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Critical Point of an Interacting Two-Dimensional Atomic Bose Gas

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We have measured the critical atom number in an array of harmonically trapped two-dimensional (2D) Bose gases of rubidium atoms at different temperatures. We found this number to be about five times higher than predicted by the semi-classical theory of Bose-Einstein condensation (BEC) in the ideal gas. This demonstrates that the conventional BEC picture is inapplicable in an interacting 2D atomic gas, in sharp contrast to the three-dimensional case. A simple heuristic model based on the Berezinskii-Kosterlitz-Thouless theory of 2D superfluidity and the local density approximation accounts well for our experimental results.

Bose-Einstein condensation (BEC) at a finite temperature is not possible in a homogeneous two-dimensional (2D) system, but an interacting Bose fluid can nevertheless become superfluid at a finite critical temperature $T_c$. This unconventional phase transition is described by the Berezinskii-Kosterlitz-Thouless (BKT) theory $\frac{1}{2}$, and does not involve any spontaneous symmetry breaking and emergence of a uniform order parameter. It is instead associated with a topological order embodied in the pairing of vortices with opposite circulations; true long-range order is destroyed by long wavelength phase fluctuations even in the superfluid state $\frac{1}{1}$.

Recent advances in producing harmonically trapped, weakly interacting (quasi-)2D atomic gases $\frac{4}{4}$, $\frac{5}{5}$, $\frac{6}{6}$, $\frac{7}{7}$, $\frac{8}{8}$, $\frac{9}{9}$, $\frac{10}{10}$, $\frac{11}{11}$, $\frac{12}{12}$, $\frac{13}{13}$, $\frac{14}{14}$ have opened the possibility for detailed studies of BKT physics in a controllable environment. There has been some theoretical debate on the nature of the superfluid transition in these systems $\frac{13}{13}$, $\frac{14}{14}$, $\frac{15}{15}$, $\frac{16}{16}$, $\frac{17}{17}$, $\frac{18}{18}$, $\frac{19}{19}$, $\frac{20}{20}$ because the harmonic confinement modifies the density of states compared to the homogeneous case. This allows for “conventional” finite temperature Bose-Einstein condensation in the ideal 2D gas $\frac{22}{22}$. Early experiments have been equally consistent with the BEC and the BKT picture of the phase transition. For example, the density profiles at very low temperatures $\frac{4}{4}$ are expected to be the same in both cases. However, recent studies of matter wave interference of independent 2D atomic clouds close to the transition have revealed both thermally activated vortices $\frac{12}{12}$, $\frac{13}{13}$, $\frac{14}{14}$ and quasi-long-range coherence properties $\frac{13}{13}$ in agreement with the BKT theory $\frac{23}{23}$.

In this Letter, we study the critical atom number in an array of 2D gases of rubidium atoms, and observe stark disagreement with the predictions of the ideal gas BEC theory. We detect the critical point by measuring (i) the onset of bimodality in the atomic density distribution and (ii) the onset of interference between independent 2D clouds. These two measurements agree with each other, and for the investigated range of temperatures $T \approx 50$–$110$ nK give critical atom numbers $N_c$ which are $\sim 5$ times higher than the ideal gas prediction for conventional Bose-Einstein condensation in our trap $\frac{20}{20}$. For comparison, in three-dimensional (3D) atomic gases, where conventional BEC occurs, the increase of the critical atom number due to repulsive interactions is typically on the order of ten percent $\frac{23}{23}$, $\frac{24}{24}$. A simple heuristic model based on the BKT theory of 2D superfluidity and the local density approximation gives good agreement with our measurements.

In $\frac{13}{13}$ we studied quasi-long-range coherence of a trapped 2D gas, which is directly related to the superfluid density $\rho_s$ $\frac{22}{22}$. In that case, signatures of the BKT transition emerge only once a significant part of the cloud becomes superfluid. Since the atomic density in the trap is not uniform, this happens slightly below the true critical temperature for the onset of superfluidity in the trap center, and the observed transition is rounded off. The present study concentrates on the exact critical point and relates to the total density at criticality $\rho_c$, which has been of long standing theoretical interest $\frac{25}{25}$, $\frac{26}{26}$.

Our experimental procedure for the preparation of cold 2D Bose gases has been described in $\frac{13}{13}$. We start with an $^{87}$Rb 3D condensate in a cylindrically symmetric magnetic trap with trapping frequencies $\omega_z = 2\pi \times 10.6$ Hz and $\omega_y = 2\pi \times 125$ Hz. To split the sample into 2D clouds we add a blue detuned one-dimensional opti-
Figure 2: Phase transition in a rubidium 2D gas. 2D clouds confined parallel to the $xy$ plane are released from an optical lattice and the density distribution is recorded by absorption imaging along $y$ after $t = 22$ ms of time of flight. The measured line densities $\tilde{\rho}(x)$ (left) and $\tilde{\rho}(y)$ (right) for an atom number just below the threshold of conventional BEC in our potential at the left panel shows the expected distribution of the 2D ideal gas and the Gaussian part of the bimodal distribution. The dashed line in the left panel shows the expected distribution of the 2D ideal gas at the threshold of conventional BEC in our potential at the same temperature ($T = 92$ nK). The dotted line in the right panel indicates the Gaussian part of the bimodal distribution.

The atomic density profiles are recorded in the $yz$ plane by resonant absorption imaging along $y$. The interaction energy $E_{\text{int}}$ at a fixed temperature. We measure the critical atom number $N_c$ corresponding to the “normal component”, and a parabolic Thomas-Fermi (TF) profile expected from superfluid hydrodynamics.

From the bimodal fits we extract the total atom number $N$ and the number of atoms within the TF part of the distribution $N_\text{c}$. The absolute detection efficiency of our imaging system was calibrated by measuring critical atom numbers for 3D BECs, taking into account interaction effects.

The absolute uncertainty of our atom number and temperature calibration is 20% and 10%, respectively.

We measure the critical atom number $N_c$ by varying the total atom number $N$ at a fixed temperature. We start with a highly degenerate sample and keep it trapped for a time $\tau$ varying between 1 and 10 s. During this time we maintain a constant temperature by applying a constant radio frequency field in the range of 10–25 kHz above the frequency corresponding to the bottom of the trap. As the hold time $\tau$ increases, $N$ gradually reduces and drops below $N_c$ due to inelastic losses.

The atomic density profiles are recorded in the $xz$ plane by resonant absorption imaging along $y$ after $t = 22$ ms of time of flight expansion. Along $z$ the profiles are Gaussian, closely corresponding to the zero point kinetic energy $h\omega_z(x = 0)/4$. Along $x$, for all $N < N_c$ a Gaussian distribution fits the data well. For $N > N_c$, the profiles exhibit a clearly bimodal shape (Fig. 2). The bimodal distributions are fitted well by the sum of a Gaussian, corresponding to the “normal component”, and a parabolic Thomas-Fermi (TF) profile expected from superfluid hydrodynamics.

The interaction energy $E_{\text{int}}/\hbar^2$ is the geometric mean of the two trapping frequencies in the plane. For comparison with our experimental results we have numerically integrated the Bose-Einstein distribution for our confining potential sketched in Fig. 1. The result depends on a ten percent level on the exact position of the lattice planes relative to the minimum of the magnetic trap potential. Since we do not fully control this position, we average over the possible configurations. We obtain the result $N_{c,\text{id}}^{\text{multi}} = p N_{c,\text{id}}$, where the effective number of planes $p$ smoothly grows from $\approx 2.2$ at 50 nK to $\approx 4.2$ at 110 nK. The resulting $N_{c,\text{id}}^{\text{multi}}(T)$ is shown in Fig. 4 as a solid line. Our measurements clearly show systematically higher $N_c$ than expected for ideal gas condensation. An empirical function...
The number of atoms in the Thomas-Fermi part of the bimodal distribution $N_0$ (●) is plotted as a function of the total atom number $N$. The solid line shows the linear fit we use to determine $N_c$, and the dashed line is its extrapolation. For comparison, the interference amplitude $|\hat{\rho}(0,k_0)|$ (●) is also displayed as a function of $N$. It shows the same threshold $N_c$ within our experimental precision. The inset shows that the temperature deduced from the Gaussian part of the fit is to a good approximation constant for all data points. Horizontal and vertical dashed lines indicate the average temperature and the critical atom number, respectively. The solid line marks the region used to determine the average temperature $T = 92(6)$ nK close to the transition. Each data point is based on 5–10 images, all error bars represent standard deviations.

$N_c = \alpha N_{\text{multi}}(T)$, with the scaling factor $\alpha$ as the only free parameter, fits the data well and gives $\alpha = 5.3(5)$, where the quoted error is statistical.

We also study the coherent fraction of the 2D gas and compare its behavior with the bimodal density profiles. We investigate the interference patterns that form after releasing the independent planar gases from the trap (Fig. 3). Fourier transforming the density profile $\rho(x,z) \rightarrow \mathcal{F}[\rho(x,z)] = \hat{\rho}(x,k_z)$ allows us to quantify the size of the coherent, i.e. interfering part of the gas as a function of $N$. The spatial frequency corresponding to the fringe period for the interference of neighboring planes is $k_0 = m/d/\hbar t$. We find that $\hat{\rho}(x,k_0)$ is well fitted by a pure Thomas-Fermi profile. Within our experimental accuracy, the radii $R_{\text{TF}}(k_0)$ of these profiles are equal to those obtained from a bimodal fit to the density. In particular, the onsets of interference and bimodality coincide (circles and diamonds in Fig. 3, respectively).

We now turn to the interpretation of our measurements in the framework of the BKT theory of 2D superfluidity. The theory predicts a universal jump of the superfluid density at the transition, from $\rho_s = 0$ to $\rho_s \lambda^2 = 4$, where $\lambda = h/\sqrt{2\pi mkT}$ is the thermal wavelength [21] (for experiments see [9, 13]). However, the total density at the critical point $\rho_c$ is not universal because it depends on the microscopic interactions. For weak interactions ($\tilde{g} < 1$), $\rho_c \lambda^2 = \ln(C/\tilde{g})$ [25], with $C = 380 \pm 3$ given by high-precision Monte Carlo calculations [26]. For our value of $\tilde{g} = 0.13$ (experimentally confirmed by measuring $R_{\text{TF}}$ as a function of $N_0$) this gives $\rho_c \lambda^2 = 8.0$.

In a harmonic trap, within the local density approximation, the transition is expected to occur when the density in the center of the cloud reaches the critical value $\rho_c$. We can heuristically relate the critical density and the corresponding critical atom number $N_{c,BKT}$ using the experimentally observed Gaussian density profiles. For a single plane with a quadratic confining potential $V(x,y)$, integrating $\rho(x,y) = \rho_c \exp(-V(x,y)/(kT))$ gives:

$$N_{c,BKT} = \rho_c \lambda^2 \left(\frac{kT}{\hbar \omega}\right)^2 = \rho_c \lambda^2 \frac{6}{\pi^2} N_{c,\text{id}}.$$  \hspace{1cm} (2)

For $\rho_c \lambda^2 = 8.0$ this gives $N_{c,BKT}/N_{c,\text{id}} = 4.9$. In a lattice configuration this ratio changes only slightly. We set the
peak density in the most populated plane to $\rho_c$, and sum the contributions of all planes $j$ using the corresponding potentials $V_j$ (here we neglect the small non-harmonic effects due to finite laser waists). The total population in the lattice is then $N_{\text{multi}}^0 = p' N_{c,\text{BKT}}$, where the effective number of planes $p'$ varies from 2.4 at 50 nK to 3.5 at 110 nK. We thus obtain $N_{c,\text{BKT}}/N_{c,\text{id}} \approx 4.7$, which is close to the experimental ratio $\alpha = 5.3$.

One could try to reproduce our observations within the self-consistent Hartree-Fock (HF) theory \cite{32, 33, 34}, by replacing $V(\mathbf{r})$ with the effective mean field potential $V(\mathbf{r}) + 2g \rho(\mathbf{r})$ and again setting the peak density to the BKT threshold $\rho_c$. For very weak interactions, $\log(1/g) \gg 1$, analytical HF calculation gives critical numbers which are only slightly larger than $N_{c,\text{id}}$ \cite{35}. This approach could in principle be implemented numerically for our value of $\tilde{g}$ and our lattice geometry. However, it has been suggested \cite{35} that treating interactions at the mean field level is insufficient for $\tilde{g} \sim 10^{-1}$, because the interactions are strong enough for the critical region to be a significant fraction of the sample. In future experiments with atomic gases $\tilde{g}$ could be varied between 1 and $10^{-2}$ using Feshbach resonances, allowing for detailed tests of the microscopic BKT theory and the possible breakdown of the mean field approximation.

In conclusion, we have shown that the ideal gas theory of Bose-Einstein condensation, which is extremely successful in 3D, cannot be used to predict the critical point in interacting 2D atomic gases, where interactions play a profound role even in the normal state. A much better prediction of the critical point is provided by the BKT theory of 2D superfluidity. We have also shown that, despite the absence of true long-range order, the low temperature state displays density profiles and local coherence largely analogous to 3D BECs.

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\begin{thebibliography}{35}
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\bibitem{31} Note that the slope $dN_0/dN \sim 0.7$ for $N > N_c$ is less than unity, contrary to what is expected for a saturated gas. This lack of saturation will be discussed elsewhere.
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