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The first coordinated observations of mid-latitude $E$-region quasi-periodic radar echoes and lower thermospheric 557.7-nm airglow

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Abstract. We present the first coordinated observations of quasi-periodic (QP) radar echoes from sporadic-$E$ ($E_s$) field-aligned irregularities (FAIs), OI 557.7-nm airglow, and neutral winds in a common volume over Shigaraki, Japan ($34.9^\circ$ N, $136.1^\circ$ E) on the night of 5 August 2002 during the SEEK-2 campaign. QP echo altitudes of 90–110 km were lower than usual by 10 km, enabling us to make a detailed comparison among QP echoes, airglow intensity, and neutral wind at around 96 km altitude. Eastward movement of the QP echo regions is consistent with the motions of neutral winds, airglow structures, and FAIs, suggesting that the electrodynamics of $E_s$-layers is fundamentally controlled by the neutral atmospheric dynamics. During the QP echo event, the echo altitudes clearly went up (down) in harmony with an airglow enhancement (subsidence) that also moved to the east. This fact suggests that the eastward-moving enhanced airglow region included an upward (downward) component of neutral winds to raise (lower) the altitude of the wind-shear node responsible for the $E_s$ formation. The airglow intensity, echo intensity, and Doppler velocity of FAIs at around 96 km altitude fluctuated with periods from 10 min to 1 h, indicating that these parameters were modulated with short-period atmospheric disturbances. Some QP echo regions below 100 km altitude contained small-scale QP structures in which very strong neutral winds exceeding 100 m/s existed. The results are compared with recent observations, theories, and simulations of QP echoes.

Keywords. Ionosphere (Ionosphere-atmosphere interactions; Ionospheric irregularities; Mid-latitude ionosphere)

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1 Introduction

Detailed behavior of 3.2-m field-aligned irregularities (FAIs) in the nighttime mid-latitude $E$-region in summer was presented for the first time by Yamamoto et al. (1991), who used the powerful middle and upper atmosphere (MU) radar at Shigaraki, Japan ($34.9^\circ$ N, $136.1^\circ$ E; $25.0^\circ$ N geomagnetic). They found spectacular radar echoes, called “quasi-periodic” (QP) echoes, that appear intermittently at altitudes above 100 km with periods of 5–20 min from post-sunset to midnight. Since then, some characteristics of the MU radar QP echoes have been revealed by Yamamoto et al. (1992, 1994, 1997) and Ogawa et al. (1995, 1998, 2002), and QP echoes similar to the MU QP echoes have been detected at other locations (e.g. Yamamoto et al., 2005). See a paper by Tsunoda et al. (2004) for radar observations of QP-like echoes prior to the finding of the QP echoes by Yamamoto et al. (1991).

QP echoes from the $E$-region are usually related to sporadic-$E$ ($E_s$) layers that accompany electron density gradients capable of inducing the gradient-drift instability under the action of the ambient electric field. A recent nonlocal theory by Seyler et al. (2004) predicts that this instability can produce electron density irregularities responsible for radar wave backscatter from both sides of a thin $E_s$-layer, superimposed on a background electron density distribution, under constant electric field and/or neutral wind. Strong polarization electric fields associated with an $E_s$-layer can rarely excite FAIs through the two-stream instability (Haldoupis et al., 1996). Neutral winds may also play a role in generating FAIs (Kagan and Kelley, 1998). From MU radar observations, Rao et al. (2000) found the new type of QP echoes (low-altitude QP echoes) at altitudes of 90–100 km that have vertical wavelengths of 0.6–1.2 km and periods of 30–90 s (see also Ogawa et al., 2002). Since then, similar low-altitude QP echoes at other locations have been reported (e.g. Urbina et
2 Experimental setup

2.1 The MU radar

The 46.5-MHz MU radar with a peak power of 1 MW has a sharp antenna beam, and can measure 3.2-m scale E-region FAIs by directing the beam northward with zenith angles of 50°–60°. Almost simultaneous observations in multi-directions are possible by using the capability of high speed beam-swinging with an interpulse period of 1.5 m/s. We used 5 beams for the observations which will be discussed in this paper. The range and time resolutions were 600 m and around 20 s, respectively. (See papers by Yamamoto et al. (1991) and Ogawa et al. (1995) for the detailed 5-beam observation technique). Figure 1 shows azimuth and zenith of each beam. The two-way half-power total beam width at these zenith angles is 4.5° and 2.3° in the vertical and horizontal planes, respectively. The 5 beams provide an azimuth coverage of ±30° about geographic north. At 100 km altitude this azimuth coverage corresponds to an east-west distance of about 155 km. All the beam bearings are within 1.8° of perpendicular to the geomagnetic field at an altitude of about 100 km. The points where the beams penetrate the 96 and 110 km altitudes are indicated in Fig. 1 by the white circles. Note that these points are nearly aligned in the geomagnetic east-west direction with separation distances of 30–50 km between neighboring beams.

Figure 2 depicts detailed configuration of the antenna pattern of beam 3 in the geographical meridian plane that looks due north (5° E of geomagnetic north) with an elevation angle of 39° (Ogawa et al., 2002). Contours of the geomagnetic aspect angle “A”, defined in the upper left corner in the figure, between the radar wave vector and the model geomagnetic field (IGRF85) are also shown. The beam bearing is exactly perpendicular to the geomagnetic field at 100 km altitude. The geomagnetic aspect angles along the beam bearing are within 90°±1° at ranges of 110–220 km, i.e. at altitudes of 70–135 km.
2.2 557.7-nm all-sky camera and Fabry-Perot interferometer

An all-sky CCD imager for measuring OI 557.7-nm airglow in the lower thermosphere has been operated at Shigaraki to obtain two-dimensional airglow intensity maps (512×512 pixels) every 5.5 min with an exposure time of 105 s (Shiokawa et al., 1999). As the 557.7-nm emission occurs mostly at altitudes of 90–100 km with its peak intensity at around 96 km (McDade et al., 1986; Iwagami et al., 2003), we assume the emission altitude of 96 km in this paper. A Fabry-Perot interferometer (FPI) has also been operated at Shigaraki to measure the Doppler shift of 557.7-nm emission at zenith angles of 50° ± 4.75° and at four azimuthal directions of N, S, E, and W, every 15 min, and to obtain neutral wind speed for the first and second fringes independently (Shiokawa et al., 2001). The difference of winds between the both fringes is estimated to be mostly less than 20 m/s. Note that the horizontal distances between the N and S points and between the E and W points are about 230 km, so that winds observed by the FPI represent those that are averaged over a large area above Shigaraki.

We operated the MU radar, all-sky imager, and FPI on the nights from 28 July to 6 August (except for 4 August) during SEEK-2 in 2002. Weather conditions, however, were bad, except for 5 and 6 August on which we could obtain simultaneous data from the radar and optical equipment. The minimum altitude of the QP echo region was 90 km on 5 August and about 100 km on 6 August, so we analyzed data only on 5 August to see the relationship among QP echoes, airglow, and neutral winds in a common volume.

3 Observations and analysis

3.1 557.7-nm airglow

Simultaneous observations of E-region FAIs, airglow, and neutral winds were successfully carried out under clear sky conditions after about 20:00 LT on the night of 5 August. Using airglow image data from the all-sky imager, we calculated a deviation of airglow intensity in each pixel of the image. Here the deviation \( \Delta I_{SS8(t)} \) (in %) is defined as \( (I_{SS8(t)} - I_0)/I_0 \), where \( I_{SS8(t)} \) is the absolute intensity (in units of Rayleighs) at time \( t \) and \( I_0 \) is the 1-h running average of \( I_{SS8(t)} \). Then the deviations were plotted on geographic coordinates by assuming an emission altitude of 96 km. An example of the deviation map is displayed in Fig. 3, where, for example, the dark airglow region extending from ENE to WSW is discerned. The detailed analysis of the consecutive deviation maps indicates that the averaged motion of the airglow patterns between 20:40 and 22:30 LT was toward SE at about 40 m/s.

The 5 beam directions of the MU radar shown in Fig. 1 are plotted in Fig. 3. A white circle on each beam represents the position where the beam penetrates the 96 km altitude. We picked up \( I_{SS8(t)} \) and \( \Delta I_{SS8(t)} \) in the pixel that is closest to the white circle. Figure 4a plots temporal variations of these parameters on all the radar beams. The background intensity of \( I_{SS8} \) is about 200 R and the maximum intensity reaches 400 R. This indicates an increase in the atomic oxygen density [O] of 26% of the background density, because \( I_{SS8} \) is approximately proportional to the cube of [O] (McDade et al., 1986). Notice in Fig. 4a that the maximum absolute intensity of 400 R on beams 1 and 5 occurs at 21:37 and 22:14 LT, respectively, indicating an apparent eastward motion of the
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Fig. 4. (a) Temporal variations of absolute intensity ($I_{558}$) and deviation of 557.7-nm airglow intensity ($\Delta I_{558}$) at around 96 km altitude (white circles in Fig. 1) on 5 MU radar beams. (b) Neutral wind speed and azimuth (clockwise from geographic north) of the wind vector determined from first and second fringes of the Fabry-Perot interferometer at Shigaraki.

bright airglow region at about 65 m/s along the dot-dashed lines at 96 km in Fig. 1. This apparent motion corresponds well to the motion of the airglow patterns in Fig. 3 (toward SE at about 40 m/s). The $\Delta I_{558}$ values (dashed curves) in Fig. 4a fluctuate with periods from 10 min (= twice the time resolution of 5.5 min for the airglow measurements) to 1 h. Interestingly, the dashed curves for beams 1 and 4 are similar to those for beams 2 and 5, respectively, with a little time difference. This fact suggests that the airglow is modulated by AGWs with a horizontal scale of about 100 km (note the separation distances of 40–50 km between neighboring radar beams), in addition to smaller scales.

Figure 4b shows temporal variations of the neutral wind speed and azimuth (clockwise from geographic north) of the wind vector that were determined from the first and second fringes of the FPI. The wind speed at around 96 km altitude was between 20 and 50 m/s while changing slowly its azimuth from the east to the south before about 22:30 LT, and from the south to the east thereafter. The wind may represent a semidiurnal tidal wind. The wind speed (20–50 m/s) and its azimuth (ENE–SSE) between 20:30 and 22:30 LT are consistent with the eastward motion of both the airglow patterns (Fig. 3) and the most bright airglow region (Fig. 4a).

3.2 E-region radar echoes

Figure 5 displays range-time-intensity (RTI) plots of signal-to-noise ratio (SNR in dB) of the MU radar echoes observed on beams 1, 3 and 5 on the night of 5 August. The altitudes shown on the right ordinate were calculated from the steered zenith angle of each radar beam. Temporal variations of absolute intensity ($I_{558}$; solid curve) at around 96 km altitude, shown in Fig. 4a, are also plotted.
by allowing the geomagnetic aspect sensitivity to be not so severe. Relying on this and consulting with Fig. 1, we estimate that the radar echo regions moved approximately to the NE before about 20:45 LT and to the SE after that. In fact, the neutral wind observed by FPI (Fig. 4b) was ENE before 20:30 LT, and thereafter changed slowly its direction toward the south. This suggests that the meridional winds control the striations of QP echoes; namely, the northward (southward) wind produces positive (negative) range rate (Pan and Tsunoda, 1999). The wind had an eastward component during 20:00–22:30 LT when the QP echoes were observed.

The temporal variations of the absolute airglow intensity ($I_{558}$) at around 96 km altitude shown in Fig. 4a are overlaid in Fig. 5. From a comparison of the radar echoes and $I_{558}$, we point out that on beam 3 (and beam 2; not shown) the altitudes where the echo intensity is high go up and down in harmony with the enhancement and subsidence of $I_{558}$, respectively, though such a behavior is not always clear on beams 1, 4 (not shown), and 5.

Figure 6 shows range-time-Doppler velocity plots of the MU radar echoes with SNRs higher than 0 dB on beams 1, 3 and 5. Plus (minus) sign of Doppler velocity represents motion away from (toward) the radar.

Figure 7 shows temporal variations of echo intensity, Doppler velocity, and deviation of 557.7-nm airglow intensity ($\Delta I_{558}$) at around 96 km altitude (white circles in Fig. 1) on beams 1, 3 and 5.

Using the $V_d$ data from all the radar beams, we estimated a vector of FAI motion ($V$) in the horizontal plane. To this end, we calculated $V_d$ averaged over 20:00–23:00 LT for altitudes above 100 km and below that for each beam. Then, assuming that $\overline{V}_d$ is a projection of $V$ along the radar beam, we obtained $V$ above 100 km ($V_1$) and below that ($V_2$). The results show that $V_1$ is 70 m/s with an azimuth of 66° (clockwise from geographic north), and $V_2$ is 53 m/s with an azimuth of 88° (eastward). $V_2$ deviates not only from the motion of the airglow structures at around 96 km (toward SE at about 40 m/s) but also from the neutral wind motions there (toward ENE–SSE at 20–50 m/s in Fig. 4b), maybe because $V_d$ below the 100-km altitude is determined by both $E \times B$ electron drift velocity and neutral winds.

Figure 7 shows temporal variations of echo intensity, $V_d$, and $\Delta I_{558}$ at around 96 km altitude (white circles in Fig. 1) on three radar beams. As is also seen in Fig. 6, $V_d$ is negative (toward the radar) for beam 1 (Fig. 7a), and mostly positive (away from the radar) for beams 3 (Fig. 7b) and 5 (Fig. 7c), again suggesting the eastward movement of FAIs. As stated before, $\Delta I_{558}$ fluctuates with periods of 10 min–1 h. Such
fluctuations are also discerned in the echo intensity and $V_d$, suggesting these parameters are modulated by common atmospheric disturbances.

The RTI plot between 20:25 and 20:45 LT in Fig. 5c is enlarged in Fig. 8a, where, as indicated by the white dashed lines, some small-scale QP structures, embedded within the positive striations, with both high echo intensity and weakly negative striation, can be seen between 20:25 and 20:37 LT at altitudes of 87–100 km. This feature means that the eastward-moving plasma blob causing the positive striation includes some small-scale plasma structures that move toward the radar. Range variations of the Doppler spectra obtained at around 20:27 and 20:30 LT are displayed in Figs. 8b and c, respectively. In each figure there are four echo bands with a range separation of about 3 km, corresponding to the altitude separation of about 2 km. The $V_d$ values at around 96 km altitude are between +40 and +80 m/s, being higher than the FPI winds (eastward ~ 30 m/s) in Fig. 4b. Interestingly, very high $V_d$ of about 120 m/s appears at 90 km altitude in Fig. 8c. Relying on that $V_d$ at this altitude is almost equal to neutral wind speed because of the high collision frequency between electrons and neutrals, we can conclude that very strong neutral winds exceeding 100 m/s existed in a local area at 90 km altitude. The solid lines in Figs. 8b and c connect the peak position of each spectrum. As can be seen, the range (altitude) profiles of $V_d$ exhibit strong shears, which may be due to AGWs with sheared wind profiles.

4 Discussion

The results are summarized as follows:

1. MU radar QP echoes were observed at altitudes of 90–110 km before midnight on 5 August 2002. The minimum altitude (90 km) of the QP echoes was lower by 10 km than that of QP echoes previously observed by the MU radar, which made it possible to compare directly the behavior of 3.2-m FAIs, 557.7-nm airglow, and neutral winds at around 96 km altitude.

2. Motion of QP echo regions had an eastward component. This motion was consistent with the motions of neutral wind, airglow structures, and FAIs. Slopes of QP echo striations in RTI plots were positive before about 20:45 LT and negative after that. This change was caused by a change in the neutral wind direction from ENE to SE.

3. During the QP echo event, airglow intensity at 96 km altitude was enhanced up to 400 R, indicating an increase in the atomic oxygen density of 26%. The motion of the enhanced airglow region had an eastward component of about 65 m/s, which was consistent with the motion of airglow structures toward SE at about 40 m/s. Echo altitudes clearly went up and down in harmony with an enhancement and subsidence of airglow intensity, respectively.

4. Airglow structures with periods of 10 min to 1 h moved to SE at about 40 m/s. Echo intensity and Doppler velocity of FAIs at around 96 km altitude also fluctuated with these periods.

5. Some QP echo regions contained small-scale QP structures with a range separation of 3 km (altitude separation of about 2 km). Doppler spectra of the small-scale structures indicate that FAI velocities at around 96 km altitude (+40–+80 m/s) were higher than FPI wind (~30 m/s), and that very strong neutral winds exceeding 100 m/s existed in a local area around 90 km altitude. Range (altitude) profiles of $V_d$ exhibit strong shears.

The present QP echo event is a rare case because it is known that MU radar QP echoes before midnight usually appear above 100 km altitude (e.g. Yamamoto et al., 1991; Ogawa et al., 2002). In fact, the MU radar observations on the nights of 28 July to 6 August (except for 4 August) during the SEEK-2 campaign showed that the QP echoes appeared more or less on all the nights, and that the altitudes of typical QP echoes before midnight were always above 100 km, except for the 5 August case.
Regarding the second result finding, the movements of QP echo regions that have been observed by the MU radar are mostly toward SW (e.g. Yamamoto et al., 1994, 1997; Ogawa et al., 1995, 2002). Such a movement has also been detected by the HF and VHF radars in Tanegashima during SEEK–2 (Saito et al., 2005) and at other locations (e.g. Hyse, 2004; Hussey et al., 2004; Tsunoda et al., 2004 and references therein). In the present case, however, the QP echo regions moved to the east. This movement is consistent with the results from the simultaneous observations of Doppler velocity, airglow structures, and neutral wind, which suggests that the $E_s$-layers, including the 3.2-m FAIs, moved together with the neutral wind. Monitoring of some FM-broadcasting circuits at Wakkanai (northernmost part of Japan) indicated that the $E_s$-layers at around 40°N on the night of 5 August moved to the east (S. Saito, private communication). Ogawa et al. (1995) reported a case study in which QP echo regions first moved to the east and then to the west. The eastward neutral wind at around 96 km, above which a thin $E_i$ is located (as inferred from Fig. 5), together with speculated westward winds above $E_s$, is suitable for the generation of $E_s$ through the well-known wind-shear mechanism.

Cosgrove and Tsunoda (2002) predicted that a mid-latitude $E_i$ at a zonal wind-shear node is unstable at night to plane wave distortions. In the Northern Hemisphere the growth rate of the instability is positive (negative) for wave phase fronts with azimuthal alignment in the NW–SE (NE–SW) direction, and zero for phase fronts aligned E–W or N–S (Tsunoda et al., 2004). Then Cosgrove and Tsunoda (2003) showed results from numerical simulations confirming this azimuth-dependent $E_i$ instability (see also Yokoyama et al., 2004a). The southwestward movement of QP echo regions with phase fronts aligned NW–SE has been documented in many papers, and seems consistent with the theory of Cosgrove and Tsunoda (2002). However, our QP echoes on 5 August moved to the east, perhaps with phase front alignment in the N–S direction, which seems unsuitable for the Cosgrove and Tsunoda theory.

The $V_d$ values were mostly less than ±150 m/s (Fig. 6). FAIs can be generated by the gradient-drift instability under the combined action of an electric field and electron density gradient, the two-stream instability due to strong electric field, or a wind-driven gradient-drift instability. $V_d$ of FAIs produced through the gradient-drift instability is slower than the $E \times B$ electron drift velocity by a factor $(1 + \psi)$, where $\psi = \nu_e/\omega_e\alpha_e$; $\nu_e(v_i)$ is the electron (ion) collision frequency and $\omega_e(\omega_i)$ is the electron (ion) cyclotron frequency (e.g. Tanaka and Venkateswaran, 1982). This factor increases rapidly with decreasing altitude; 0.01 at 108 km, 0.1 at 100 km, 1.0 at 93 km, and 10 at 87 km. Hence, $V_d$ observed at altitudes lower than, say, 100 km does not always represent the electron drift velocity: see Ogawa et al. (2002) in detail. The very high $V_d$ values exceeding 100 m/s at 90 km altitude (item 5) were observed only on beam 5 but not on other beams (Fig. 6) and are perhaps caused by localized, high-speed neutral winds. Very strong neutral winds exceeding 100 m/s at around 90 km altitude were also observed by Larsen (2000b). Our FPI, however, which is most sensitive to the 557.7-nm airglow at around 96 km and can observe winds averaged over a large area above Shigaraki, could not detect such a strong wind in a local area at 90 km. Urbina et al. (2004) believe that the gradient-drift instability mechanism does not play a primary role in the generation of low-altitude (~92 km) radar echo regions drifted at about 100 m/s although a secondary gradient-drift process is probably the cause of Bragg-scale electron density waves.

The eastward movement of the isolated bright airglow region is interesting (item 3). We suppose that the eastward neutral winds, which accompany vertical wind shears to form $E_s$, transported an enhanced $[O]$ from the west to the east, and that this region included an upward (downward) wind component to raise (lower) the altitude of the wind-shear node, namely, the QP echo altitudes.

Regarding the fourth result finding, past QP echo observations demonstrated that the radar echo intensity and Doppler velocity of FAIs usually fluctuates in space and time, perhaps due to AGWs in the neutral atmosphere (e.g. Ogawa et al., 1995). In this paper we have first showed that the temporal and spatial fluctuations of the echo intensity and Doppler velocity associated with QP echoes have counterparts in those of the 557.7-nm airglow in the lower thermosphere. Short-period AGWs causing these fluctuations may come from the lower atmosphere, as has first been demonstrated in numerical simulations by Horinouchi (2004). Moreover, recent simulations by Yokoyama et al. (2004b) have shown that AGWs from below modulate an $E_i$ to produce wave-like patterns (QP structures) of plasma density, and that southwestward-propagating gravity waves can create the QP echo structures whose phase fronts align from NW to SE, consistent with Cosgrove and Tsunoda’s theory (2002). The Yokoyama et al. (2004b) simulations suggest that eastward-propagating AGWs from below are saturated/dissipated at a critical level, located below a wind-shear node, due to the eastward wind. These eastward-propagating AGWs may account for the south-eastward propagating airglow patterns at 96 km altitude below the $E_i$-layer. Analyzing OI and OH airglows and $E$-region FAI data from Tanegashima (about 700 km southwest of Shigaraki) during SEEK–2, Onoma et al. (2005) conclude that the gravity waves seen in the airglow images during FAI events propagated mostly southeastward, which is consistent with our case.

Another candidate for the generation of QP-like echo patterns is the Kelvin-Helmholtz (K-H) instability caused by a sheared neutral wind profile (e.g. Larsen, 2000a, 2000b; Bernhardt, 2002), which does not always require a preferred phase front alignment of echo regions. Bernhardt (2002) have demonstrated from numerical simulations how K-H neutral turbulence evolves nonlinearly to modulate an $E_i$-layer and plasma density in and around $E_s$. The vortex structures (“K-H billows”) seen in the simulations may account for the fluctuations of airglow intensity, QP echo intensity, QP echo altitudes, and Doppler velocity of FAIs. Sripathi et al. (2003) suggested the K-H instability as a possible source
for the daytime QP echoes in the lower $E$-region. Short-period AGWs that interact with an $E_t$ may form the QP structures (e.g., Kagan, 2002). We need more investigations to know which process, K-H billows, AGWs from below, etc., is the primary cause for the fluctuations of the parameters that we observed.

Our fifth result finding strongly suggests that two kinds of atmospheric waves, one causing the positive QP striations and the other causing the negative striations in an RTI plot, might exist in the same region. At this stage we do not know the generation mechanism of these low-altitude echoes, though some hints have been given by Urbina et al. (2004).

5 Conclusions

On the night of 5 August 2002 during the SEEK–2 campaign, we made, for the first time, coordinated observations of QP radar echoes, lower thermospheric OI 557.7-nm airglow, and neutral winds in a common volume over Shigaraki. The QP echoes appeared at altitudes of 90–110 km, which were lower by 10 km than usual QP echo altitudes, enabling us to study the relationship among FAIs, short-period atmospheric disturbances, and neutral winds at around 96 km altitude. The QP echo regions moved toward the east, contrary to the southwestward movements that have been often observed by the MU radar and radars at other locations. Temporal and spatial fluctuations of echo intensity and Doppler velocity associated with the QP echoes had counterparts in those of the 557.7-nm airglow at around 96 km in the lower thermosphere.

The eastward movement of the QP echo regions was consistent with the motions of neutral winds, airglow structures, and FAIs, a very reasonable result expected when the electrodynamics of the $E_t$ is mainly controlled by the neutral atmospheric dynamics. During the QP echo event, 1) airglow intensity at 96 km altitude was enhanced up to 400 R (increase of 26% in [O]), 2) the enhanced airglow region moved to the east, contrary to the southwestward movements that have been often observed by the MU radar and radars at other locations. Temporal and spatial fluctuations of echo intensity and Doppler velocity associated with the QP echoes had counterparts in those of the 557.7-nm airglow at around 96 km in the lower thermosphere.

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