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EXPERIMENTAL STUDY ON THE ELECTRIC FIELD OF DISLOCATIONS IN COLOURED ALKALI HALIDE CRYSTALS

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Abstract. — Previously the «dislocation photoconduction spectrum» in the F band was measured by us in X-ray coloured KCl crystals. The spectrum was interpreted by assuming that the internal electric field of moving dislocations perturbed the F-centres producing a quadratic Stark-effect in them. Comparing our results with Chiarotti, Grassano et al., we estimated the internal electric field in our crystal. This work was carried out on several other alkali halides as well.

In ionic crystals edge dislocations have an excess electric charge which is neutralized by the field of ions surrounding the dislocation. If these crystals are deformed in a suitable way an electric signal can be observed due to the movement of dislocations. According to some measurements, at deformation potentials of several thousand volts may appear on the surface of the crystals [1].

In the last few years we have been investigating the interaction of edge dislocations and point defects in X-irradiated alkali halide crystals approaching the problem from the side of dislocations. From our experimental results we obtained information on the magnitude of the internal dislocation electric field itself.

The dislocation photoconduction spectrum. — Our X-ray coloured, usually nominally pure alkali halide crystals were deformed by four-point bending and the exponential decay of the developed voltage upon illuminating the sample with light of constant intensity in the F band was measured [2]. According to our experimental results the quantity a characterizing the decay varies fairly linearly with the intensity of incident light and it is not in simple correlation with the amount of deformation and the F-center concentration. Plotting a/2 against the wavelength (\lambda) of the illuminating light we obtained the dislocation photoconduction spectrum (DPS) [3]. In principle this should follow the shape of the F optical absorption band and its shape should be determined by the electric field of dislocations. But as it can be seen in figure 1 there is a marked difference between the DPS and the F optical absorption band [3-4].

The main differences:

1) The DPS is broadened compared to the F optical absorption band.
2) The DPS exhibits local maxima.

Fig. 1. — Dislocation photoconduction spectrum of a NaCl crystal (curve — - - -) and its F optical absorption band (curve — — —).

In the present we are dealing mainly with the first feature of the DPS, for the second one we had given an explanation elsewhere [5].

For the interpretation of our spectra one had to assume that they were due to electrons released by the incident light from F-centers in the vicinity of dislocations, i.e. interacting with the electric field of dislocations. The fact that our spectra are essentially independent from the bulk concentration of F-centers in the range of \(10^{12}-10^{17}\) cm\(^{-3}\) supports this idea too. Thus our method enables us to detect solely those F-centers which are interacting with dislocations.

Stark-effect of the F-center. — It is well known that a strong electric field causes a quadratic Stark-effect in the F-center. According to the theoretical considerations of Henry, Schnatterly and Slichter the effect can be easily analyzed by the method of moments of the absorption line shape [6].
The absorption band is characterized by the line shape function \( f(E) \) (\( E \) is the photon energy), which is related to the absorption coefficient \( \alpha \) by

\[
\alpha(E) = CE f(E),
\]

\( C \) is a constant.

The application of an electric field changes the line shape function to \( g(E) \). According to the theory the field leaves the zeroth and first moments unchanged

\[
A = \int f(E) \, dE = \int g(E) \, dE, \Delta A = 0
\]

\[
\langle \Delta E \rangle = 0
\]

and causes a change of the second moment \( \langle \Delta E^2 \rangle \) which for light polarized in the field direction is proportional to the square of the local field strength \( F \).

The change of the absorption coefficient at the peak of the F band is related to the change of the second moment by the approximative relation:

\[
\frac{\Delta \alpha}{\alpha} \approx -\frac{\langle \Delta E^2 \rangle}{\langle 2E^2 \rangle}.
\]

Therefore

\[
\left| \frac{\Delta \alpha}{\alpha} \right| \propto F^2 \text{too}.
\]

The experimental results of Chiarotti, Grassano et al. [7] and those of Rhyner and Cameron [8] supported the HSS theory and plotting \( \Delta \alpha/\alpha \) against

![Fig. 2. Difference curves \( g(E) - f'(E) \) for a) NaCl; b) KCl; c) KBr; d) KI crystals.](image-url)
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the square of the field strength they obtained the expected straight line.

Results and discussion. — We assumed, that in our case we were dealing with the Stark-effect of the F-center caused by the dislocation electric field, and as an approximation the HSS theory is applicable to our measurements too. According to the theory the zeroth and first moments of the F band perturbed by the dislocation electric field remain unchanged, and based on this the \( g(E) - f(E) \) difference curves were constructed, where \( f(E) \) is the line shape function of the unperturbed F band and \( g(E) \) is the line shape function of the DPS. Figures 2a, b, c, d show these for NaCl, KCl, KBr and KI crystals.

From the difference curves and the F band the value of \( \Delta \varepsilon/\varepsilon \) in the peak of the F band was determined. Than extrapolating the experimental results obtained for \( \Delta \varepsilon/\varepsilon(F^2) \) with an external electric field [7] we obtained the following local electric field values for our crystals.

\[
\begin{array}{|c|c|c|}
\hline
\text{Crystal} & \Delta \varepsilon/\varepsilon (*) & [\text{V/cm}] \\
\hline
\text{NaCl} & 0.367 & 1.05 \times 10^7 \\
\text{KCl} & 0.381 & 5.0 \times 10^6 \\
\text{KBr} & 0.359 & 3.5 \times 10^6 \\
\text{KI} & 0.242 & 5.2 \times 10^6 \\
\hline
\end{array}
\]

(*) At the peak of the F band.

The values given in the table were obtained from measurements carried out on one sample in case of all crystals. The scattering of the internal electric field values obtained on the same sample at different measurement was within 8 % as we could ascertain for the case of KCl. Measurements on different samples of the same alkali halide gave field strength values varying only within 20 %. Extremely pure OH-free crystals and crystals coloured additively gave also essentially the same results (Fig. 3, 4, 5).

Our results should be considered approximative ones, due to the following:

1) The field in the vicinity of dislocations is inhomogeneous, we determined only a resultant average value.

2) The condition of illumination with light of constant intensity is not satisfied through the whole
thickness of the sample. This error is diminished probably by the light scattering within the sample and reflection from internal surfaces.

3) The local maxima of the DPS explainable by lattice distortion around dislocations, appear simultaneously with the electric field effect making the results more complex and the interpretation more difficult.

4) Our measurements were carried out at room temperature while those carried out with an external electric field at low temperature. This, due to the variation of $<E^2>$ with temperature should modify our values obtained for $F$ by a factor of 2.

5) One expects that impurities should play an important role in the DPS, this problem is now under investigation.

References