} z_m \\ \quad + g_3(x_m, y_m, z_m, t) \end{cases} \quad (12)$$

$$\text{Slave} \begin{cases} \frac{d^{\alpha_1} x_s}{dt^{\alpha_1}} = b_{11} x_s + b_{12} y_s + b_{13} z_s \\ \quad + g_1(x_m, y_m, z_m, t) + u_1(t) \\ \frac{d^{\alpha_2} y_s}{dt^{\alpha_2}} = b_{21} x_s + b_{22} y_s + b_{23} z_s \\ \quad + g_2(x_m, y_m, z_m, t) + u_2(t) \\ \frac{d^{\alpha_3} z_s}{dt^{\alpha_3}} = b_{31} x_s + b_{32} y_s + b_{33} z_s \\ \quad + g_3(x_m, y_m, z_m, t) + u_3(t) \end{cases} \quad (13)$$

where $\frac{d^{\alpha_i}}{dt^{\alpha_i}}$ is the fractional differential operator in Caputo sense, $0 < \alpha_i \leq 1$, g_i is a nonlinear function and b_{ij} are real constants, for $i, j = 1, 2, 3$. Subscripts m and s stand for the master system and slave system, respectively, $\mathbf{u}(t) = [u_1(t), u_2(t), u_3(t)]^T$ is the linear controller to be designed such that these two chaotic systems synchronize. Defining the synchronization error as,

$$e_1(t) = x_s(t) - x_m(t)$$

$$e_2(t) = y_s(t) - y_m(t)$$

$$e_3(t) = z_s(t) - z_m(t)$$

we get the error system,

$$\begin{cases} \frac{d^{\alpha_1} e_1}{dt^{\alpha_1}} = b_{11}e_1 + b_{12}e_2 + b_{13}e_3 + u_1(t) \\ \frac{d^{\alpha_2} e_2}{dt^{\alpha_2}} = b_{21}e_1 + b_{22}e_2 + b_{23}e_3 + u_2(t) \\ \frac{d^{\alpha_3} e_3}{dt^{\alpha_3}} = b_{31}e_1 + b_{32}e_2 + b_{33}e_3 + u_3(t) \end{cases} \quad (14)$$

Our aim is to determine the controller $\mathbf{u}(t)$, given below, such that the drive system (12) and the response system (13) are synchronized ($\|\mathbf{e}(t)\| \rightarrow 0$, as $t \rightarrow +\infty$).

$$\begin{aligned} \mathbf{u}(t) &= \begin{pmatrix} u_1(t) \\ u_2(t) \\ u_3(t) \end{pmatrix} \\ &= \begin{pmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{pmatrix} \cdot \begin{pmatrix} e_1(t) \\ e_2(t) \\ e_3(t) \end{pmatrix} \end{aligned} \quad (15)$$

where a_{ij} are real constants. For this purpose, we design the controller $\mathbf{u}(t) = [u_1(t), u_2(t), u_3(t)]^T$ such that the solutions of the error system $e_i(t)$ go to 0, $i = 1, 2, 3$, as t goes to $+\infty$. In this paper, three cases are studied to design the controller $\mathbf{u}(t)$.

Case 1. If $\alpha = \alpha_1 = \alpha_2 = \alpha_3$ then, using Theorem 2, the drive system (12) and the response system (13) are synchronized when $|\arg(\text{spec}(A_u))| > \alpha\pi/2$, where,

$$A_u = \begin{pmatrix} a_{11} + b_{11} & a_{12} + b_{12} & a_{13} + b_{13} \\ a_{21} + b_{21} & a_{22} + b_{22} & a_{23} + b_{23} \\ a_{31} + b_{31} & a_{32} + b_{32} & a_{33} + b_{33} \end{pmatrix} \quad (16)$$

In this case, we choose the controller $\mathbf{u}(t)$ such that $|\arg(\lambda)| > \alpha\pi/2$ for every eigenvalue λ of the matrix A_u . Of course, if all eigenvalues of A_u are negative real numbers then both systems achieve synchronization.

Case 2. If α_i 's are rational numbers such that $\alpha_i = k_i/m_i$, $\gamma = 1/m$ where m is the least common multiple of the denominators m_i of α_i 's, $k_i, m_i \in \mathbb{N}$,

for $i = 1, 2, 3$, then, using Theorem 3, the drive system (12) and the response system (13) are synchronized when all roots λ of the characteristic equation $\det(\text{diag}(\lambda^{m\alpha_1}, \lambda^{m\alpha_2}, \lambda^{m\alpha_3}) - A_u) = 0$ satisfy $|\arg(\lambda)| > \gamma\pi/2$. In this case, we choose the controller $\mathbf{u}(t)$ such that all the roots of the equation $\det(\text{diag}(\lambda^{r_1}, \lambda^{r_2}, \lambda^{r_3}) - A_u) = 0$ satisfy $|\arg(\lambda)| > \gamma\pi/2$, where $r_i = mk_i/m_i$.

Case 3. In particular, if $\alpha_1, \alpha_2, \alpha_3$ are any positive real numbers between 0 and 1 then, using Theorem 4, the drive system (12) and the response system (13) are synchronized when A_u is an upper or lower triangular matrix and all eigenvalues of A_u are negative real numbers. In this case we choose the controller $\mathbf{u}(t)$ such that A_u becomes an upper or lower triangular matrix and all roots of the characteristic equation $(\lambda - a_{11} - b_{11})(\lambda - a_{22} - b_{22})(\lambda - a_{33} - b_{33}) = 0$ are negative real numbers.

To demonstrate these techniques, some examples of synchronization for two identical fractional-order chaotic systems are discussed in the following sections.

5 Synchronization of fractional-order Chen systems

Here, we consider Chen system, which has been found by Chen and Ueta [1999] in 1999. Its fractional version reads as,

$$\begin{cases} \frac{d^{\alpha_1} x}{dt^{\alpha_1}} = a(y - x) \\ \frac{d^{\alpha_2} y}{dt^{\alpha_2}} = (c - a)x - xz + cy \\ \frac{d^{\alpha_3} z}{dt^{\alpha_3}} = xy - bz \end{cases} \quad (17)$$

where $\alpha = (\alpha_1, \alpha_2, \alpha_3)$ is subject to $0 < \alpha_1, \alpha_2, \alpha_3 \leq 1$. When $\alpha = (1, 1, 1)$, system (17) is the original integer order Chen system, which is chaotic, for example, when $(a, b, c) = (35, 3, 28)$, see [Chen & Ueta, 1999]. Simulations are performed to obtain chaotic behavior of the fractional-order Chen system for different fractional orders α . Simulation results demonstrate that chaos indeed exists in the fractional-order Chen system with order less than

3 (i.e. $\alpha_1 + \alpha_2 + \alpha_3 < 3$). For example, chaotic attractors are found in [Yan & Li, 2007] when $\alpha = (0.95, 0.95, 0.95)$ and $(a, b, c) = (35, 3, 28)$. In [Li & Chen, 2004b] chaotic behaviors are found when $\alpha = (0.9, 0.9, 0.9)$ and $(a, b, c) = (35, 3, 28)$.

Moreover, in [Wang *et al.*, 2006], it is found that for the parameters $\alpha = (0.985, 0.99, 0.98)$ and $(a, b, c) = (35, 3, 28)$ the fractional-order Chen system can display chaotic attractors. The chaotic attractors of fractional Chen system (17) for some fractional orders are shown in Figs. 1-4.

Figure 1

Figure 2

Figure 3

Figure 4

According to our approach, the drive and the response fractional-order Chen systems are described as,

$$\begin{pmatrix} \frac{d^{\alpha_1} x_m}{dt^{\alpha_1}} \\ \frac{d^{\alpha_2} y_m}{dt^{\alpha_2}} \\ \frac{d^{\alpha_3} z_m}{dt^{\alpha_3}} \end{pmatrix} = \begin{pmatrix} -a & a & 0 \\ c-a & c & 0 \\ 0 & 0 & -b \end{pmatrix} \begin{pmatrix} x_m \\ y_m \\ z_m \end{pmatrix} + \begin{pmatrix} 0 \\ -x_m z_m \\ x_m y_m \end{pmatrix} \quad (18)$$

$$\begin{pmatrix} \frac{d^{\alpha_1} x_s}{dt^{\alpha_1}} \\ \frac{d^{\alpha_2} y_s}{dt^{\alpha_2}} \\ \frac{d^{\alpha_3} z_s}{dt^{\alpha_3}} \end{pmatrix} = \begin{pmatrix} -a & a & 0 \\ c-a & c & 0 \\ 0 & 0 & -b \end{pmatrix} \begin{pmatrix} x_s \\ y_s \\ z_s \end{pmatrix} + \begin{pmatrix} 0 \\ -x_m z_m \\ x_m y_m \end{pmatrix} + \begin{pmatrix} u_1(t) \\ u_2(t) \\ u_3(t) \end{pmatrix} \quad (19)$$

where the controller $\mathbf{u}(t) = [u_1(t), u_2(t), u_3(t)]^T$ is defined in (15). Then the error system gives,

$$\frac{d^\alpha}{dt^\alpha} \mathbf{e}(t) = B_u \mathbf{e}(t), \quad (20)$$

where $\mathbf{e}(t) = (e_1(t), e_2(t), e_3(t))$ and,

$$B_u = \begin{pmatrix} a_{11} - a & a_{12} + a & a_{13} \\ a_{21} + c - a & a_{22} + c & a_{23} \\ a_{31} & a_{32} & a_{33} - b \end{pmatrix}.$$

In view of case 1, if $\alpha_1 = \alpha_2 = \alpha_3 = \alpha$, then the drive system (18) and the response system (19) are synchronized if all roots of the equation $\det(\lambda I - B_u) = 0$ lie in the region $|\arg(\lambda)| > \alpha\pi/2$, see examples 1 and 2 below.

Also, in view of case 2, if $\alpha_1 = k_1/m_1$, $\alpha_2 = k_2/m_2$ and $\alpha_3 = k_3/m_3$ are rational numbers, then both systems are synchronized if all roots of the equation $\det(\text{diag}(\lambda^{r_1}, \lambda^{r_2}, \lambda^{r_3}) - B_u) = 0$ lie in the region $|\arg(\lambda)| > \pi/2m$, where m is the least common multiple of $\{m_1, m_2, m_3\}$ and $r_i = m\alpha_i$, see example 3 below.

Moreover, in particular and according to case 3, as in examples 3 and 4 below, if $\alpha_1, \alpha_2, \alpha_3$ are any positive real numbers between 0 and 1, then synchronization of the drive system (18) and the response system (19) is sufficiently achieved under the control laws,

$$\begin{cases} a_{12} = -a \\ a_{13} = 0 \\ a_{23} = 0 \\ a_{11} < a \\ a_{22} < -c \\ a_{33} < b \end{cases} \quad (21)$$

or,

$$\begin{cases} a_{21} = a - c \\ a_{31} = 0 \\ a_{32} = 0 \\ a_{11} < a \\ a_{22} < -c \\ a_{33} < b \end{cases} \quad (22)$$

Example 1 Taking $\alpha = (0.95, 0.95, 0.95)$, $(a, b, c) = (35, 3, 28)$ and $\mathbf{u}(t) = [(37, -33, 0), (2, -31, 0), (0, 0, -5)] \mathbf{e}(t)$, then the roots of the equation $\det(\lambda I - B_u) = 0$ are,

$$\begin{aligned} \lambda_1 &= -5, \\ \lambda_2 &= \frac{1}{2}(-1 + \sqrt{15}i), \\ \lambda_3 &= \frac{1}{2}(-1 - \sqrt{15}i), \end{aligned}$$

where $i^2 = -1$. Of course, all roots $\lambda_1, \lambda_2, \lambda_3$ have negative real parts, and so they lie in the region $|\arg(\lambda)| > (0.95)\pi/2$. Therefore, the drive system (18) and the response system (19) are synchronized.

The error functions evolution, in this case, is shown in Fig. 5 (a).

Example 2 If we take $\alpha = (0.9, 0.9, 0.9)$, $(a, b, c) = (35, 3, 28)$ and $\mathbf{u}(t) = [(32, -35, 0), (7, -26, -20), (0, 30, 2)]\mathbf{e}(t)$, then the roots of the equation $\det(\lambda I - B_u) = 0$ are,

$$\begin{aligned}\lambda_1 &= -3, \\ \lambda_2 &= \frac{1}{2}(1 + \sqrt{2391}i), \\ \lambda_3 &= \frac{1}{2}(1 - \sqrt{2391}i).\end{aligned}$$

Even the eigenvalues λ_2 and λ_3 have positive real parts, systems (18) and (19) are synchronized since $|\arg(\lambda_{2,3})| > (0.9)\pi/2$. The error functions evolution, in this case, is shown in Fig. 5 (b).

Example 3 If we take $\alpha = (0.985, 0.99, 0.98)$, $(a, b, c) = (35, 3, 28)$ and $\mathbf{u}(t) = [(25, 0, 0), (7, -33, 5), (0, 0, 1)]\mathbf{e}(t)$, we have $r_1 = 985$, $r_2 = 990$, $r_3 = 980$ and $m = 1000$ and then the equation $\det(\text{diag}(\lambda^{985}, \lambda^{990}, \lambda^{980}) - B_u) = 0$ can be written as,

$$(\lambda^{985} + 10)(\lambda^{990} + 5)(\lambda^{980} + 2) = 0.$$

Using simple calculations, we can show that all roots of above equation lie in the region $|\arg(\lambda)| > \pi/2000$. Therefore, both systems (18) and (19) are synchronized. Also, we can conclude that the systems are synchronized since the controller $\mathbf{u}(t)$ satisfies the control law (22). The error functions evolution, in this case, is shown in Fig. 5 (c).

Example 4 If we take $\alpha = (0.985, 0.99, 0.98)$, $(a, b, c) = (35, 3, 28)$ and $\mathbf{u}(t) = [(27, -35, 0), (8, -30, 0), (5, 0, 2.5)]\mathbf{e}(t)$, then $\mathbf{u}(t)$ satisfies the control law (21), and so both systems (18) and (19) are synchronized. The error functions evolution, in this case, is shown in Fig. 5 (d).

Figure 5

From Fig. 5, it is obvious that the components of the error system (20) decay towards zero as t goes to $+\infty$. Therefore, we can numerically conclude that

the designed linear controller can effectively control the chaotic fractional-order Chen system to achieve synchronization between the drive system (18) and the response system (19).

6 Synchronization of fractional-order Rössler systems

Now, we consider the Rössler system, which has been found by Rössler [1976] in 1976. Its fractional version reads as,

$$\begin{cases} \frac{d^{\alpha_1} x}{dt^{\alpha_1}} = -(y + z) \\ \frac{d^{\alpha_2} y}{dt^{\alpha_2}} = x + ay \\ \frac{d^{\alpha_3} z}{dt^{\alpha_3}} = z(x - c) + b \end{cases} \quad (23)$$

where $\alpha = (\alpha_1, \alpha_2, \alpha_3)$ is subject to $0 < \alpha_1, \alpha_2, \alpha_3 \leq 1$. When $\alpha = (1, 1, 1)$, system (23) is the original integer order Rössler system, which is chaotic, for example, when $(a, b, c) = (0.2, 0.2, 5)$, see Fig 6. Simulations are performed to obtain chaotic behavior of the fractional-order Rössler system and the results demonstrate that chaos indeed exists with order less than 3 (i.e. $\alpha_1 + \alpha_2 + \alpha_3 < 3$). For example, chaotic attractors are found in [Li & Chen, 2004a] when $\alpha = (0.9, 0.9, 0.9)$ and $(a, b, c) = (0.4, 0.2, 10)$. Also, in [Peng, 2007] it is found that for the parameters $\alpha = (1, 1, 0.8)$ and $(a, b, c) = (0.2, 0.2, 5)$ the fractional-order Rössler system can display chaotic behaviors. Moreover, we find that chaotic behaviors can be displayed for the fractional-order Rössler system when $(a, b, c) = (0.2, 0.2, 5)$ and $\alpha = (0.95, 0.95, 0.95)$ or $\alpha = (0.9, 0.9, 0.9)$, see Figs. 7 and 8.

Figure 6

Figure 7

Figure 8

According to our approach, the drive and the response fractional-order Rössler systems are described as,

$$\begin{pmatrix} \frac{d^{\alpha_1} x_m}{dt^{\alpha_1}} \\ \frac{d^{\alpha_2} y_m}{dt^{\alpha_2}} \\ \frac{d^{\alpha_3} z_m}{dt^{\alpha_3}} \end{pmatrix} = \begin{pmatrix} 0 & -1 & -1 \\ 1 & a & 0 \\ 0 & 0 & -c \end{pmatrix} \begin{pmatrix} x_m \\ y_m \\ z_m \end{pmatrix} + \begin{pmatrix} 0 \\ 0 \\ x_m z_m + b \end{pmatrix} \quad (24)$$

$$\begin{pmatrix} \frac{d^{\alpha_1} x_s}{dt^{\alpha_1}} \\ \frac{d^{\alpha_2} y_s}{dt^{\alpha_2}} \\ \frac{d^{\alpha_3} z_s}{dt^{\alpha_3}} \end{pmatrix} = \begin{pmatrix} 0 & -1 & -1 \\ 1 & a & 0 \\ 0 & 0 & -c \end{pmatrix} \begin{pmatrix} x_s \\ y_s \\ z_s \end{pmatrix} + \begin{pmatrix} 0 \\ 0 \\ x_m z_m + b \end{pmatrix} + \begin{pmatrix} u_1(t) \\ u_2(t) \\ u_3(t) \end{pmatrix} \quad (25)$$

where the controller $\mathbf{u}(t) = [u_1(t), u_2(t), u_3(t)]^T$ is defined in (15). Then the error system gives,

$$\frac{d^\alpha}{dt^\alpha} \mathbf{e}(t) = C_u \mathbf{e}(t), \quad (26)$$

where,

$$C_u = \begin{pmatrix} a_{11} & a_{12} - 1 & a_{13} - 1 \\ a_{21} + 1 & a_{22} + a & a_{23} \\ a_{31} & a_{32} & a_{33} - c \end{pmatrix}.$$

In view of case 1, if $\alpha_1 = \alpha_2 = \alpha_3 = \alpha$, then the drive system (24) and the response system (25) are synchronized if all roots of the equation $\det(\lambda I - C_u) = 0$ lie in the region $|\arg(\lambda)| > \alpha\pi/2$, see example 5 below.

Also, in view of case 2, if $\alpha_1 = k_1/m_1$, $\alpha_2 = k_2/m_2$ and $\alpha_3 = k_3/m_3$ are rational numbers, then both systems are synchronized if all roots of the equation $\det(\text{diag}(\lambda^{r_1}, \lambda^{r_2}, \lambda^{r_3}) - C_u) = 0$ lie in the region $|\arg(\lambda)| > \pi/2m$, where m is the least common multiple of $\{m_1, m_2, m_3\}$ and $r_i = m\alpha_i$, see example 6 below.

Moreover, in particular and according to case 3, if $\alpha_1, \alpha_2, \alpha_3$ are any positive real numbers between 0 and 1, then synchronization of the drive system (24) and the response system (25) is sufficiently achieved under the control laws,

or,

$$\begin{cases} a_{12} = -1 \\ a_{13} = 1 \\ a_{23} = 0 \\ a_{11} < 0 \\ a_{22} < -a \\ a_{33} < c \end{cases} \quad (27)$$

$$\begin{cases} a_{21} = -1 \\ a_{31} = 0 \\ a_{32} = 0 \\ a_{11} < 0 \\ a_{22} < -a \\ a_{33} < c \end{cases} \quad (28)$$

Example 5 Taking $\alpha = (0.9, 0.9, 0.9)$, $(a, b, c) = (0.4, 0.2, 10)$ and $\mathbf{u}(t) = [(2, 9, 1), (-18, -2.4, 0), (0, 0, 7)]\mathbf{e}(t)$, then the roots of the equation $\det(\lambda I - C_u) = 0$ are,

$$\lambda_1 = -3, \quad \lambda_2 = \sqrt{132}i, \quad \lambda_3 = -\sqrt{132}i.$$

Since $|\arg(\lambda_{1,2,3})| > (0.9)\pi/2$, the drive system (24) and the response system (25) are synchronized. The error functions evolution, in this case, is shown in Fig. 9 (a).

Example 6 If we take $\alpha = (1, 1, 0.8)$, $(a, b, c) = (0.2, 0.2, 5)$ and $\mathbf{u}(t) = [(-2, -2, 1), (6, -1.2, 0), (0, 0, -2)]\mathbf{e}(t)$, we have $r_1 = 10$, $r_2 = 10$, $r_3 = 8$ and $m = 10$ and then the equation $\det(\text{diag}(\lambda^{10}, \lambda^{10}, \lambda^8) - C_u) = 0$ can be written as,

$$(\lambda^8 + 2)(\lambda^{20} + 3\lambda^{10} + 20) = 0.$$

Simply, we can show that all roots of above equation lie in the region $|\arg(\lambda)| > \pi/20$. Therefore, both systems (24) and (25) are synchronized. The error functions evolution, in this case, is shown in Fig. 9 (b).

Figure 9

From Fig. 9, for the given parameters, we can conclude that the components of the error system (26) decay towards zero as t goes to $+\infty$, and so the drive system (24) and the response system (25) are synchronized.

7 Synchronization of fractional-order cubic Chua systems

Now, we consider the Chua's equations [Chua & Lin, 1990; Chua *et al.*, 1993] with cubic nonlinearity (i.e. cubic Chua system), see for example [Huang *et al.*, 1996]. Its fractional version reads as,

$$\begin{cases} \frac{d^{\alpha_1} x}{dt^{\alpha_1}} = a(y + bx + cx^3) \\ \frac{d^{\alpha_2} y}{dt^{\alpha_2}} = x - y + z \\ \frac{d^{\alpha_3} z}{dt^{\alpha_3}} = -\beta y - \gamma z \end{cases} \quad (29)$$

where $\alpha = (\alpha_1, \alpha_2, \alpha_3)$ is subject to $0 < \alpha_1, \alpha_2, \alpha_3 \leq 1$. We fix the parameters as,

$$(a, b, c) = (9.5, 0.15, -0.3), (\beta, \gamma) = (14, 0.02), \quad (30)$$

and we perform numerical simulations to obtain chaotic behaviors of the fractional-order cubic Chua system. First of all, the integer case $\alpha = (1, 1, 1)$ is studied and chaotic attractors are found as shown in Fig. 10. Then, in the fractional-order case, with $\alpha = (0.98, 0.98, 0.98)$ or $\alpha = (0.96, 0.96, 0.96)$, chaotic behaviors are also obtained as shown in Figs. 11 and 12.

Figure 10

Figure 11

Figure 12

According to our approach, the drive and the response systems for the fractional-order cubic Chua system are described as,

$$\begin{aligned} \begin{pmatrix} \frac{d^{\alpha_1} x_m}{dt^{\alpha_1}} \\ \frac{d^{\alpha_2} y_m}{dt^{\alpha_2}} \\ \frac{d^{\alpha_3} z_m}{dt^{\alpha_3}} \end{pmatrix} &= \begin{pmatrix} -ab & a & 0 \\ 1 & -1 & 1 \\ 0 & -\beta & -\gamma \end{pmatrix} \begin{pmatrix} x_m \\ y_m \\ z_m \end{pmatrix} \\ &+ \begin{pmatrix} -acx_m^3 \\ 0 \\ 0 \end{pmatrix} \end{aligned} \quad (31)$$

$$\begin{aligned} \begin{pmatrix} \frac{d^{\alpha_1} x_s}{dt^{\alpha_1}} \\ \frac{d^{\alpha_2} y_s}{dt^{\alpha_2}} \\ \frac{d^{\alpha_3} z_s}{dt^{\alpha_3}} \end{pmatrix} &= \begin{pmatrix} -ab & a & 0 \\ 1 & -1 & 1 \\ 0 & -\beta & -\gamma \end{pmatrix} \begin{pmatrix} x_s \\ y_s \\ z_s \end{pmatrix} \\ &+ \begin{pmatrix} -acx_m^3 \\ 0 \\ 0 \end{pmatrix} + \begin{pmatrix} u_1(t) \\ u_2(t) \\ u_3(t) \end{pmatrix} \end{aligned} \quad (32)$$

Then, the error system gives,

$$\frac{d^\alpha}{dt^\alpha} \mathbf{e}(t) = D_u \mathbf{e}(t), \quad (33)$$

where,

$$D_u = \begin{pmatrix} a_{11} - ab & a_{12} + a & a_{13} \\ a_{21} + 1 & a_{22} - 1 & a_{23} + 1 \\ a_{31} & a_{32} - \beta & a_{33} - \gamma \end{pmatrix}.$$

Similar to the previous systems, in view of case 1, if $\alpha_1 = \alpha_2 = \alpha_3 = \alpha$, then the drive system (31) and the response system (32) are synchronized if all roots of the equation $\det(\lambda I - D_u) = 0$ lie in the region $|\arg(\lambda)| > \alpha\pi/2$.

Also, in view of case 2, if $\alpha_1 = k_1/m_1$, $\alpha_2 = k_2/m_2$ and $\alpha_3 = k_3/m_3$ are rational numbers, then both systems are synchronized if all roots of the equation $\det(\text{diag}(\lambda^{r_1}, \lambda^{r_2}, \lambda^{r_3}) - D_u) = 0$ lie in the region $|\arg(\lambda)| > \pi/2m$, where m is the least common multiple of $\{m_1, m_2, m_3\}$ and $r_i = m\alpha_i$.

Moreover, in particular and according to case 3, as in examples 7 and 8, if $\alpha_1, \alpha_2, \alpha_3$ are any positive real numbers between 0 and 1, then synchronization of the drive system (31) and the response system (32) is sufficiently achieved under the control laws,

$$\begin{cases} a_{12} = -a \\ a_{13} = 0 \\ a_{23} = -1 \\ a_{11} < ab \\ a_{22} < 1 \\ a_{33} < \gamma \end{cases} \quad (34)$$

or,

$$\left\{ \begin{array}{l} a_{21} = -1 \\ a_{31} = 0 \\ a_{32} = \beta \\ a_{11} < ab \\ a_{22} < 1 \\ a_{33} < \gamma \end{array} \right. \quad (35)$$

Example 7 Taking $\alpha = (0.98, 0.98, 0.98)$, the same parameters as given in Eq. (30) and $\mathbf{u}(t) = [(-2.5, -9.5, 0), (2, -3, -1), (5, 3, -2)]\mathbf{e}(t)$, then $\mathbf{u}(t)$ satisfies the control law (34), and so both systems (31) and (32) are synchronized. The error functions evolution, in this case, is shown in Fig. 13 (a).

Example 8 If we take $\alpha = (0.96, 0.96, 0.96)$, the same parameters as given in Eq. (30) and $\mathbf{u}(t) = [(-4, 3, 5), (-1, 0.2, 1), (0, 14, -3)]\mathbf{e}(t)$, then $\mathbf{u}(t)$ satisfies the control law (35), and so both systems (31) and (32) are synchronized. The error functions evolution, in this case, is shown in Fig. 13 (b).

Figure 13

From Fig. 13, for the given parameters, we can conclude that the components of the error system (33) decay towards zero as $t \rightarrow +\infty$, and so the drive system (31) and the response system (32) are synchronized.

8 Conclusion

In this paper, we present theoretical results for drive-response synchronization between fractional-order systems, based on stability results of linear fractional-order systems. Suitable sufficient conditions for achieving synchronization of two identical fractional-order systems via a suitable linear controller applied to the response system are derived. We briefly discuss the issue of controlling fractional-order chaotic Chen, Rössler and modified Chua systems to realize synchronization with linear error feedback control. For this purpose, we design the controller such that the components of the error system decay towards zero as t goes to $+\infty$. Numerical simulations are given to verify the effectiveness of the proposed

synchronization scheme and to show that some typical chaotic systems can be synchronized. We emphasize that this scheme of synchronization can be extended to other chaotic fractional-order systems.

References

- [1] Yamada T., Fujisaka H. [1983] "Stability theory of synchronized motion in coupled-oscillator systems," *Progress of Theoretical Physics*, **70** 1240-1248.
- [2] Pecora L.M., Carrol T.L. [1990] "Synchronization in chaotic systems," *Phys. Rev. Lett.*, **64** 821-824.
- [3] Ott E., Grebogi C., Yorke J.A. [1990] "Controlling chaos," *Phys. Rev. Lett.*, **64** 1196-1199.
- [4] Chen G., Dong X. [1998] From chaos to order: methodologies, perspectives and applications, (World Scientific, Singapore).
- [5] Derivière S., Aziz-Alaoui M.A. [2003] "Estimation of attractors and synchronization of generalized Lorenz systems," *Dynamics of Continuous, Discrete and Impulsive Systems Series B: Applications and Algorithms*, **10**(6) 833-852.
- [6] Aziz-Alaoui M.A. [2006] "Synchronization of chaos," *Encyclopedia of Mathematical Physics*, Elsevier, 213-226.
- [7] Petráš I. [2008] "A note on the fractional-order Chua's system," *Chaos, Solitons & Fractals*, **38**(1) 140-147.
- [8] Podlubny I. [1999] Fractional Differential Equations, (Academic Press, New York).
- [9] Hilfer R. [2000] Applications of Fractional Calculus in Physics, (World Scientific, New Jersey).
- [10] Ge Z.M., Ou C.Y. [2007] "Chaos in a fractional order modified Duffing system," *Chaos, Solitons & Fractals*, **34**(2) 262-291.
- [11] Hartley T., Lorenzo C., Qammer H. [1995] "Chaos in a fractional order Chua's system," *IEEE Trans. CAS-I*, **42** 485-490.
- [12] Li C., Chen G. [2004a] "Chaos and hyperchaos in fractional order Rössler equations," *Physica A*, **341** 55-61.
- [13] Li C., Peng G. [2004] "Chaos in Chen's system with a fractional order," *Chaos, Solitons & Fractals*, **22**(2) 443-450.
- [14] Li C., Chen G. [2004b] "Chaos in the fractional order Chen system and its control," *Chaos, Solitons & Fractals*, **22**(3) 549-554.

- [15] Lu J.G., Chen G. [2006] "A note on the fractional-order Chen system," *Chaos, Solitons & Fractals*, **27**(3) 685-688.
- [16] Grigorenko I., Grigorenko E. [2003] "Chaotic dynamics of the fractional Lorenz system," *Phys. Rev. Lett.*, **91**(3) 034101.
- [17] Arneodo A., Couillet P., Spiegel E., Tresser C. [1985] "Asymptotic chaos," *Physica D*, **14**(3) 327-347.
- [18] Deng W.H., Li C.P. [2005] "Chaos synchronization of the fractional Lü system," *Physica A*, **353** 61-72.
- [19] Li C.G., Liao X.X., Yu J.B. [2003] "Synchronization of fractional order chaotic systems," *Phys. Rev. E*, **68** 067203.
- [20] Li C., Zhou T. [2005] "Synchronization in fractional-order differential systems," *Physica D*, **212**(1-2) 111-125.
- [21] Zhou S., Li H., Zhu Z., Li C. [2008] "Chaos control and synchronization in a fractional neuron network system," *Chaos, Solitons & Fractals*, **36**(4) 973-984.
- [22] Peng G. [2007] "Synchronization of fractional order chaotic systems," *Phys. Lett. A*, **363**(5-6) 426-432.
- [23] Matignon D. [1996] "Stability results of fractional differential equations with applications to control processing", in *Proceeding of IMACS, IEEE-SMC*, Lille, France, 963-968.
- [24] Deng W., Li C., Lü J. [2007] "Stability analysis of linear fractional differential system with multiple time delays," *Nonlinear Dyn.*, **48** 409-416.
- [25] Sheu L.J., Chen H.K., Chen J.H., Tam L.M. [2007] "Chaos in a new system with fractional order," *Chaos, Solitons & Fractals*, **31**(5) 1203-1212.
- [26] Yan J., Li C. [2007] "On chaos synchronization of fractional differential equations," *Chaos, Solitons & Fractals*, **32**(2) 725-735.
- [27] Li C., Yan J. [2007] "The synchronization of three fractional differential systems," *Chaos, Solitons & Fractals*, **32**(2) 751-757.
- [28] Wang J., Xiong X., Zhang Y. [2006] "Extending synchronization scheme to chaotic fractional-order Chen systems," *Physica A*, **370**(2) 279-285.
- [29] Li C.P., Deng W.H., Xu D. [2006] "Chaos synchronization of the Chua system with a fractional order," *Physica A*, **360**(2) 171-185.
- [30] Zhu H., Zhou S., Zhang J. "Chaos and synchronization of the fractional-order Chua's system," *Chaos, Solitons & Fractals*, in press.
- [31] Wu X., Li J., Chen G.J. [2008] "Chaos in the fractional order unified system and its synchronization," *Journal of the Franklin Institute*, **345**(4) 392-401.
- [32] Yu Y., Li H. [2008] "The synchronization of fractional-order Rössler hyperchaotic systems," *Physica A*, **387**(5-6) 1393-1403.
- [33] Caputo M. [1967] "Linear models of dissipation whose Q is almost frequency independent, Part II," *J. Roy. Astr. Soc.*, **13** 529-539.
- [34] Gorenflo R., Mainardi F. [1997] "Fractional calculus: Integral and differential equations of fractional order," in the book *Fractals and fractional calculus*, (eds.: Carpinteri and Mainardi), New York.
- [35] Oldham K.B., Spanier J. [1974] *The Fractional Calculus*, (Academic Press, New York).
- [36] Chen C., Ueta T. [1999] "Yet another chaotic attractor," *I. J. Bifurc. Chaos*, **9**(7) 1465-1466.
- [37] Rössler O.E. [1976] "An equation for continuous chaos," *Physics Letters*, **57A**(5) 397-398.
- [38] Chua L.O., Lin G.-N. [1990] "Canonical realization of Chua's circuit family," *IEEE transactions on Circuits and Systems*, **37** 885-902.
- [39] Chua L.O., Makoto I., Kocarev L., Eckert K. [1993] "Chaos synchronization in Chua's circuit," *J. Circuits, Systems and Computers*, **3**(1) 93-108.
- [40] Huang A., Pivka L., Wu C.W. [1996] "Chua's equation with cubic nonlinearity," *I. J. Bifurc. Chaos*, **6** 2175-2222.

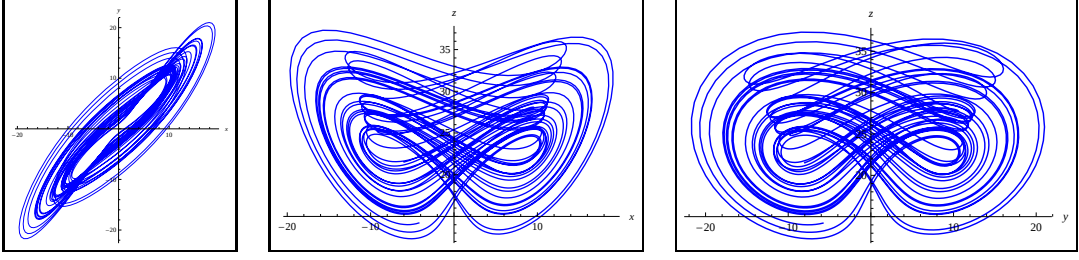


Figure 1: Chaotic attractor of the Chen system (17), when $\alpha = (1, 1, 1)$ and $(a, b, c) = (35, 3, 28)$.

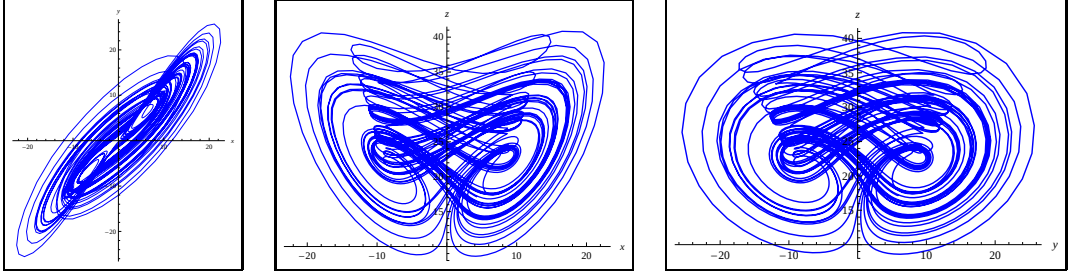


Figure 2: Chaotic attractor of the Chen system (17), when $\alpha = (0.9, 0.9, 0.9)$ and $(a, b, c) = (35, 3, 28)$.

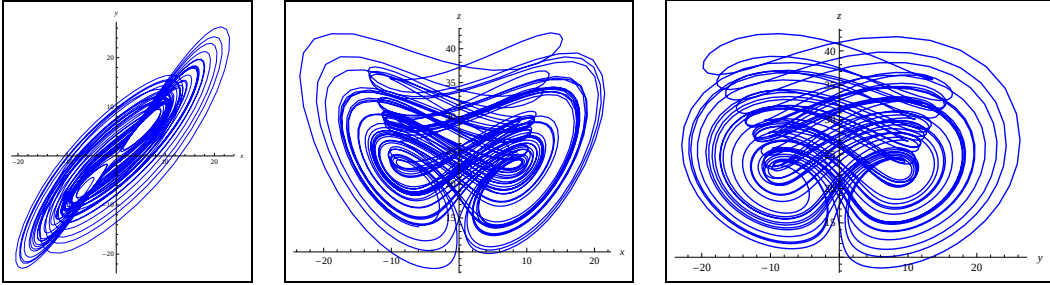


Figure 3: Chaotic attractor of the Chen system (17), when $\alpha = (0.85, 0.85, 0.85)$ and $(a, b, c) = (35, 3, 28)$.

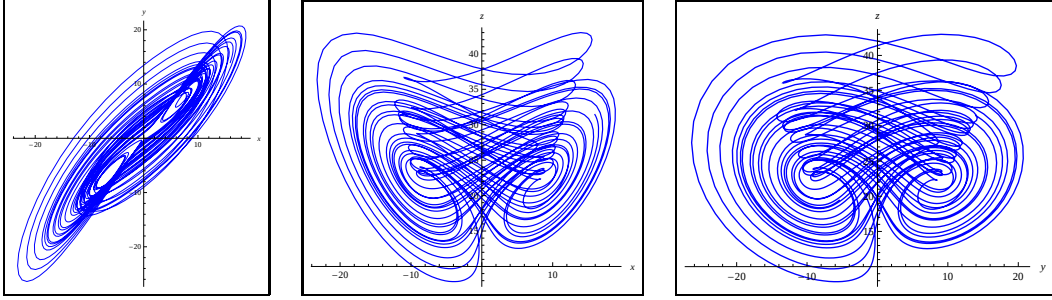


Figure 4: Chaotic attractor of the Chen system (17), when $\alpha = (0.985, 0.99, 0.98)$ and $(a, b, c) = (35, 3, 28)$.

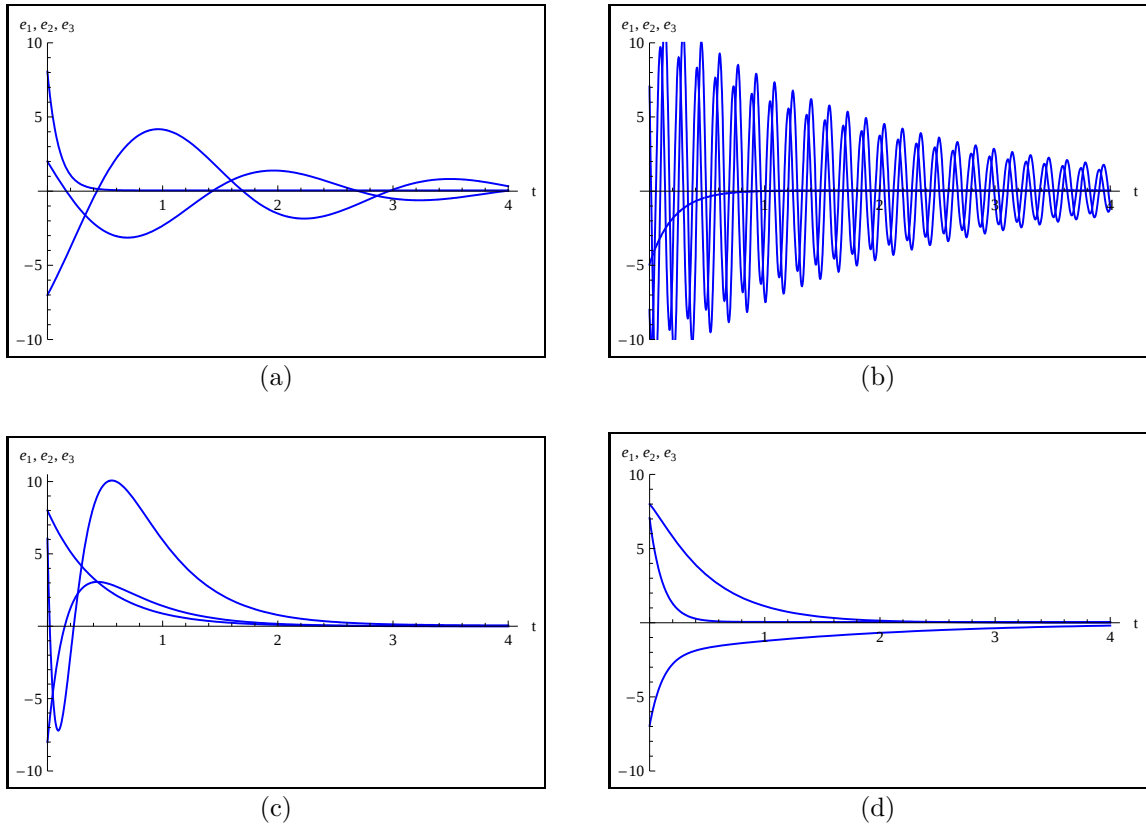


Figure 5: Synchronization of the fractional-order Chen system (17). It is obvious that the evolution of the error functions $e_1(t)$, $e_2(t)$, $e_3(t) \rightarrow 0$ as $t \rightarrow +\infty$.

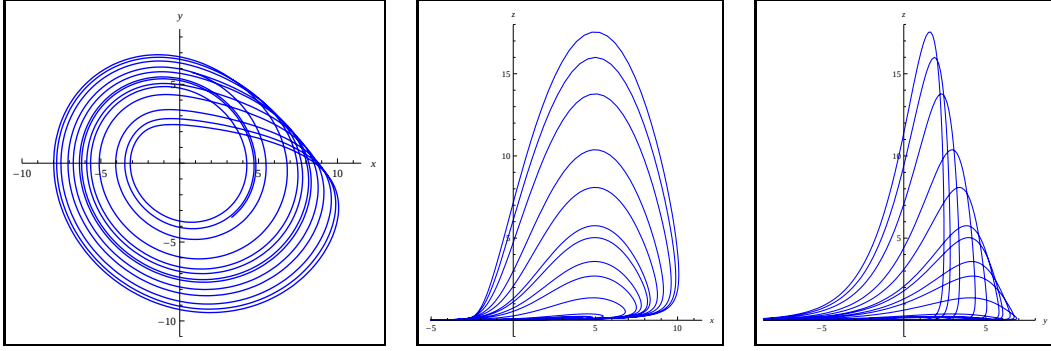


Figure 6: Chaotic attractor of the Rössler system (23), when $\alpha = (1, 1, 1)$ and $(a, b, c) = (0.2, 0.2, 5)$.

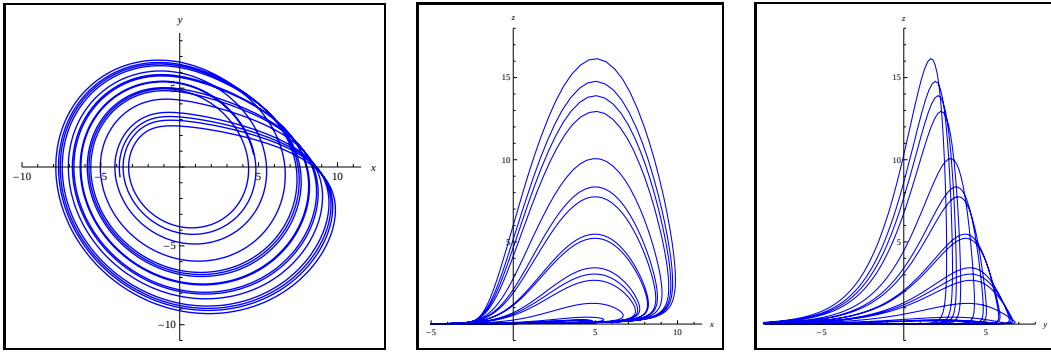


Figure 7: Chaotic attractor of the Rössler system (23), when $\alpha = (0.95, 0.95, 0.95)$ and $(a, b, c) = (0.2, 0.2, 5)$.

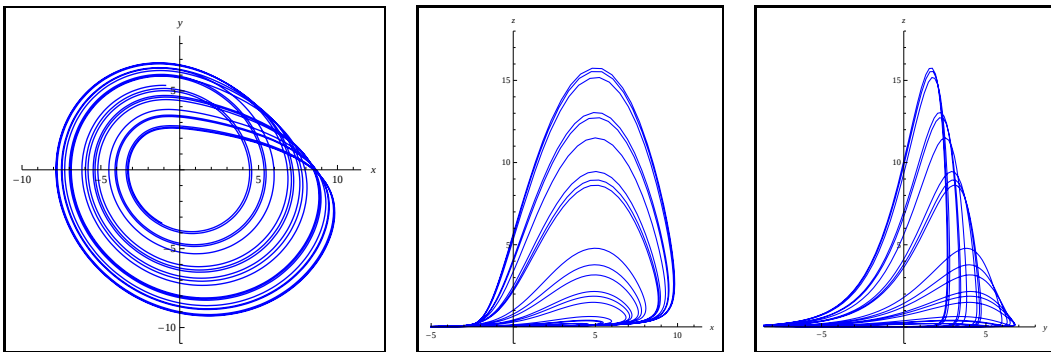


Figure 8: Chaotic attractor of the Rössler system (23), when $\alpha = (0.9, 0.9, 0.9)$ and $(a, b, c) = (0.2, 0.2, 5)$.

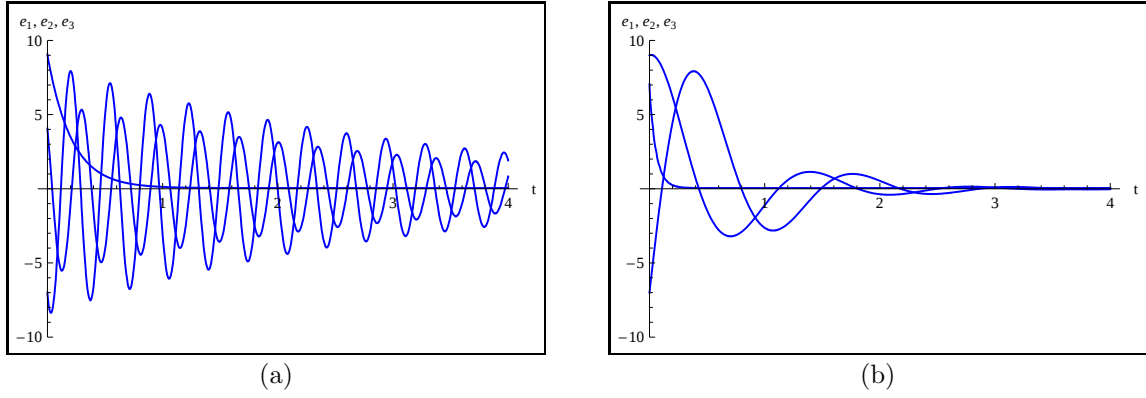


Figure 9: Synchronization of the fractional-order Rössler system (23). It is obvious that the evolution of the error functions $e_1(t)$, $e_2(t)$, $e_3(t) \rightarrow 0$ as $t \rightarrow +\infty$.

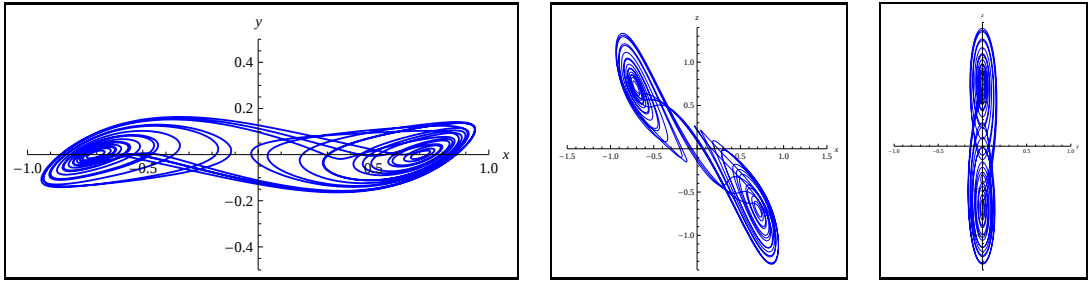


Figure 10: Chaotic attractor of the Chua cubic system (29), when $\alpha = (1, 1, 1)$ and the parameters given in Eq. (30).

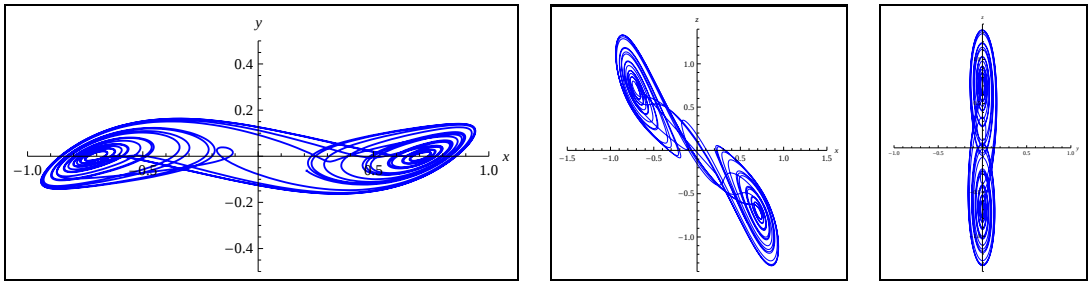


Figure 11: Chaotic attractor of the fractional Chua cubic system (29), when $\alpha = (0.98, 0.98, 0.98)$ and the parameters given in Eq. (30).

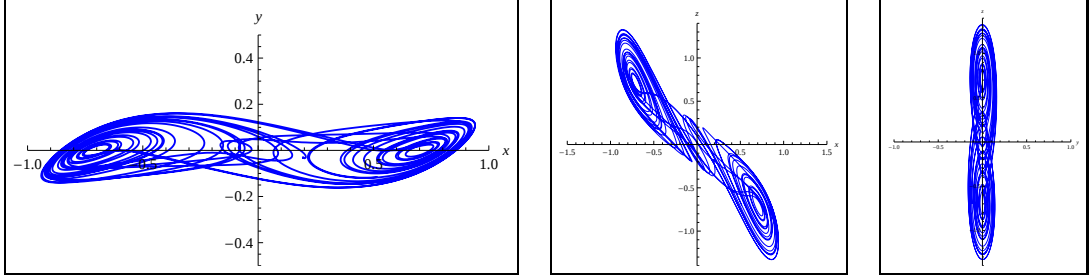


Figure 12: Chaotic attractor of the fractional Chua cubic system (29), when $\alpha = (0.96, 0.96, 0.96)$ and the parameters given in Eq. (30).

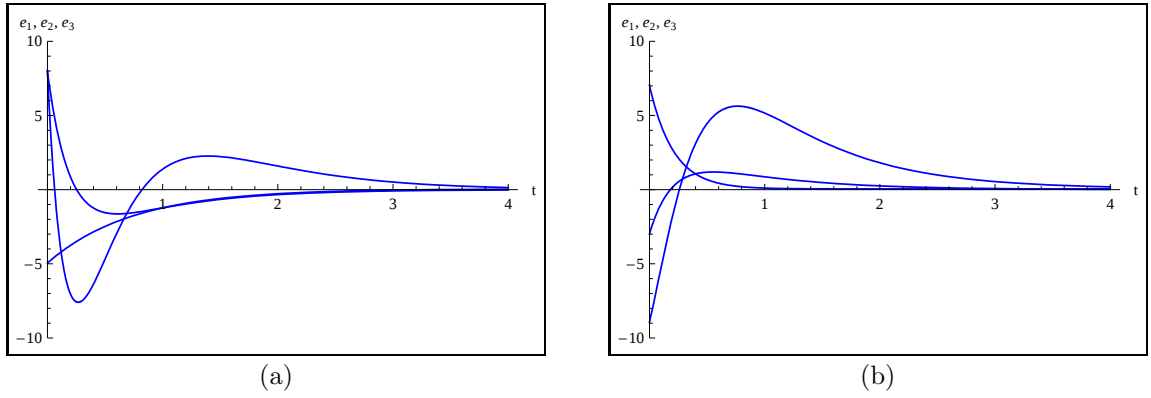


Figure 13: Synchronization of the fractional-order Chua cubic system (29). It is obvious that the evolution of the error functions $e_1(t)$, $e_2(t)$, $e_3(t) \rightarrow 0$ as $t \rightarrow +\infty$.