Local limit theorems for Brownian additive functionals and penalisation of Brownian paths, IX
Bernard Roynette, Marc Yor

To cite this version:
Bernard Roynette, Marc Yor. Local limit theorems for Brownian additive functionals and penalisation of Brownian paths, IX. Prépublication IECN 2008/10. 2008. <hal-00257593>
Local limit theorems for Brownian additive functionals and penalisation of Brownian paths, IX

B. Roynette\(^{(1)}\), M. Yor\(^{(2),(3)}\)

12/02/08

\(^{(1)}\) Institut Elie Cartan, Université Henri Poincaré, B.P. 239, 54506 Vandoeuvre les Nancy Cedex

\(^{(2)}\) Laboratoire de Probabilités et Modèles Aléatoires, Université Paris VI et VII, 4 place Jussieu - Case 188 F - 75252 Paris Cedex 05

\(^{(3)}\) Institut Universitaire de France

**Abstract** We obtain a local limit theorem for the laws of a class of Brownian additive functionals and we apply this result to a penalisation problem. We study precisely the case of the additive functional : \( A_t^- := \int_0^t 1_{X_s < 0} ds, t \geq 0 \). On the other hand, we describe Feynman-Kac type penalisation results for long Brownian bridges thus completing some similar previous study for standard Brownian motion (see [RVY, I]).

**Key words** Limit theorems for additive functionals, Feynman-Kac functionals, long Brownian bridges.

**2000 Mathematics Subject Classification** 60 F 17, 60 G 44, 60 J 25, 60 J 35, 60 J 55, 60 J 57, 60 J 60, 60 J 65.
0 Notations and introduction

0.1 Notations

- \( \Omega, (X_t, \mathcal{F}_t)_{t \geq 0}, \mathcal{F}_\infty = \vee_{t \geq 0} \mathcal{F}_t, P_x(x \in \mathbb{R}) \) denotes the canonical realisation of the one-dimensional Wiener process. \( \Omega = C([0, \infty[ \rightarrow \mathbb{R}) \) is the space of continuous functions on \([0, \infty[, (X_t, t \geq 0)\) the coordinates process on this space, \( (\mathcal{F}_t, t \geq 0) \) its natural filtration and \( (P_x, x \in \mathbb{R}) \) the family of Wiener measures on \((\Omega, \mathcal{F}_\infty)\), with \( P_x(X_0 = x) = 1 \). When \( x = 0 \), we write simply \( P \) for \( P_0 \).

- We denote by \( (L^t_x, t \geq 0, x \in \mathbb{R}) \) the jointly continuous family of the local times of \((X_t, t \geq 0)\). We denote \( (L_t, t \geq 0) \) for \((L^0_t, t \geq 0)\), the (continuous) local time process at level 0 and by \( (\tau_l, l \geq 0) \) its right-continuous inverse:

\[
\tau_l := \inf\{s > 0; L_s > l\} \quad (l \geq 0)
\]

- To \( q \) a positive Radon measure on \( \mathbb{R} \), \( q \neq 0 \), we associate the continuous additive functional:

\[
A^q_t := \int_{\mathbb{R}} L^t x q(dx)
\]

When \( q \) admits a density with respect to Lebesgue measure, we keep the former notation by still writing \( q \) for the density; we have:

\[
A^q_t = \int_0^t q(X_s)ds
\]

from the occupation density formula for Brownian motion.

Throughout the following, we shall assume that \( q \) satisfies one of the three following hypotheses:

H1. (The integrable case) \( \int_{\mathbb{R}} (1 + |x|) q(dx) < \infty \).

H2. (The left unilateral case) \( \int_{-\infty}^0 (1 + |x|) q(dx) < \infty \) and there exists \( \alpha < 1 \) such that \( \lim_{x \to \infty} x^{2\alpha} q^{(a)}(x) > b > 0 \) where \( q^{(a)} \) denotes the absolutely continuous part of \( q \).

H3. (The right unilateral case) \( \int_{0}^\infty (1 + |x|) q(dx) < \infty \) and there exists \( \alpha < 1 \) such that \( \lim_{x \to -\infty} |x|^{2\alpha} q^{(a)}(x) > b > 0 \).

Of course, if the pair \((X_t, A^q_t), t \geq 0\) satisfies H2 (resp. H3), then the pair \((-X_t, A^q_t), t \geq 0\) satisfies H3 (resp. H2).
0.2 Introduction

0.2.1

In [RVY, I], we obtained the following results:

i) Under $H_1$, $H_2$ or $H_3$, for any $\lambda > 0$:

$$\lim_{t \to \infty} \sqrt{t} E_x \left[ e \left( -\frac{\lambda}{2} A^q_t \right) \right] = \varphi_{\lambda q}(x)$$  \hspace{1cm} (0.4)

(Later, we shall give other presentations of $\varphi_{\lambda q}$).

ii) For any $s \geq 0$ and $\Lambda_s \in \mathcal{F}_s$:

$$\lim_{t \to \infty} \frac{E \left[ 1_{\Lambda_s} e \left( -\frac{\lambda}{2} A^q_t \right) \right]}{E \left[ e \left( -\frac{\lambda}{2} A^q_t \right) \right]} := Q^{(\lambda q)}(\Lambda_s)$$  \hspace{1cm} (0.5)

where formula (0.5) induces a probability $Q^{(\lambda q)}$ on $(\Omega, \mathcal{F}_\infty)$ (see [RVY, I] for more details about this penalisation result).

The first part of this work consists in:

- Using the result i) to obtain a limit theorem relative to the law of the additive functional $(A^q_t, t \geq 0)$. This is the content of Theorem 1.1.

- Obtaining a penalisation result, which is more general than (0.5) i.e., by replacing the exponential function $x \to e^{-\frac{\lambda}{2} x^2}$ by a more general function. This is the content of Theorem 1.5.

0.2.2

In Section 2 of this work, we study in detail the situation where $q = 1_{-\infty,0}$ i.e.,

$A^q_t := A^-_t := \int_0^t 1_{X_s < 0} ds$. In particular, we prove a penalisation theorem for long Brownian bridges: this is the content of Theorem 2.1

0.2.3

Section 3 of this work is devoted to the study of Feynman-Kac penalisation for long Brownian bridges, which generalizes what we have done in [RVY, I] for standard Brownian motion. This is the content of Theorem 3.1.

To summarize, this work extends, in the above directions, our preceding work [RVY, I].
1 A local limit theorem for the laws of some Brownian additive functionals and a penalisation result

1.1 A local limit theorem

Theorem 1.1. Let \( q \) satisfy one of the hypotheses H1, H2 or H3, and let \( (A^q_t, t \geq 0) \) be defined by (0.2) (or (0.3)). Then, for every \( x \in \mathbb{R} \), there exists a positive, \( \sigma \)-finite measure \( \nu_x \), carried by \( \mathbb{R}_+ \), such that:

\[
\sqrt{t} P_x(A^q_t \in dz) \xrightarrow{t \to \infty} \nu_x(dz)
\]

The convergence in (1.1) is understood in the following sense: for any function \( f : \mathbb{R}_+ \to \mathbb{R}_+ \) Borel, and sub-exponential i.e. there exist two positive constants \( C_1 \) and \( C_2 \) such that:

\[
0 \leq f(x) \leq C_1 e^{-C_2x}
\]

then

\[
\sqrt{t} E_x[f(A^q_t)] \xrightarrow{t \to \infty} \int_{\mathbb{R}_+} f(z)\nu_x(dz)
\]

The measure \( \nu_x \) is characterized by:

\[
\int_0^\infty e^{-\frac{y}{2}}\nu_x(dy) = \varphi_{\lambda q}(x)
\]

(1.2)

1.2 Proof of Theorem 1.1

We first begin with some precisions, taken from [RVY, I], (see also S. Kotani, [K]) about \( \varphi_{\lambda q} \), which was defined from (0.3) but admits at least another characterization, namely:

\( \varphi_{\lambda q} \) is the unique solution of the Sturm-Liouville equation:

\[
\varphi''(dx) = \lambda \varphi(x)q(dx)
\]

(1.3)

This equation is taken in the sense of Schwartz distributions, and subject to the following boundary conditions:

Under H1. : \( \varphi'(+\infty) = -\varphi'(-\infty) = \sqrt{\frac{2}{\pi}} \)

(1.4)

Under H2. : \( \varphi'(-\infty) = -\sqrt{\frac{2}{\pi}} \) and \( \varphi(+\infty) = 0 \)

(1.5)

Under H3. : \( \varphi'(+\infty) = \sqrt{\frac{2}{\pi}} \) and \( \varphi(-\infty) = 0 \)

(1.6)

Theorem 1.1. is now an immediate consequence of the next lemma.
Lemma 1.2. Under either of the hypotheses H1, H2, or H3, the function \( \lambda \rightarrow \varphi_{\lambda q}(x) (\lambda > 0) \) is, for any real \( x \), completely monotone, i.e., it satisfies:

\[
(-1)^n \frac{\partial^n}{\partial \lambda^n} \varphi_{\lambda q}(x) \geq 0
\]  

(1.7)

Consequently, there exists a positive, \( \sigma \)-finite measure \( \nu_x \), carried by \( \mathbb{R}_+ \), such that:

\[
\varphi_{\lambda q}(x) = \int_0^\infty e^{-\frac{\lambda}{2} z} \nu_x(dz)
\]  

(1.8)

We shall give two proofs for Lemma 1.2.

1.3 A first proof of Lemma 1.2

We define, for every \( f : \mathbb{R}_+ \rightarrow \mathbb{R} \) and every real \( h \neq 0 \):

\[
D_h f(\lambda) := \frac{f(\lambda + h) - f(\lambda)}{h}
\]  

(1.9)

For \( f(\lambda) := \exp -\frac{\lambda}{2} A^q_t \), we get:

\[
(D_h)^n(f)(\lambda) = e^{-\frac{\lambda A^q}{2}} \left( e^{-\frac{A^q h}{2}} - 1 \right)^n
\]

and, hence for all \( h \neq 0 \):

\[
(-1)^n(D_h)^n(f)(\lambda) \geq 0
\]  

(1.10)

Consequently, taking the expectation of the LHS in (1.10), we obtain:

\[
\sqrt{t} (-1)^n E_x \left[ (D_h)^n \left( \exp -\frac{\lambda A^q}{2} \right) \right] \geq 0
\]  

(1.11)

Hence, from (0.4):

\[
\sqrt{t} (-1)^n E_x \left[ (D_h)^n \left( \exp -\frac{\lambda A^q}{2} \right) \right] \overset{t \rightarrow \infty}{\longrightarrow} (-1)^n(D_h)^n(\varphi_{\bullet q}(x))
\]

Thus:

\[
(-1)^n(D_h)^n(\varphi_{\bullet q}(x))(\lambda) \geq 0
\]  

(1.12)

Letting \( h \rightarrow 0 \) in (1.12), and using the fact that \( D_h f \overset{h \rightarrow 0}{\longrightarrow} f' \), we get:

\[
(-1)^n \frac{\partial^n}{\partial \lambda^n} \left( \varphi_{\lambda q}(x) \right) \geq 0
\]  

(1.13)
1.4 A second proof of Lemma 1.2

We shall only give this second proof under the hypothesis H1 and for \(x = 0\). In [R VY, I], Proposition 4.13, formula (4.43), we have obtained the following explicit formula for \(\varphi_{\lambda q}(0)\):

\[
\varphi_{\lambda q}(0) = \frac{1}{\sqrt{2\pi}} \int_0^\infty \left[ Q^{(0)}_l(e^{-\lambda < Y, q^- >}) \cdot Q^{(2)}_l(e^{-\lambda < Y, q^+ >}) + Q^{(2)}_l(e^{-\lambda < Y, q^- >})Q^{(0)}_l(e^{-\lambda < Y, q^+ >}) \right] dl \tag{1.14}
\]

where, in this formula (1.14), the process \((Y_x, x \geq 0)\) is, under \(Q^{(0)}_l\) (resp. under \(Q^{(2)}_l\)), a squared Bessel process with dimension 0, (resp. 2), starting from \(l\), and we denote:

\[
< Y, q^+ > = \int_0^\infty Y_x q(dx) ; \quad < Y, q^- > = \int_{-\infty}^0 Y_{-x} q(dx) \tag{1.15}
\]

It is then clear from (1.14) that \(\lambda \rightarrow \varphi_{\lambda q}(0)\) is the Laplace transform of a positive measure, as an integral, with respect to the parameter \(l\) of the product of two Laplace transforms of positive measures (indexed by \(l\)).

We shall now give some examples for which the measure \(\nu_x\) may be computed explicitly. We recall that \(\nu_x\) is characterized by:

\[
\int_0^\infty e^{-\frac{z^2}{2}} \nu_x(dz) = \varphi_{\lambda q}(x) \quad \text{where} \quad \varphi_{\lambda q}(x) \quad \text{is given by} \quad (1.3) \ldots (1.6).
\]

1.5 Computation of \(\nu_x\) for \(q(dy) = \delta_0(dy)\)

In this case, the hypothesis H1 is verified and \(A^q_t = L_t\) is the local time at level 0

\[
\varphi_{\lambda q}(x) = \sqrt{\frac{2}{\pi}} \left( \frac{2}{\lambda} + |x| \right) \quad (\text{cf} \quad \text{[R VY, I], Ex. 4.8, p. 199-200})
\]

\[
= \int_0^\infty e^{-\frac{z^2}{2}} \left( \sqrt{\frac{2}{\pi}} 1_{z \geq 0} dz + \sqrt{\frac{2}{\pi}} |x| \delta_0(dz) \right) \tag{1.16}
\]

Thus:

\[
\nu_x(dz) = \sqrt{\frac{2}{\pi}} 1_{[0,\infty)}(z) dz + \sqrt{\frac{2}{\pi}} |x| \delta_0(dz) \tag{1.17}
\]

1.6 Computation of \(\nu_x\) for \(q(dy) = \delta_a(dy) + \delta_b(dy)\) with \((a < b)\)

In this case, the hypothesis H1 is satisfied and \(A^q_t = L^a_t + L^b_t\) where \((L^a_t, \ t \geq 0)\) resp. \((L^b_t, \ t \geq 0)\) denotes the local time at level \(a\), resp. at level \(b\). We know (see [R VY, I], Ex. ...
that
\[
\varphi_{\lambda q}(x) = \begin{cases} 
\frac{\sqrt{2}}{\pi} \left( \frac{1}{\lambda} + x - b \right) & \text{if } x > b \\
\frac{\sqrt{2}}{\pi} \frac{1}{\lambda} & \text{if } x \in [a, b] \\
\frac{\sqrt{2}}{\pi} \left( \frac{1}{\lambda} + a - x \right) & \text{if } x < a
\end{cases}
\]  
(1.18)

\[
= \sqrt{\frac{2}{\pi}} \int_0^\infty e^{-\frac{1}{2}z} \left\{ \frac{1}{2} dz + (x - b)\delta_{0}(dz) + (a - x)\delta_{0}(dz) \right\}
\]

Hence :
\[
\nu_x(dz) = \sqrt{\frac{2}{\pi}} \left\{ \frac{1}{2} 1_{[0,\infty[}(z)dz + (x - b)^{+}\delta_{0}(dz) + (a - x)^{+}\delta_{0}(dz) \right\}
\]  
(1.19)

1.7 Computation of \( \nu_x, \) for \( q(y) = e^{2y} \)

In this case, the hypothesis H2 is satisfied and \( A_q^t = \int_0^t e^{2X_s} ds. \)

To begin with, we show :
\[
\varphi_{\lambda q}(x) = \sqrt{\frac{2}{\pi}} K_0(\sqrt{\lambda} e^x)
\]  
(1.20)

where \( K_0 \) denotes the Bessel-Mc Donald function with index 0 (see [Leb], p. 108).

Let \( \psi(x) := \sqrt{\frac{2}{\pi}} K_0(\sqrt{\lambda} e^x). \) To check (1.20), it suffices to see that :
\[
\psi''(x) = \lambda e^{2x} \psi(x), \quad \psi(x) \xrightarrow{x \to \infty} 0, \quad \psi'(x) \xrightarrow{x \to -\infty} -\sqrt{\frac{2}{\pi}}
\]  
(1.21)

Now (1.21) follows from (see [Leb], p. 110) :
\[
K_0' = -K_1, \quad -K_1'(z) = \frac{1}{z} K_1(z) + K_0(z)
\]

and
\[
\psi(x) \xrightarrow{x \to -\infty} \sqrt{\frac{2}{\pi}} \left( \frac{\pi}{2\sqrt{\lambda} e^{x}} \right) e^{-\sqrt{\lambda} e^x} \xrightarrow{x \to -\infty} 0 \quad ([\text{Leb}], \text{p.123})
\]
\[
\psi'(x) = -\sqrt{\frac{2}{\pi}} \sqrt{\lambda} e^{x} K_1(\sqrt{\lambda} e^x) \sim -\sqrt{\frac{2}{\pi}} \sqrt{\lambda} e^{x} \frac{1}{2} \sqrt{\lambda} e^{x} \xrightarrow{x \to -\infty} -\sqrt{\frac{2}{\pi}} \quad ([\text{Leb}], \text{p.111})
\]

This proves (1.20). But, we also have :
\[
K_0(\sqrt{\lambda} e^x) = \frac{1}{2} \int_0^\infty e^{-\frac{\lambda z^2}{4t}} \frac{dt}{t} \quad (\text{cf [Leb], p. 119})
\]
\[
= \frac{1}{2} \int_0^\infty e^{-\frac{2z^2}{4u}} \frac{du}{u}
\]
Hence:
\[
\nu_x(dz) = \frac{1}{\sqrt{2\pi}} e^{-\frac{x^2}{2}} \delta_0(z) dz
\]
(1.22)

1.8 Computation of \( \nu_x \) for \( q_0(dx) = 1_{[-\infty,0]}(x)dx \)

Here, it is the hypothesis H3 which is satisfied, and
\[
A_t^{q_0} = \int_0^t 1_{[-\infty,0]}(X_s) ds
\]

By scaling, one has, under \( P_0 : A_t^{q_0} \overset{law}{=} t A_1^{q_0} \), and it is well known that under \( P_0, A_t^{q_0} \) follows the arc sine law, i.e., the beta \( \left( \frac{1}{2}, \frac{1}{2} \right) \) law. We shall recall the law of \( A_1^{q_0} \) under \( P_0 \), (see Subsection 2.1 below), which will allow to obtain the following result:

\[
\nu_x(dz) = x + \sqrt{\frac{2}{\pi}} \delta_0(dz) + \frac{1}{\sqrt{2\pi}} e^{-\frac{x^2}{2}} \delta_0(z) dz
\]
(1.23)

For the moment, we shall prove (1.23) without using the explicit law of \( A_t^{q_0} \). For this purpose, we already observe that:

\[
\varphi_{\lambda q_0}(x) = \sqrt{\frac{2}{\pi}} \left\{ e^{x\sqrt{\lambda}} \frac{1}{\sqrt{\lambda}} 1_{x \leq 0} + (x + \frac{1}{\sqrt{\lambda}}) 1_{x > 0} \right\}
\]
(1.24)

Indeed we have:

\[
\varphi_{\lambda q_0}'(x) = \lambda 1_{-\infty,0}(x) \varphi_{\lambda q}(x), \quad \varphi_{\lambda q_0}(+\infty) = \sqrt{\frac{2}{\pi}}, \quad \varphi_{\lambda q_0}(-\infty) = 0.
\]

Then, it remains to see that:

\[
\int_0^\infty e^{-\frac{x^2}{2}} \nu_x(dz) = \varphi_{\lambda q_0}(x)
\]
(1.25)

where \( \nu_x \) is defined via (1.23) and \( \varphi_{\lambda q_0}(x) \) by (1.24). Now, for \( x > 0 \), one has:

\[
\int_0^\infty e^{-\frac{x^2}{2}} \nu_x(dz) = x + \sqrt{\frac{2}{\pi}} + \frac{1}{\pi} \int_0^\infty e^{-\frac{x^2}{2}} \frac{dz}{\sqrt{2}} = x + \sqrt{\frac{2}{\pi}} + \sqrt{\frac{2}{\pi}} \frac{1}{\sqrt{\lambda}} \Gamma(1/2) = x + \sqrt{\frac{\pi}{2}} + \sqrt{\frac{2}{\pi}} \frac{1}{\sqrt{\lambda}} = \varphi_{\lambda q_0}(x)
\]
wheras for \( x < 0 \):

\[
\int_0^\infty e^{-\frac{x^2}{2}} \nu_x(dz) = \frac{1}{\pi} \int_0^\infty e^{-\frac{x^2}{2} - \frac{z^2}{2}} \frac{dz}{\sqrt{2}} = \frac{2}{\pi} K_{1/2}(\pi \sqrt{\lambda}) \left( \frac{x^2}{\lambda} \right)^{1/4} \quad \text{(see [Leb], p. 119)}
\]

However, one has: \( K_{1/2}(\pi \sqrt{\lambda}) = \left( \frac{\pi}{2|x|\sqrt{\lambda}} \right)^{1/2} e^{-|x|\sqrt{\lambda}} \). Hence:

\[
\int_0^\infty e^{-\frac{x^2}{2}} \nu_x(dz) = \frac{2}{\pi} \left( \frac{x^2}{\lambda} \right)^{1/4} \left( \frac{\pi}{2|x|\sqrt{\lambda}} \right)^{1/2} e^{-|x|\sqrt{\lambda}} = \sqrt{\frac{2}{\pi \lambda}} e^{-|x|\sqrt{\lambda}} = \varphi_{\lambda q}(x)
\]
1.9 Computation of $\nu_x$ when $q(y) = 1_{[a,b]}(y)$ \quad (a < b)

The hypothesis H1 is satisfied and $A_t^q = \int_0^t 1_{[a,b]}(X_s) \, ds$. We shall prove that:

$$\nu_x^{(a,b)}(dz) = \begin{cases} 
\sqrt{\frac{2}{\pi}} (x-b) \delta_0(dz) + \frac{1}{\pi \sqrt{z}} 1_{[0,\infty]}(z) dz \left( 1 + 2 \sum_{n=1}^{\infty} e^{-\frac{n^2(a-b)^2}{2z}} \right) & \text{if } x > b \\
\sqrt{\frac{2}{\pi}} (a-x) \delta_0(dz) + \frac{1}{\pi \sqrt{z}} 1_{[0,\infty]}(z) dz \left( 1 + 2 \sum_{n=1}^{\infty} e^{-\frac{n^2(a-b)^2}{2z}} \right) & \text{if } x < a
\end{cases} \quad (1.26)$$

Here, the explicit form of $\varphi_{\lambda q}^{(a,b)}(x)$ is (see [RVY, I], Ex. 4.7, p. 199):

$$\varphi_{\lambda q}^{(a,b)}(x) = \begin{cases} 
\sqrt{\frac{2}{\pi}} \left( \frac{1}{\sqrt{\lambda} \tanh (\sqrt{\lambda} \frac{b-a}{2})} + x-b \right) & \text{if } x > b \\
\sqrt{\frac{2}{\pi}} \left( \frac{1}{\sqrt{\lambda} \tanh (\sqrt{\lambda} \frac{b-a}{2})} + a-x \right) & \text{if } x < a
\end{cases} \quad (1.27)$$

It now remains to prove that:

$$\int_0^\infty e^{-\frac{x^2}{2}z} \nu_x^{(a,b)}(dz) = \varphi_{\lambda q}^{(a,b)}(x) \quad (1.28)$$

where $\nu_x^{(a,b)}$ is defined via (1.26) and $\varphi_{\lambda q}^{(a,b)}$ via (1.27). But, (1.28) follows, after some elementary computations from the identities, for every real $u$ and $v > 0$:

$$\frac{\cosh(\sqrt{\lambda}u)}{\sqrt{\lambda} \sinh(\sqrt{\lambda}v)} = \sum_{n=0}^{\infty} \int_0^\infty dh \left( e^{-\sqrt{\lambda}(h+(2n+1)v-u)} + e^{-\sqrt{\lambda}(h+(2n+1)v+u)} \right) \quad (1.29)$$

$$= \sum_{n=0}^{\infty} \int_0^\infty dh \int_0^\infty ds \left( H_{h+(2n+1)v-u}(s) + H_{h+(2n+1)v+u}(s) \right) e^{-\lambda s} \quad (1.30)$$

with

$$H_{a}(u) := \frac{a}{2\sqrt{\pi} u^3} e^{-a^2/4u} = \frac{1}{\pi u} \frac{\partial}{\partial u} (e^{-u^2/\pi}) \quad (a > 0)$$

Passing from (1.29) to (1.30) is obtained by using the elementary formula:

$$e^{-\sqrt{\lambda}u} = \int_0^\infty e^{-\lambda u} H_a(u) du = \int_0^\infty e^{-\lambda u} \frac{a}{2\sqrt{\pi} u^3} e^{-u^2/\pi} du \quad (1.31)$$
We now show (1.29).

**Remark 1.9**

(i) If in formula (1.26), we take : $b = 0$, and we let $a$ tend to $-\infty$, we obtain:

$$
\lim_{a \to -\infty} \nu^{a,0}(dz) = \begin{cases} 
\sqrt{\frac{2}{\pi}} \frac{1}{\sqrt{x}} & \text{if } x > 0 \\
\frac{1}{\pi \sqrt{z}} & \text{if } x \leq 0
\end{cases}
$$

We note that the RHS of (1.32) is nothing else but the measure $\nu_x$ associated with $q_0(y) = 1_{[-\infty,0]}$ (see (1.23)). This may be interpreted as "a continuity property" of $\varphi^{a,b}$, as $a \to -\infty$.

(ii) In the same spirit, but taking up now the computation from subsection 1.9, where we choose for $q$ the function : $q^{(c)}(y) = \frac{1}{2c} \mathbb{1}_{[-c,c]}(y)$, we have:

$$
\int_0^\infty e^{-\frac{\lambda v}{2}} \nu_x^{(c)}(dz) \xrightarrow{c \to 0} \int_0^\infty e^{-\frac{\lambda v}{2}} \nu_x(dz) = \sqrt{\frac{2}{\pi}} \left( \frac{2}{\lambda} + |x| \right)
$$

where $\nu_x$ is the measure associated to $q(dz) = \delta_0(dz)$ (see (1.16)). In other terms, since:

$$
\frac{1}{2c} \int_0^t \mathbb{1}_{[-c,c]}(X_s)ds \xrightarrow{c \to 0} L_t \text{ a.s.,}
$$

we witness there also a "continuity property of $\nu_x^{(c)}$" as $c \to 0$. Let us show (1.33) for $x = 0$; from (1.27):

$$
\int_0^\infty e^{-\frac{\lambda v}{2}} \nu_0^{(c)}(dz) = \sqrt{\frac{2}{\pi}} \left( \frac{\cosh \left( \frac{\lambda}{2c} c \right)}{\sqrt{\frac{\lambda}{2c} \sinh \left( \frac{\lambda}{2c} c \right)} \cosh \left( \frac{\lambda}{2c} c \right) \sinh \left( \frac{\lambda}{2c} c \right)} \right) \xrightarrow{c \to 0} \sqrt{\frac{2}{\pi}} \cdot \frac{2}{\lambda}
$$
and for \( x \neq 0 \), and \( c \) small enough, we obtain from (1.27) that :

\[
\int_0^\infty e^{-\frac{\lambda z^2}{2c}} \nu_x(c)(dz) = \sqrt{\frac{2}{\pi}} \left\{ \frac{1}{\sqrt{\frac{\lambda}{2c}} \tanh \left( \sqrt{\frac{\lambda}{2c}} \cdot c \right)} + |x - c| \right\} \xrightarrow{c \to 0} \sqrt{\frac{2}{\lambda}} \sqrt{\frac{2}{\pi}} \left( \frac{2}{\lambda} + |x| \right)
\]

### 1.10 Computation of \( \nu_x \) when \( q(y) = 1_{[0,\infty]}(y)^\alpha, \alpha > 0 \)

The hypothesis \( H2 \) is satisfied, and we have :

\[
A'_q = \int_0^t 1_{(X_s > 0)} X_s^\alpha ds.
\]

We now show the existence of a constant \( C_\alpha > 0 \) such that :

\[
\nu_0(dz) = \frac{C_\alpha}{z^{2+\alpha}} 1_{[0,\infty]}(z)dz \quad (1.34)
\]

Indeed, thanks to the scaling property, we have :

\[
E_0 \left( e^{-\frac{1}{2} \int_0^t 1_{X_s > 0} X_s^\alpha ds} \right) = E_0 \left( e^{\left( -\frac{1}{2} t^{1+\alpha/2} A'_q \right)} \right)
= E_0 \left( \exp \left( -\frac{1}{2} A'_q \lambda^{\frac{\alpha}{2+\alpha}} t \right) \right) \quad (1.35)
\]

Thus, multiplying (1.35) by \( \sqrt{t} \) and letting \( t \) tend to \(+\infty\), we obtain :

\[
\varphi_{\lambda q}(0) = \frac{1}{\lambda^{\frac{1}{2+\alpha}}} \varphi_{1,q}(0) = \frac{1}{\lambda^{\frac{1}{2+\alpha}}} c'_\alpha = c_\alpha \int_0^\infty e^{-\frac{z^2}{2}} \frac{dz}{z^{\frac{1}{2+\alpha}}}
\]

The same computations, performed this time with \( x \neq 0 \), lead to :

\[
\varphi_{\lambda q}(x) = \frac{1}{\lambda^{\frac{1}{2+\alpha}}} \varphi_{1,q}(x \lambda^{\frac{1}{2+\alpha}}), \text{ i.e. } \nu_x = \frac{1}{\lambda^{\frac{1}{2+\alpha}}} \nu_{\lambda}(x \lambda^{\frac{1}{2+\alpha}}), \nu_{\lambda}
\]

where \( \nu_{\lambda} \) is the image of \( \nu \) by the application \( z \rightarrow \lambda z \).

### Question 1.4

We know (see [R,Y], chap. X) that, if \( q \) is an integrable function, then :

\[
\frac{1}{\sqrt{t}} \int_0^t q(x + X_s)ds \xrightarrow{\text{law}} \left( \int q(x + y)dy \right) |N| = \left( \int q(y)dy \right) |N| \quad (1.36)
\]

where \( N \) is a standard Gaussian variable, and on the LHS of (1.36), \( (X_s, s \geq 0) \) is a Brownian motion starting from 0. Let \( g \) denote the density of the r.v. \( \bar{q}|N| \) with \( \bar{q} = \int q(y)dy \) that is :

\[
g(z) = \frac{1}{\bar{q}} \sqrt{\frac{2}{\pi}} e^{-\frac{z^2}{2\bar{q}}} 1_{[0,\infty]}(z)
\]
Let us now consider the supplementary hypothesis \( \tilde{H} \), which seems reasonable enough in view of (1.36) that the density \( g_t(x, \cdot) \) of the r.v. \( \frac{1}{\sqrt{t}} \int_0^t q(x + X_s) \, ds \) converges, as \( t \to \infty \), uniformly on every compact, towards \( g \).

However, this would imply that, for every function \( h \), which is continuous with compact support, one would have:

\[
\sqrt{t} \, E_x[h(A^q_t)] = \sqrt{t} \, E_x[h \left( \frac{A^q_t}{\sqrt{t}} \right)] = \sqrt{t} \int_0^\infty h(z) g_t(x, z) \, dz = \int_0^\infty h(y) g_t(x, \frac{y}{\sqrt{t}}) \, dy \to \infty \int_0^\infty h(y) g(0) \, dy
\]

But, from Theorem 1.1., we know that:

\[
\sqrt{t} \, E_x[h(A^q_t)] \to \int_{\mathbb{R}_+} h(x) \nu_x(dz)
\]

Thus, this would imply that the measure \( \nu_x(dz) \) would be equal to:

\[
\frac{1}{\sqrt{2\pi}} \int_{0,\infty}^{1/2} dz \quad (1.37)
\]

so that, the measure \( \nu_x \) would not depend on \( x \), and would be proportional to Lebesgue measure on \( \mathbb{R}_+ \). But clearly, this is not the case for either of the examples in subsections 1.4 to 1.9. Consequently, the hypothesis \( \tilde{H} \) is not satisfied for the corresponding \( q \)'s. It would be of interest to know for which \( q \)'s, if any, it is satisfied.

### 1.11 Penalties by \( h(A^q_t) \)

Let \( q \) satisfy one of the previous hypotheses H1, H2 or H3, and denote, as before:

\[
A^q_t = \int_\mathbb{R} L^q_t q(dx) = \int_0^t q(X_s) \, ds \quad \text{if } q \text{ admits a density}. \]

Let now \( h : \mathbb{R}_+ \to \mathbb{R}_+ \) such that:

\[
\sqrt{t} \, E_x[h(A^q_t)] \to \int_0^\infty h(z) \nu_x(dz) \quad (1.38)
\]

Then, (1.38) is satisfied, from Theorem 1.1., as soon as \( h \) is sub-exponential (for example if \( h \) is continuous, with compact support). We shall now study the penalisation of Wiener measure by the functional \( h(A^q_t) \), i.e.: we shall study the limit, as \( t \to \infty \), of:

\[
\frac{E_x(1_{A^q_t} h(A^q_t))}{E_x(h(A^q_t))} \quad (s \geq 0, \Lambda_s \in \mathcal{F}_s) \quad (1.39)
\]

We have already made this study in two situations:

1) \( q(dy) = \delta_0(dy) \) then \( A^q_t = L_t \) (cf [RVY, II])

2) \( A^q_t = \int_\mathbb{R} L^q_t q(dy) \) and \( h(u) = \exp\left(-\frac{\lambda}{2} u \right) \) (cf [RVY, I]).
This time, Theorem 1.1. allows us to obtain:

**Theorem 1.5.** : Let $q$, $A^q$ and $h$ as above. Then:

1) For every $s \geq 0$, and every $\Lambda_s \in \mathcal{F}_s$:

\[
\lim_{t \to \infty} \frac{E_x(1_{\Lambda_s} h(A^q_t))}{E_x(h(A^q_t))} \quad \text{exists} \tag{1.40}
\]

2) This limit equals $E_x(1_{\Lambda_s} M^{h,q}_s) := Q^{h,q}(\Lambda_s)$, where

\[
M^{h,q}_s := \frac{\int_{\mathbb{R}^+} \nu_{X_s}(dz) h(z + A^q_s)}{\int_{\mathbb{R}^+} \nu_z(dz) h(z)} \tag{1.41}
\]

Furthermore, $(M^{h,q}_s, s \geq 0)$ is a positive martingale. In the case when $h(u) := e^{-\frac{\lambda}{2} u}$ $(u, \lambda \geq 0)$, we then obtain:

\[
M^{h,q}_s = \frac{\varphi_{\lambda q}(X_s)}{\varphi_{\lambda q}(x)} \exp\left(-\frac{\lambda}{2} A^q_s\right) \tag{1.42}
\]

**Proof of Theorem 1.5.**

We have:

\[
\frac{E_x(1_{\Lambda_s} h(A^q_t))}{E_x(h(A^q_t))} = \frac{E_x(1_{\Lambda_s} E_b(h(a + A^q_{t-s})))}{E_x(h(A^q_t))}
\]

from the Markov property, where $b = X_s$ and $a = A^q_s$. Thus, from Theorem 1.1. :

\[
E_x(h(A^q_t)) \sim \frac{1}{\sqrt{t}} \int_{0}^{\infty} \nu_x(dz) h(z)
\]

and

\[
E_b(h(a + A^q_{t-s})) \sim \frac{1}{\sqrt{t-s}} \int_{0}^{\infty} \nu_y(dz) h(a+z)
\]

Hence :

\[
\frac{E_x(1_{\Lambda_s} h(A^q_t))}{E_x(h(A^q_t))} \xrightarrow{t \to \infty} \frac{\sqrt{t}}{\sqrt{t-s}} \frac{E_x(1_{\Lambda_s} \int_{\mathbb{R}^+} \nu_{X_s}(dz) h(z + A^q_s))}{\int_{\mathbb{R}^+} h(z) \nu_z(dz)} \xrightarrow{t \to \infty} E_x(1_{\Lambda_s} M^{h,q}_s)
\]

In the preceding lines, we have been a little careless concerning the exchange of limit and expectation. Likewise, although it is easy to see that $(M^{h,q}_s, s \geq 0)$ is a local martingale, some care is needed in order to show that it is a true martingale. However, all this is correct as soon as $h$ is sub-exponential. We leave details to the reader.
2 A detailed study for \( q_0 = 1_{[-\infty,0]} \), \( A_t^- := \int_0^t 1(X_s<0)ds \)

Throughout this section, we choose \( q_0 = 1_{[-\infty,0]} \). Thus, the hypothesis H3 is now satisfied. We shall study this situation in detail, which we are able to do as we know (see \( [Y] \)) the law of \( A_t^{q_0} = \int_0^t q_0(X_s)ds \) under \( P_x \), for every real \( x \) (see (2.5) and (2.7) below). We shall, successively:
- compute explicitly the measure \( \nu_x \) starting from the knowledge of the law of \( A_t^{q_0} \) and we shall recover the result of subsection 1.8 above;
- study the penalisation, not only of the process \((X_t, t \geq 0)\) by \( h(A_t^{q_0}) \), but also the penalisation of the "long bridges" by this functional;
- describe precisely the behavior of the canonical process under the probability \( Q^{h,q_0} \), where \( Q^{h,q_0} \) is defined via:

\[
Q^{h,q_0}(\Lambda_s) = E(1_{\Lambda_s}M_s^{h,q_0}) \quad (s \geq 0, \Lambda_s \in \mathcal{F}_s) \quad (2.1)
\]

2.1 The law of \( A_t^- \) and the computation of \( \nu_x \)

To simplify notation, we denote :

\[
A_t^- = \int_0^t 1_{(X_s<0)}ds = \int_0^t q_0(X_s)ds \quad (2.2)
\]

We recall the following result, which is found in \( [Y] \). For any \( f : [0,1] \rightarrow \mathbb{R}_+ \), Borel, sub-exponential (see Theorem 1.1) and any \( y > 0 \):

\[
E_0\left[f\left(\int_0^1 1_{(X_s<y)}ds\right)\right] = \int_0^1 \frac{du}{\pi\sqrt{u(1-u)}} e^{-\frac{y^2}{2u}} f(u) + f(1) \sqrt{\frac{2}{\pi}} \int_0^y e^{-\frac{\alpha^2}{2}} d\alpha \quad (2.3)
\]

whereas, for any \( y < 0 \), we use :

\[
\int_0^1 1_{(X_s<y)}ds \overset{\text{law}}{=} \int_0^1 1_{(X_s>-y)}ds \; \overset{\text{law}}{=} \; 1 - \int_0^1 1_{(X_s<-y)}ds \quad (2.4)
\]

and by the scaling property :

\[
E_x\left[f\left(\int_0^1 1_{(X_s<0)}ds\right)\right] = E_0\left(f\left(\int_0^1 1_{(X_s<-x)}ds\right)\right) = E_0\left(f\left(t \int_0^1 1_{(X_s<-\frac{x}{\sqrt{t}})}ds\right)\right)
\]

hence, from (2.3) and (2.4), if \( x \leq 0 \):

\[
E_x[f(A_t^-)] = \int_0^t \frac{du}{\pi\sqrt{v(t-v)}} e^{-\frac{2}{\pi v}} f(v) + f(t) \sqrt{\frac{2}{\pi}} \int_0^{\left|\frac{v}{\sqrt{t}}\right|} e^{-\frac{\alpha^2}{2}} d\alpha \quad (2.5)
\]

\[
\lim_{t \to \infty} \frac{1}{\pi\sqrt{t}} \int_0^\infty \frac{du}{\sqrt{v}} e^{-\frac{2}{\pi v}} f(v) \quad (2.6)
\]
whereas, if $x > 0$:

$$E_x[f(A_t^-)] = f(0) \sqrt{\frac{2}{\pi}} \int_0^x e^{-\frac{y^2}{2}} \, dy + \int_0^t \frac{dv}{\pi \sqrt{v(t-v)}} e^{-\frac{v}{2v(t-v)}} f(v)$$  \hspace{1cm} (2.7)

$$\sim_{t \to \infty} f(0) \sqrt{\frac{2}{\pi}} \int_0^x e^{-\frac{y^2}{2}} \, dy + \frac{1}{\pi \sqrt{t}} \int_0^\infty \frac{dv}{\sqrt{v}} f(v)$$  \hspace{1cm} (2.8)

Thus, we obtain, for $f : \mathbb{R}_+ \to \mathbb{R}_+$, Borel, sub-exponential (and then $\int_0^\infty \frac{dv}{\sqrt{v}} f(v) < \infty$)

$$\sqrt{t} E_x[f(A_t^-)] \sim_{t \to \infty} \int_0^\infty f(z) \nu_x(\,dz\,$$

with

$$\nu_x(\,dz\, = x_+ \sqrt{\frac{2}{\pi}} \delta_0(\,dz\,) + \frac{1}{\pi} e^{-\frac{x^2}{2}} 1_{[0,\infty)}(z) \frac{dz}{\sqrt{x}}$$

which is precisely (1.23).

2.2 Penalisation by $h(A_t^-)$. A study of ”long bridges” and of the $Q_h$–process

We recall that, from (2.5), the density of $A_t^-$ under $P_0$, which we denote by $p_{A_t^-}$, equals:

$$p_{A_t^-}(y) = \frac{1}{\pi} \frac{1}{\sqrt{y(t-y)}} 1_{[0,t]}(y) \quad (\text{: the arc sine law}).$$  \hspace{1cm} (2.9)

Throughout the following, $h$ denotes a function from $\mathbb{R}_+$ to $\mathbb{R}_+$ such that:

$$\int_0^\infty \frac{dy}{\sqrt{y}} h(y) < \infty$$

and we assume, without loss of generality, that:

$$\int_0^\infty \frac{dy}{\sqrt{y}} h(y) = 1$$  \hspace{1cm} (2.10)

**Theorem 2.1**

1) For every $s \geq 0$ and every $\Lambda_s \in \mathcal{F}_s$:

$$\lim_{t \to \infty} E_0(1_{\Lambda_s} | A_t^- = a) = Q^{(a)}(\Lambda_s)$$  \hspace{1cm} (2.11)

with

$$Q^{(a)}(\Lambda_s) := \sqrt{\frac{2}{\pi}} \frac{1_{a < s}}{\sqrt{s-a}} E_0(1_{\Lambda_s} X_s^+ | A_s^- = a)$$

$$+ E[1_{\Lambda_s} \sqrt{\frac{a}{a-A_s^-}} 1_{(X_s^- < a)} e^{-\frac{(X_s^-)^2}{2(a-A_s^-)}}]$$  \hspace{1cm} (2.12)

(recall that $X_s^+ = 0 \lor X_s$, $X_s^- = -(X_s \land 0)$ and $A_s^- = \int_0^s 1_{X_u < 0} \, du$)
2) For every function \( h \) which satisfies (2.10), for every \( s \geq 0 \) and any \( \Lambda_s \in \mathcal{F}_s \):

\[
\lim_{t \to \infty} \frac{E_0(1_{\Lambda_s} h(A_t^-))}{E_0(h(A_t^-))} = E(1_{\Lambda_s} M_s^h) \tag{2.13}
\]

where \( (M_s^h, s \geq 0) \) is the positive martingale given by:

\[
M_s^h := \sqrt{2\pi} X_s^+ h(A_s^-) + \int_0^\infty \frac{dy}{\sqrt{y}} e^{-\frac{(X_s-y)^2}{2y}} h(A_s^- + y) \tag{2.14}
\]

(note that \( M_0^h = 1 \)).

3) Formula (2.13) induces a probability \( Q^h \) on \((\Omega, \mathcal{F}_\infty)\), which admits the following disintegration:

\[
Q^h(\Lambda) = \int_0^\infty Q^{(a)}(\Lambda) \frac{h(a)}{\sqrt{a}} da \quad (\Lambda \in \mathcal{F}_\infty) \tag{2.15}
\]

where \( Q^{(a)} \) is given by (2.12).

4) Under \( Q^h \), the canonical process \((X_t, t \geq 0)\) satisfies:

i) \( A_\infty \) is finite a.s., and admits as density \( \frac{h(y)}{\sqrt{y}} 1_{y>0} \) \tag{2.16}

ii) let \( g = \inf\{t ; A_t^- = A_\infty \} = \sup\{t ; X_t \leq 0 \} \) \tag{2.17}

Then \( Q^h(g < \infty) = 1 \)

iii) the processes \((X_t, t \leq g)\) and \((X_{g+t}, t \geq 0)\) are independent ;

iv) the process \((X_{g+t}, t \geq 0)\) is a 3-dimensional Bessel process starting from 0.

Moreover, while proving Theorem 2.1, we shall give a precise description of the process \((X_t ; t \leq g)\).

2.3 Proof of Theorem 2.1

2.3.1 Proof of point 1) in Theorem 2.1

For this purpose, we choose a function \( h \), which is Borel, positive, and satisfies (2.10). We first write:

\[
E_0(1_{\Lambda_s} h(A_t^-)) = \int_0^t E_0(1_{\Lambda_s} |A_t^- = a)p_{A_t^-}(a)h(a)da \tag{2.18}
\]

Then, conditioning with respect to \( \mathcal{F}_s \), we obtain:

\[
E_0(1_{\Lambda_s} h(A_t^-)) = E_0 \left( 1_{\Lambda_s} E_0 \left( h \left( a + \int_0^{t-s} 1_{(X_u < -x)} du \right) \right) \right) \tag{2.19}
\]
with \( a = A_s^- \) and \( x = X_s \). Using now (2.5) and (2.6), we obtain:

\[
E_0(1_{\Lambda_s}h(A_s^-)) = E_0(1_{\Lambda_s}1_{x<0} \left( \int_0^{t-s} \frac{dv}{\pi \sqrt{v(t-s-v)}} e^{-\frac{x^2}{2v}} h(a+v) \right) + h(a+t-s) \psi \left( \frac{|x|}{\sqrt{t-s}} \right)) + E_0(1_{\Lambda_s}1_{x>0} h(a)) \]

\[
\int_0^{t-s} \frac{dv}{\pi \sqrt{v(t-s-v)}} \int_0^{v+s} da \ p_{A_s^-}(a-v) \ E_0(1_{\Lambda_s}1_{x<0} e^{-\frac{x^2}{2v}} |A_s^- = a-v) h(a)
\]

\[
+ \int_{t-s}^{t} \ p_{A_s^-}(a-s-t) E_0(1_{\Lambda_s}1_{x<0} \psi \left( \frac{|X_s|}{\sqrt{t-s}} \right) |A_s^- = a+s-t) h(a) da
\]

Similarly:

\[
(2)_t = \int_0^{s} p_{A_s^-}(a) E_0(1_{\Lambda_s}1_{x>0} \psi \left( \frac{|X_s|}{\sqrt{t-s}} \right) |A_s^- = a) h(a) da
\]

\[
+ \int_0^{s} \int_0^{t-s} \frac{dv}{\pi \sqrt{v(t-s-v)}} \ E_0(1_{\Lambda_s}1_{x>0} e^{-\frac{x^2}{2v}} |A_s^- = a) h(a+v)
\]

Then, comparing (2.18), (2.22), (2.24) and identifying the "coefficient of \( h(a) \)”, it follows that:

\[
E_0(1_{\Lambda_s} |A_s^- = a) = (1)_{t} + (2)_{t}
\]
with:

\[(\dagger)\]
\[
\int_0^t \frac{1}{p_{A_s^-(a)}} \frac{dv}{\pi \sqrt{v(t-s-v)}} 1_{\nu t} \leq a < v + p_{A_s^-(a)}(a-v) E_0(1_{A_s^-(a-v)} e^{-\frac{x^2}{2\pi} |A_s^- = a-v})
\]
\[
+ \int_0^t \frac{1}{p_{A_s^-(a)}} 1_{t-s < a < v} p_{A_s^-}(a-v) E_0(1_{A_s^- e^{-\frac{x^2}{2(a-v)}} |A_s^- = a-v})
\]
\[
\rightarrow \int_0^t \frac{1}{\pi \sqrt{a-w}} p_{A_s^-}(w) E_0(1_{A_s^- e^{-\frac{x^2}{2(a-w)}} |A_s^- = w}) dw
\]
\[
= E_0(1_{A_s^- e^{-\frac{x^2}{2(a-c_s)}} \sqrt{a-A_s^-} 1_{A_s^- < a}) (2.25)
\]

since \(p_{A_s^-}(a) = \frac{1}{\pi \sqrt{a(t-a)}} 1_{[0,t]}(a)\). Similarly, one has:

\[(\ddagger)\]
\[
\int_0^t \frac{a}{\leq a < v + p_{A_s^-(a)}(a-v) E_0(1_{A_s^- e^{-\frac{x^2}{2(a-v)}} |A_s^- = a-v})
\]
\[
+ \int_0^t \frac{1}{p_{A_s^-(a)}} 1_{s < v} E_0(1_{A_s^- e^{-\frac{x^2}{2(v-a)}} |A_s^- = a})
\]
\[
\rightarrow \int_0^t \frac{1}{\pi \sqrt{v-a}} p_{A_s^-}(a) E_0(1_{A_s^- e^{-\frac{x^2}{2(a-v)}} |A_s^- = a-v})
\]
\[
= E_0(1_{A_s^- e^{-\frac{x^2}{2(a-c_s)}} \sqrt{a-A_s^-} 1_{A_s^- < a})
\]

Hence, point 1 of Theorem 2.1 follows.

2.3.2 Proof of point 2 in Theorem 2.1

In fact, this point has already been shown while proving Theorem 1.5. With the help of the formula (1.41) of \(M^h\) and the explicit computation of \(\nu_x\) (see formula (1.23)), we obtain:

\[
M^h = \int_0^\infty \nu_X(dv) h(A_s^- + y) = \int_0^\infty h(A_s^- + y) \left[ X_s^+ \sqrt{\frac{2}{\pi}} \delta_0(dy) + \frac{1}{\pi} e^{-\frac{(X_s^+)^2}{2y}} dy \right]
\]
\[
= \sqrt{2\pi} X_s^+ \delta_0(dy) + \int_0^\infty e^{-\frac{(X_s^+)^2}{2y}} h(A_s^- + y) \quad (2.26)
\]

Now, clearly, this point 1 of Theorem 2.1 which we just proved implies also point 2 of the same Theorem 2.1. Indeed, we have:

\[
E_0(1_{A_s^-}(a)) \int_0^t \frac{E_0(1_{A_s^-} = a) h(a) p_{A_s^-}(a) da}{E_0(h(A_s^-))} \int_0^t h(a) p_{A_s^-}(a) da
\]
From the above point 1, and with the help of the explicit form of \( p_{A_t^-}(a) \) as given by (2.9) the above quantity converges, as \( t \to \infty \), towards :

\[
\int_0^\infty \frac{h(a)}{\sqrt{a}} Q^{(a)}(\Lambda_s) da = \int_0^\infty \frac{h(a)}{\sqrt{a}} Q^{(a)}(\Lambda_s) da
\]

(2.27)

since we assumed : \( \int_0^\infty \frac{h(a)}{\sqrt{a}} da = 1. \)

It now remains to compute :

\[
\int_0^\infty \frac{h(a)}{\sqrt{a}} Q^{(a)}(\Lambda_s) da = \sqrt{2} \int_0^\infty \frac{1}{\sqrt{a(s-a)}} E_0(1_{\Lambda_s} X_s^+ | A_s^- = a) h(a) da \\
+ \int_0^\infty \frac{h(a)}{\sqrt{a}} E_0(1_{\Lambda_s} \sqrt{\frac{a}{a-A_s^-}} 1_{A_s^- < a} e^{-\frac{(X_s^+)^2}{2(a-A_s^-)}}) da \quad \text{(from (2.12))}
\]

\[
= \sqrt{2\pi} \int_0^\infty p_{A_s^-}(a) E_0(1_{\Lambda_s} X_s^+ | A_s^- = a) h(a) da \\
+ \int_0^\infty dy \sqrt{\frac{y}{y_a}} E_0(1_{\Lambda_s} e^{-\frac{(X_s^+)^2}{2y}} h(A_s^- + y))
\]

(after the change of variable \( a - A_s^- = y \))

\[
= \sqrt{2\pi} E_0(1_{\Lambda_s} X_s^+ h(A_s^-)) + \int_0^\infty E_0(1_{\Lambda_s} e^{-\frac{(X_s^+)^2}{2y}} h(A_s^- + y)) dy \\
= E_0(1_{\Lambda_s} M_{\sigma_a}^h) \quad \text{(from (2.14))}
\]

(2.29)

We now remark that point 3 in Theorem 2.1 states precisely formula (2.29) we just established.

### 2.3.3 Proofs of points 4i) and 4ii) in Theorem 2.1

**a)** From formula (2.15) and from Doob’s optional sampling theorem, we deduce :

\[
Q^h(A_{\sigma_a}^- > a) = E[M_{\sigma_a}^h], \quad \text{with} \quad \sigma_a := \inf \{ t : A_t^- > a \}
\]

(2.30)

But :

\[
M_{\sigma_a}^h = \sqrt{2\pi} h(a) X_{\sigma_a}^+ + \int_0^\infty dy \sqrt{\frac{y_a}{y}} e^{-\frac{(X_{\sigma_a}^-)^2}{2y}} h(a + y)
\]

\[
= \int_0^\infty dy \sqrt{\frac{y_a}{y}} e^{-\frac{(X_{\sigma_a}^-)^2}{2y}} h(a + y) \quad \text{since} \quad X_{\sigma_a}^+ = 0
\]

We recall that the process \((X_{\sigma_a}^- : a \geq 0)\) is distributed as the reflecting Brownian motion \((|X_a|, a \geq 0)\), where \((X_a, a > 0)\) is a standard Brownian motion starting from 0 (see, e.g., [K.S], Th. 3.1, p. 419). Hence, we obtain :

\[
E[M_{\sigma_a}^h] = \int_0^\infty dy \sqrt{y} h(a + y) E(e^{-\frac{y^2}{2y}})
\]

\[
= \int_0^\infty dy \sqrt{y} h(a + y) \sqrt{\frac{y}{y + a}} = \int_0^\infty \frac{dy}{\sqrt{y + a}} h(a + y) = \int_a^\infty \frac{dy}{\sqrt{y}} h(y)
\]
b) We now remark that it is easy to recover the law of $A_\infty^-$ under $Q^h$ from points 1 and 2 in Theorem 2.1. We may already prove that, under $Q^{(a)}$, one has $A_\infty^- = a$ a.s. Indeed, this follows from:

\[
Q^{(a)}(A_s^- > b) = \sqrt{\frac{\alpha}{\pi}} \frac{1_{a<s}}{\sqrt{s-a}} E_0 X_s^+ 1_{A_s^- > b} |A_s^- = a \rangle
\]

Hence, passing to the limit as $s \to \infty$, if $b > a$, $Q^{(a)}(A_\infty^- > b) = 0$.

On the other hand, it is clear that $E_0(1_{A_{\infty}^- < a}|A_t^- = a) = 1$ (if $t > s$), hence, passing to the limit as $t \to \infty$, and then, letting $s \to \infty$ we obtain:

\[
Q^{(a)}(A_\infty^- \leq a) = 1
\]

Finally, from (2.15), we get:

\[
Q^h(A_\infty^- \leq b) = \int_0^\infty \frac{h(a)}{\sqrt{a}} Q^{(a)}(A_\infty^- \leq b) da = \int_0^b \frac{h(a)}{\sqrt{a}} da.
\]

### 2.3.4 Computation of Azéma’s supermartingale $Z_t := Q^h(g > t|\mathcal{F}_t)$

Our proof of points 4iii) and 4iv) in Theorem 2.1 is based on the theory of enlargements of filtration (cf [J] or [JY]). In order to apply this theory, we need to calculate Azéma’s supermartingale $Q^h(g > t|\mathcal{F}_t)$. We start with this computation.

Let $g = \inf\{t \geq 0; A_t^- = A_\infty^-=\sup\{t \geq 0; X_t \leq 0\}$

\[
(2.31)
\]

**Lemma 2.2** The following explicit formula holds:

\[
Z_t := Q^h(g > t|\mathcal{F}_t) = 1_{(X_t < 0)} + 1_{(X_t > 0)} \int_0^\infty \frac{dv}{\sqrt{v}} h(A_t^- + v) \frac{M_t^h}{M_t^h}
\]

\[
(2.32)
\]

**Proof of Lemma 2.2** We note that, for $A_t \in \mathcal{F}_t$:

\[
Q^h(1_{g > t} 1_{A_t}) = Q^h(1_{A_t} 1_{X_t < 0}) + Q^h(1_{A_t} 1_{X_t > 0} 1_{d_t < \infty})
\]

(\text{where } d_t \text{ denotes the first return time to 0 after time } t)

\[
= Q^h(1_{A_t} 1_{X_t < 0}) + E(1_{A_t} 1_{X_t > 0} M_{d_t}^h)
\]

We have:

\[
M_{d_t}^h = \sqrt{2\pi} h(A_{d_t}^-) X_{d_t}^+ + \int_0^\infty \frac{dv}{\sqrt{v}} h(A_{d_t}^- + v) = \int_0^\infty \frac{dw}{\sqrt{w}} h(A_t^- + v) \quad \text{(from (2.14))}
\]

since $X_{d_t} = 0$ and $A_{d_t}^- = A_t^-$ on the set $\{X_t > 0\}$.

Hence:

\[
Q^h(1_{g > t} 1_{A_t}) = Q^h\left(1_{A_t} \left(1_{X_t < 0} + 1_{X_t > 0} \int_0^\infty \frac{dw}{\sqrt{w}} h(A_t^- + w) \frac{M_t^h}{M_t^h}\right)\right)
\]

This proves (2.32) and Lemma 2.2.
2.3.5 Proof of \( Q^h(g < \infty) = 1 \)

We deduce from (2.32) that:

\[
Q[g < t] = 1 - Q[g > t]
\]

\[
= 1 - E\left[1_{X_t < 0} M_t^h + 1_{X_t > 0} \int_0^\infty \frac{dv}{\sqrt{v}} h(A_t - v)\right]
\]

\[
= \sqrt{2\pi} E[X_t^+ h(A_t^-)]
\]

(from 2.26) and since \( (M_t^h, t \geq 0) \) is a martingale s.t. \( E(M_t^h) = 1 \)

\[
= \sqrt{2\pi} E\left(\int_0^t h(A_s^-) dL_s\right)
\]

from Itô-Tanaka formula

\[
\xrightarrow{t \to \infty} \sqrt{\frac{\pi}{2}} E\left(\int_0^\infty h(A_s^-) dL_s\right) = \sqrt{\frac{\pi}{2}} E\left(\int_0^\infty h(a) dL_a\right)
\]

(where \( (\sigma_a, a \geq 0) \) denotes the right continuous inverse of \( (A_t^-, t \geq 0) \))

\[
= 2 \sqrt{\frac{\pi}{2}} E\left(\int_0^\infty h(a) dL_a\right)
\]

(since \( (X_{\sigma_a}, a \geq 0) \) is distributed as \( (|X_a|, a \geq 0; \text{cf point a}) \) of Subsection 2.3.3)

\[
= 2 \sqrt{\frac{\pi}{2}} \int_0^\infty h(a) E(dL_a) = \int_0^\infty \frac{h(a)}{\sqrt{a}} da = 1
\]

since \( E(L_a) = \sqrt{a} \cdot \sqrt{\frac{2}{\pi}} \).

2.3.6 Description of the canonical process \((X_t, t \geq 0)\) under \( Q^h \)

For this purpose, we shall use the technique of enlargement of filtrations. Thus, let \((G_t, t \geq 0)\) denote the smallest filtration which makes \( g \) a \((G_t, t \geq 0)\) stopping time, and which contains \((F_t, t \geq 0)\).

The application of Girsanov’s Theorem and (2.14) imply the existence of a \((F_t, Q^h)\) Brownian motion \((\beta_t, t \geq 0)\) such that, under \( Q^h \):

\[
X_t = \beta_t + \int_0^t \left\{ \sqrt{2\pi} h(A_s^-) 1_{X_s > 0} - \left( \int_0^\infty \frac{dw}{w^{3/2}} e^{-\frac{(X_s^-)^2}{2w}} h(A_s^- + w)\right) X_s^- \right\} ds \quad (2.33)
\]

We now apply the enlargement formulae (cf [J], [JY], [MY]). We first observe that:

\[
dZ_t = -\int_0^\infty \frac{dw}{\sqrt{w} h(A_t^- + w)} \left( \sqrt{2\pi} h(A_t^-) 1_{X_t > 0} dX_t \right) + dV_t \quad (2.34)
\]

where \((V_t, t \geq 0)\) has bounded variations and therefore:

\[
d < Z, X >_t = -\int_0^\infty \frac{dw}{\sqrt{w} h(A_t^- + w)} \sqrt{2\pi} h(A_t^-) 1_{(X_t > 0)} dt \quad (2.35)
\]
Thus, there exists a \((G_t, t \geq 0), Q^h\) Brownian motion \((\tilde{\beta}_t, t \geq 0)\) such that:
\[
\begin{align*}
\frac{dX_t}{dt} &= \frac{d\tilde{\beta}_t}{dt} + \frac{1}{M_t^h} \left\{ \sqrt{2\pi} h(A_t^-) 1_{X_t < 0} - \left( \int_0^{\infty} \frac{dw}{w^{3/2}} e^{-w \frac{(X_t^-)^2}{2}} h(A_t^- + w) \right) X_t^- \right\} dt \\
&\quad + 1_{t < g} \left[ - \frac{f_{\infty}}{\sqrt{\pi} M_t^h} h(A_t^-) + \int_0^{\infty} \frac{dw}{w^{3/2}} (2\sqrt{2\pi} h(A_t^-) 1_{X_t > 0}) \right] dt \\
&\quad - 1_{t > g} \left[ - \frac{f_{\infty}}{\sqrt{\pi} M_t^h} h(A_t^-) + \int_0^{\infty} \frac{dw}{w^{3/2}} (2\sqrt{2\pi} h(A_t^-) X_t^+) \right] dt
\end{align*}
\]
This yields, after some simplifications:
\[
X_t = \tilde{\beta}_t - \int_0^{t \wedge g} \frac{1}{M_s} \left( \int_0^{\infty} \frac{dw}{w^{3/2}} e^{-w \frac{(X^-_s)^2}{2}} h(A_s^- + w) \right) X_s^- ds + \int_0^t \frac{ds}{L_s} X_s.
\] (2.36)

since, after \(g, X_t^- = 0\), hence \(X_t^+ = X_t\).

Points 4 iii) and iv) now follow immediately from (2.36).

**Remark 2.3** When \(h(x) = e^{-\frac{x^2}{2}}\) \((\lambda > 0, x \geq 0)\), the equation (2.36) simplifies as:
\[
\begin{align*}
X_t = \tilde{\beta}_t &- \int_0^{t \wedge g} \sqrt{X_s^-} \lambda^2 \left( \frac{X^-_s}{\lambda} \right)^{1/4} 2 K_{-1/2}(\sqrt{\lambda} X^-_s) ds + \int_0^t \frac{ds}{L_s} X_s.
\end{align*}
\] (2.37)

This formula (2.37) follows from:
\[
\begin{align*}
\int_0^{\infty} \frac{dw}{w^{3/2}} e^{-w \frac{(X^-_s)^2}{2}} \lambda^2 \frac{1}{2} w dw &= (\frac{(X^-_s)^2}{\lambda})^{1/4} 2 K_{-1/2}(\sqrt{\lambda} X^-_s) \\
\int_0^{\infty} \frac{dw}{w^{1/2}} e^{-w \frac{(X^-_s)^2}{2}} \frac{1}{\lambda} w dw &= (\frac{(X^-_s)^2}{\lambda})^{1/4} 2 K_{1/2}(\sqrt{\lambda} X^-_s)
\end{align*}
\]
and from: \(K_{-1/2}(z) = K_{1/2}(z) = \left( \frac{\pi}{2z} \right)^{1/2} e^{-z}\) ([Leb], p. 112 and p. 119).

### 2.3.7 Markovian limit process

Theorem 2.1 shows that the process \((X_t, t \geq 0)\) is not Markovian under \(Q^h\), whereas the 2-dimensional process \((\tilde{X}_t, \tilde{A}_t), t \geq 0\) is Markovian.

Indeed, \(g\) is not a \((\mathcal{F}_t, t \geq 0)\) stopping time and the dynamics of \((X_t)\) is not the same before and after \(g\). On the other hand, we know (see [RVY, I]) that if \(h(x) = e^{-\frac{x^2}{2}}\) \((\lambda, x \geq 0)\), then the \(Q^h\)-process is Markovian. It is the diffusion with infinitesimal generator \(L^h\):
\[
L^h f(x) = \frac{1}{2} f''(x) + \frac{\varphi'}{\varphi}(x) f'(x), \quad f \in C^2_0
\]
where \(\varphi\) denotes the unique solution of \(\varphi'' = \lambda \varphi, \varphi(-\infty) = 0; \varphi'(+\infty) = \sqrt{\frac{2}{\pi}}\). In this case, the solution of this equation (see (1.24)) takes the explicit form:
\[
\varphi_\lambda(x) = \sqrt{\frac{2}{\pi}} \left\{ e^{\sqrt{\lambda} x} \frac{1}{\sqrt{\lambda}} 1_{x \leq 0} + \left( x + \frac{1}{\sqrt{\lambda}} \right) 1_{x > 0} \right\}
\] (2.38)
Under \( Q^h \), we obtain:

\[
X_t = B_t + \int_0^t \frac{du}{X_u^{+} + \frac{1}{\sqrt{\lambda}}} \quad \text{(compare with (2.37))} \tag{2.39}
\]

where \((B_t, t \geq 0)\) is a \(((\mathcal{F}_t, t \geq 0), Q^h)\) Brownian motion. The martingale \((M^h_s, s \geq 0)\) is equal to:

\[
M^h_s = \varphi_{\lambda}(X_s) \exp\left(-\frac{\lambda}{2} \int_0^s 1_{[-\infty,0]}(X_u) \, du\right) \tag{2.40}
\]

This example motivated us to raise the question: which are the functions \( h \) such that the \( Q^h \) process is Markovian? The answer is given by the following:

**Proposition 2.4**: Let \( h \) be regular, bounded, satisfying equation (2.10) i.e., \( \int_0^\infty \frac{dy}{\sqrt{y}} h(y) = 1 \) and such that the process \((X_t, t \geq 0)\) is Markov under \( Q^h \). Then, there exists \( \lambda \geq 0 \) such that

\[
h(x) = \sqrt{\frac{2\pi}{\lambda^2}} e^{-\frac{x^2}{\lambda^2}} (x \geq 0).
\]

**Proof of Proposition 2.4**

To answer this question, we come back to equation (2.33). The problem is to find under which conditions the drift term:

\[
\frac{\sqrt{2\pi} h(A_t^{-}) 1_{X_t > 0} - \left( \int_0^\infty \frac{dw}{w^{3/2}} e^{-\frac{(X_t^{-})^2}{2w}} h(A_t^{-} + w) \right) X_t^{-}}{\sqrt{2\pi} h(A_t^{-}) X_t^{+} + \int_0^\infty \frac{dw}{w^{1/2}} e^{-\frac{(X_t^{-})^2}{2w}} h(A_t^{-} + w)} \tag{2.41}
\]

does not depend on \( A_t^{-} \). Considering this expression when \( X_t < 0 \), the problem amounts to study the functions \( h \) for which:

\[
x \int_0^\infty \frac{dw}{w^{3/2}} e^{-\frac{x^2}{2w}} h(a + w) \int_0^\infty \frac{dw}{w^{1/2}} e^{-\frac{x^2}{2w}} h(a + w) := k(x) \tag{2.42}
\]

does not depend on \( a \). (2.42) may be written:

\[
\frac{\partial}{\partial x} \log(\theta(x, a)) = -k(x)
\]

where we have denoted: \( \theta(x, a) := \int_0^\infty \frac{dw}{\sqrt{w}} e^{-\frac{x^2}{2w}} h(a + w) \).

Hence, by integration we obtain the existence of two functions \( \varphi_1 \) and \( \varphi_2 \) such that:

\[
\int_0^\infty \frac{x}{\sqrt{2\pi w^3}} e^{-\frac{x^2}{2w}} h(a + w) \, dw = \varphi_1(a) \varphi_2(x) \tag{2.43}
\]

Letting \( x \to 0 \) in (2.43), we obtain \( h(a) = \varphi_1(a) \varphi_2(0) \). Note that the LHS in (2.43) writes \( E[h(a + T_x)] \), where \((T_x, x \geq 0)\) is the \( \frac{1}{2} \)-stable subordinator of Brownian first hitting times. Hence we have:

\[
E[h(a + T_x)] = P_x(h)(a) = E[\varphi_1(a + T_x) \varphi_2(0)] = \varphi_1(a) \varphi_2(0) \tag{2.44}
\]

23
where \((P_x, x \geq 0)\) denotes the semi-group associated with the subordinator \((T_x, x \geq 0)\), whose infinitesimal generator is \(\left(\frac{\partial^2}{\partial x^2}\right)^{\frac{1}{2}}\). In other terms, from (2.43), we get:

\[
P_x \varphi_1(a) = \frac{\varphi_2(x)}{\varphi_2(0)} \varphi_1(a)
\]

(2.45)

\(\varphi_1\) is an eigenfunction of \(P_x\), and consequently an eigenfunction of \(\frac{\partial^2}{\partial x^2}\). \(\varphi_1\) being positive and bounded:
\[
\varphi_1(a) = c e^{-\frac{\lambda}{2} a} \quad (a, \lambda \geq 0) \quad \text{and} \quad h(a) = c e^{-\frac{\lambda}{2}} \varphi_2(0) = c' e^{-\frac{\lambda}{2}}.
\]

3 A local limit theorem for a class of additive functionals of the ”long Brownian bridges”

3.1 Statement of Theorem 3.1

In this section, our aim is to obtain results similar to those in Section 1, but, now, Brownian motion \((X_s, s \geq 0)\) is being replaced by the Brownian bridge with length \(t\), with \(t \to \infty\). \(q\) denotes a function from \(\mathbb{R}\) to \(\mathbb{R}^+\), which is Borel, and such that:

\[
0 < \int_{-\infty}^{\infty} (1 + x^2) q(x) \, dx < \infty.
\]

(3.1)

We let:

\[
A^q_t := \int_0^t q(X_s) \, ds
\]

(3.2)

**Theorem 3.1**

1) For every \(x\) and \(y \in \mathbb{R}\), and \(\mu > 0\):

\[
E_x \left( \exp \left( -\frac{\mu}{2} \int_0^t q(X_s) \, ds \right) \bigg| X_t = y \right) \sim \frac{\pi}{2} \frac{\varphi_{\mu q}(x) \varphi_{\mu q}(y)}{t}
\]

where \(\varphi_{\mu q}\) denotes the unique solution of:

\[
\varphi'' = (\mu q) \cdot \varphi, \quad \lim_{x \to +\infty} \varphi'(x) = -\lim_{x \to -\infty} \varphi'(x) = \sqrt{\frac{2}{\pi}}
\]

(3.4)

2) \(\lim_{t \to \infty} t \ P_x \ (A^q_t \in dz | X_t = y) = \nu_x * \nu_y (dz)\)

(3.5)

where \(\nu_x\) and \(\nu_y\) have been defined in Theorem 1.1. The convergence in (3.5) has the same meaning as in Theorem 1.1.

3.2 Proof of Theorem 3.1

Without loss of generality, we shall assume that \(\mu = 1\).
3.2.1 An auxiliary lemma

**Lemma 3.2** There exists a constant $C > 0$, depending only on $q$, such that:

$$E_x \left( \exp \left( -\frac{1}{2} \int_0^t q(X_s) \, ds \right) \big| X_t = y \right) \leq C \, e^{\frac{(x-y)^2}{2t}} \frac{1 + |x| (1 + |y|)}{1 + t}$$  \hspace{1cm} (3.6)

3.2.2 Proof of lemma 3.2

1) As an intermediary result, we already show that:

$$E_x \left( \exp \left( -\frac{1}{2} \int_0^t q(X_s) \, ds \right) \big| X_t = y \right) \leq C \, e^{\frac{(x-y)^2}{2t}} \frac{1 + |x|}{\sqrt{1 + t}}$$  \hspace{1cm} (3.7)

for a constant $C$ which does not depend on $x, y, t$.

To prove (3.7), we condition with respect to $X_{t/2}$, and we get:

$$E_x \left( \exp \left( -\frac{1}{2} \int_0^t q(X_s) \, ds \right) \big| X_{t/2} = y \right) = \mathbb{E}_x \left( \exp -\frac{1}{2} A_{t/2}^q \big| X_{t/2} = c \right)$$

$$E_x \left( \exp -\frac{1}{2} A_{t/2}^q \big| X_{t/2} = y \right) \cdot e^{-\frac{(x-y)^2}{2t}} \leq \mathbb{E}_x \left( \exp -\frac{1}{2} A_{t/2}^q \big| X_{t/2} = c \right)$$

$$\leq \mathbb{E}_x \left( \exp -\frac{1}{2} A_{t/2}^q \right) \leq C \frac{1 + |x|}{\sqrt{1 + t}}$$

In (3.8), we majorize $E_x \left( \exp -\frac{1}{2} A_{t/2}^q \big| X_{t/2} = y \right)$ by 1, and we get:

$$E_x \left( \exp -\frac{1}{2} A_{t/2}^q \big| X_{t/2} = y \right) \cdot e^{-\frac{(x-y)^2}{2t}} \leq \frac{1}{\sqrt{\pi t}} \int_{-\infty}^{\infty} E_x \left( \exp -\frac{1}{2} A_{t/2}^q \big| X_{t/2} = c \right) e^{-\frac{(x-c)^2}{2t}} \, dc$$

from Lemma 4.3 in [RVY, I]. Thus, we have obtained (3.7).

2) Then, plugging the estimate (3.7) in (3.8), we obtain:

$$E_x \left( \exp \left( -\frac{1}{2} A_{t/2}^q \big| X_{t/2} = y \right) \right) \leq e^{\frac{(x-y)^2}{2t}} \frac{C (1 + |y|)}{\sqrt{1 + t}} \int_{-\infty}^{\infty} E_x \left( \exp -\frac{1}{2} A_{t/2}^q \big| X_{t/2} = c \right) e^{-\frac{(x-c)^2}{2t}} \, dc$$

$$\leq \frac{C (1 + |y|)}{\sqrt{1 + t}} \cdot e^{\frac{(x-y)^2}{2t}}$$

since:

$$E_x \left( \exp -\frac{1}{2} A_{t/2}^q \big| X_{t/2} = y \right) = E_y \left( \exp -\frac{1}{2} A_{t/2}^q \big| X_{t/2} = c \right)$$

and

$$e^{\frac{(x-y)^2}{2t}} \cdot \frac{e^{\frac{(1 + |y|)}{2}}}{\sqrt{1 + t}} \cdot e^{\frac{(x-y)^2}{2t}} \leq e^{\frac{(x-y)^2}{2t}} \frac{C (1 + |y|)}{1 + t}$$

by applying once again Lemma 4.3 in [RVY, I].
3.2.3 Another auxiliary lemma

Lemma 3.3 Let \( Z(t, x, y) := E_x \left( \exp - \frac{1}{2} A_t^0 | X_t = y \right) \). We also denote by \( U(t, x, y) \) the solution of:

\[
\begin{align*}
\frac{\partial U}{\partial t}(t, x, y) - \frac{1}{2} \frac{\partial^2 U}{\partial x^2}(t, x, y) + \frac{1}{2} U(t, x, y) q(x) &= 0 \\
U(0, \bullet, y) &= \delta_y
\end{align*}
\]

(3.9)

Then:

\( Z(t, x, y) = \sqrt{2\pi t} \exp \left( \frac{(x-y)^2}{2t} \right) U(t, x, y) \).

In particular, it follows from Lemma 3.2, that:

\[
U(t, x, y) \leq C \left( 1 + |x| \right) \left( 1 + |y| \right)^3 (t \geq 1)
\]

(3.10)

3.2.4 Proof of Lemma 3.3

We know that, for every regular function \( f \):

\[
Z^f(t, x) := E_x \left[ \exp \left( - \frac{1}{2} A_t^0 \right) f(X_t) \right]
\]

is solution of:

\[
\frac{\partial Z^f}{\partial t} - \frac{1}{2} \frac{\partial^2 Z^f}{\partial x^2} + \frac{1}{2} Z^f \cdot q = 0, \quad Z^f(0, x) = f(x)
\]

(3.11)

It suffices, in order to obtain Lemma 3.3, to write:

\[
Z(t, x, y) = \lim_{\varepsilon \downarrow 0} \frac{E_x \left[ \left( \exp \left( - \frac{1}{2} A_t^0 \right) \right) f_\varepsilon(X_t) \right]}{E_x(f_\varepsilon(X_t))}
\]

where \( f_\varepsilon \) is a family of functions which converges weakly towards \( \delta_y \), and to use (3.11).

3.2.5 We define, for every \( \lambda > 0 \)

\[
A(\lambda, x, y) = \int_0^\infty e^{-\lambda t} U(t, x, y) \, dt
\]

(3.12)

Since \( Z(t, x, y) \) is a decreasing function of \( t \), we deduce the following equivalences from the Tauberian theorem:

\[
\begin{align*}
i) \quad Z(t, x, y) &\sim_{t \to \infty} \frac{\pi}{2} \frac{\varphi_q(x) \varphi_q(y)}{t} \\
ii) \quad U(t, x, y) &\sim_{t \to \infty} \frac{\sqrt{\pi}}{2\sqrt{2}} \frac{\varphi_q(x) \varphi_q(y)}{t^{3/2}} \\
iii) \quad \left| \frac{\partial}{\partial \lambda} A(\lambda, x, y) \right| &\sim_{\lambda \to 0} \frac{1}{2\sqrt{2}} \frac{\varphi_q(x) \varphi_q(y)}{\sqrt{\lambda}}
\end{align*}
\]

(3.13)
We shall now show (3.13). We already deduce from Lemmas 3.2 and 3.3 that:

\[ \lim_{\lambda \to 0} A(\lambda, x, y) = \int_0^\infty U(t, x, y) \, dt < \infty \quad (3.14) \]

\[ A(\lambda, x, y) \leq C (1 + |x|)(1 + |y|) \quad (3.15) \]

\[ \left| \frac{\partial}{\partial \lambda} A(\lambda, x, y) \right| \leq \frac{C}{\sqrt{\lambda}} (1 + |x|)(1 + |y|) \quad (3.16) \]

To prove (3.13) we shall show that:

\[ \psi(x, y) := \lim_{\lambda \to 0} \sqrt{\lambda} \frac{\partial}{\partial \lambda} A(\lambda, x, y) \]

satisfies the Sturm-Liouville equation (for any fixed \( y \)):

\[ \frac{\partial^2}{\partial x^2} \psi = \psi q, \quad \text{with adequate limit conditions in } x = \pm \infty \quad (3.17) \]

### 3.2.6 We get, from (3.9)

\[ U(t, x, y) = \frac{1}{\sqrt{2\pi t}} e^{-\frac{(x-y)^2}{2t}} - \frac{1}{2} \int_0^t ds \int_{-\infty}^\infty \frac{e^{-\frac{(x-z)^2}{2(t-s)}}}{\sqrt{2\pi (t-s)}} U(s, z, y) q(z) dz \quad (3.18) \]

Thus, after taking the Laplace transform in the variable \( t \) of the two sides of (3.18), we obtain:

\[ A(\lambda, x, y) = g_\lambda(x, y) - \frac{1}{2} \int_{-\infty}^\infty g_\lambda(x, z) A(\lambda, z, y) q(z) dz \quad (3.19) \]

where \( g_\lambda \) denotes the density of the resolvent kernel of Brownian motion:

\[ g_\lambda(x, z) = \frac{1}{\sqrt{2\lambda}} e^{-|x-z|\sqrt{2\lambda}} \]

We write (3.19) in the form:

\[ A(\lambda, x, y) = G_\lambda \left[ \delta_y - \frac{1}{2} (A(\lambda, \bullet, y) q(\bullet)) (x) \right] \quad (3.20) \]

with for any Radon measure \( \mu(dz) \):

\[ G_\lambda \mu(x) := \int_{-\infty}^\infty g_\lambda(x, z) \mu(dz) \quad (3.21) \]

and we use the resolvent equation:

\[ \frac{\partial^2}{\partial x^2} G_\lambda \mu = -2\mu + 2\lambda G_\lambda \mu, \]

to obtain:

\[ \frac{\partial^2}{\partial x^2} A(\lambda, x, y) = 2\lambda A(\lambda, x, y) - \left[ 2\delta_y - A(\lambda, x, y) q(x) \right] \quad (3.22) \]

As a consequence, differentiating with respect to \( \lambda \), then multiplying by \( \sqrt{\lambda} \), we obtain:

\[ \frac{\partial^2}{\partial x^2} \left( \sqrt{\lambda} \frac{\partial A}{\partial \lambda} (\lambda, x, y) \right) - \sqrt{\lambda} \frac{\partial A}{\partial \lambda} (\lambda, x, y) q(x) = 2\sqrt{\lambda} A(\lambda, x, y) + 2\lambda^{3/2} \frac{\partial A}{\partial \lambda} (\lambda, x, y) \quad (3.23) \]

Hence, from (3.16) and (3.15), and denoting \( \tilde{A}(\lambda, x, y) := \sqrt{\lambda} \frac{\partial A}{\partial \lambda} (\lambda, x, y) \), it follows that:

\[ \left| \frac{\partial^2}{\partial x^2} \tilde{A}(\lambda, x, y) - \tilde{A}(\lambda, x, y) q(x) \right| \leq C \sqrt{\lambda} (1 + |x|)(1 + |y|) \quad (\lambda \to 0) \quad (3.24) \]

(3.24) is the first step to prove that \( \tilde{A}(\lambda, x, y) \) converges, as \( \lambda \to 0 \), to a solution of the Sturm-Liouville equation (3.17).
3.2.7 We now examine the limit conditions in \( x = \pm \infty \)

We come back to equation (3.19) which we differentiate with respect to \( \lambda \), then we multiply by \( \lambda \):

\[
\sqrt{\lambda} \tilde{A}(\lambda, x, y) = -\frac{1}{2} A(\lambda, x, y) - \frac{1}{2} \int_{-\infty}^{\infty} e^{-|x-z|\sqrt{\lambda}} |x-z| \left( \delta_y (dz) - A(\lambda, z, y) q(z) dz \right) \\
- \frac{1}{\sqrt{2}} \int_{-\infty}^{\infty} e^{-|x-z|\sqrt{\lambda}} A(\lambda, x, z) q(z) dz
\]  
(3.25)

From (3.16) and (3.14), respectively we deduce that:

\[
\sqrt{\lambda} \tilde{A}(\lambda, x, y) \xrightarrow{\lambda \to 0} 0 \quad \text{and} \quad A(\lambda, x, y) \quad \text{converges as} \quad \lambda \to 0.
\]

hence, from (3.25), since \( \int_{-\infty}^{\infty} (1 + x^2) q(dz) < \infty \):

\[
\lim_{\lambda \to 0} \int_{-\infty}^{\infty} e^{-|x-z|\sqrt{\lambda}} \tilde{A}(\lambda, x, z) q(z) dz \quad \text{exists.} \quad (3.26)
\]

On the other hand, differentiating (3.25) with respect to \( x \), we obtain:

\[
\frac{\partial \tilde{A}}{\partial x}(\lambda, x, y) = \frac{\partial B}{\partial x} = \frac{1}{2} \left\{ \int_{-\infty}^{\infty} e^{-|x-z|\sqrt{\lambda}} \tilde{A}(\lambda, z, y) q(z) dz + \int_{x}^{\infty} e^{-|x-z|\sqrt{\lambda}} \tilde{A}(\lambda, z, y) q(z) dz \right\}
\]  
(3.27)

with

\[
B := -\frac{1}{\sqrt{2}} \left\{ A + \int_{-\infty}^{\infty} e^{-|x-z|\sqrt{\lambda}} (A(\lambda, z, y) q(z) dz - \delta_y (dz)) \right\}
\]  
(3.28)

We deduce from (3.27), (3.26) and (3.28) that:

\[
\lim_{\lambda \to 0} \frac{\partial}{\partial x} \tilde{A}(\lambda, x, y) = - \lim_{\lambda \to 0} \frac{\partial}{\partial x} \tilde{A}(\lambda, x, y) = C(y) \quad (3.29)
\]
(cf [RVY, I], p. 194-197 for similar computations). Thus, from the equivalence between i), ii) and iii) which we recalled in (3.13), we get:

\[
E_x \left[ \exp - \frac{1}{2} \int_{0}^{t} q(X_s) ds \left| X_t = y \right. \right] \xrightarrow{t \to \infty} \frac{\psi(x, y)}{t}
\]  
(3.30)

where \( \psi \) is solution to:

\[
\frac{\partial^2 \psi}{\partial x^2}(x, y) = \psi(x, y) q(x), \quad \lim_{x \to +\infty} \frac{\partial \psi}{\partial x}(x, y) = - \lim_{x \to -\infty} \frac{\partial \psi}{\partial x}(x, y) = C(y)
\]  
(3.31)

Thus, from the definition of \( \varphi_q \) (see (3.4)), we get:

\[
\psi(x, y) = C(y) \sqrt{\frac{\pi}{2}} \varphi_q(x)
\]

Now, since \( Z(t, x, y) \) is symmetric in \( x \) and \( y \):

\[
\psi(x, y) = K \varphi(x) \varphi(y)
\]  
(3.32)
It remains to determine the value of $K$. For this purpose, we write:

$$
\varphi_q(x) = E_x \left( \exp \left( - \frac{1}{2} \int_0^t ds \, q(X_s) \right) \right) \varphi_q(X_t)
$$

(since $\varphi_q(X_t)$ exp$\left( - \frac{1}{2} A^q_t \right)$, $t \geq 0$ is a martingale)

$$
= \int_{-\infty}^\infty E_x \left( \exp \left( - \frac{1}{2} A^q_t \right) | X_t = y \right) e^{-\frac{(x-y)^2}{2t}} \varphi_q(y) dy
$$

$$
\sim_{t \to \infty} K \int_{-\infty}^\infty \frac{\varphi_q(x) \varphi_q(y)}{t} e^{-\frac{(x-y)^2}{2}} \varphi_q(y) dy
$$

$$
= \frac{K \varphi_q(x)}{t} E_x(\varphi^2_q(X_t)) \sim_{t \to \infty} \frac{K \varphi_q(x)}{2 \pi t}
$$

since $\varphi_q(z) \sim_{|z| \to \infty} \sqrt{\frac{2}{\pi} |z|}$. Hence $K \frac{2}{\pi} = 1$, that is: $K = \frac{\pi}{2}$.

Thus, we have obtained point 1 of Theorem 3.1.

3.2.8 Point 2 of Theorem 3.1 may be proven

with the help of (3.3), exactly as Theorem 1.1.

**Remark 3.2** Under our hypothesis H1 on $q$, there is the equivalence:

$$
Z(t, x, y) \equiv E_x \left( \exp \left( - \frac{1}{2} \int_0^t q(X_s) ds \right) | X_t = y \right) \sim_{t \to \infty} \frac{\pi}{2t} \varphi_q(x) \varphi_q(y)
$$

(3.33)

Intuitively, we may think of the bridge of duration $t$ going from $x$ to $y$ as "resembling", as $t \to \infty$, to the concatenation of two Brownian motions each being defined on a time interval $[0, \frac{t}{2}]$, with the first one starting from $x$ and the second one, after time reversal, starting from $y$, these two parts being independent. If this were true, then:

$$
Z(t, x, y) = E_x \left( \exp \left( - \frac{1}{2} A^q_t \right) | X_t = y \right) = E_x \left( \exp \left( - \frac{1}{2} A^q_{t/2} \right) \right) \cdot E_y \left( \exp \left( - \frac{1}{2} A^q_{t/2} \right) \right)
$$

$$
\sim_{t \to \infty} \frac{\varphi_q(x)}{\sqrt{t/2}} \cdot \frac{\varphi_q(y)}{\sqrt{t/2}} = \frac{4}{\pi} \left( \frac{\pi \varphi_q(x) \varphi_q(y)}{t} \right)
$$

Thus, comparing with (3.33) the factor $\frac{4}{\pi}$ which we just obtained measures, in some sense, the default of independence of these two Brownian components.

**Remark 3.3** Theorem 3.1 allows to "penalize long Brownian Bridges". More precisely, for every $s \geq 0$ and $\Lambda_s \in \mathcal{F}_s$:

$$
\frac{E_x \left( 1_{\Lambda_s} \exp \left( - \frac{1}{2} A^q_t \right) | X_t = y \right)}{E_x \left( \exp \left( - \frac{1}{2} A^q_t \right) | X_t = y \right)} \sim_{t \to \infty} E_x \left( 1_{\Lambda_s M^q_s} \right)
$$

(3.34)

29
with \( M^\varphi_s := \frac{\varphi_q(X_s)}{\varphi_q(x)} \exp\left( -\frac{1}{2} A^q_s \right) \). et \((M^\varphi_s, s \geq 0)\) is a positive martingale. In other terms, comparing with Theorem 5.1 in [RVY, I], the penalisation is the same for "long bridges" as for Brownian motion itself. Once more (see [RVY, III]), we obtain that a long bridge of duration \( t \), as \( t \to \infty \), behaves as a standard Brownian motion.

### 3.2.9 Finally, we show (3.34)

\[
\frac{E_x \left( 1_{\Lambda_s} \exp\left( -\frac{1}{2} A^q_t \right) | X_t = y \right) }{E_x \left( \exp\left( -\frac{1}{2} A^q_t \right) | X_t = y \right) } \xrightarrow{t \to \infty} \frac{E_x \left( 1_{\Lambda_s} \exp\left( -\frac{1}{2} A^q_t \right) \cdot \frac{\varphi_q(X_s)\varphi_q(y)}{\varphi_q(x)\varphi_q(y)} \right) }{t - s } \xrightarrow{t \to \infty} E_x \left( 1_{\Lambda_s} M^\varphi_s \right).
\]
References


