Topological correlations and asymptotic freedom in cellular aggregates
Faez Miri, Christophe Oguey

To cite this version:

HAL Id: hal-00171704
https://hal.archives-ouvertes.fr/hal-00171704
Submitted on 13 Sep 2007

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L’archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d’enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.
Topological correlations and asymptotic freedom in cellular aggregates

F. Miri\textsuperscript{a}, C. Oguey\textsuperscript{b,*}

\textsuperscript{a}IASBS, Zanjan, Iran
\textsuperscript{b}LPTM, Univ. de Cergy-Pontoise, France\textsuperscript{1}

Abstract

In random cellular systems, both observation and maximum entropy inference give a specific form to the topological pair correlation: it is bi-affine in the cells number of edges with coefficients depending on the distance between the two cells of the pair. Assuming this form for the pair correlations, we make explicit the conditions of statistical independence at large distance. When, on average, the defects do not contribute, the layer population and the enclosed topological charge both increase polynomially with distance. In dimension 2, the exponent of the leading terms depend on sum rules satisfied, or not, by the maximum entropy coefficients.

Key words: Foam, cells, space partition, correlation decay, sum rules, topological charge
PACS: 89.75.Fb, 82.70.Rr, 87.18.Hf, 02.10.Ox

1 Introduction

Foams are random but the precise probability distributions describing their structure are still largely unknown. Leaving aside the geometrical and energetic details, obviously dependent on the system considered, we focus on the topological characteristics, believed to be more universal. Only a few investigations have been devoted to correlations beyond nearest neighbours. Pair correlations at arbitrary distance were analysed by Fortes-Pina [1], Rivier \textit{et al} [2,3], Szeto \textit{et al} [4], and measured in [4].

\textsuperscript{*} Corresponding author. Address: LPTM, Univ. de Cergy-Pontoise, 2, ave A. Chauvin, F–95302 Cergy-Pontoise
\textit{Email addresses: miri@iasbs.ac.ir, oguey@ptm.u-cergy.fr} (C. Oguey).
\textsuperscript{1} LPTM is CNRS UMR 8089

Preprint submitted to Elsevier Science 13 September 2007
Maximum entropy is one of the few methods able to predict some aspects of the probability distributions. Maximum entropy (maxent) arguments yield a specific form for the pair correlations at arbitrary distance [2,5].

Here, we examine the asymptotic behaviour of pair correlations and how the independence of occurrences at large distances constrains the parameters in the maxent formulae.

The present analysis is mostly devoted to bi-dimensional foams, subject to extensive theoretical and experimental investigations [1,2,3,4,5,6,7,8,9,10,11,12,13,14,15,16,17,18,19,20] and where the progresses are ahead of those on 3D foams.

2 Foam statistics

A foam $F$ divides space into $N = |F|$ polygonal cells. Here $F$ is viewed as a set of cells and $|F|$ represents the number of elements in the set.

2.1 One cell statistics

The local observable is $n$, the number of sides of each cell (“polygonality”), or the topological charge $q = 6 - n$. The fraction of $n$-sided cells is $p(n) = N(n)/N$.

The generic and physically stable vertex coordination (or degree) is $z = 3$. Euler’s theorem implies:

$$\langle q \rangle = \langle 6 - n \rangle = \sum_n (6 - n) p(n) = 6\chi/N \quad (1)$$

where $\chi$ characterises the embedding space: $\chi = 2$ for the sphere, 0 for the torus$^2$.

In the large $N$ limit, $\langle q \rangle \to 0$. The second moment is $\mu_2 = \langle q^2 \rangle = \langle (n - 6)^2 \rangle$. 

2
2.2 Two cells, correlations

The topological distance $j$ between cells is measured in nearest neighbour steps. The $j$th layer around a given cell $o$, $\text{lay}(j\mid o)$, is the set of cells at distance $j$ from $o$ (Fig. 2). It has population $K_j(o) = |\text{lay}(j\mid o)|$. The average over $n$-sided central cells is $\langle K_j(n) \rangle$ and the overall average is $\langle \langle K_j \rangle \rangle$.

As stated here, equation (1) only holds if the foam fills the 2D space, without boundary. Otherwise, the edges on the boundary must be counted only once in $\langle n \rangle N = 2E - E_{\text{boundary}}$, implying a additional boundary term in (1) [21]. $E$ is the total number of edges, internal and on the boundary. In particular, this is so for a bounded cluster in the plane, where $\chi = 1$. In any reasonable case, $\langle q \rangle = \langle 6 - n \rangle \to 0$ as $N \to \infty$. 

---

Fig. 1. A pair of cells, $(k, n) = (6, 5)$, at distance $j$.

Fig. 2. Stratification of a Poisson-Voronoi froth into layers.
The joint distribution \( p^{(2)}(k, n) \) — the probability that a \((k, n)\)-sided pair of cells occur at mutual distance \( j \) — and the corresponding marginal distribution \( s_j(n) = \sum_k p^{(2)}(k, n) \) — the probability that a cell is at distance \( j \) from an \( n \)-sided one\(^3\) — satisfy \( s_j(n) = \langle K_j(n) \rangle p(n) \).

The correlator \( A_j(k, n) \) and correlation function \( g_j(k, n) \) are defined by

\[
p^{(2)}_j(k, n) = A_j(k, n) \frac{p(k) p(n)}{\langle \langle K_j \rangle \rangle} = g_j(k, n) s_j(k) s_j(n).
\]

Both account for the statistical dependence of the concurrent occurrence of a \( k \) and a \( n \)-sided cell at distance \( j \). They only differ in the way they are normalised: the correlation function is 1 whereas the correlator is the mean population \( \langle \langle K_j \rangle \rangle \) in independent situations. The correlator is related to the layer population by \([1,2]\)

\[
\sum_k p(k) A_j(k, n) = \langle K_j(n) \rangle.
\]

A similar identity follows from counting the edges (‘polygonality’ \( P_j \) of layer \( j \)) \([1]\):

\[
\sum_k k p(k) A_j(k, n) = \langle P_j(n) \rangle.
\]

The polygonality of a set is the sum of the individual polygonalities \( n \). Because charges are additive, and simply related to \( n \) by \( q = 6 - n \), it is more natural to consider the total charge\(^4\) of layer \( j \), \( Q(\text{lay}(j|o)) = 6K_j(o) - P_j(o) \). Here we are getting closer to \([2]\). Then averaging over \( n \)-sided central cells \( o \) yields, according to (4), (5),

\[
\langle Q(\text{lay}(j|n)) \rangle = \sum_k (6 - k) p(k) A_j(k, n).
\]

3 Maximum entropy

Maximum entropy arguments (maxent) and the recursion relation to be described in sec. 5 give the following expressions, also observed in numerical

\(^3\) This is also the probability that a \( n \)-sided cell is at distance \( j \) from an other one.

\(^4\) The charge of a set of cells is the sum of the individual charges \( q = 6 - n \).
\[ A_j(k, n) = \sigma_j(k - 6)(n - 6) + a_j(n + k - 12) + b_j, \quad (7) \]
\[ \langle K_j(n) \rangle = a_j(n - 6) + b_j, \quad (8) \]

\[ \sigma_j, a_j, b_j \text{ are real parameters for } j = 1, 2, \ldots. \text{ The contributions from defects have been neglected. In the infinite foam limit, the average of (8) gives } \langle \langle K_j \rangle \rangle = b_j. \]

### 4 Asymptotic freedom

In normal systems of statistical physics, spatially distant events become uncorrelated. In foams, this was first measured by [4]. As \( j \to \infty \),

\[ p_j^{(2)}(k, n) \rightarrow s_j(k) s_j(n) \Leftrightarrow g_j(k, n) \rightarrow 1. \quad (9) \]

With

\[ g_j(k, n) = \frac{A_j(k, n)\langle\langle K_j \rangle\rangle}{\langle K_j(k)\rangle\langle K_j(n)\rangle} \]
\[ = 1 + \left( \frac{\sigma_j}{b_j} - \left( \frac{a_j}{b_j} \right)^2 \right) (k - 6)(n - 6) \quad (10) \]
and maxent, asymptotic de-correlation holds if and only if

\[ \frac{\sigma_j}{b_j} - \left( \frac{a_j}{b_j} \right)^2 \rightarrow 0. \quad (12) \]

A sufficient condition is that the ratios \( \frac{\sigma_j}{b_j} \to 0 \) and \( \frac{a_j}{b_j} \to 0 \), meaning that, in both \( A_j \) and \( K_j \), the \( n \) dependent terms would be dominated by the constant one \( (b_j) \) at large \( j \), a sensible result. In this limit, (8) implies \( \langle K_j(n) \rangle \to \langle \langle K_j \rangle \rangle \) and then \( s_j(n) \to p(n) \); the effect of the central condition vanishes, another manifestation of asymptotic de-correlation. In (9), \( s_j \) can then be replaced by \( p \).

### 5 Recursion equation

The following equation was derived in [1] and in [2,22]:

\[ \Delta \langle K_j(n) \rangle + \langle q_j(n) \rangle \langle K_j(n) \rangle = \langle I_j(n) \rangle \simeq 0. \quad (13) \]
The curly bracket quantities \( \prec q_j(n) \) (resp. \( \prec m_j(n) \)) are the cellular charge (resp. sidedness) averaged over the neighbours at topological distance \( j \) to an \( n \)-side cell. They are defined by \( \prec q_j(n) \equiv 6 - \prec m_j(n) \equiv \langle Q(\text{lay}(j|n)) \rangle / \langle K_j(n) \rangle \). \( \Delta K_j = K_{j+1} - 2K_j + K_{j-1} \) is the discrete laplacian.

The right hand side is due to the presence of defects, the cells of layer \( j \) which have no edge in common with the next layer, \( j+1 \) (Fig. 3). It is assumed that, on average, this contribution vanishes: \( \langle I_j(n) \rangle = 0 \).

Using (6) for the total charge \( \langle Q(\text{lay}(j|n)) \rangle = \prec q_j(n) \rangle \langle K_j(n) \rangle \) allows to put the recursion relation (13) in the following form:

\[
\Delta \langle K_j(n) \rangle = \sum_k (k - 6) p(k) A_j(k,n).
\]

(14)

Fig. 3. Defects (left) and an inclusion (right), marked hatched.

6 Asymptotic behaviour and sum rules

With the maxent form (7),(8) of the correlator \( A_j(k,n) \) and population \( \langle K_j(n) \rangle \), the recursion relation implies the following system, with \( \alpha_j = a_j \) and \( \beta_j = b_j - 6a_j \):

\[
\Delta \begin{pmatrix} \alpha_j \\ \beta_j \end{pmatrix} = \mu_2 \begin{pmatrix} \sigma_j \\ \alpha_j - 6\sigma_j \end{pmatrix},
\]

(15)

\[
\begin{pmatrix} \alpha_0 \\ \beta_0 \end{pmatrix} = 0, \quad \begin{pmatrix} \alpha_1 \\ \beta_1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}.
\]

(16)
The initial conditions follow from $\langle K_0(n) \rangle = 0$ and $\langle K_1(n) \rangle = n$; $\mu_2$ was defined in sec. 2.1.

The solutions involve $f_k = -\mu_2 \sigma_k$, conjectured to decrease fast at large $k$. Then, $\sum_{k=1}^{j-1} f_k \to S$ as $j \to \infty$, so that $f/S$ defines a normalised distribution (like probabilities except that some $f_k$ may be negative). Let $M_n$ be the moments of this distribution:

$$S = \sum_{k \geq 1} f_k, \quad M_n = \frac{1}{S} \sum_{k \geq 1} k^n f_k. \quad (17)$$

The solutions to (15), (16) are

$$\alpha_j = \sum_{k=1}^{j-1} k f_k + j (1 - \sum_{k=1}^{j-1} f_k), \quad \beta_j = \frac{\mu_2}{6} \left[ (j+1)j(j-1)(1-\sum f_k) + (3j^2 - 1) \sum k f_k 
- 3j \sum k^2 f_k + \sum k^3 f_k \right] + 6 \sum (j-k) f_k, \quad (18)$$

where the sums run from $k = 1$ to $j - 1$. As $j \to \infty$, they behave like

$$\alpha_j \to j(1 - S) + S M_1, \quad \beta_j \to \frac{\mu_2}{6} S \left[ (j+1)j(j-1)(1/S - 1) + (3j^2 - 1) M_1 
- 3j M_2 + M_3 \right] + 6S(j - M_1), \quad (20)$$

The leading order will be lower if the $\{f_k = -\mu_2 \sigma_k\}$ satisfy certain sum rules. Indeed, if $S = 1$, the estimates become

$$\alpha_j \to M_1, \quad \beta_j \to \frac{\mu_2}{6} \left[ (3j^2 - 1) M_1 - 3j M_2 + M_3 \right] + 6(j - M_1), \quad (22)$$

so that the average population $\langle K_j(n) \rangle$ would not grow faster than $j^2$.

The so-called Euclidean scaling \cite{3,23}, $\langle K_j(n) \rangle \sim j$ in 2D, requires a second condition:

$$M_1 = -\mu_2 \sum_{k \geq 1} k \sigma_k = 0.$$
7 Topological charge

The purpose of this section is to relate the moments $S$ and $M_k$ to the average charge contained in a ball or in a layer. It does not enforce the precise values found in the previous section, as we first hoped; but it shows that the sum rules and asymptotes, for charges and populations, are compatible.

7.1 Charge enclosed in a ball

Let $B_j(o)$ be the union of the cells at distance at most $j$ from an arbitrarily chosen cell $o$. The total charge enclosed in the cluster $B_j(o)$ satisfies, averaging over $n$-sided origins $o$,

$$\langle Q(B_j) \rangle = \left( 6 - n - \mu_2 \frac{j(j+1)}{2} \right) \left[ 1 - \sum_{k=1}^{j} f_k \right]$$

$$+ \frac{\mu_2}{2} \left[ \sum_{k=1}^{j} k^2 f_k - (2j+1) \sum_{k=1}^{j} kf_k \right].$$

(24)

where $\langle Q(B_j) \rangle$ stands for $\langle Q(B_j(n)) \rangle$. This expression follows from inserting the solutions $\alpha_j, \beta_j$ into (7, 8) and either summing the layer charges (6) from 0 to $j$ or integrating the difference equation (13). It is consistent with (21).

Indeed, if the individual charges are normal random variables, averaging to zero by Euler equation (1), their sum should behave like fluctuations:

$$|\langle Q(B_j(n)) \rangle| \leq \text{const} \langle |B_j(n)| \rangle^{1/2}.$$  

(25)

Now, if $S \neq 1$, the population $\langle |B_j(n)| \rangle \propto \sum_{l=0}^{j} \langle K_l(n) \rangle \sim j^4$ and $\langle Q(B_j(n)) \rangle \sim j^2$, as in (24) without further condition.

If, on the other hand, the charge fluctuations are less, $\langle Q(B_j(n)) \rangle = o(j^2)$, then $S = 1$ must hold, predicting a layer population $\sim j^2$, according to (23).

A possible interpretation of $S = 1$ can be deduced from (24). Define the average $\langle Q_j \rangle = \sum_n p(n) \langle Q(B_j(n)) \rangle$. It coincides with the mean charge of a ball centred at a neutral cell: $\langle Q(B_j(6)) \rangle$. Then, using (24), we can calculate the mean excess charge in the ball due to conditioning on $n$-sided central cells:

$$\langle Q(B_j(n)) \rangle - \langle Q_j \rangle = (6 - n) \left[ 1 - \sum_{k=1}^{j} f_k \right].$$

(26)
As $j \to \infty$, this excess charge becomes $(6-n)(1-S)$; then $S = 1$ means global neutrality: the central charge $q = 6-n$ is exactly screened by the excess charge $(6-n)S$ in the layers around.

### 7.2 Layer charge and Aboav-Weaire’s law

Inserting the maxent formula (8) for $\langle K_j(n) \rangle$ into (13) and using (15) gives an expression for the average charge in layer $j$:

$$\langle Q(\text{lay}(j|n)) \rangle = (n-6)f_j - \mu_2 \alpha_j,$$

or, dividing by the mean layer population,

$$\langle q_j(n) \rangle = \frac{\langle Q(\text{lay}(j|n)) \rangle}{\langle K_j(n) \rangle} = \frac{(n-6)f_j - \mu_2 \alpha_j}{\alpha_j n + \beta_j}.$$  \hspace{1cm} (28)

This is a fractional linear function of $n$ as in Aboav-Weaire’s law [24], which is in fact the first case, $j = 1$, of this sequence labelled by $j$. Because double contacts are negligible in foams, the first layer population is just the number if sides of the central cell: $K_1(n) = \langle K_1(n) \rangle = n$ and $\langle q_1(n) \rangle = \langle q_1(n) \rangle = 6 - \langle m_1(n) \rangle$, where $\langle m_1(n) \rangle$ is the mean number of sides (the polygonality) of the first neighbours of $n$-sided cells. So, with (16), equation (28) for $j = 1$ specialises to

$$\langle q_1(n) \rangle = \frac{(n-6)f_1 - \mu_2}{n} = f_1 - \frac{6f_1 + \mu_2}{n}.$$  \hspace{1cm} (29)

Compared to the usual form of Aboav-Weaire’s law [24,25]: $\langle m_1(n) \rangle = 6 - a + (6a + \mu_2)/n$, eq. (29) gives an interpretation of the parameter $a$ as the $nk$ coefficient in the pair correlator: $a = f_1 = -\mu_2 \sigma_1$.

### 8 Conclusion

To summarise, the major part of this article is a brief review of what is known, so far, on correlations in foams beyond first neighbours. This includes the definitions and sum rules in sec 2, the (bi)affine form of the correlations in sec. 3, the recursion equation in sec. 5 and the solutions in sec. 6. Except for the affine ansatz, controversially [25] justified by maximum entropy arguments [2,5], and the related hypothesis that defects do not contribute on average, which we assumed from the start, our purpose has been to restrain approximations or ad
hoc substitutions to a minimum. Some approximations (factorisations, etc.)
give interesting perspectives and will be treated in an other article [26].

The new results are mainly contained in sections 4, 6, 7. First, we have shown
that the maxent affine ansatz is compatible with asymptotic freedom, and how
the correlation decay affects the affine coefficients.

Next, the decay of the correlations, or the rate of increase of the layer popula-
tions, at large distance are tightly related to sum rules for the coefficients $\sigma_j$ in
equ. (7), specifying the screening of (topological) charges. Charge neutrality
means that a given charge is surrounded by a cloud of total opposite charge.
In electrostatics, charge neutrality results from energy bounds and shielding
the long range Coulomb field. This strong Debye screening implies subnormal
charge fluctuations [27].

In foams, as far as we can see, no electric field constrains the topological
charge. Global neutrality is a consequence of Euler equation (1). Neutrality
seems also true at an intermediate scale. The charge fluctuations result from
a strange compromise between statistical disorder and geometrical constraints
[28]. In sec. 6, the order of the asymptotic polynomial behaviour of the layer
population is related to specific sum rules satisfied by the correlation coeffi-
cients. In turn, these sum rules command the overall charge (fluctuations) in
large domains.

The conclusion is that both populations and charges are consistently related.
The next question is what do we need to get more precise, or more predictive,
estimates.

Acknowledgement

We would like to deeply thank the referees for a number of interesting, helpful
questions and comments.

References

   (1998)
   (1998); H.M. Ohlenbush, N. Rivier, T. Aste, B. Dubertret, DIMACS Series in
[22] N. Rivier, T. Aste,