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A novel “*CO*rrrelated Ion and Neutral Time OfFlight” Method: event-by-event detection of neutral and charged fragments in Collision Induced Dissociation (CID) of mass selected ions

C. Teyssier¹, R. Fillol¹, H. Abdoul-Carime¹, B. Farizon¹, M. Farizon¹ and T.D. Märk²

Abstract

A new mass spectrometric MS-MS method based on Time Of Flight measurements performed on an event-by-event detection technique is presented. This “*CO*rrrelated Ion and Neutral Time of Flight (*COINTOF*)” method allows to explore CID fragmentation processes by directly identifying not only all ions and neutral fragments produced but also their arrival time correlations within each single fragmentation event from a dissociating molecular ion. This constitutes a new step in the characterization of molecular ions. The method will be illustrated here for a prototypical case involving Collision Induced Dissociation (CID) of protonated water clusters $\text{H}^+(\text{H}_2\text{O})_{n=1-5}$ upon collisions with argon atoms.

¹*Université de Lyon, F-69003, Lyon, France; Université Lyon 1, Villeurbanne; CNRS/IN2P3, UMR5822, Institut de Physique Nucléaire de Lyon; F-69622, Villeurbanne.*

²*Institut für Ionenphysik und Angewandte Physik, Leopold Franzens Universität Innsbruck, Technikerstrasse 25, A-6020 Innsbruck, Austria*

I. Introduction

Mass spectrometry is nowadays a standard technique used in many fundamental and applied areas of science. As an analytical technique, it is used in diverse fields such as for the detection of drugs¹ in forensic applications, for the controlling of processes² in food production, or for the analysis during microchips manufacturing processes³ in the field of nanotechnology. One analytical method in mass spectrometry involves Collision Induced Dissociation (CID) where a primary molecular ion dissociates into secondary fragments after collision with a gas target. Indeed, the CID technique is also a powerful tool in physics and chemistry to shed light on the nature of ion collision reactions present in many fundamental and technical phenomena (e.g., interstellar media, planetary ionospheres, plasmas). Moreover, it is widely applied in analytical chemistry to elucidate the identity and the nature of molecular and cluster ions via MS/MS spectrometry⁴. Although so far limited to the detection and identification of the produced ion fragments in recent studies it has been extended to the analysis of ions formed in single events⁵ or to the detection of neutrals species⁶. Nevertheless directly monitoring all ionized and neutral fragments produced and their correlations within a single fragmentation event – as will be described in the following- is a major step towards the complete description of the physics and chemistry of the reaction.

Here a new type of mass spectrometry technique, named COINTOF (CorrelatedIon and Neutral Time Offlights) mass spectrometry developed in our laboratory⁷ is presented. This method allows the exploration of fragmentation processes by recording all ionized and neutral products in a single dissociation event of an individual precursor ion. Detailed information on the physics of the dissociation process is obtained by a statistical analysis of an ensemble of

recorded events. This new method is illustrated in the following for the case of CID of protonated water cluster ions.

II. Principle of the COINTOF method

The principle of the method is presented in Figure 1. As the COINTOF technique is a part of an MS-MS technique, a prerequisite to COINTOF is the production of a mass and energy selected beam of parent (precursor) ions. This production of an intense mass and energy selected ion beam has been thoroughly described elsewhere⁸. Briefly, protonated water clusters are produced by electron impact in a supersonic nozzle expansion of water vapor. They are accelerated to an energy between 2 to 8 keV and mass selected by a double focusing sector field mass spectrometer⁹. For the present CID experiment, the energy and mass-selected $H^+(H_2O)_n$ beam crosses an effusive argon gas jet whose density is controlled in such a way as to insure single collision conditions in the CID region. It is noteworthy however, that the precursor ions can collide with other type of particles (protons, ions, electrons or photons) at the collision point. The newly designed time of flight COINTOF apparatus is mounted right after the collision region along the direction of the incident beam as shown in Fig. 1. A single dissociative collision of a $H^+(H_2O)_n$ parent ion with an Ar atom will in general produce a daughter ion and one or more neutral fragments. This daughter ion is then extracted from the collision area and accelerated through 1.8kV before reaching a field free time of flight TOF zone. Neutral fragments emerging from the collision zone however will not be affected by this acceleration field and remain after the collision close to the initial velocity of the parent ion (except for small changes due to additional kinetic energy released in the collision process itself, see below). All the fragments (ionized and neutral) are then detected by the one Micro Channel Plate detector (MCP, Hamamatsu) with the charged fragment arriving significantly

earlier than the neutral species. The MCP set-up is designed to ensure the collection/detection of all the fragments by this detector. The signals produced are sampled by a fast Analog to Digital Converter (ADC8 GHW/10bits from Acqiris®, Agilent) and it is recorded without dead time during a given arrival time window. The intensity of the initial parent ion beam is controlled in such a way as to ensure that only the fragments of one single parent ion can reach the detector during the acquisition time window. Therefore, the correlation between the arrival times of the charged and neutral fragments resulting from the dissociation of one single parent ion can be recorded. This correlation is determined consecutively for a large number of dissociation events.

III. Description of the COINTOF detector

The COINTOF detector is based on the Wiley-McLaren type¹⁰ time of flight principle. It consists of three separate zones with different values of electric fields (see Figure 1). The first zone is the extraction zone and the second the acceleration zone. The collision induced dissociation takes place at the center of the extraction zone that is the crossing point between the cluster ion beam and the neutral target beam. In the present CID experiments argon or helium is used as a neutral target gas and the neutral gas jet has a maximum extent of 2mm at the crossing point¹¹. Since all fragments (neutral and ionized) produced after dissociation have a velocity close to the velocity of the parent ion, the use of an extraction voltage is *a priori* not a mandatory requirement. Nevertheless the charged products are extracted by an electric field applied and after reaching the acceleration zone accelerated further by an acceleration field produced by a series of electrodes which provide a homogenous electric field in this zone.

Therefore any charged species will be accelerated by the same electric field and thus acquire on top of the initial velocity an additional velocity characteristic of its mass/charge ratio.

After this acceleration zone, all species travel through a field free time of flight zone before reaching the MCP detector. This TOF zone consists of a drift tube of 150 mm length held at the acceleration potential. Detection is carried out with a MCP detector operated with a difference voltage of 1700 V and positioned 40 mm behind this field free zone. In the present case we used a pair of MicroChannel Plates (MCP-F4293-07-Hamamatsu) with an effective area of 20 mm in diameter and with channel diameters of 12 μm mounted in a *Chevron* like fashion. The MCP signal is collected on a single anode. It should be noted that this anode can be replaced by wire anodes, to determine the exact position of impact of the collision products on the detector.

One of the key points in the COINTOF technique is that the parent ions, in the present case protonated water cluster ions, have an energy prior to collision of 2 keV up to 8 keV and this high energy ensures that after fragmentation both ionized and neutral products are detectable on the detector because both will lead upon collision at the entrance plate of the MCP to the emission of secondary electrons which are then amplified by the channel plates. It must be underlined though that neutral products are not accelerated after the interaction zone while the charged species receive an additional energy due to acceleration in the electric field applied. The velocity of the neutral products is close to the velocity of the parent ion. Nevertheless the neutral products from the parent ion reach the detector with sufficient energy (at least 2 keV) to be detected and they are well separated in their arrival time from the ionized fragments which are arriving much earlier due to the additional energy drawn from the accelerating field. It is clear that the present technique is limited to studies where the

initial collision energy is at least 2 keV, because below this energy the efficiency for emission of secondary electrons at the entrance plate of the MCP is too small.

IV. Simulation of the COINTOF geometry

The necessary geometry of the present time-of-flight set-up and the corresponding trajectories of the particles has been calculated using the charged particle optics simulation software, SIMION®¹². Figure 2 shows the present COINTOF detector simulated by SIMION® showing the equipotential lines (in blue) in the three different zones and voltages applied. As mentioned above the produced ions and neutrals will have after the collision the velocity of the parent ion that is they will have a kinetic energy in proportion of their mass to the mass of the parent ion. These fragment kinetic energies will be modified in case of ionized fragments by the accelerating field and in addition by the kinetic energy released (KER) during the fragmentation process. This KER may amount to up to several 100 meV¹³ (i.e., typical H bond energies). This energy is however negligible to a first approximation in comparison to the kinetic energy of the fragments (fraction of keV) and the energy gained by the accelerating field (1.8 keV). Nevertheless, this additional KER will lead to an additional broadening of the measured distribution of the time difference between the arrival times of the ionized and the neutral products. On the other hand, the mean time difference between ionized fragment and the neutral fragments is directly related to the mass of the ionized (and accelerated) fragment.

SIMION® simulations were carried out for various ion masses and a range of energies. These simulations provide the arrival times on the MCP detector for each ion and corresponding neutral fragment separately. Table 1 presents results for fragments from different parent ions all at

the same parent velocity of $130 \times 10^3 \text{ m/s}$. As mentioned above, the time of flight for the neutral fragments is independent of their mass since the applied electric fields have no influence on their flight times, in contrast to that of the ions. The differences in times between charged and neutral (last column) depend on the mass of the charged fragment and thus allow the identification of the charged fragment for each dissociation event.

V. The acquisition set up

The output signal generated by the impact of neutral and ionized fragments on the MCP detector is connected to an acquisition device via preamplifiers (FAMPS8, HandelsRoentdek GmbH). During the measurements carried out to validate the present method, two acquisition chains are installed in parallel as shown in Figure 3:

- The "TDC multihit" acquisition line consists of a constant fraction discriminator, a TDC (Time Digital Converter) module and a fast bus (ECL) for transferring the time values in a memory. The unit is managed by standard CAMAC. The TDC module allows multi-hits to record with a time resolution of 0.5 ns, however with a dead time of 15 ns between successive pulses. It is programmed in the "common stop" mode. In this mode the time of arrival of the various fragments produced by the dissociation of a given precursor is saved and labeled as the arrival times corresponding to the dissociation of the corresponding precursor ion. This acquisition line provides a direct COIN TOF mass spectrum of the charged fragments obtained in continuous operation without an external trigger (CID-recorded COIN TOF spectrum). Figure 3 shows such a CID mass spectrum for the case of 8 keV protonated water tetramer

parent ions interacting with argon target gas. It is noteworthy that the COINTOF mass spectrum differs from that obtained traditionally, where the charged species are plotted as a function of the mass to charge ratio. In the COINTOF mass spectrum, the ions are plotted as a function of the difference between the arrival instant of the neutral and the ion, ΔT . Since the neutral species remain at the velocity of the precursor (to a first approximation), they can be used as an intrinsic time reference within a single dissociation event. Thus the lighter the ion species, the larger is ΔT , as shown in fig.3 (lower right side) where the four mass peaks refer to $H^+(H_2O)_3$, $H^+(H_2O)_2$, $H^+(H_2O)$ and H^+ from left to right respectively.

- The "ADC-signal SID" acquisition line uses a 10-bit analog to digital converter, coupled to a memory 8GHz Quick Access 1.2 MB (Acqiris map). The PCI card is connected to a standard PC computer. Recording of the detected signal is performed on an event by event basis. Within a single event, in contrast to the TDC line, there is no dead time. For each dissociation the signal is recorded individually as shown in Fig.3. By summing up these individual signals this acquisition line can also produce a COINTOF mass spectrum of charged fragments as in case of using the TDC-multihit acquisition line but in this mode more information can be deduced from the data. By choosing an appropriate time window, it is possible to record only specific dissociation events.

From these data we can immediately deduce the number of products for a given precursor ion. The precursor ions arrive one after the other in the interaction zone at a rate of typically one ion per 1ms (i.e., 1000 cps). On the other hand, the acquisition time window is about two to four microseconds. Therefore on average there is no overlap between successive dissociation events and the number of peaks can be directly related to the reaction under

study. In the example shown in Figure 3, three products are detected for the tetramer parent ion, i.e., a charged product and two neutral fragments. Indeed, the charged product from the precursor, which has been accelerated, hits the detector before the two neutral products from the same precursor. Moreover, the acquisition set up includes a series of signal triggering modules thus allowing us to select a specific (adjustable) acquisition time after the first (neutral) product of interaction hits the detector. Moreover the duration of the acquisition window is adjustable. This allows making specific measurements, for example, using acquisition strategies reducing the acquisition time.

The present acquisition set-up also allows the recording of the amplitude of the detected signal as shown in Figure 4 for an example utilizing the ADC converter line. The amplitude of the detected signal is digitized into 10 bits at a rate of 8 gigahertz (125ps). For each event, the shape of the detection pulse is also recorded. The acquisition device is organised in such a way as to label the set of numbers obtained for a given precursor as an event associated with the interaction of this one precursor with the target. This labelling step allows to distinguish information acquired (arrival time and magnitude of the peaks) between different precursors, thereby preserving the correlation between information obtained for each event. It should be noted that without an external time reference, this approach enables us to work with a continuous beam of precursor ions of strong intensity and to have a high rate of repetition. Typically, we carry out experiments with MCP intensities of 500 counts/s. The acquisition time window is about a couple of μ s thus allowing us to record event after event.

VI. CID of protonated water cluster ions explored by the COINTOF Method: a test case

The COINTOF method can be applied to a wide range of molecular ions. As a test case, we present here results obtained for collision induced dissociation of protonated water cluster $(\text{H}_2\text{O})_{n=2-5}\text{H}^+$ with an argon target beam. All parent cluster ions studied here are accelerated to energies corresponding to the velocity of 130.10^3 m/s and mass selected by an ExB sector field mass spectrometer. Collision induced dissociation of these cluster ions at this velocity leads to the production of $(\text{H}_2\text{O})_{n-1}\text{H}^+, \dots, (\text{H}_2\text{O})\text{H}^+$ and H^+ charged fragments which are usually accompanied by the loss of one or several neutral water molecules from the precursor ion.

Two prototypical examples of recorded single dissociation event signals are given in Fig. 5 for CID of the protonated dimer. Each arrival time window is associated with the dissociation of one single $\text{H}^+(\text{H}_2\text{O})_2$. Due to the acceleration zone after the collision, the first peak in each time window is attributed to the charged fragment. The arrival time for the neutral fragments is always larger than any of that of the charged fragments. Moreover, the difference ΔT between the arrival time of the charged fragment(s) and the neutral one(s) is directly related to the mass to charge ratio(s) of the ionized fragments involved.

In Figure 5A, ΔT is characteristic of the $\text{H}^+(\text{H}_2\text{O})$ ion and by comparison with the ion trajectory simulations the underlying dissociation event corresponds to



In contrast, three peaks are observed in the time arrival window shown in Figure 5B, corresponding to one charged fragment and two neutral fragments. The difference in arrival times between the charged and neutral fragments is characteristic of the H^+ ion. Thus this arrival time spectrum indicates the dissociation of the protonated dimer in three fragments, the proton and two neutral water molecules via reaction



Note that the arrival time difference between the two neutral fragments (caused by different KER values of the two fragments) may be smaller than the typical dead time of discriminators used in standard arrival time measurements with a time to digital converter. Nevertheless, in the present case it is accessible by the use of a fast ADC.

The example presented in Figure 6 involving CID of the protonated water pentamer illustrates various steps in a particular analysis. In the left upper insert we show the spectrum of the charged fragments for CID of 8 keV $(\text{H}_2\text{O})_5\text{H}^+$ with argon as target gas. Each ΔT (ns) peak is associated to a given mass to charge ratio(s) for a specific charged fragment produced in the collision event (see Table 1), corresponding to the detection of $\text{H}^+(\text{H}_2\text{O})$ (19 amu - $\Delta T \approx 490$ ns), $\text{H}^+(\text{H}_2\text{O})_2$ (37 amu - $\Delta T \approx 312$ ns), $\text{H}^+(\text{H}_2\text{O})_3$ (55 amu - $\Delta T \approx 230$ ns) and $\text{H}^+(\text{H}_2\text{O})_4$ (73 amu - $\Delta T \approx 182$ ns) fragments. In the lower insert right hand side, a single event for the dissociation channel of $(\text{H}_2\text{O})_5\text{H}^+$ into $(\text{H}_2\text{O})_3\text{H}^+ + 2(\text{H}_2\text{O})$ is reported; three peaks are observed, the first one corresponding to the ionized fragment and the two close peaks at higher arrival times corresponding to the two neutral water units. Figure 6 shows also the most stable protonated water pentamer structure from theoretical calculations of Miyazaki et al.¹⁴. After collision the excess energy can lead to various dissociation channels with two neutral subunits that can arise from any site of the $\text{H}^+(\text{H}_2\text{O})$ central ion. Nevertheless, the measurement of the time difference between the two neutral water subunits (≈ 31.3 ns) will be a key information regarding the structure of the decaying parent ion.

In addition CID experiments have been performed here with protonated water cluster ions $\text{H}^+(\text{H}_2\text{O})_n$ for sizes n from 1 up to 5 all at a velocity of $130 \cdot 10^3$ m/s. In Figure 7 we show the mass spectra of the respective charged fragments obtained by the COINTOF method. The

arrival time of the first neutral fragment is used as a time reference to identify the charged fragments for each dissociating precursor ion $\text{H}^+(\text{H}_2\text{O})_n$ using the results obtained with the SIMION simulation (Table 1). The red asterisks in Figure 7 correspond to the values obtained with the SIMION calculations. On the right hand side of Figure 7 the various dissociation channels present in the mass spectra are listed for each incident precursor ion $\text{H}^+(\text{H}_2\text{O})_n$. As expected the charged fragments observed correspond to the loss of one up to n H_2O molecules. Nevertheless differences are observed concerning the abundance of the various $\text{H}^+(\text{H}_2\text{O})_p$ productions formed by these collisions. The most surprising result is the relative large abundance of the H_3O^+ fragment ion among the fragment ions produced in the protonated water pentamer case. This may be understood however by the important role of the H_3O^+ ion (Eigen ion) as the central core of protonated water clusters as discussed in numerous papers and shown experimentally by Headricks et al.¹⁵ in infrared spectra of $\text{H}^+(\text{H}_2\text{O})_n$ clusters for specific small cluster sizes. The present results on the water pentamer ions suggest that the parent ion structure seems in accordance with a H_3O^+ ion (Eigen) core surrounded by H_2O molecules as calculated by Miyazaki et al.¹⁴ (see Figure 6).

VII. Conclusion

In conclusion, the novel CID mass spectrometry technique, COINTOF, presented here provides compared to earlier studies a more detailed description of a collision induced dissociation process by directly probing not only all ionized and neutral fragments produced but also their correlations in time within each single dissociative event. This COINTOF method is implemented on DIAM¹, a recently designed experimental setup at University Lyon 1 to investigate processes resulting from the irradiation of various molecular nanosystems. This new MS/MS spectrometer detects neutral and charged fragments arising from the

dissociation of a single precursor ion. The set up does not need any prior external time reference and can be operated in a continuous mode. The "event by event" recording allows to analyze CID reactions beyond the sole average production rate of charged fragments. This technique allows the correlated measurements of the time of flights of charged and neutral fragments produced within the same dissociation event of a mass and energy pre-selected molecular parent ion. The power of this COINTOF technique is illustrated here showing examples from collision induced dissociation of protonated water cluster ions with atomic gas targets (e.g. argon).

FIGURES

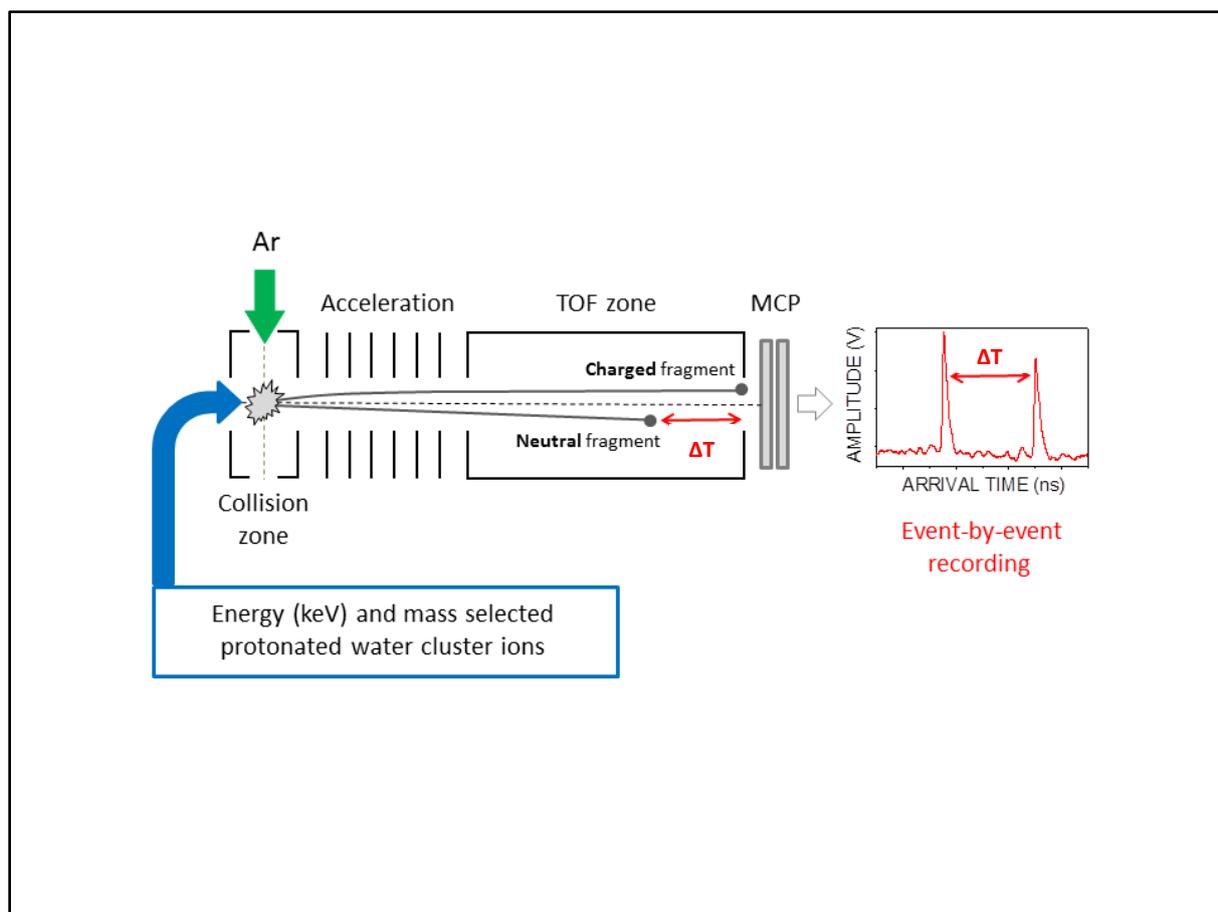


Figure 1: Principle of the COINTOF method. The mass and energy selected precursor ions collide, here, with Ar atoms. The ionized fragments are extracted from the collision zone and accelerated before reaching a field free zone. The neutral species are not influenced by these electric fields. All the fragments resulting from the dissociation of a single precursor ion hit the detector, and the associated signals are recorded.

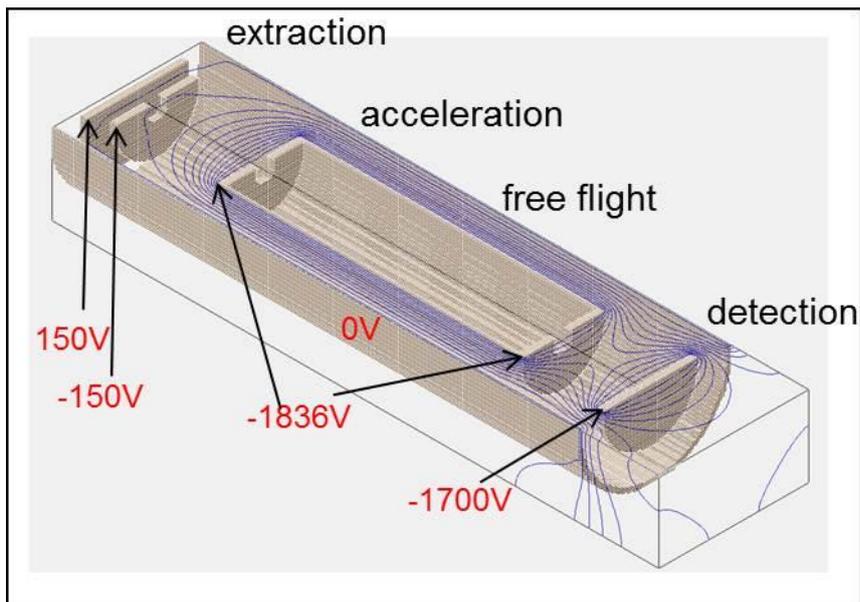


Figure2:

Simulation of the COINTOF geometry by SIMION[®] software. Three zones corresponding to the extraction, the acceleration and the free flight part are represented with their typical applied voltages. The equipotential lines are drawn in blue.

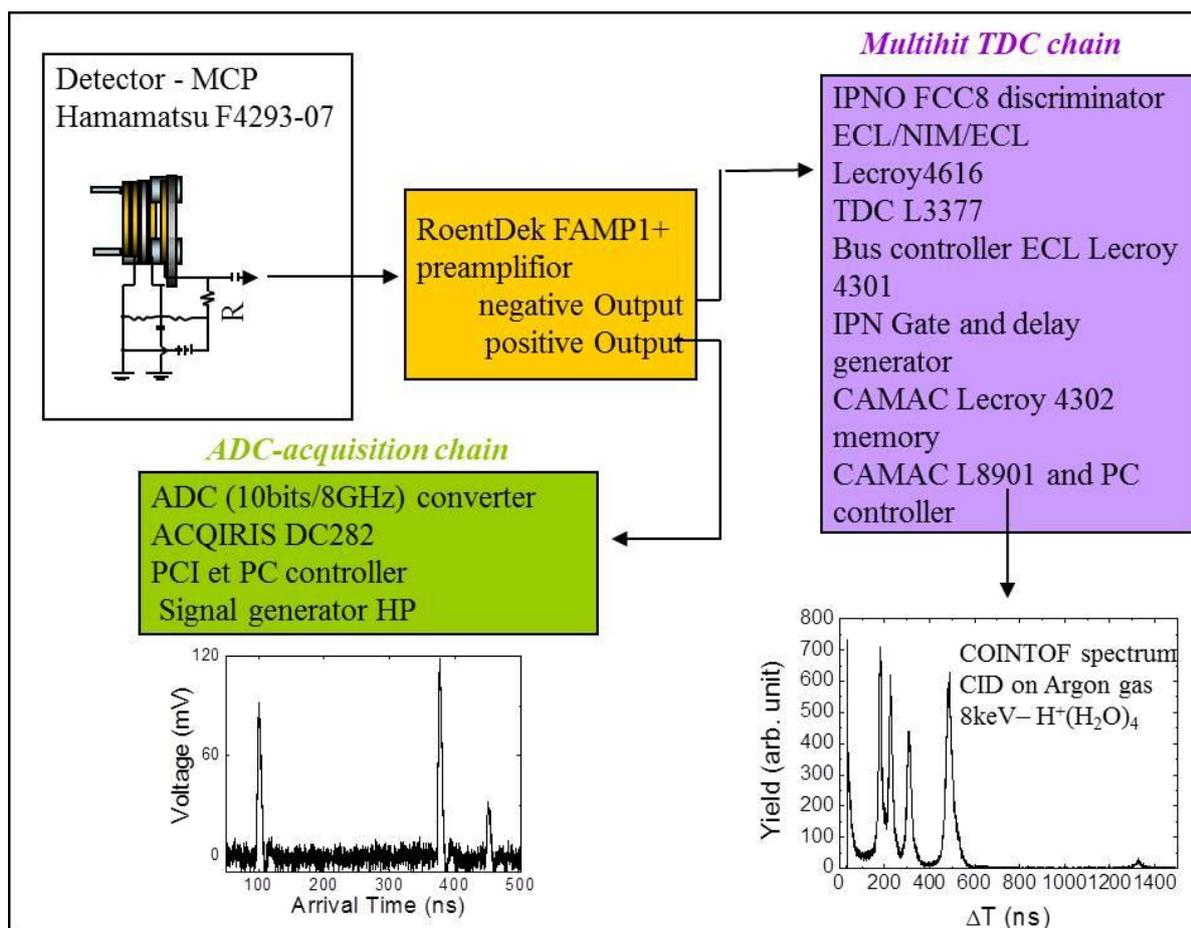


Figure 3: Principle of the acquisition set-ups used. Two acquisition lines have been used in parallel. The TDC-multihit chain provides the COINTOF mass spectra (right hand side). The ADC (10bits/8GHz) acquisition chain is used to record an arrival time spectrum for a single dissociation event in a specific dissociation channel (left hand side).

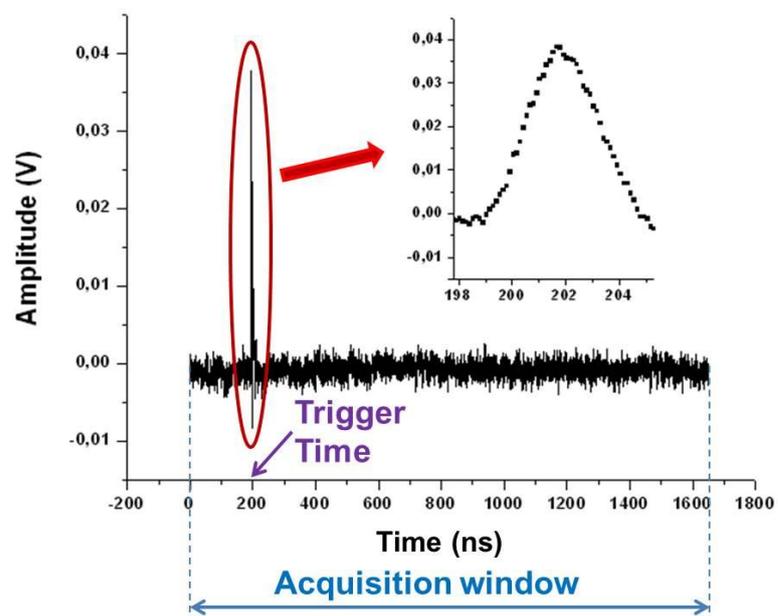


Figure 4: Signal produced from the impact of one particle on the detector by the fast ADC (8GHz/10bits) acquisition line.

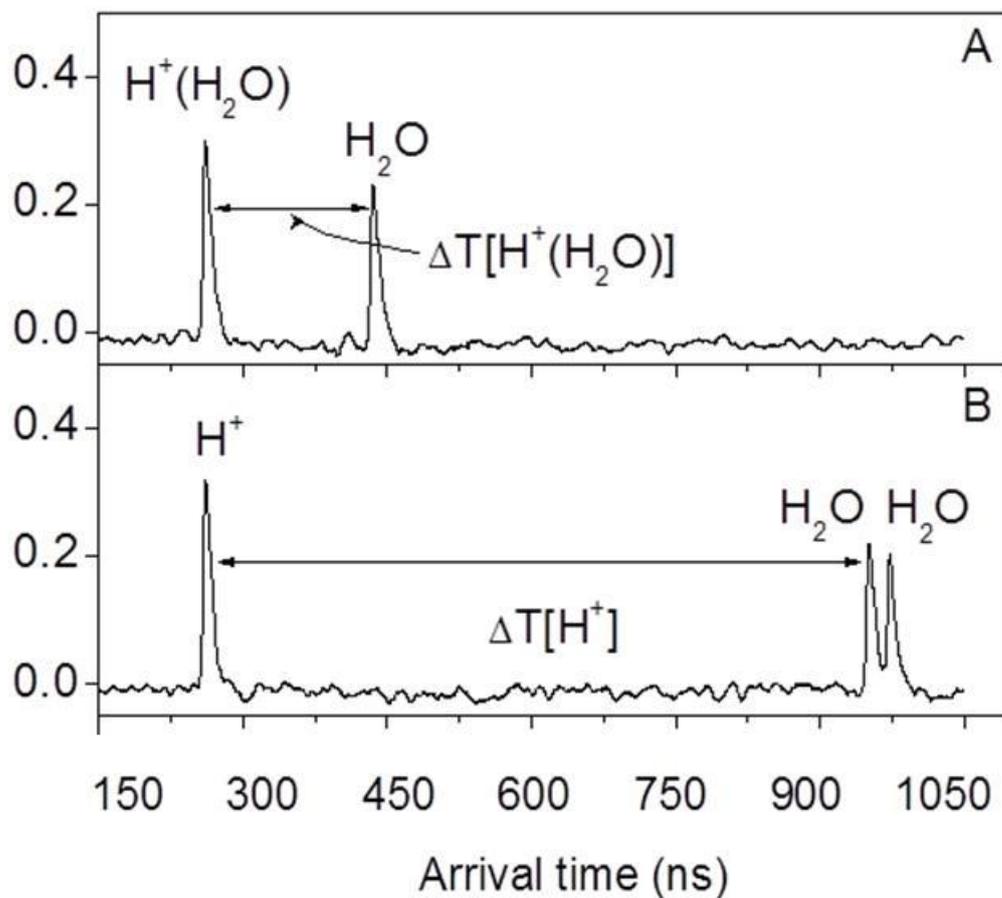


Figure 5: COINTOF arrival time spectra for 8keV $H^+(H_2O)_2$ ions interacting with an Ar target beam:

(A) Detection of a $H^+(H_2O)$ fragment ion and one neutral fragment

(B) Detection of a H^+ fragment ion and two neutral fragments

The time differences between the ion and the neutral(s) are in accordance with those obtained from ion trajectory simulation.

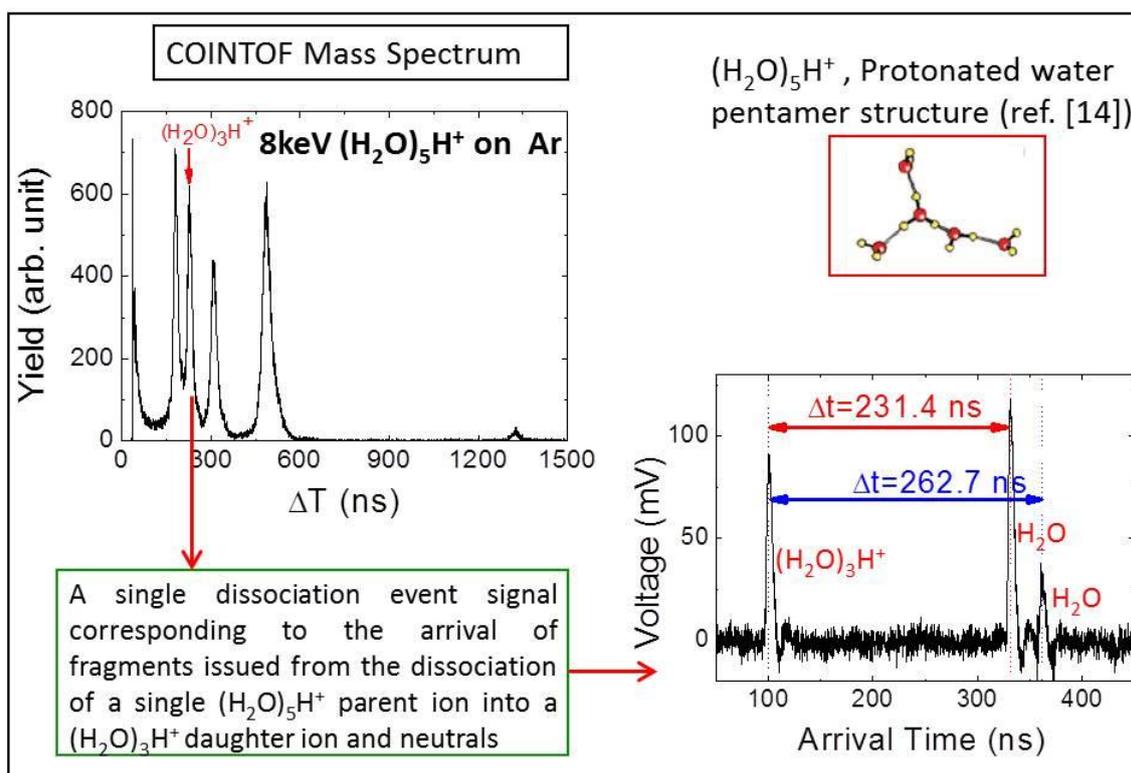


Figure 6: CID of the protonated water pentamer cluster ion

Left upper side: COINTOF mass spectrum of the charged fragments in the Collision Induced Dissociation of 8keV $(\text{H}_2\text{O})_5\text{H}^+$ ions with an argon gas beam

Right upper side: The most stable protonated water pentamer structure from theoretical calculations of Miyazaki et al.¹⁴.

Right lower side: A recorded single event signal for the dissociation channel $(\text{H}_2\text{O})_5\text{H}^+$ into $(\text{H}_2\text{O})_3\text{H}^+ + 2(\text{H}_2\text{O})$: three peaks are observed corresponding to the three fragments.

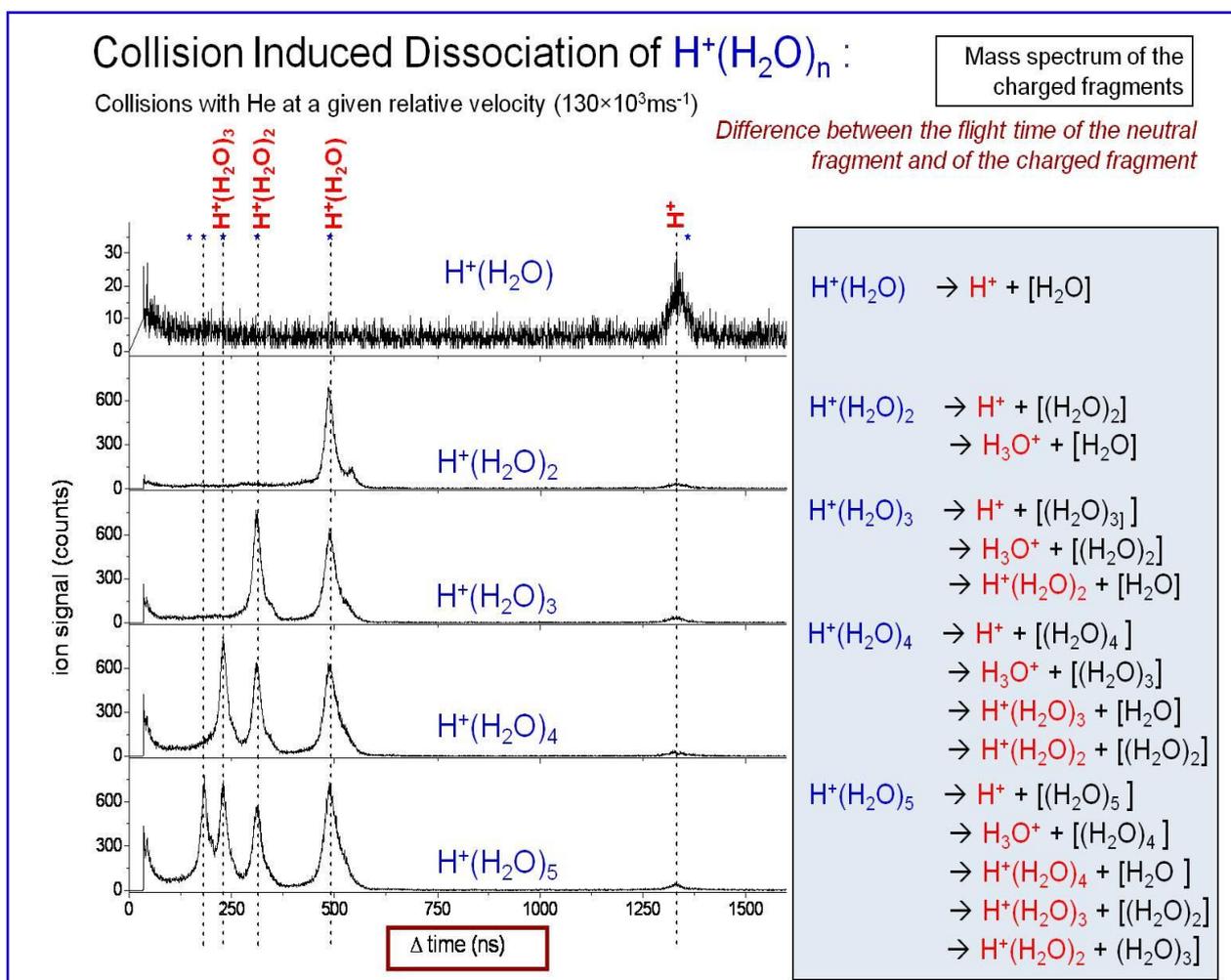


Figure 7: COINTOF mass spectra of the charged fragments for $H^+(H_2O)_n$ parent ions with n from 1 to 5. The arrival time of the first neutral fragment is used as a time reference to identify the charged fragments for each of the dissociating precursor ions $H^+(H_2O)_n$. The red asterisks correspond to the predictions from SIMION calculations (Table 1). All the parent ions are initially accelerated to the same velocity of $1.3 \cdot 10^7$ cm/s ($130 \cdot 10^3$ m/s). The various dissociation channels observed are given on the right hand side.

TABLES

Parent mass of protonated water cluster ions (amu)	Parent energy (eV)	Fragment charged mass (amu)	Charged fragment KE (eV)	Charged fragment arrival time (μs)	Neutral fragment arrival time (μs)	Arrival time difference ΔT (μs)
19	1670.33	1	87.91210	0.522104	1.90407	1.381966
37	3252.747	19	1670.329	1.41434	1.90407	0.48973
		1	87.91208	0.522104		1.381966
55	4835.165	37	3252.747	1.59219	1.90407	0.31188
		19	1670.329	1.41434		0.48973
		1	87.91209	0.522104		1.381966
73	6417.582	55	4835.164	1.67453	1.90407	0.22954
		37	3252.747	1.59219		0.31188
		19	1670.329	1.41434		0.48973
		1	87.91208	0.522104		1.381966
91	8000	73	6417.582	1.72225	1.90407	0.18182
		55	4835.164	1.67453		0.22954
		37	3252.747	1.59219		0.31188
		19	1670.329	1.41434		0.48973
		1	87.91208	0.522104		1.381966

Table 1: Arrival time difference between the charged fragment ion and a neutral fragment for different initial kinetic energies (amounting to the same initial velocity of 130.10^3 m/s for all parent ions studied).

REFERENCES

- ¹ S. Balabanova, J. Homoki, Z. Rechtsmed 98, 235 (1987)
- ² R. Flamini, A. Panighel, Mass. Spectrom. Rev. 25, 741 (2006)
- ³ T. Sikanen et al., Mass Spectrom. Rev. 29, 351 (2010)
- ⁴ B.T.Chait, Science, 314, 65 (2006)
- ⁵ F. Gobet, B. Farizon, M. Farizon, M.J. Gaillard, S. Louc, N. Goncalves, M. Barbatti, H. Luna, G. Jalbert, N.V. de Castro Faria, M.C. Bacchus Montabonel, J.P. Buchet, M. Carré, T.D. Märk, Phys. Rev. Lett. 86, 4263-4266 (2001)
- ⁶ M.Ohkubo et al, Int.J.MassSpectrom. 299, 94 (2011)
- ⁷ M. Farizon, B. Farizon, H. Abdoul-Carime, G. Bruny, S. Eden, S. Feil, C. Montano , Method for tandem time of flight analysis and analysis appliance using said method (EN). / Procédé d'analyse à temps de vol en tandem et appareil d'analyse en faisant application (FR). Université Claude Bernard Lyon 1, Dépôt de brevet n° 09 58991 du 15 Décembre 2009, PCT/FR2010/052733 patent ; Jul, 7 2011: WO 2011/080455.
- ⁸ G. Bruny, S. Eden, S. Feil, R. Fillol, K. El Farkh, M.M. Harb, C. Teyssier, S.Ouaskit, H. Abdoul-Carime, B. Farizon, M. Farizon and T.D. Märk, Review of Scientific Instruments 83 (2012) 013305
- ⁹ Z. P.Wang et al., Int. J. Mass Spectrom. 285, 143-148 (2009)
- ¹⁰ W.C. Wiley, I.H. McLaren, Time of Flight Mass Spectrometer with Improved Resolution, Review of Scientific, 26, 1150 (1955)
- ¹¹ B. Farizon, M. Farizon, M.J. Gaillard, E. Gerlicet S. Ouaskit, Nucl. Instr. and Meth. in Phys. Res. B 101, 287 (1995)
- ¹² Dahl DA.SIMION® 3D Version 7.0, in 50th ASMS Conference on Mass Spectrometry and Allied Topics, June 2000, Long Beach, CA; 717.
- ¹³ J. Laskin, C. Lifshitz, J. Mass Spectr. 36459 (2001)
- ¹⁴ M.Miyazaki et al. Science 304, 1134 (2004)
- ¹⁵ J. Headrick, E. G. Diken, R. S. Walters, et al., Science 308, 1765 (2005)