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To cite this version:

Patrice Bertail, Emmanuelle Gautherat, Hugo Harari-Kermadec. Exponential bounds for multivariate self-normalized sums. Revised version. 2008. <hal-00333999>

HAL Id: hal-00333999
https://hal.archives-ouvertes.fr/hal-00333999
Submitted on 24 Oct 2008

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Exponential bounds for multivariate self-normalized sums

Patrice Bertail, Emmanuelle Gautherat & Hugo Harari-Kermadec

9th April 2008

Abstract

In a non-parametric framework, we establish some non-asymptotic bounds for self-normalized sums and quadratic forms in the multivariate case for symmetric and general random variables. This bounds are entirely explicit and essentially depends in the general case on the kurtosis of the Euclidean norm of the standardized random variables.

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AMS Keywords: Primary 62G15 ; secondary 62E17, 62H15.
Exponential inequalities; self-normalization; multivariate; Hoeffding inequality.

1 Introduction

Let \( Z, Z_1, ..., Z_n \) be i.i.d. random centered vectors from a probability space \((\Omega, \mathcal{A}, \Pr)\) to \((\mathbb{R}^q, \mathcal{B}, \mathbb{P})\). We denote \( E \) the expectation under \( \mathbb{P} \). In the following we put \( \overline{Z}_n = n^{-1} \sum_{i=1}^{n} Z_i \). Define \( S \) a square root of the matrix \( S^2 = E(ZZ') \) and similarly \( S_n \) a square root of \( S_n^2 = n^{-1} \sum_{i=1}^{n} Z_i Z_i' \). We assume in the following that \( S^2 \) and \( S_n^2 \) are full rank. For further use, we define \( \gamma_r = E(\|S^{-1}Z\|_2^r) \), \( r > 0 \), where \( \| \cdot \|_2 \) is the Euclidean norm on \( \mathbb{R}^q \). Now consider the self normalized sum

\[
\left( n^{1/2} S_n^{-1} \overline{Z}_n \right)_{i=1}^{n} \left( \sum_{i=1}^{n} Z_i Z_i' \right)^{-1/2} \sum_{i=1}^{n} Z_i \tag{1}
\]

and its Euclidean norm

\[
n \overline{Z}_n \cdot S_n^{-2} \overline{Z}_n \tag{2}
\]

Self-normalized sums have recently given rise to an important literature : see for instance Jing and Wang (1999), Chistyakov and Götze (2003) or Bercu et al. (2002) for self-normalized processes. It has been proved that non-asymptotic exponential bounds can be obtained for these quantities under very weak conditions on the underlying moments of the variables \( Z_i \). Unfortunately, except in the symmetric case, these bounds established in the real case \( (q = 1) \) are not universal and depend on the skewness \( \gamma_3 = E|S^{-1}Z|^3 \) or even an higher moments for instance \( \gamma_{10/3} = E|S^{-1}Z|^{10/3} \), see Jing and Wang (1999). Actually, uniform bounds in \( \mathbb{P} \) are impossible to obtain, otherwise this would contradict Bahadur and Savage’s Theorem, see Bahadur and Savage (1956), Romano and Wolf (2000). Recall that the behaviour of self-normalized sums is closely linked to the behaviour of the statistics of Student, which is the basic
asymptotic root for constructing confidence intervals (see Remark 2 below). Moreover, available bounds are not explicit and only valid for \( n \geq n_0, n_0 \) large and unknown. To our knowledge, non-asymptotic exponential bounds with explicit constants are only available for symmetric distribution Hoeffding (1963), Efron (1969), Pinelis (1994), in the unidimensional case \( (q = 1) \). In this paper, we obtain generalizations of these bounds for \( (2) \) in the multivariate case by using a multivariate extension of the symmetrization method developed by Panchenko (2003) as well as arguments taken from the literature on self-normalized process, see Bercu et al. (2002). Our bounds are explicit but depend on the kurtosis \( \gamma_4 \) of the Euclidean norm of \( S_n^{-1} Z \) rather than on the skewness. They hold for any value of the parameter size \( q \). One technical difficulty in the multidimensional case is to obtain an explicit exponential bound for the smallest eigenvalue of the empirical variance which allows to control the deviation of \( S_n^2 \) from \( S^2 \), a result which has its own interest.

2 Exponential bounds for self-normalized sums

Some bounds for self-normalized sums may be quite easily obtained in the symmetric case (that is for random variables having a symmetric distribution) and are well-known in the unidimensional case. In non-symmetric and/or multidimensional case these bounds are new and not trivial to prove. One of the main tools for obtaining exponential inequalities in various setting is the famous Hoeffding inequality (see Hoeffding (1963)) yielding that for independent real random variables (r.v.) \( Y_i, i = 1, \ldots, n \), with finite support say \([0, 1]\), we have

\[
\Pr \left( \frac{1}{n} \left( \sum_{i=1}^{n} Y_i \right)^2 \geq t \right) \leq 2 \exp \left( -\frac{t^2}{2} \right).
\]

A direct application of this inequality to self-normalized sums (via a randomization step introducing Rademacher r.v.’s) yields (see Efron (1969), Eaton and Efron (1970)) that, for \( n \) independent random variables \( Z_i \), symmetric about 0, and not necessarily bounded (nor identically distributed), we have

\[
\Pr \left( \frac{\left( \sum_{i=1}^{n} Z_i \right)^2}{\sum_{i=1}^{n} Z_i^2} \geq t \right) \leq 2 \exp \left( -\frac{t}{2} \right).
\]

In the general non-symmetric case, the master result of Jing and Wang (1999) for \( q = 1 \) states that if \( \gamma_{10/3} < \infty \), then for some \( A \in \mathbb{R} \) and some \( a \in [0, 1] \),

\[
\Pr \left( \frac{\left( \sum_{i=1}^{n} Z_i \right)^2}{\sum_{i=1}^{n} Z_i^2} \geq t \right) \leq 2 F_1(t) + A \gamma_{10/3} n^{-1/2} e^{-at/2},
\]

where \( F_q(t) \) is the survival function of a \( \chi^2(q) \) distribution defined by \( F_q(t) = \int_{t}^{\infty} f_q(y) dy \) with \( f_q(y) = \frac{1}{2^{\frac{q}{2}} \Gamma(q/2)} y^{q/2-1} e^{-\frac{y}{2}} \) and \( \Gamma(p) = \int_{0}^{\infty} y^{p-1} e^{-y} dy \).

However the constants \( A \) and \( a \) are not explicit and, despite of its great interest to understand the large deviation behaviour of self normalized sums, the bound is of no direct practical use. In the non-symmetric case our bounds are worse than \( (4) \) as far as the control of the approximation by a \( \chi^2(q) \) distribution are concerned, but entirely explicit.

Theorem 1 Let \( Z, (Z_{11})_{1 \leq i \leq n} \), be an i.i.d. sample in \( \mathbb{R}^q \) with probability \( P \). Suppose that \( S^2 \) is of rank \( q \). Then the following inequalities hold, for finite \( n > q \) and for \( t < nq \),

a) if \( Z \) has a symmetric distribution, then, without any moment assumption,

\[
\Pr \left( n S_n^{q-2} \mathbf{Z} \geq t \right) \leq 2 q e^{-\frac{t^2}{2}},
\]

b) if \( Z \) has a non-symmetric distribution, then, with any moment assumption,
b) for general distribution of $Z$, with $\gamma_4 < \infty$, for any $a > 1$,

$$\Pr \left( nZ_n S_n^{-2} Z_n \geq t \right) \leq 2q e^{\frac{1}{2q} - \frac{q}{q_1(q+1)}} + C(q) n^3q_4^{-\frac{q}{q_1(q+1)}} \left( 1 - \frac{1}{q_1} \right)^2$$

$$\leq 2q e^{\frac{1}{2q} - \frac{q}{q_1(q+1)}} + K(q) n^3q_4^{-\frac{q}{q_1(q+1)}} \left( 1 - \frac{1}{q_1} \right)^2$$

with $q_1 = \frac{q-1}{q+1}$ and

$$C(q) = \frac{(2e\pi)^{2q}(q + 1)}{2^{2/(q+1)}(q - 1)^{3q}} \quad \text{and} \quad K(q) = \frac{C(q)}{q^{2q}} \leq 8.$$ 

Moreover for $nq \leq t$, we have

$$\Pr \left( nZ_n S_n^{-2} Z_n \geq t \right) = 0.$$ 

The proof is postponed to Appendix (1). Part a) in the symmetric multidimensional case follows by an easy but crude extension of Hoeffding (1963) or Efron (1969), Eaton and Efron (1970). The exponential inequality (5) is classical in the unidimensional case. Other type of inequalities with suboptimal rate in the exponential term have also been obtained by Major (2004).

In the general multidimensional framework, the main difficulty is actually to keep the self-normalized structure when symmetrizing the original sum. We first establish the inequality in the symmetric case by an appropriate diagonalization of the estimated covariance matrix, which reduces the problem to $q$-unidimensional inequalities. The next step is to use a multidimensional version of Panchenko’s symmetrization lemma (see Panchenko (2003)). However this symmetrization lemma destroys partly the self-normalized structure (the normalization is then $S_n^2 + S^2$ instead of the expected $S_n^2$), which can be retrieved by obtaining a lower tail control of the distance between $S_n^2$ and $S^2$. This is done by studying the behavior of the smallest eigenvalue of the normalizing empirical variance. The second term in the right hand side of inequality (6) is essentially due to this control.

However, for $q > 1$, the bound of part a) is clearly not optimal. A better bound, which has not exactly an exponential form, has been obtained by Pinelis (1994) following previous works by Eaton (1974). Pinelis’ result gives a much more precise evaluation of the tail for moderate $q$. It essentially says that in the symmetric case the tail of the self-normalized sum can essentially be bounded by the tail of a $\chi^2(q)$ distribution. Notice that this tail $\Phi_q$ satisfies the following approximation (see Abramovitch and Stegun (1970), p. 941, result 26.4.12)

$$\Phi_q(t) \sim \frac{1}{\Gamma(\frac{q}{2})} \left( \frac{t}{2} \right)^{\frac{q}{2} - 1} \exp(-\frac{t}{2}).$$

This bounds gives the right behavior of the tail (in $q$) as $t$ grows, which is not the case for a). However, in the unidimensional case a) still gives a better approximation than Pinelis (1994). a) can still be used in the multidimensional case to get crude but exponential bounds. We expect however Pinelis’ inequality to give much better bounds for moderate $q$ and moderate sample size $n$ in the symmetric case. For these reason, we will extend the results of Theorem 1 by using a $\chi^2(q)$ type of control. This essentially consists in extending lemma 1 of Panchenko (2003) to non exponential bound.

**Theorem 2** The following inequalities hold, for finite $n > q$ and for $t < nq$:

a) (Pinelis 1994) if $Z$ has a symmetric distribution, without any moment assumption, then we have

$$\Pr \left( nZ_n S_n^{-2} Z_n \geq t \right) \leq \frac{2q^3}{9} \Phi_q(t),$$

$$\Pr \left( nZ_n S_n^{-2} Z_n \geq t \right) \leq \frac{2q^3}{9} \Phi_q(t).$$
b) for general distribution of $Z$ with kurtosis $\gamma_4 < \infty$, for any $a > 1$ and for $t \geq 2q(1+a)$ and $\tilde{q} = \frac{q-1}{q+1}$ we have

$$\Pr \left( n Z_n S_n^{-2} Z_n \geq t \right)$$

$$\leq \frac{2 e^3}{9 \Gamma \left( \frac{q}{2} + 1 \right)} \left( \frac{t - q(1+a)}{2(1+a)} \right)^{\frac{q}{2}} e^{-\frac{t-q(1+a)}{2(1+a)}} + C(q) \left( \frac{n^3}{\gamma_4} \right)^{\frac{q}{4}} e^{-\frac{n(1-\frac{1}{q})^2}{\gamma_4(n+1)}}$$

$$\leq \frac{2 e^3}{9 \Gamma \left( \frac{q}{2} + 1 \right)} \left( \frac{t - q(1+a)}{2(1+a)} \right)^{\frac{q}{2}} e^{-\frac{t-q(1+a)}{2(1+a)}} + K(q) n^{3q} e^{-\frac{n(1-\frac{1}{q})^2}{\gamma_4(n+1)}} \quad (8)$$

For $t \geq nq$, we have $\Pr \left( n Z_n S_n^{-2} Z_n \geq t \right) = 0$.

Remark 1 Notice that the constant $K(q)$ does not increase with large $q$ as it can be seen on Figure 1. A close examination of the bounds shows that essentially $\gamma_4(q+1)$ has to be small compared to $n$ for practical use of these bounds. Of course practically $\gamma_4$ is not known, however one may use an estimator or an upper bound for this quantity to get some insight on a given estimation problem.

Remark 2 It can be tempting to compare our bounds with some more classical results in statistics. We recall that, in an unidimensional framework, the studentized ratio is given by $T_n = \frac{Z_n}{\sqrt{n} S_n}$ where $S_n$ is the unbiased estimator of the variance $\hat{S}_n = \frac{1}{n} \sum_{i=1}^{n} (Z_i - \bar{Z}_n)^2$. In a Gaussian framework, $\bar{T}_n$ has a Student distribution with $(n-1)$ degrees of freedom. In opposition, our self-normalized sum is defined by $T_n = \frac{\sqrt{n}}{n-1} \left( \frac{1}{n-1} \sum_{i=1}^{n} Z_i^2 \right)^{-1/2} \bar{Z}_n$. It is related to $\bar{T}_n$ by the relation $T_n = f_n(\bar{T}_n)$ with $f_n(x) = \sqrt{\frac{n}{n-1}} \left( 1 + \frac{x^2}{n-1} \right)^{-1/2}$. As a consequence, one gets in an unidimensional symmetric case, for $t > 0$,

$$\Pr(\bar{T}_n \geq t) \leq \exp \left\{ -\frac{1}{2} \frac{n}{n-1} \frac{t^2}{1 + \frac{t^2}{n-1}} \right\}.$$}

For large $n$ we recover an sub-gaussian type of inequality. At fixed $n$, this inequality is noninformative for $t \to \infty$ since the right-hand side tends to 1. Recall that, in a Gaussian framework, the tail $\Pr(\bar{T}_n > t)$ is of order $O(\frac{1}{\sqrt{n}})$ as $t \to \infty$.

Remark 3 In the best case, past studies give some bounds for $n$ sufficiently large, without an exact value for "sufficiently large". Here, the bounds are valid and explicit for any $n > q$.

These bounds are motivated by some statistical applications to the construction of non-asymptotic confidence intervals with conservative coverage probability in a semi-parametric setting. Self-normalized sums appear naturally in the context of empirical likelihood and its generalization to Cressie-Read divergences, see Harari-Kermadec (2006), Owen (2001). In particular, Bertail et al. (2005) shows how the bounds obtained here may be used to construct explicit non asymptotic confidence regions, even when $q$ depends on $n$.  

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**Figure 1:** Value of $K(q)$ as a function of $q$.
A Proofs of the main results

A.1 Some lemmas

The first lemma is a direct extension of Panchenko, 2003, Corollary 1 to the multidimensional case, which will be used both in theorem 1 and 2.

**Lemma 1** Let $J_q$ be the unit sphere of $\mathbb{R}^q$, $J_q = \{ \lambda \in \mathbb{R}^q, \| \lambda \|_2 = 1 \}$. Let $Z^{(n)} = (Z_i)_{1 \leq i \leq n}$ and $Y^{(n)} = (Y_i)_{1 \leq i \leq n}$ be i.i.d. centered random vectors in $\mathbb{R}^q$ with $Z^{(n)}$ independent of $Y^{(n)}$. We denote, for any random vector $W = (W_i)_{1 \leq i \leq n}$, $S_{n,W}^2 = \frac{1}{n} \sum_{i=1}^n W_i W_i'$.

If there exists $D > 0$ and $d > 0$ such that, for all $t \geq 0$,

$$\Pr \left( \sup_{\lambda \in J_q} \left( \sqrt{n} \lambda \left( \frac{Z_n - \bar{Y}_n}{\sqrt{\lambda' S_{n,(Z^{(n)} - Y^{(n)})} \lambda}} \right) \right) \geq \sqrt{t} \right) \leq D e^{-dt},$$

then, for all $t \geq 0$,

$$\Pr \left( \sup_{\lambda \in J_q} \frac{\sqrt{n} \lambda' Z_n}{\sqrt{\lambda' S_{n} \lambda + \lambda' S^{2} \lambda}} \geq \sqrt{t} \right) \leq D e^{d - dt}. \quad (9)$$

**Proof:** This proof follows Lemma 1 of Panchenko (2003) with some adaptations to the multidimensional case. Denote

$$A_n(Z^{(n)}) = \sup_{\lambda \in J_q} \sup_{b > 0} \left\{ E \left[ 4b (\lambda' (Z_n - \bar{Y}_n) - b \lambda' S_{n,Z^{(n)} - Y^{(n)}} \lambda) \right] \right\}$$

and

$$C_n(Z^{(n)}, Y^{(n)}) = \sup_{\lambda \in J_q} \sup_{b > 0} \left\{ 4b (\lambda' (Z_n - \bar{Y}_n) - b \lambda' S_{n,Z^{(n)} - Y^{(n)}} \lambda) \right\}.$$

By Jensen inequality, we have $\Pr$-almost surely

$$A_n(Z^{(n)}) \leq E[C_n(Z^{(n)}, Y^{(n)}) | Z^{(n)}]$$

and, for any convex function $\Phi$, by Jensen inequality, we also get

$$\Phi(A_n(Z^{(n)})) \leq E[\Phi(C_n(Z^{(n)}, Y^{(n)}) | Z^{(n)})].$$

We obtain

$$E(\Phi(A_n(Z^{(n)}))) \leq E(\Phi(C_n(Z^{(n)}, Y^{(n)}))). \quad (10)$$

Now remark that

$$A_n(Z^{(n)}) = \sup_{\lambda \in J_q} \sup_{b > 0} \left\{ 4b (\lambda' Z_n - b \lambda' S_{n} \lambda - b \lambda' S^{2} \lambda) \right\}$$

and

$$C_n(Z^{(n)}, Y^{(n)}) = \sup_{\lambda \in J_q} \left( \frac{\lambda' (Z_n - \bar{Y}_n)}{\sqrt{\lambda' S_{n,Z - Y} \lambda}} \right)^2.$$

Now, notice that $\sup_{\lambda \in J_q} \frac{\lambda' Z_n}{\sqrt{\lambda' S_{n,Z}}} > 0$ and apply the arguments of the proof of Panchenko (2003)’s Corollary 1 applied to inequality (10) to obtain the result. ■

The next lemma allows to establish an non exponential version of the preceding lemmas. Indeed a consequence of this lemma is that, if the tail of the symmetrized version in (A.1) is controlled by a chi-square tail, then the non symmetrized version may be controlled by an exponential multiplied by a polynomial. The rate in the exponential is asymptotically correct.
Lemma 2 Let \( \nu \) and \( \xi \), be two r.v.’s, satisfying \( E(\xi) \leq E(\nu) \) and such that, there exists a constant \( C > 0 \), such that, for \( t > 0 \),
\[
\Pr(\nu > t) \leq CF_q(t)
\]
then, for \( t \geq 2q \), we have
\[
\Pr(\xi > t) \leq C \left( \frac{(t - q)}{2} \right)^q e^{-\frac{(t - q)}{\Gamma(q/2 + 1)}}.
\]
and for \( t > q \), we have
\[
\Pr(\xi > t) \leq CF_{q+2}(t - q).
\]

Proof: We follow the lines of the proof of Panchenko’s lemma, with function \( \Phi \) given by \( \Phi(x) = \max(x - t + q; 0) \) for some \( t > q \). Remark that \( \Phi(0) = 0 \) and \( \Phi(t) = q \), then we have
\[
\Pr(\xi \geq t) \leq \frac{1}{\Phi(t)} \left( \Phi(0) + \int_0^{+\infty} \Phi'(x) \Pr(\nu \geq x)dx \right)
\]
\[
\leq \frac{C}{q} \int_{t-q}^{+\infty} F_q(x)dx.
\]
By integration by parts, we have
\[
\int_{t-q}^{+\infty} F_q(x)dx = \int_{t-q}^{+\infty} xf_q(x)dx - (t - q) \int_{t-q}^{+\infty} f_q(x)dx.
\]
It follows by straightforward calculations that, for \( t > q \),
\[
\Pr(\xi \geq t) \leq \frac{C}{q} \int_{t-q}^{+\infty} F_q(x)dx = C \left( F_{q+2}(t - q) - \frac{t - q}{q} F_q(t - q) \right).
\]
For \( t \geq 2q \), and using the recurrence relation 26.4.8 of Abramovitch and Stegun (1970), page 941.
\[
\Pr(\xi \geq t) \leq C \left( F_{q+2}(t - q) - F_q(t - q) \right) = \left( \frac{(t - q)}{2} \right)^{q/2} C e^{-\frac{(t - q)}{\Gamma(q/2 + 1)}}.
\]
Moreover, for \( t > q \) we have \( \Pr(\xi \geq t) \leq CF_{q+2}(t - q) \). \( \blacksquare \)

We now extend a result of Barbe and Bertail (2004), which controls the behavior of the smallest eigenvalue of the empirical variance. In the following, for a given symmetric matrix \( A \), we denote \( \mu_1(A) \) its smallest eigenvalue.

Lemma 3 Let \((Z_i)_{1 \leq i \leq n}\) be i.i.d. random vectors in \( \mathbb{R}^q \) with common mean 0. Assume \( 1 \leq \tilde{\gamma}_4 = E(\|Z_i\|^4_2) < +\infty \). Then, for any \( 1 \leq q < n \) and \( 0 < t \leq \mu_1(S_n^2) \),
\[
\Pr \left( \mu_1(S_n^2) \leq t \right) \leq C(q) \frac{n^{3q} \mu_1(S^2)^{2q}}{\tilde{\gamma}_4^q} \exp \left( - \frac{n(\mu_1(S^2) - t)^2}{\tilde{\gamma}_4(q + 1)} \right) \wedge 1,
\]
with
\[
C(q) = \pi^{2q}(q + 1)e^{2q}(q - 1)^{-3q}2^{2q} \frac{\delta^q}{\tilde{\gamma}_4}
\]
\[
\leq 4\pi^2(q + 1)e^q(q - 1)^{-3q}.
\]
**Proof**: The proof of this result is adapted from Barbe and Bertail (2004) and makes use of some idea of Bercu et al. (2002).

We first have by a truncation argument and applying Markov’s inequality on the last term in the inequality (see the proof of Barbe and Bertail, 2004, Lemma 4), for every $M > 0$,

$$
\Pr \left( \mu_1 \left( \sum_{i=1}^{n} Z_i Z'_i \right) \leq nt \right) \leq \Pr \left( \inf_{v \in J_q} \sum_{i=1}^{n} (v' Z_i)^2 \leq nt, \sup_{i=1,...,n} ||Z_i||_2 \leq M \right) + n \frac{\gamma}{M^4} \quad (13)
$$

We call $I$ the first term on the right hand side of this inequality.

Notice that by symmetry of the sphere, we can always work with the northern hemisphere of the sphere rather than the sphere. In the following, we denote by $\mathcal{N}_q$ the northern hemisphere of the sphere.

Notice that, if the supremum of the $||Z_i||_2$ is smaller than $M$, then for $u, v$ in $\mathcal{N}_q$, we have

$$
\left| \sum_{i=1}^{n} (v' Z_i)^2 - \sum_{i=1}^{n} (u' Z_i)^2 \right| \leq 2n ||u - v||_2 M^2.
$$

Thus if $u$ and $v$ are apart of $t\eta/(2M^2)$ then $| \sum_{i=1}^{n} (v' Z_i)^2 - \sum_{i=1}^{n} (u' Z_i)^2 | \leq \eta nt$. Now let $N(\mathcal{N}_q, \varepsilon)$ be the smallest number of caps of radius $\varepsilon$ centered at some points on $\mathcal{N}_q$ (for the $||.||_2$ norm) needed to cover $\mathcal{N}_q$. Following the same arguments as Barbe and Bertail (2004), we have, for any $\eta > 0$,

$$
I \leq N \left( N_q, \frac{t\eta}{2M^2} \right) \max_{u \in \mathcal{N}_q} \Pr \left( \sum_{i=1}^{n} (u' Z_i)^2 \leq (1 + \eta)nt \right).
$$

The proof is now divided in three steps, i) control of $N(\mathcal{N}_q, \frac{t\eta}{2M^2})$, ii) control of the maximum over $\mathcal{N}_q$ of the last expression in $I$, iii) optimization over all the free parameters.

i) On the one hand, we have, for some constant $b(q) > 0$,

$$
N(\mathcal{N}_q, \varepsilon) \leq b(q) \varepsilon^{-(q-1)} \vee 1. \quad (14)
$$

For instance, we may choose $b(q) = \pi^{-q}$.

Indeed, following Barbe and Bertail (2004), the northern hemisphere can be parameterized in polar coordinates, realizing a diffeomorphism with $J_{q-1} \times [0, \pi]$. Now proceed by induction, notice that for $q = 2$, $\mathcal{N}_q$, the half circle can be covered by $[\pi/2\varepsilon] \vee 1 + 1 \leq 2[(\pi/2\varepsilon) \vee 1] \leq \pi/\varepsilon \vee 1$ caps of diameter $2\varepsilon$, that is, we can choose the caps with their center on a $\varepsilon$-grid on the circle. Now, by induction we can cover the cylinder $J_{q-1} \times [0, \pi]$ with $[\pi/2\varepsilon \pi^{q-2}/\varepsilon^{q-2}] \vee 1 + 1 \leq \pi^{q-1}/\varepsilon^{q-1}$ intersecting cylinders which in turn can be mapped to region belonging to caps of radius $\varepsilon$, covering the whole sphere (this is still a covering because the mapping from the cylinder to the sphere is contractive).

ii) On the other hand, for all $t > 0$, we have by exponentiation and Markov’s inequality, and independence of $(Z_i)_{1 \leq i \leq n}$, for any $\lambda > 0$

$$
\max_{u \in \mathcal{N}_q} \Pr \left( \sum_{i=1}^{n} u' Z_i Z'_i u \leq nt \right) \leq e^{\lambda M} \max_{u \in \mathcal{N}_q} \left( E \left[ e^{-\lambda u' Z_i Z'_i u} \right] \right)^n.
$$
Now, using the classical inequalities, $\log(x) \leq x - 1$ and $e^{-x} - 1 \leq -x + x^2/2$, both valid for $x > 0$, we have

$$\max_{w \in \mathcal{N}_q} \left( E \left[ e^{-\lambda w' Z_i Z_i' u} \right] \right)^n = \max_{w \in \mathcal{N}_q} \left\{ n \log \left( E \left[ e^{-\lambda w' Z_i Z_i' u} \right] \right) \right\} \leq \max_{w \in \mathcal{N}_q} \left\{ nE \left[ e^{-\lambda w' Z_i Z_i' u - 1} \right] \right\} \leq \max_{w \in \mathcal{N}_q} \left\{ n \left( -\lambda w' S^2 u + \frac{\lambda^2}{2} \gamma_4 \right) \right\} = \exp \left( \frac{\lambda^2}{2} n \gamma_4 - \lambda n \mu_1(S^2) \right).$$

(16)

iii) From (16) and (14), we deduce that, for any $t > 0$, $\lambda > 0$, $\eta > 0$,

$$I \leq b(q) \left( \frac{2M^2}{q-1} \right)^{q-1} e^{\lambda n(t - \mu_1(S^2)) + n\lambda^2 \gamma_4/2}.$$

Optimizing the expression $\exp(-(q - 1) \log(\eta) + \eta t)$ in $\eta > 0$, yields immediately, for any $t > 0$, any $M > 0$, any $\lambda > 0$

$$I \leq b(q) \left( \frac{2n \mu_1(S^2)}{\gamma_4(q - 1)} \right)^{q-1} \exp \left( -\frac{n}{\gamma_4} \left( \mu_1(S^2) - t \right)^2 \right) + n \gamma_4 \frac{M^2}{M^2}.$$

We now optimize in $M^2 > 0$ and the optimum is attained at

$$M^2_s = \left( \frac{2n \gamma_4}{(q - 1)b(q)} \right)^{\frac{q+1}{q-1}} \left( \frac{2n \mu_1(S^2)}{\gamma_4} \right)^{-\frac{q(q-1)}{q+1}} \exp \left( -\frac{n}{\gamma_4} \left( \mu_1(S^2) - t \right)^2 \right),$$

yielding the bound

$$\Pr \left( \mu_1 \left( \sum_{i=1}^{n} Z_i Z_i' \right) \leq t \right) \leq \hat{C}(q) n^{\frac{3q+1}{q+1}} \mu_1(S^2)^{-\frac{2(q-1)}{q+1}} \gamma_4^{-\frac{q(q-1)}{q+1}} \exp \left( -\frac{n}{\gamma_4} \left( \mu_1(S^2) - t \right)^2 \right),$$

with

$$\hat{C}(q) = b(q) \gamma_4 (q+1) e^{\frac{2(q-1)}{q+1}} (q - 1)^{-\frac{3q-1}{q+1}} 2^{\frac{2q-1}{q+1}}.$$

Using the constant $b(q) = \pi^{q-1}$ we get the expression of $C(q)$, which is bounded by the simpler bound (for large $q$ this bound will be sufficient) $4\pi^2 (q+1)e^{2(q - 1)} - \frac{3q-1}{q+1} 2^{\frac{2q-1}{q+1}}$, using the fact that $\gamma_4 \geq 1$.

The result of the Lemma follows by using this inequality combined with inequality 13. ■

### A.2 Proof of Theorem 1

**Proof**: Notice that we have always $Z_n' S_n^{-2} Z_n \leq q$. Indeed, there exists an orthogonal transformation $O_n$ and a diagonal matrix $\Lambda_n^2 := \text{diag}(\mu_j)_{1 \leq j \leq q}$ with $\mu_j > 0$ being the eigenvalues of $S_n^2$, such that $S_n^2 = O_n \Lambda_n^2 O_n$. Now put $Y_{i,n} := [Y_{i,j,n}]_{1 \leq j \leq q} = O_n Z_i$. It is easy to see that by construction the empirical variance of the $Y_{i,n}$ is

$$\frac{1}{n} \sum_{i=1}^{n} Y_{i,n} Y_{i,n}' = \frac{1}{n} \sum_{i=1}^{n} O_n Z_i Z_i' O_n' = O_n S_n^2 O_n' = \Lambda_n^2.$$
It also follows from this equality that, for all \( j = 1, \ldots, q \), \( \frac{1}{n} \sum_{i=1}^{n} Y_{i,j,n}^2 = \hat{\mu}_j \), and
\[
\bar{Z}_n \sigma_n^{-2} \bar{Z}_n = \bar{Y}_n \Lambda_n^{-2} \bar{Y}_n = \sum_{j=1}^{q} \left( \frac{1}{n} \sum_{i=1}^{n} Y_{i,j,n} \right)^2 / \hat{\mu}_j.
\]
This quantity is lower than \( q \) by Cauchy-Schwartz inequality. So, it follows that, for all \( t > qn \)
\[
\Pr \left( n \bar{Z}_n \sigma_n^{-2} \bar{Z}_n \geq t \right) = 0.
\]

\textbf{a)} In the symmetric and unidimensional framework \((q = 1)\), this bound follows from Hoeffding inequality
(see Efron (1969)). Consider now the symmetric multidimensional framework \((q > 1)\). Let \( \sigma_i, 1 \leq i \leq n \) be
Rademacher random variables, independent from \((Z_i)_{1 \leq i \leq n} \). Let \( \sigma_i = -1 \) if \( \sigma_i = 1 \) if \( 1/2 \). We denote
\[
\sigma_n(Z) = \left( \frac{1}{\sqrt{n}} \sum_{i=1}^{n} \sigma_i Z_i \right)
\]
and remark that \( \tilde{S}_n^2 = \frac{1}{n} \sum_{i=1}^{n} \sigma_i Z_i^2 \sigma_i \). Since the \( Z_i \)'s have a symmetric
distribution, meaning that \(-Z_i \) has the same distribution as \( Z_i \), we make use of a first symmetrization step:
\[
\Pr \left( n \bar{Z}_n \sigma_n^{-2} \bar{Z}_n \geq t \right) = \Pr(\sigma_n(Z)' \sigma_n(Z) \geq t).
\]
Now, we have
\[
\sigma_n(Z)' \sigma_n(Z) = \sigma_n(Y)' \Lambda_n^{-2} \sigma_n(Y)
\]
\[
= \sum_{j=1}^{q} \left( \sum_{i=1}^{n} \sigma_i Y_{i,j,n} \right)^2 / \sum_{i=1}^{n} Y_{i,j,n}^2.
\]
It follows that, for \( t > 0, \)
\[
\Pr(\sigma_n(Z)' \sigma_n(Z) \geq t) \leq \sum_{j=1}^{q} \Pr \left( \frac{\sum_{i=1}^{n} \sigma_i Y_{i,j,n}}{\sqrt{\sum_{i=1}^{n} Y_{i,j,n}^2}} \geq \sqrt{t/q} \right)
\]
\[
\leq 2 \sum_{j=1}^{q} E \Pr \left( \frac{\sum_{i=1}^{n} \sigma_i Y_{i,j,n}}{\sqrt{\sum_{i=1}^{n} Y_{i,j,n}^2}} \geq \sqrt{t/q} \left( Z_i \right)_{1 \leq i \leq n} \right).
\]
Apply now Hoeffding inequality to each unidimensional self-normalized term in this sum to conclude.

\textbf{b)} The \( Z_i \)'s are not anymore symmetric. Our first step is to control the probability \( \Pr( n \bar{Z}_n \sigma_n^{-2} \bar{Z}_n \geq t ) \).
Define
\[
B_n = \sup_{\lambda \in J_n} \left\{ \frac{\lambda \bar{Z}_n}{\sqrt{\lambda^2 S_n^2}} \right\}
\]
\[
D_n = \sup_{\lambda \in J_n} \left\{ \sqrt{1 + \frac{\lambda^2 S_n^2}{\lambda^2 S_n^2}} \right\}.
\]
First of all, remark that the following events are equivalent
\[
\left\{ n \bar{Z}_n \sigma_n^{-2} \bar{Z}_n \geq t \right\} = \left\{ B_n \geq \sqrt{\frac{t}{n}} \right\}.
\]
and notice that
\[
\Pr \left( B_n \geq \sqrt{\frac{t}{n}} \right) \leq \inf_{a > 1} \left\{ \Pr \left( B_n D_n^{-1} \geq \sqrt{\frac{t}{n(1 + a)}} \right) + \Pr(D_n \geq \sqrt{1 + a}) \right\}.
\]
The control of the first term on the right side is obtained in two steps. First apply part a) of Theorem 1 to \( n^{1/2} \sup_{\lambda \in J_n} \frac{\lambda \bar{Z}_n - \bar{Y}_n}{\sqrt{\lambda^2 S_n^2}} \). Then, by application of Lemma 1 and the previous remark, we get
\[ \sqrt{n} B_n D_n^{-1} \leq n^{1/2} \sup_{\lambda \in J_q} \frac{\lambda' Z_n}{\sqrt{\lambda' S_n^2 \lambda}} \]

we have for all \( t > 0 \),

\[ \Pr \left( B_n D_n^{-1} \geq \sqrt{\frac{t}{n(1 + a)}} \right) \leq 2q e^{\frac{1}{2q(1 + a)}}. \]

For \( a \leq 0 \), the control of the second term is trivial and useless. Whereas, for all \( a > 0 \), and all \( t > 0 \) we have

\[ \{ D_n \geq \sqrt{1 + a} \} = \left\{ \sup_{\lambda \in J_q} \left( 1 + \frac{\lambda' S_n^2 \lambda}{\lambda' S_n^2 \lambda} \right) \geq 1 + a \right\} \]

\[ = \left\{ \inf_{\lambda \in J_q} \left( \lambda' S_n^{-1} S_n^2 S_n^{-1} \lambda \right) \leq \frac{1}{a} \right\} \subset \left\{ \mu_1 S_n^2 S_n^{-1} \leq \frac{1}{a} \right\}. \]

We now use Lemma 3 applied to the r.v.’s \((S_n^{-1} Z_i)_{1 \leq i \leq n}\) with covariance matrix equal to \( I d_q \). It is easy to check that \( \gamma_4 = \tilde{\gamma}_4 \). For all \( 1 < a \), we have,

\[ \Pr(D_n > \sqrt{1 + a}) \leq C(q) \left( \frac{n^3}{\gamma_4} \right)^\frac{\tilde{q}}{4} e^{-\frac{n^3}{\gamma_4} \frac{1}{(1 - \frac{a}{4q})^2}}. \]

Since \( \inf_{a > -1} \leq \inf_{a > 1} \), we conclude that, for any \( t > 0 \),

\[ \Pr \left( B_n > \sqrt{\frac{t}{n}} \right) \leq \inf_{a > 1} \left\{ 2q e^{-\frac{n^3}{\gamma_4} \frac{1}{(1 - \frac{a}{4q})^2}} + C(q) \left( \frac{n^3}{\gamma_4} \right)^\frac{\tilde{q}}{4} e^{-\frac{n^3}{\gamma_4} \frac{1}{(1 - \frac{a}{4q})^2}} \right\}. \]

We achieve the proof by noticing that \( \gamma_4 \geq q^2 \) from Jensen’s inequality and \( E(||S_n^{-1} Z||^2_2) = q \). \( \blacksquare \)

### A.3 Proof of Theorem 2.

Part a) is proved in Pinelis (1994). Now, the proof of part b) follows the same lines as the Theorem 1 combining Lemmas 1, 2 and 3.

### References


