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# Regular Solutions for Landau-Lifschitz Equation in $\mathbb{R}^3$

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**ABSTRACT** - In this paper we prove local existence, global existence with small data and uniqueness of regular solutions for Landau-Lifschitz equations. Furthermore we establish local existence and uniqueness for a system coupling Maxwell and Landau-Lifschitz equations arising from Micromagnetism theory.

**AMS Subject Classification.** 35K15, 35Q60

## 1 Introduction

In Micromagnetism theory (see [2]), the behaviour of a ferromagnet is represented by an unitary vector field  $u$  called magnetic moment. The variations of  $u$  are governed by Landau-Lifshitz equation

$$\begin{cases} \frac{\partial u}{\partial t} = u \wedge \Delta u - u \wedge (u \wedge \Delta u) \\ u(0, \cdot) = u_0(\cdot) \end{cases} \quad (1.1)$$

where we assume that  $|u_0| = 1$ .

Existence and non uniqueness for weak solutions of (1.1) are proved by F. Alouges and A. Soyeur in [1]. Furthermore, P.L. Sulem, C. Sulem and C. Bardos establish in [6] local existence of regular solutions for the equation

$$\frac{\partial u}{\partial t} = u \wedge \Delta u. \quad (1.2)$$

Concerning equation (1.1) we prove the following theorems.

**Theorem 1.** *Assume that*

$$|u_0| = 1, \quad \nabla u_0 \in \mathbb{H}^1(\mathbb{R}^3).$$

*Then there exists  $T > 0$ , there exists an unique  $u$  such that*

- (i)  $u \in \mathbb{L}^\infty((0, T) \times \mathbb{R}^3), \quad |u| = 1,$
- (ii)  $\nabla u \in L^\infty((0, T); \mathbb{H}^1(\mathbb{R}^3)) \cap L^2((0, T); \mathbb{H}^2(\mathbb{R}^3)),$
- (iii)  $u$  satisfies (1.1).

**Theorem 2.** *There exists  $\delta > 0$  such that if  $|u_0| = 1$  and if  $\nabla u_0 \in \mathbb{H}^1(\mathbb{R}^3)$  with  $\|\nabla u_0\|_{\mathbb{H}^1(\mathbb{R}^3)} < \delta$ , then the solution  $u$  given by Theorem 1 exists for  $T = +\infty$ .*

The term  $-u \wedge (u \wedge \Delta u)$  in (1.1) is in fact a dissipation term. For this reason, global existence with small data is valid for (1.1) and not for (1.2) (see [6]).

On the other hand, the propagation of electromagnetic waves in the ferromagnet is governed by a system coupling Landau-Lifschitz and Maxwell equations.

$$\frac{\partial u}{\partial t} = u \wedge (\Delta u + H) - u \wedge (u \wedge (\Delta u + H)),$$

$$\frac{\partial B}{\partial t} + \operatorname{curl} E = 0,$$

$$\frac{\partial E}{\partial t} - \operatorname{curl} H = 0,$$

$$B = H + u,$$

$$u(0, \cdot) = u_0(\cdot), \quad B(0, \cdot) = B_0(\cdot), \quad E(0, \cdot) = E_0(\cdot),$$

where we assume that  $|u_0| = 1$  and  $\operatorname{div} B_0 = 0$ .

Instead of working with  $(u, H, E)$ , we will write the system with the unknowns  $(u, B, E)$ .

$$\frac{\partial u}{\partial t} = u \wedge (\Delta u + B) - u \wedge (u \wedge (\Delta u + B)) \tag{1.3}$$

$$\frac{\partial B}{\partial t} = -\operatorname{curl} E \tag{1.4}$$

$$\frac{\partial E}{\partial t} = \operatorname{curl} B - \operatorname{curl} u \tag{1.5}$$

$$u(0, \cdot) = u_0(\cdot), \quad B(0, \cdot) = B_0(\cdot), \quad E(0, \cdot) = E_0(\cdot) \tag{1.6}$$

and we still assume that

$$|u_0| = 1 \text{ and } \operatorname{div} B_0 = 0 \text{ in } \mathbb{R}^3. \tag{1.7}$$

We will establish the following theorem.

**Theorem 3.** Let  $(u_0, E_0, B_0)$  such that

$$\begin{aligned} E_0 \in \mathbb{H}^1(\mathbb{R}^3), \quad B_0 \in \mathbb{H}^1(\mathbb{R}^3), \quad \operatorname{div} B_0 = 0, \\ |u_0| = 1, \quad \nabla u_0 \in \mathbb{H}^1(\mathbb{R}^3). \end{aligned}$$

Then there exists  $T > 0$ , there exists an unique  $(u, E, B)$  such that

- (i)  $|u| = 1, \quad \nabla u \in L^\infty((0, T); \mathbb{H}^1(\mathbb{R}^3)) \cap L^2((0, T); \mathbb{H}^2(\mathbb{R}^3)),$
- (ii)  $E$  and  $B$  belong to  $L^\infty((0, T); \mathbb{H}^1(\mathbb{R}^3)),$
- (iii)  $(u, E, B)$  satisfies (1.3)-(1.6).

The existence of weak solutions for the system (1.3)-(1.6) is proved in [7] and in [3] in the case of a bounded domain.

In [4], J.L. Joly, G. Métivier and J. Rauch prove the existence of solutions for a system similar to (1.3)-(1.6) but without  $\Delta u$  in (1.3).

The asymptotic behaviour of weak solutions of (1.3)-(1.6) in a bounded domain is studied in [3].

The proof of Theorems 1, 2 and 3 is based on a semi-discretization used in [1] and [6].

In Part 2, we describe the discretization process. Part 3 is devoted to the proof of Theorems 1 and 2. Theorem 3 is proved in the last part.

## 2 Discretization space, notations

We fix  $h > 0$  and we set  $x_j^h = jh$  for  $j \in \mathbb{Z}$ .

For  $\alpha = (i, j, k) \in \mathbb{Z}^3$ , we note  $X_\alpha^h = (x_i^h, x_j^h, x_k^h)$  and

$$C_\alpha^h = \left\{ (x, y, z), x_i^h \leq x < x_i^h + h, x_j^h \leq y < x_j^h + h, x_k^h \leq z < x_k^h + h \right\}.$$

We denote

$$Z_h^3 = \left\{ X_\alpha^h \in \mathbb{R}^3, \alpha \in \mathbb{Z}^3 \right\}.$$

In the sequel of this part, in order to simplify the notations, we will omit the exponent  $h$ .

We consider the following operators defined for  $u : Z^3 \rightarrow \mathbb{R}^3$  :

$$\tau_1^+ u(x_i, x_j, x_k) = u(x_{i+1}, x_j, x_k)$$

$$D_1^+ u = \frac{1}{h} (\tau_1^+ u - u)$$

$$\tau_1^- = (\tau_1^+)^{-1}, \quad D_1^- = \tau_1^- \circ D_1^+$$

In the same way, we denote  $\tau_2^+$ ,  $D_2^+$ ,  $\tau_2^-$ ,  $D_2^-$  the same operations concerning the second variable, and  $\tau_3^+$ ,  $D_3^+$ ,  $\tau_3^-$ ,  $D_3^-$  for the third variable.

We set

$$\tilde{\Delta} = \sum_{i=1}^3 D_i^- D_i^+,$$

$$\widetilde{\text{div}} u = \sum_{i=1}^3 D_i^+ u^i$$

$$\widetilde{\text{curl}} u = (D_2^+ u^3 - D_3^+ u^2, D_3^+ u^1 - D_1^+ u^3, D_1^+ u^2 - D_2^+ u^1).$$

We denote

$$\int_{Z^3} u = \sum_{\alpha \in \mathbb{Z}^3} h^3 u(X_\alpha).$$

and we use the following classical notations

$$\|u\|_{l^p} = \left( \int_{Z^3} |u|^p \right)^{\frac{1}{p}}, \quad \|D^+ u\|_{l^p} = \left( \sum_{i=1}^3 \|D_i^+ u\|_{l^p}^p \right)^{\frac{1}{p}},$$

$$\|u\|_{w^{1,p}} = (\|u\|_{l^p}^p + \|D^+ u\|_{l^p}^p)^{\frac{1}{p}}, \quad \|u\|_{h^1} = \|u\|_{w^{1,2}},$$

$$\|u\|_{h^2} = \left( \|u\|_{h^1}^2 + \sum_{ij} \|D_i^+ D_j^+ u\|_{l^2}^2 \right)^{\frac{1}{2}}, \quad \|u\|_{l^\infty} = \sup_{\alpha \in \mathbb{Z}^3} |u(X_\alpha)|.$$

We remark that

$$\int_{Z^3} D_i^+ u \cdot v = - \int_{Z^3} u \cdot D_i^- v,$$

furthermore,

$$D_i^+ (uv) = D_i^+ u \tau_i^+ v + u D_i^+ v.$$

We recall now the following discrete Sobolev inequalities.

**Lemma 1.** *There exists a constant  $C$  independent of  $h$  such that for all  $u : Z^3 \rightarrow \mathbb{R}^3$ ,*

$$\text{if } \|u\|_{h^1} < +\infty, \text{ then } \|u\|_{l^6(Z^3)} \leq C \|D^+ u\|_{l^2(Z^3)}, \quad (2.1)$$

Discrete versions of Sobolev inequalities are established in [5] using interpolation procedures.

The same interpolation process is used in this paper (see 3.3 and 4.3) and is presented in Section 3.3.

### 3 Proof of theorems 1 and 2

#### 1. Discretization

For  $h > 0$ , let  $u_0^h$  defined on the mesh  $Z_h^3$  such that

- $|u_0^h| = 1$  on  $Z_h^3$ ,
- $r_h u_0^h$  tends to  $u_0$  in  $L_{loc}^2(\mathbb{R}^3)$
- $\alpha \|D^+ u_0^h\|_{h^1(Z_h^3)} \leq \|\nabla u_0\|_{H^1(\mathbb{R}^3)} \leq \frac{1}{\alpha} \|D^+ u_0^h\|_{h^1(Z_h^3)}$

where  $r_h$  is the interpolating operator defined in [5] and in Section 3.3 of this paper, and where  $\alpha$  does not depend on  $h$ .

Now we fix  $h > 0$  and we solve

$$\begin{cases} \frac{du^h}{dt} = u^h \wedge \tilde{\Delta} u^h - u^h \wedge (u^h \wedge \tilde{\Delta} u^h) \\ u^h(t=0) = u_0^h \end{cases} \quad (3.1)$$

The map  $u \mapsto u \wedge \tilde{\Delta} u - u \wedge (u \wedge \tilde{\Delta} u)$  is locally Lipschitz in  $l^\infty(Z_h^3)$ , so there exists an unique solution of (3.1) with Cauchy-Lipschitz Theorem.

In order to simplify the notation, we will omit the exponent  $h$  in the computations of the following subsection.

#### 2. Estimates

We multiply (3.1) by  $u$  and we obtain that

$$\frac{d}{dt} |u|^2 = 0,$$

hence  $|u| = 1$  on  $Z^3$ , as it is the case for  $u_0$ .

With the above remark we can modify the form of the equation. We first note that

$$u \wedge (u \wedge \tilde{\Delta} u) = (u \cdot \tilde{\Delta} u)u - \tilde{\Delta} u.$$

Furthermore, writting that  $\tilde{\Delta} |u|^2 = 0$  as  $|u| = 1$ , we obtain that

$$2u \cdot \tilde{\Delta} u + |D^+ u|^2 + |D^- u|^2 = 0.$$

Hence, Equation (3.1) takes the form

$$\frac{du}{dt} = u \wedge \tilde{\Delta}u + \tilde{\Delta}u + \frac{1}{2}(|D^+u|^2 + |D^-u|^2)u \quad (3.2)$$

**First estimate.** We multiply (3.2) by  $\tilde{\Delta}u$  and after summation on  $Z^3$ , we get

$$-\frac{1}{2} \frac{d}{dt} \|D^+u\|_{l^2}^2 = \|\tilde{\Delta}u\|_{l^2}^2 + \frac{1}{2} \int_{Z^3} (|D^-u|^2 + |D^+u|^2) u \cdot \tilde{\Delta}u,$$

so we obtain

$$\frac{d}{dt} \|D^+u\|_{l^2}^2 + 2\|\tilde{\Delta}u\|_{l^2}^2 \leq 2\|D^+u\|_{l^4}^2 \|\tilde{\Delta}u\|_{l^2} \quad (3.3)$$

**Second estimate.** We multiply now (3.2) by  $\tilde{\Delta}^2u$  and after summation on  $Z^3$  we obtain

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|\tilde{\Delta}u\|_{l^2}^2 &= \int_{Z^3} u \wedge \tilde{\Delta}u \cdot \tilde{\Delta}^2u - \|D^+\tilde{\Delta}u\|_{l^2}^2 \\ &+ \frac{1}{2} \int_{Z^3} (|D^+u|^2 + |D^-u|^2) u \cdot \tilde{\Delta}^2u. \end{aligned}$$

We remark now that

$$\begin{aligned} \left| \int_{Z^3} u \wedge \tilde{\Delta}u \cdot \tilde{\Delta}^2u \right| &= \left| \sum_i \int_{Z^3} D_i^+(u \wedge \tilde{\Delta}u) \cdot D_i^+\tilde{\Delta}u \right| \\ &= \left| \sum_i \int_{Z^3} D_i^+u \wedge \tilde{\Delta}u \cdot D_i^+\tilde{\Delta}u \right| \\ &\leq \|D^+\tilde{\Delta}u\|_{l^2} \|\tilde{\Delta}u\|_{l^3} \|D^+u\|_{l^6} \end{aligned} \quad (3.4)$$

Furthermore,

$$\begin{aligned} \left| \frac{1}{2} \int_{Z^3} (|D^-u|^2 + |D^+u|^2) u \cdot \tilde{\Delta}^2u \right| &= \left| \frac{1}{2} \sum_i \int_{Z^3} D_i^+ \left( (|D^-u|^2 + |D^+u|^2)u \right) \cdot D_i^+\tilde{\Delta}u \right| \\ &\leq \frac{1}{2} \|D^+ \left( (|D^-u|^2 + |D^+u|^2)u \right)\|_{l^2} \|D^+\tilde{\Delta}u\|_{l^2}. \end{aligned}$$

Now we compute  $D_i^+ (|D^+u|^2)$  and we obtain that

$$\|D^+ \left( (|D^-u|^2 + |D^+u|^2)u \right)\|_{l^2} \leq C \left( \sum_{ij} \|D_i^+ D_j^+ u \cdot D_j^+ u\|_{l^2} + \| |D^+u|^3 \|_{l^2} \right)$$

as  $|u| = 1$ .

Thus there exists a constant  $K$  independant of  $h$  such that

$$\frac{d}{dt} \|\tilde{\Delta}u\|_{l^2}^2 + 2\|D^+\tilde{\Delta}u\|_{l^2}^2 \leq +K \left( \|\tilde{\Delta}u\|_{l^3} \|D^+u\|_{l^6} + \|D^+u\|_{l^6}^3 \right) \|D^+\tilde{\Delta}u\|_{l^2} \quad (3.5)$$

With Lemma 2.1 and by interpolation there exists a universal constant  $C$  such that

$$\begin{aligned} \|D^+u\|_{l^6} &\leq C \|\tilde{\Delta}u\|_{l^2} \\ \|\tilde{\Delta}u\|_{l^3} &\leq C \|\tilde{\Delta}u\|_{l^2}^{\frac{1}{2}} \|D^+\tilde{\Delta}u\|_{l^2}^{\frac{1}{2}} \end{aligned} \quad (3.6)$$

From (3.5) and (3.6) we deduce

$$\frac{d}{dt}\|\tilde{\Delta}u\|_{l^2}^2 + 2\|D^+\tilde{\Delta}u\|_{l^2}^2 \leq K \left( \|D^+\tilde{\Delta}u\|_{l^2}^{\frac{3}{2}} \|\tilde{\Delta}u\|_{l^2}^{\frac{3}{2}} + \|D^+\tilde{\Delta}u\|_{l^2} \|\tilde{\Delta}u\|_{l^2}^3 \right) \quad (3.7)$$

**Estimate for Theorem 1.**

We absorb  $\|D^+\tilde{\Delta}u\|_{l^2}$  in (3.7) and we obtain

$$\frac{d}{dt}\|\tilde{\Delta}u\|_{l^2}^2 + \|D^+\tilde{\Delta}u\|_{l^2}^2 \leq K\|\tilde{\Delta}u\|_{l^2}^6 \quad (3.8)$$

On the other hand, from (3.3), by interpolation in  $\mathbb{L}^4$ , we derive

$$\frac{d}{dt}\|D^+u\|_{l^2}^2 + \|\tilde{\Delta}u\|_{l^2}^2 \leq K\|D^+u\|_{l^2} \|D^+u\|_{l^6}^3 \leq K\|D^+u\|_{l^2} \|\tilde{\Delta}u\|_{l^2}^3 \quad (3.9)$$

Combining (3.8) and (3.9) we obtain

$$\frac{d}{dt}\|D^+u\|_{h^1}^2 + \|\tilde{\Delta}u\|_{h^1}^2 \leq K \left( 1 + \|D^+u\|_{h^1}^6 \right) \quad (3.10)$$

We set now  $g(t) = \|D^+u\|_{h^1}^2$  and we have

$$\frac{dg}{dt} \leq K(1 + g^3),$$

hence there exist  $T > 0$  and  $K$  independant of  $h$  such that

$$\begin{cases} \|D^+u\|_{L^\infty(0,T;h^1)} \leq K, \\ \|\tilde{\Delta}u\|_{L^2(0,T;h^1)} \leq K, \\ \left\| \frac{du}{dt} \right\|_{L^\infty(0,T;l^2)} \leq K. \end{cases} \quad (3.11)$$

The last estimate is obtained using (3.1) and the previous estimate concerning  $D^+u$ .

**Estimate for Theorem 2.**

We absorb  $\|D^+\tilde{\Delta}u\|_{l^2}$  only in the first term of the right hand-side of (3.7) writting

$$\|D^+\tilde{\Delta}u\|_{l^2}^{\frac{3}{2}} \|\tilde{\Delta}u\|_{l^2}^{\frac{3}{2}} \leq \frac{1}{2} \left( \|D^+\tilde{\Delta}u\|_{l^2}^2 + \|D^+\tilde{\Delta}u\|_{l^2} \|\tilde{\Delta}u\|_{l^2}^3 \right),$$

and we obtain

$$\frac{d}{dt}\|\tilde{\Delta}u\|_{l^2}^2 + \|D^+\tilde{\Delta}u\|_{l^2}^2 \leq K\|\tilde{\Delta}u\|_{l^2}^3 \|D^+\tilde{\Delta}u\|_{l^2} \quad (3.12)$$

Combining (3.9) and (3.12) we derive

$$\frac{d}{dt}\|D^+u\|_{h^1}^2 + \|\tilde{\Delta}u\|_{h^1}^2 \leq K \left( \|\tilde{\Delta}u\|_{l^2}^3 \|D^+\tilde{\Delta}u\|_{l^2} + \|D^+u\|_{l^2}^2 \|\tilde{\Delta}u\|_{l^2}^2 \right)$$

Hence there exists a constant  $K$  independant of  $h$  such that

$$\frac{d}{dt}\|D^+u\|_{h^1}^2 + \|\tilde{\Delta}u\|_{h^1}^2 \leq K\|D^+u\|_{h^1}^2 \|\tilde{\Delta}u\|_{h^1}^2$$

thus

$$\frac{d}{dt} \|D^+ u\|_{h^1}^2 + \|\tilde{\Delta} u\|_{h^1}^2 (1 - K \|D^+ u\|_{h^1}^2) \leq 0 \quad (3.13)$$

We set now  $\delta = \frac{1}{\sqrt{K}}$  and we suppose that  $\|D^+ u_0\|_{h^1} < \delta$ .

We claim that for all  $t \geq 0$ ,  $\|D^+ u(t)\|_{h^1} < \delta$ .

If it is not the case, then let  $t_1$  be the first  $t > 0$  such that  $\|D^+ u(t)\|_{h^1} \geq \delta$ .

For all  $t < t_1$ ,

$$1 - K \|D^+ u\|_{h^1}^2 \geq 0,$$

hence, for all  $t < t_1$ ,

$$\frac{d}{dt} \|D^+ u\|_{h^1}^2 \leq 0, \text{ so } \|D^+ u\|_{h^1}^2(t_1) \leq \|D^+ u(0)\|_{h^1}^2 < \delta$$

which leads to a contradiction.

Therefore, if  $\|D^+ u_0\| < \delta$ ,

$$\|D^+ u\|_{L^\infty(0, +\infty; h^1)} < \delta \quad (3.14)$$

and from (3.13) we deduce that there exists  $K$  such that

$$\|\tilde{\Delta} u\|_{L^2(0, +\infty; h^1)} \leq K \quad (3.15)$$

### 3. Limit when $h$ goes to zero

Let us prove Theorem 1.

In the preceding subsection, for all  $h > 0$  we have constructed a solution  $u^h$  of (3.1) defined on the mesh  $Z_h^3$  which satisfies (3.1) and (3.11).

We extend  $u^h$  to the whole space using an interpolation process described in [5], p. 224.

We introduce the following interpolating operators :

For  $X = (x_1, x_2, x_3) \in C_\alpha^h$ , if we note  $X_\alpha^h = (x_1^h, x_2^h, x_3^h)$ , we set

- $r_h u^h(X) = u^h(X_\alpha^h),$
- $p_h u^h(X) = u^h(X_\alpha^h) + \sum_{i=1}^3 D_i^+ u^h(X_\alpha^h)(x_i - x_i^h)$   
 $+ \sum_{1 \leq i < j \leq 3} D_i^+ D_j^+(X_\alpha^h)(x_i - x_i^h)(x_j - x_j^h) + D_1^+ D_2^+ D_3^+ u^h(X_\alpha^h) \prod_{i=1}^3 (x_i - x_i^h),$
- $q_h^k u^h(X) = u^h(X_\alpha^h) + \sum_{i \neq k} D_i^+ u^h(X_\alpha^h)(x_i - x_i^h)$   
 $+ \sum_{\substack{1 \leq i < j \leq 3 \\ i, j \neq k}} D_i^+ D_j^+(X_\alpha^h)(x_i - x_i^h)(x_j - x_j^h).$

We recall that

$$\frac{\partial}{\partial x_i} (p_h u^h) = q_h^i (D_i^+ u^h).$$

Furthermore we have the following proposition proved in [5].



**Proposition 1.** *If one of the interpolates  $p_h u^h$ ,  $q_h u^h$ , or  $r_h u^h$  converges strongly (resp. weakly) in  $L^2$  when  $h$  goes to zero, then the two others also converge to the same limit in  $L^2$  strongly (resp. weakly).*

The estimate (3.11) gives that there exists  $K > 0$  such that for all  $h > 0$ ,

$$\|r_h(D_i^+ D_j^+ D_k^+ u^h)\|_{L^2((0,T) \times \mathbb{R}^3)} \leq K$$

$$\|r_h(D_i^+ D_j^+ u^h)\|_{L^2((0,T) \times \mathbb{R}^3)} \leq K$$

$$\|r_h(D_i^+ u^h)\|_{L^2((0,T) \times \mathbb{R}^3)} \leq K$$

Furthermore  $r_h u^h$  is bounded in  $L^2_{loc}$  independently of  $h > 0$ . Thus, up to subsequences, we deduce that when  $h$  goes to zero,

$$r_h u^h \rightharpoonup u \text{ in } L^2_{loc} \text{ weakly,}$$

$$r_h(D_i^+ u^h) \rightharpoonup v_i \text{ in } L^2 \text{ weakly,}$$

$$r_h(D_i^+ D_j^+ u^h) \rightharpoonup w_{ij} \text{ in } L^2 \text{ weakly,}$$

$$r_h(D_i^+ D_j^+ D_k^+ u^h) \rightharpoonup \omega_{ijk} \text{ in } L^2 \text{ weakly,}$$

$$r_h\left(\frac{du^h}{dt}\right) \rightharpoonup f \text{ in } L^2 \text{ weakly.}$$

Now with Proposition 1.,  $q_h^i(D_i^+(D_j^+ D_k^+ u^h))$  and  $r_h(D_i^+ D_j^+ D_k^+ u^h)$  have the same limit  $\omega_{ijk}$ , and since  $q_h^i(D_i^+(D_j^+ D_k^+ u^h)) = \frac{\partial}{\partial x_i}(p_h(D_j^+ D_k^+ u^h))$ , as  $p_h(D_j^+ D_k^+ u^h)$  tends to  $w_{jk}$  in  $L^2$  weak, we deduce by uniqueness of the limit in  $\mathcal{D}'$  that

$$\omega_{ijk} = \frac{\partial}{\partial x_i} w_{jk}.$$

With the same raisonnement we deduce that

$$v_i = \frac{\partial u}{\partial x_i}, \quad w_{ij} = \frac{\partial^2 u}{\partial x_i \partial x_j}, \quad \omega_{ijk} = \frac{\partial^3 u}{\partial x_i \partial x_j \partial x_k}.$$

In addition, we have  $f = \frac{\partial u}{\partial t}$  and, since  $\frac{\partial}{\partial t} p_h u^h \rightharpoonup \frac{\partial u}{\partial t}$  and  $\frac{\partial}{\partial x_i} p_h u^h \rightharpoonup \frac{\partial u}{\partial x_i}$  in  $L^2$  weakly, since  $p_h u^h \rightharpoonup u$  in  $L^2_{loc}$  weakly, we deduce that

$$p_h u^h \rightarrow u \text{ in } L^p_{loc} \text{ strongly for } 2 \leq p < 6,$$

using the compactness of Sobolev embeddings in bounded domains.

In order to prove that  $u$  satisfies (1.1), we take  $\varphi \in \mathcal{D}((0, T) \times \mathbb{R}^3)$  and we introduce  $\Omega$  such that  $\varphi$  is zero outside of  $\Omega$ . We have

$$\int_{\Omega} r_h\left(\frac{du^h}{dt}\right) \cdot \varphi = \int_{\Omega} r_h u^h \wedge r^h \tilde{\Delta} u^h \cdot \varphi - \int_{\Omega} r^h u^h \wedge (r^h u^h \wedge r^h \tilde{\Delta} u^h) \cdot \varphi \quad (3.16)$$

Now,

$$r_h \frac{du^h}{dt} \rightharpoonup \frac{\partial u}{\partial t} \quad \text{in } L^2(\mathbb{R}^3) \text{ weakly,}$$

$$r_h \tilde{\Delta} u^h \rightharpoonup \Delta u \quad \text{in } L^2(\mathbb{R}^3) \text{ weakly,}$$

$$r_h u^h \rightarrow u \quad \text{in } L^2(\Omega) \text{ strongly (for the first term),}$$

$$r_h u^h \rightarrow u \quad \text{in } L^4(\Omega) \text{ strongly (for the second term).}$$

Thus we can take the limit in (3.16) and we obtain that  $u$  satisfies (1.1) in  $\mathcal{D}'(\mathbb{R}^3)$ .

Furthermore, by lower semicontinuity of the different norms, we obtain from (3.11) that  $u$  satisfies

$$\nabla u \in L^\infty((0, T); \mathcal{H}^1(\mathbb{R}^3)) \cap L^2((0, T); \mathcal{H}^2(\mathbb{R}^3)).$$

Finally, since  $r_h u^h \rightarrow u$  in  $L^2_{loc}$  strongly, by extracting a subsequence,  $r_h u^h \rightarrow u$  a.e., hence  $|u| = 1$  as it is the case for  $r_h u^h$ .

Let us prove now the uniqueness of the solution of (1.1) satisfying (i) and (ii) in Theorem 1.

Let  $\tilde{u}$  be another solution, and let  $\bar{u} = u - \tilde{u}$ .

We have

$$\frac{\partial \bar{u}}{\partial t} = \Delta \bar{u} + \bar{u} \wedge \Delta u + \tilde{u} \wedge \Delta \bar{u} + |\nabla u|^2 \bar{u} - \nabla \bar{u} \cdot (\nabla u + \nabla \tilde{u}) \tilde{u} \quad (3.17)$$

We multiply (3.17) by  $\bar{u}$  and we obtain

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|\bar{u}\|_{L^2}^2 + \|\nabla \bar{u}\|_{L^2}^2 &\leq \|\nabla \bar{u}\|_{L^2} \|\bar{u}\|_{L^2} \|\nabla \tilde{u}\|_{L^\infty} + \|\bar{u}\|_{L^2}^2 \|\nabla u\|_{L^\infty}^2 \\ &\quad + \|\bar{u}\|_{L^2} \|\nabla \bar{u}\|_{L^2} (\|\nabla u\|_{L^\infty} + \|\nabla \tilde{u}\|_{L^\infty}). \end{aligned}$$

We absorb  $\|\nabla \bar{u}\|_{L^2}$  in the left hand-side of the inequality and we obtain

$$\frac{d}{dt} \|\bar{u}\|_{L^2}^2 + \|\nabla \bar{u}\|_{L^2}^2 \leq K \|\bar{u}\|_{L^2} (\|\nabla u\|_{L^\infty} + \|\nabla \tilde{u}\|_{L^\infty}).$$

Now since  $\nabla u$  and  $\nabla \tilde{u}$  belong to  $L^1(0, T; L^\infty)$  (with Sobolev injections), we can use Gronwall Lemma to conclude that  $\bar{u} = 0$ .

Therefore Theorem 1 is proved.

In the same way we prove Theorem 2, starting from Estimates (3.14) and (3.15).

1. Discretization.

For all  $h > 0$  we consider  $(u_0^h, E_0^h, B_0^h)$  defined on  $Z_h^3$  such that

- $|u_0^h| = 1, \quad r_h u_0^h \xrightarrow{h \rightarrow 0} u_0$  in  $L_{loc}^2(\mathbb{R}^3)$ ,
- $\alpha \|D^+ u_0^h\|_{h^1(Z_h^3)} \leq \|\nabla u_0\|_{H^1(\mathbb{R}^3)} \leq \frac{1}{\alpha} \|D^+ u_0^h\|_{h^1(Z_h^3)}$ ,
- $B_0^h \in l^\infty(Z_h^3), \quad r_h B_0^h \xrightarrow{h \rightarrow 0} B_0$  in  $L^2(\mathbb{R}^3)$ ,
- $\alpha \|B_0^h\|_{h^1(Z_h^3)} \leq \|\nabla B_0\|_{H^1(\mathbb{R}^3)} \leq \frac{1}{\alpha} \|B_0^h\|_{h^1(Z_h^3)}$ ,
- $H_0^h \in l^\infty(Z_h^3), \quad r_h H_0^h \xrightarrow{h \rightarrow 0} H_0$  in  $L^2(\mathbb{R}^3)$ ,
- $\alpha \|H_0^h\|_{h^1(Z_h^3)} \leq \|\nabla H_0\|_{H^1(\mathbb{R}^3)} \leq \frac{1}{\alpha} \|H_0^h\|_{h^1(Z_h^3)}$ ,

where  $\alpha$  does not depend on  $h$ .

We remark that  $B_0^h$  and  $E_0^h$  are bounded in  $l^\infty(Z_h^3)$  but not uniformly in  $h$ .

Now, for  $h$  fixed, we solve the system

$$\frac{du^h}{dt} = u^h \wedge (\tilde{\Delta} u^h + B^h) - u^h \wedge (u^h \wedge (\tilde{\Delta} u^h + B^h)) \quad (4.1)$$

$$\frac{dB^h}{dt} = -\widetilde{\text{curl}} E^h \quad (4.2)$$

$$\frac{dE^h}{dt} = \widetilde{\text{curl}} B^h - \widetilde{\text{curl}} u^h \quad (4.3)$$

$$u^h(t=0) = u_0^h, \quad B^h(t=0) = B_0^h, \quad E^h(t=0) = E_0^h \quad (4.4)$$

Using Cauchy-Lipschitz theorem, there exists a local solution of (4.1)-(4.4).

Multiplying (4.1) by  $u$  we prove that  $|u^h| = 1$  hence, we can write (4.1) on the form

$$\frac{du^h}{dt} = u^h \wedge (\tilde{\Delta} u^h + B^h) + \tilde{\Delta} u^h + \frac{1}{2}(|D^+ u^h|^2 + |D^- u^h|^2)u^h - u^h \wedge (u^h \wedge B^h) \quad (4.5)$$

Furthermore, we can eliminate  $E$  in (4.2)-(4.3) to obtain

$$\frac{d^2 B^h}{dt^2} - \tilde{\Delta} B^h = \widetilde{\text{curl}} \widetilde{\text{curl}} u^h \quad (4.6)$$

as  $\widetilde{\text{div}} B^h = 0$ .

In order to simplify the notations we will omit the exponent  $h$  in the computations of the following section.

## 2. Estimates.

**First Estimate.** We multiply (4.5) by  $\tilde{\Delta}u$  and after summation on  $Z^3$ , we get

$$-\frac{d}{dt}\|D^+u\|_{l^2}^2 = \|\tilde{\Delta}u\|_{l^2}^2 + \int_{Z^3} ((u \wedge B) - u \wedge (u \wedge B)) \cdot \tilde{\Delta}u + \frac{1}{2} \int_{Z^3} (|D^+u|^2 + |D^-u|^2) u \cdot \tilde{\Delta}u.$$

Hence

$$\frac{d}{dt}\|D^+u\|_{l^2}^2 + \|\tilde{\Delta}u\|_{l^2}^2 \leq C \left( \|B\|_{l^2} \|\tilde{\Delta}u\|_{l^2} + \|D^+u\|_{l^2} \|\tilde{\Delta}u\|_{l^2}^2 \right) \quad (4.7)$$

**Second Estimate.** We multiply (4.5) by  $\tilde{\Delta}^2u$  and we obtain

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|\tilde{\Delta}u\|_{l^2}^2 + \|D^+\tilde{\Delta}u\|_{l^2}^2 &= \int_{Z^3} u \wedge \tilde{\Delta}u \cdot \tilde{\Delta}^2u \\ &+ \frac{1}{2} \int_{Z^3} (|D^+u|^2 + |D^-u|^2) u \cdot \tilde{\Delta}^2u \\ &- \int_{Z^3} \sum_i D_i^+(u \wedge B) \cdot D_i^+\tilde{\Delta}u + \int_{Z^3} \sum_i D_i^+(u \wedge (u \wedge B)) \cdot D_i^+\tilde{\Delta}u. \end{aligned}$$

Now

$$\int_{Z^3} \sum_i D_i^+(u \wedge B) \cdot D_i^+\tilde{\Delta}u = \sum_i \int_{Z^3} \left( D_i^+u \wedge B + \tau_i^+u \wedge D_i^+B \right) \cdot D_i^+\tilde{\Delta}u,$$

thus

$$\left| \int_{Z^3} \sum_i D_i^+(u \wedge B) \cdot D_i^+\tilde{\Delta}u \right| \leq K \|D^+\tilde{\Delta}u\|_{l^2} \left( \|D^+u\|_{l^4} \|B\|_{l^4} + \|D^+B\|_{l^2} \right).$$

In the same way,

$$\left| \int_{Z^3} \sum_i D_i^+(u \wedge (u \wedge B)) \cdot D_i^+\tilde{\Delta}u \right| \leq K \|D^+\tilde{\Delta}u\|_{l^2} \left( \|D^+u\|_{l^4} \|B\|_{l^4} + \|D^+B\|_{l^2} \right).$$

We treat the first two terms as in part 3 and we get

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|\tilde{\Delta}u\|_{l^2}^2 + \|D^+\tilde{\Delta}u\|_{l^2}^2 &\leq K \|D^+\tilde{\Delta}u\|_{l^2} \left( \|D^+u\|_{l^4} \|B\|_{l^4} + \|D^+B\|_{l^2} \right) \\ &+ K \left( \|D^+\tilde{\Delta}u\|_{l^2}^{\frac{3}{2}} \|\tilde{\Delta}u\|_{l^2}^{\frac{3}{2}} + \|D^+\tilde{\Delta}u\|_{l^2} \|\tilde{\Delta}u\|_{l^2}^3 \right) \end{aligned} \quad (4.8)$$

Now, by interpolation and discrete Sobolev embeddings, and absorbing  $\|D^+\tilde{\Delta}u\|_{l^2}$  in the right hand-side of (4.8), we obtain

$$\begin{aligned} \frac{d}{dt} \|\tilde{\Delta}u\|_{l^2}^2 + \|D^+\tilde{\Delta}u\|_{l^2}^2 &\leq K \left( \|D^+u\|_{l^2} \|\tilde{\Delta}u\|_{l^2} \|B\|_{l^2} \|D^+B\|_{l^2} \right. \\ &\left. + \|D^+B\|_{l^2}^2 + \|\tilde{\Delta}u\|_{l^2}^6 \right) \end{aligned} \quad (4.9)$$

**Third estimate.** We multiply (4.2) by  $B$  and (4.3) by  $E$  and we obtain, as  $\|\widetilde{\text{curl}} u\|_{l^2} \leq C \|D^+u\|_{l^2}$ ,

$$\frac{d}{dt} \left( \|B\|_{l^2}^2 + \|E\|_{l^2}^2 \right) \leq K \|D^+u\|_{l^2} \|E\|_{l^2} \quad (4.10)$$

**Fourth estimate.** We multiply (4.6) by  $\frac{dB}{dt}$  and we get, since  $\|\widetilde{\text{curl}} \widetilde{\text{curl}} u\|_{l^2} \leq K\|\tilde{\Delta}u\|_{l^2}$ ,

$$\frac{d}{dt} \left( \left\| \frac{dB}{dt} \right\|_{l^2}^2 + \|D^+ B\|_{l^2}^2 \right) \leq K\|\tilde{\Delta}u\|_{l^2} \left\| \frac{dB}{dt} \right\|_{l^2} \quad (4.11)$$

Combining (4.7), (4.9), (4.10) and (4.11), if we denote

$$\mathcal{E}(t) = \left( \|D^+ u\|_{l^2}^2 + \|\tilde{\Delta}u\|_{l^2}^2 + \|E\|_{l^2}^2 + \|B\|_{l^2}^2 + \left\| \frac{dB}{dt} \right\|_{l^2}^2 + \|D^+ B\|_{l^2}^2 \right) (t),$$

we obtain

$$\frac{d\mathcal{E}}{dt} + \|\tilde{\Delta}u\|_{h^1}^2(t) \leq K(1 + \mathcal{E}^3).$$

Therefore, there exist  $T$  and  $K$  independent of  $h$  such that

$$\begin{aligned} \|D^+ u\|_{L^\infty(0,T;h^1)} &\leq K, \quad \|\tilde{\Delta}u\|_{L^2(0,T;h^1)} \leq K \\ \left\| \frac{du}{dt} \right\|_{L^\infty(0,T;l^2)} &\leq K \end{aligned} \quad (4.12)$$

$$\|B\|_{L^\infty(0,T;h^1)} \leq K, \quad \|E\|_{L^\infty(0,T;h^1)} \leq K$$

We note that we can estimate  $\|D^+ E\|_{l^2}$  since

$$\begin{aligned} \|D^+ E\|_{l^2} &= \|\widetilde{\text{div}} E\|_{l^2}^2 + \|\widetilde{\text{curl}} E\|_{l^2}^2 \\ &= \|\widetilde{\text{div}} E_0\|_{l^2}^2 + \left\| \frac{dB}{dt} \right\|_{l^2}^2. \end{aligned}$$

### 3. Limit when $h$ goes to zero and uniqueness.

As in Part 3, we extend the discrete solution  $(u^h, E^h, B^h)$  to the whole space and with the same arguments, we can take the limit when  $h$  goes to zero, using (4.12).

The limit  $(u, E, B)$  satisfies the properties (i), (ii), and (iii) announced in theorem 3.

Now let us prove the uniqueness of the regular solution for (1.3)-(1.6).

Let us consider  $(\tilde{u}, \tilde{E}, \tilde{B})$  another regular solution for (1.3)-(1.6). We set

$$\bar{u} = u - \tilde{u}, \quad \bar{E} = E - \tilde{E}, \quad \bar{B} = B - \tilde{B}.$$

We have

$$\frac{\partial \bar{B}}{\partial t} = -\text{curl } \bar{E} \quad (4.13)$$

$$\frac{\partial \bar{E}}{\partial t} = \text{curl } \bar{B} - \text{curl } \bar{u} \quad (4.14)$$

$$\begin{aligned} \frac{\partial \bar{u}}{\partial t} &= \Delta \bar{u} + \bar{u} \wedge \Delta u + \tilde{u} \wedge \Delta \bar{u} + \bar{u} \wedge B - \tilde{u} \wedge \bar{B} \\ &\quad + |\nabla u|^2 \bar{u} - \nabla \bar{u} \cdot (\nabla u + \nabla \tilde{u}) \tilde{u} \end{aligned} \quad (4.15)$$

$$-\bar{u} \wedge (u \wedge B) - \tilde{u} \wedge (\bar{u} \wedge B + \tilde{u} \wedge \bar{B}).$$

We multiply (4.15) by  $\bar{u}$  and we obtain

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|\bar{u}\|_{L^2}^2 + \|\nabla \bar{u}\|_{L^2}^2 &= \int \tilde{u} \wedge \Delta \bar{u} \cdot \bar{u} - \int \tilde{u} \wedge \bar{B} \cdot \bar{u} \\ &+ \int |\nabla u|^2 |\bar{u}|^2 - \int \nabla \bar{u} \cdot (\nabla u + \nabla \tilde{u})(\tilde{u} \cdot \bar{u}) \\ &- \int \tilde{u} \wedge (\bar{u} \wedge B) \cdot \bar{u} - \int \tilde{u} \wedge (\tilde{u} \wedge \bar{B}) \cdot \bar{u}. \end{aligned}$$

Hence, using that  $\|\bar{u}\|_{L^4}^2 \leq C \|\bar{u}\|_{L^2} \|\nabla \bar{u}\|_{L^2}$ , and absorbing the term  $\|\nabla \bar{u}\|_{L^2}$ , we obtain that

$$\frac{d}{dt} \|\bar{u}\|_{L^2}^2 + \|\nabla \bar{u}\|_{L^2}^2 \leq K \|\bar{u}\|_{L^2}^2 \left( \|\nabla \tilde{u}\|_{L^\infty}^2 + \|\nabla u\|_{L^\infty}^2 + \|B\|_{L^2}^2 + 1 \right) + \|\bar{B}\|_{L^2}^2 \quad (4.16)$$

Furthermore, multiplying (4.13) by  $\bar{B}$  and (4.14) by  $\bar{E}$ , we obtain

$$\frac{d}{dt} \left( \|\bar{E}\|_{L^2}^2 + \|\bar{B}\|_{L^2}^2 \right) \leq 2 \|\nabla \bar{u}\|_{L^2} \|\bar{E}\|_{L^2} \quad (4.17)$$

We set

$$\mathcal{E}(t) = \left( \|\bar{E}\|_{L^2}^2 + \|\bar{B}\|_{L^2}^2 + \|\bar{u}\|_{L^2}^2 \right) (t)$$

and

$$f(t) = \left( 1 + \|B\|_{L^2}^2 + \|\nabla u\|_{L^\infty}^2 + \|\nabla \tilde{u}\|_{L^\infty}^2 \right) (t).$$

Combining (4.16) and (4.17) and absorbing  $\|\nabla \bar{u}\|_{L^2}$ , we obtain

$$\frac{d}{dt} \mathcal{E}(t) \leq K f(t) \mathcal{E}(t).$$

We remark that  $f(t) \in L^1(0, T)$ , and with Gronwall Lemma, we conclude that  $\mathcal{E} = 0$ .

Therefore, Theorem 3 is proved.

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