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EIGENVALUE ESTIMATES FOR THE DIRAC OPERATOR AND HARMONIC 1-FORMS OF CONSTANT LENGTH

ANDREI MOROIANU AND LIVIU ORNEA

ABSTRACT.

We prove that on a compact n -dimensional spin manifold admitting a non-trivial harmonic 1-form of constant length, every eigenvalue λ of the Dirac operator satisfies the inequality $\lambda^2 \geq \frac{n-1}{4(n-2)} \inf_M \text{Scal}$. In the limiting case the universal cover of the manifold is isometric to $\mathbb{R} \times N$ where N is a manifold admitting Killing spinors.

Estimations de valeurs propres pour l'opérateur de Dirac et 1-formes harmoniques de longueur constante

RÉSUMÉ.

Nous démontrons que toute valeur propre λ de l'opérateur de Dirac d'une variété spinorielle compacte, de dimension n , qui admet une 1-forme harmonique non-triviale de longueur constante vérifie l'inégalité $\lambda^2 \geq \frac{n-1}{4(n-2)} \inf_M \text{Scal}$. Dans le cas limite le revêtement universel de la variété est isométrique à $\mathbb{R} \times N$ où N est une variété admettant des spineurs de Killing.

1. INTRODUCTION

Let (M^n, g) be a compact spin manifold of (real) dimension n . For positive scalar curvature $\text{Scal} := \text{Scal}(M, g)$, every eigenvalue λ of the Dirac operator D satisfy the well-known Friedrich inequality [4]:

$$(1) \quad \lambda^2 \geq \frac{n}{4(n-1)} \inf_M \text{Scal}.$$

There exist manifolds on which the inequality is indeed sharp: the limiting case is equivalent to the existence of a Killing spinor, *i.e.* a spinor Ψ satisfying the equation

$$(2) \quad \nabla_X \Psi + \frac{\lambda}{n} X \cdot \Psi = 0.$$

Key words and phrases. Spin manifold, eigenvalues of Dirac operator, harmonic form, parallel form, Killing spinor.

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On the other hand, Hijazi noticed that a manifold admitting a parallel k -form, $k \neq 0, n$, carries no Killing spinors [5]. Consequently, Friedrich's inequality cannot be sharp on manifolds with geometric structures which support parallel forms. Various authors have obtained sharp improvements of (1) on Kähler and quaternionic Kähler manifolds (see [6], [7], [8]).

Another recent improvement of Friedrich inequality in this latter direction was found by Alexandrov, Grantcharov and Ivanov. They proved in [1] that the existence of a parallel 1-form on M^n , $n \geq 3$, implies the inequality

$$(3) \quad \lambda^2 \geq \frac{n-1}{4(n-2)} \inf_M \text{Scal}.$$

The universal covering space of the manifolds appearing in the limiting case was also described.

In this note we generalize the above result showing that (3) can be derived only from the existence of a *harmonic* 1-form with constant length:

Theorem 1. *Inequality (3) holds on any compact spin manifold (M^n, g) , $n \geq 3$, admitting a non-trivial harmonic 1-form θ of constant length. The limiting case is obtained if and only if θ is parallel and the eigenspinor Ψ corresponding to the smallest eigenvalue of the Dirac operator satisfies the following Killing type equation:*

$$(4) \quad \nabla_X \Psi + \frac{\lambda}{n-1} (X \cdot \Psi - \langle X, \theta \rangle \theta \cdot \Psi) = 0,$$

after rescaling θ to unit length.

The condition for the norm of the 1-form to be constant is essential, in the sense that the topological constraint alone – the existence of a non-trivial harmonic 1-form – does not allow any improvement of Friedrich's inequality.

Indeed, motivated by a conjecture appearing in an earlier version of this note, Bär and Dahl [3] have constructed, on any compact spin manifold M^n and for every positive real number ϵ , a metric g_ϵ on M with the property that $\text{Scal}_{g_\epsilon} \geq n(n-1)$ and such that the first eigenvalue of the Dirac operator satisfies $\lambda_1^2(D_\epsilon) \leq \frac{n^2}{4} + \epsilon$. This construction clearly shows that no improvement of Friedrich's inequality can be obtained under purely topological restrictions.

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2. THE MAIN INEQUALITY

Let θ be a 1-form of unit length on a spin manifold (M^n, g) and let Ψ be an arbitrary spinor field on M . We identify 1-forms with vector fields by means of the scalar product that we denote with $\langle \cdot, \cdot \rangle$.

Consider the following “twistor-like” operator $T : TM \otimes \Sigma M \rightarrow \Sigma M$

$$T_X \Psi = \nabla_X \Psi + \frac{1}{n-1} X \cdot D\Psi - \frac{1}{n-1} \langle X, \theta \rangle \theta \cdot D\Psi - \langle X, \theta \rangle \nabla_\theta \Psi,$$

where the dot \cdot denotes Clifford multiplication. A simple calculation yields

$$(5) \quad |T\Psi|^2 = |\nabla\Psi|^2 - \frac{1}{n-1} |D\Psi|^2 - |\nabla_\theta \Psi|^2 + \frac{2}{n-1} \langle D\Psi, \theta \cdot \nabla_\theta \Psi \rangle.$$

From now on we will suppose that θ is harmonic, M is compact with volume element $d\mu$ and has positive scalar curvature Scal , and Ψ is an eigenspinor of the Dirac operator D of M corresponding to the least eigenvalue (in absolute value), say λ . We let $\{e_i\}$, $i = 1, \dots, n$ denote a local orthonormal frame on M .

The harmonicity of θ implies the following useful relation:

$$(6) \quad D(\theta \cdot \Psi) = -\theta \cdot D\Psi - 2\nabla_\theta \Psi.$$

Indeed, one may write:

$$\begin{aligned} D(\theta \cdot \Psi) &= \sum e_i \cdot \nabla_{e_i} (\theta \cdot \Psi) = \sum e_i \cdot (\nabla_{e_i} \theta) \cdot \Psi + e_i \cdot \theta \cdot \nabla_{e_i} \Psi \\ &= (d\theta + \delta\theta) \cdot \Psi + e_i \cdot \theta \cdot \nabla_{e_i} \Psi \\ &= -\theta \cdot \sum e_i \cdot \nabla_{e_i} \Psi - 2\langle e_i, \theta \rangle \nabla_{e_i} \Psi \\ &= -\theta \cdot D\Psi - 2\nabla_\theta \Psi. \end{aligned}$$

Taking the square norm in (6) yields

$$(7) \quad |D(\theta \cdot \Psi)|^2 = |\theta \cdot D\Psi|^2 + 4|\nabla_\theta \Psi|^2 - 4\langle D\Psi, \theta \cdot \nabla_\theta \Psi \rangle,$$

By integration over M in (5), using (7) to express the last term in the right hand side of (5), and the Lichnerowicz formula $D^2 = \nabla^* \nabla + \frac{1}{4} \text{Scal}$, we get

$$(8) \quad \int_M |T\Psi|^2 d\mu = \int_M \left\{ \frac{n-2}{n-1} |D\Psi|^2 - \frac{1}{4} \text{Scal} |\Psi|^2 - \frac{n-3}{n-1} |\nabla_\theta \Psi|^2 - \frac{1}{2(n-1)} (|D(\theta \cdot \Psi)|^2 - |\theta \cdot D\Psi|^2) \right\} d\mu.$$

The term in the last bracket of the integrand is clearly positive since, from the choice of λ to be minimal, we have from the classical Rayleigh inequality

$$\lambda^2 \leq \frac{\int_M |D\Phi|^2 d\mu}{\int_M |\Phi|^2 d\mu}$$

for every Φ . In particular, for $\Phi = \theta \cdot \Psi$ this reads

$$\int_M |D(\theta \cdot \Psi)|^2 d\mu \geq \lambda^2 \int_M |\theta \cdot \Psi|^2 d\mu = \lambda^2 \int_M |\Psi|^2 d\mu = \int_M |D\Psi|^2 d\mu = \int_M |\theta \cdot D\Psi|^2 d\mu.$$

Thus (8) gives

$$\int_M \left(\frac{n-2}{n-1} \lambda^2 - \frac{1}{4} \text{Scal} \right) |\Psi|^2 d\mu = \int_M |T\Psi|^2 + \frac{n-3}{n-1} |\nabla_\theta \Psi|^2 + \frac{1}{2(n-1)} (|D(\theta \cdot \Psi)|^2 - |\theta \cdot D\Psi|^2) d\mu \geq 0,$$

which immediately implies the first statement of Theorem 1.

3. THE LIMITING CASE

Suppose now that equality is reached in (3) for the eigenvalue λ with corresponding eigenspinor Ψ . Then $T\Psi = 0$. Contracting with e_i :

$$\sum e_i \cdot \nabla_{e_i} \Psi + \frac{\lambda}{n-1} \sum e_i \cdot e_i \cdot \Psi - \frac{\lambda}{n-1} \sum e_i \cdot \langle e_i, \theta \rangle \theta \cdot \Psi - e_i \cdot \langle e_i, \theta \rangle \nabla_{\theta} \Psi = 0,$$

gives $\theta \cdot \nabla_{\theta} \Psi = 0$, so $\nabla_{\theta} \Psi = 0$ (for $n > 3$ this follows directly from the vanishing of the integral $\int_M \frac{n-3}{n-1} |\nabla_{\theta} \Psi|^2 d\mu$). Thus Ψ satisfies the Killing type equation

$$(9) \quad \nabla_X \Psi = aX \cdot \Psi - a\langle X, \theta \rangle \theta \cdot \Psi, \quad a = -\frac{\lambda}{n-1}.$$

In order to show that θ is parallel, we first compute the spin curvature operator $\mathcal{R}_{Y,X} = [\nabla_Y, \nabla_X] - \nabla_{[Y,X]}$ acting on Ψ . We have successively :

$$\begin{aligned} \frac{1}{a} \nabla_Y \nabla_X \Psi &= \nabla_Y X \cdot \Psi + aX \cdot (Y - \langle Y, \theta \rangle \theta) \cdot \Psi - \langle \nabla_Y X, \theta \rangle \theta \cdot \Psi \\ &\quad - \langle X, \theta \rangle \nabla_Y \theta \cdot \Psi - \langle X, \nabla_Y \theta \rangle \theta \cdot \Psi - a\langle X, \theta \rangle \theta \cdot (Y - \langle Y, \theta \rangle \theta) \cdot \Psi \\ \frac{1}{a} \mathcal{R}_{Y,X} \cdot \Psi &= a(X \cdot Y - Y \cdot X) \cdot \Psi - a(\langle Y, \theta \rangle X \cdot \theta - \langle X, \theta \rangle Y \cdot \theta) \cdot \Psi \\ &\quad + (\langle Y, \nabla_X \theta \rangle - \langle X, \nabla_Y \theta \rangle) \theta \cdot \Psi + (\langle Y, \theta \rangle \nabla_X \theta - \langle X, \theta \rangle \nabla_Y \theta) \cdot \Psi \\ &\quad + a(\langle Y, \theta \rangle \theta \cdot X - \langle X, \theta \rangle \theta \cdot Y) \cdot \Psi \\ &= a(X \cdot Y - Y \cdot X) \cdot \Psi + 2a(\langle Y, \theta \rangle \theta \cdot X - \langle X, \theta \rangle \theta \cdot Y) \cdot \Psi \\ &\quad + (\langle Y, \nabla_X \theta \rangle - \langle X, \nabla_Y \theta \rangle) \theta \cdot \Psi + (\langle Y, \theta \rangle \nabla_X \theta - \langle X, \theta \rangle \nabla_Y \theta) \cdot \Psi. \end{aligned}$$

Using again the harmonicity of θ we easily derive :

$$(10) \quad \begin{aligned} \frac{1}{2a} Ric(X) \cdot \Psi &= \frac{1}{a} \sum e_i \cdot \mathcal{R}_{e_i, X} \Psi = \\ &= 2(n-2)a(X - \langle X, \theta \rangle \theta) \cdot \Psi + 2\theta \cdot \nabla_X \theta \cdot \Psi - 2\langle X, \nabla_{\theta} \theta \rangle \Psi. \end{aligned}$$

But $\nabla_{\theta} \theta = 0$ because θ is unitary and closed. Indeed, for any vector field Y :

$$\langle Y, \nabla_{\theta} \theta \rangle = d\theta(\theta, Y) - \langle \theta, Y \rangle = 0.$$

Hence, taking $X = \theta$ in (10), we obtain $Ric(\theta) = 0$. Then, as θ is harmonic, the Bochner formula assures that θ is parallel. This completes the proof of Theorem 1.

Since the 1-form θ has to be parallel in the limiting case, we can apply Theorem 3.1 in [1] to determine the universal cover of M . In fact, as proved in *loc. cit.*, this is isometric to a Riemannian product $\mathbb{R} \times N$, where N is a spin manifold carrying a real Killing spinor, hence can be described by Bär's classification [2]. Finally, M turns out to be a suspension of an isometry of N/Γ (a finite quotient of N) over the circle.

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