

Mathematical derivation of viscous shallow-water equations with zero surface tension

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February 12, 2010

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Abstract. The purpose of this paper is to derive rigorously the so called viscous shallow water equations given for instance page 958-959 in [A. ORON, S.H. DAVIS, S.G. BANKOFF, *Rev. Mod. Phys.*, 69 (1997), 931–980]. Such a system of equations is similar to compressible Navier-Stokes equations for a barotropic fluid with a non-constant viscosity. To do that, we consider a layer of incompressible and Newtonian fluid which is relatively thin, assuming *no surface tension* at the free surface. The motion of the fluid is described by 3d Navier-Stokes equations with constant viscosity and free surface. We prove that for a set of suitable initial data (asymptotically close to “shallow water initial data”), the Cauchy problem for these equations is well-posed, and the solution converges to the solution of viscous shallow water equations. More precisely, we build the solution of the full problem as a perturbation of the strong solution to the viscous shallow water equations. The method of proof is based on a Lagrangian change of variable that fixes the fluid domain and we have to prove the well-posedness in thin domains: we have to pay a special attention to constants in classical Sobolev inequalities and regularity in Stokes problem.

Keywords: Navier-Stokes, shallow water, thin domain, free surface, asymptotic analysis, Sobolev spaces.

AMS subject classification: 35Q30, 35R35, 76A20, 76B45, 76D08.

1 Introduction

This paper deals with a free boundary problem of nonstationary Navier-Stokes equations in thin domain justifying the derivation of the viscous shallow water equations without surface tension. For information, such a system may be used to describe long-scale evolution of thin liquid films : See Equations (4.12a)-(4.12d) written in [18]. Depending on the geometry, there are two kinds of free boundary problems for the motion of a viscous fluid. The first one describes the motion of an isolated mass of fluid bounded by a free boundary and the second one describes the motion of a fluid occupying a semi-infinite domain in \mathbb{R}^n bounded above by a free surface and below by a fixed part of the boundary. We will consider in the present paper the second problem and refer the interested reader to [18], [12] for applications for thin liquid films and rivers. First mathematical results for free boundary problems were local existence theorems: A first local existence theorem published in 1977 by V.A. SOLONNIKOV, in [19] ; a local existence theorem for the second problem by J.T. BEALE, in 1980, in [3]. The following years brought other local existence theorems for equations of motion of incompressible fluids: see for instance [1], [2], [22], [25], [26]. The next step in investigating free boundary problems for incompressible Navier-Stokes equations was to obtain global existence theorems initiated by J.T. BEALE in his paper [4] published in 1984. This paper was devoted to the motion of a fluid contained in a three dimensional infinite ocean. The first global existence result concerning the motion of a fixed mass of a fluid bounded by a free surface appeared, in the paper [20], by V.A. SOLONNIKOV in 1986. The method used to prove the global existence results in the two papers mentioned above are completely different. J.T. BEALE examined the surface waves after transforming it to the equilibrium domain through an adequate change of variable (dilatation in the vertical variable) while V.A. SOLONNIKOV applied Lagrangian coordinates and this way transformed the considered drop problem to the initial domain. Other global existence theorem can be found. All global existence results for incompressible fluids are obtained for initial data sufficiently close to an equilibrium solution. Note that A. TANI, in [22] [23], used the strategy developed by V.A. SOLONNIKOV for free surface problem to get a local existence result through the use of Lagrangian coordinates to fix the domain and the method of successive approximations. In con-

trast to J.T. BEALE, the regularity of the solution obtained by A. TANI is sharp. Note that in J.T. BEALE existence of arbitrary $T > 0$ but for sufficiently small initial data in dependence on T is also proved. The result is obtained in the same class of functions than the local existence. To the author's knowledge, in all these studies Dirichlet boundary conditions at the bottom are considered instead of Navier-Slip boundary conditions. Only recent papers with some rigid-lid assumptions consider Navier-Stokes equations in thin domains with Navier-Slip conditions, see for instance [9], [14]. In this paper, we will adopt the Lagrangian transform strategy in order to prove existence and stability of the viscous shallow water solution. To do that, we have to prove the well-posedness in thin domains with a special attention to constants in classical Sobolev inequalities and regularity in Stokes problem. We also prove the stability result in the Lagrangian form as in the paper written by D. HOFF, see [13], where uniqueness of weak solutions of the Navier-Stokes equations of multidimensional, compressible flow is proved using the Lagrangian formulation. We remark also that, using the Lagrangian formulation, no surface tension term is needed compared to [8]. Remark that there exists, at least, two different shallow-water type systems mainly derived from two different bottom boundary conditions choices for Navier-Stokes equations with free surface: friction boundary conditions, no-slip boundary conditions. The system based on friction boundary conditions assumption is studied in the present paper and was mentioned in [18]. The system based on no-slip assumption has been formally derived in [5] and recently justified in [8] assuming non zero surface tension. Reader interested by shallow-water equations and related topics introduction is referred to the recent Handbook [6].

The paper will be divided in four Sections. The first one presents the Navier-Stokes equations with free surface and the viscous shallow-water approximation. The second section is devoted to the linearized Navier-Stokes equations around the approximate solution built from a solution of the viscous shallow-water equations. We give anisotropic estimates and Korn's inequality in thin domains and then prove the well posedness of the linearized Navier-Stokes equations. The main theorem and its corollary are proved in Section 4 that means well posedness and convergence of the full Navier-Stokes equations in Lagrangian form. On the last section, we draw some conclusions on this paper and present some perspectives to this work.

2 Navier-Stokes equations with free surface and viscous shallow-water approximation

In this section, we write $3d$ Navier-Stokes equations for a thin layer of an incompressible Newtonian fluid with a free surface. We then show how to obtain an *approximate* solution of these equations in the shallow water scaling: for that purpose, we will use smooth solution (see W. WANG and C.J. XU [27]) of a viscous shallow water model (6) derived for instance in A. ORON, S.H. DAVIS and S.G. BANKOFF [18] (see also J.-F. GERBEAU and B. PERTHAME [12]). We construct formally a second order approximate solution (with respect to the aspect ratio ε which measures the relative shallowness of the fluid layer) which will help to solve the Cauchy problem for Navier-Stokes equations with free surface and provides also convergence to viscous shallow-water equations. For this purpose, we formulate the equations in Lagrangian coordinates.

2.1 The equations and an existence theorem for the associated viscous shallow-water system

In this paper, we consider the motion of a thin layer of Newtonian fluid with a constant density, set to 1. The differential operators $\nabla_x, \operatorname{div}_x$ will concern the horizontal coordinates $x = (x_1, x_2)$. In the following, the fluid velocity u is denoted (u_H, u_V) where $u_H \in \mathbb{R}^2$ (resp $u_V \in \mathbb{R}$) is the horizontal (resp. vertical) fluid velocity. Navier-Stokes equations are then written, in their non dimensional form as

$$\begin{aligned} \partial_t u + u \cdot \nabla u + \frac{1}{F_0^2} \nabla p &= -\frac{e_3}{F_0^2} + \frac{1}{\operatorname{Re}} \operatorname{div}(D(u)), \\ \operatorname{div} u &= 0, \end{aligned} \tag{1}$$

for all $t > 0$ and $(x, z) \in \Omega_t$ the fluid domain

$$\Omega_t = \{(x, z) / x \in \mathbb{X}, \quad z \in (0, h(x, t))\}, \quad \mathbb{X} = \mathbb{T}^n, \mathbb{R}^n, \quad n = 1, 2.$$

The nondimensional numbers F_0, Re are respectively the Froude and Reynolds numbers defined as

$$F_0^2 = \frac{U^2}{gL}, \quad \operatorname{Re} = \frac{LU}{\nu},$$

with U a characteristic velocity of the fluid and L a characteristic wavelength. Navier-Stokes equations are completed by boundary conditions. The impermeability condition at the free surface reads

$$\partial_t h + u_H|_{z=h} \cdot \nabla_x h = u_V|_{z=h}, \tag{2}$$

which means that the fluid velocity is parallel to the free surface. We also assume a Navier slip condition at the bottom

$$u_V|_{z=0} = 0, \quad \partial_z u_H|_{z=0} = \gamma u_H|_{z=0}, \quad (3)$$

and at the free surface, we assume the continuity of the fluid stress

$$\left(\frac{1}{\text{Re}} D(u) - \frac{p}{F_0^2} Id \right) |_{z=h} \vec{n} = -\frac{p_{atm}}{F_0^2} \vec{n}. \quad (4)$$

Here $D(u) = \nabla u + \nabla u^T$ is the total deformation tensor of the fluid and can be written as

$$D(u) = \begin{pmatrix} D_x(u_H) & \partial_z u_H + \nabla_x u_V \\ (\partial_z u_H + \nabla_x u_V)^T & 2\partial_z u_V \end{pmatrix}, \quad D_x(u_H) = \nabla_x u_H + (\nabla_x u_H)^T.$$

The vector \vec{n} is the normal to the free surface:

$$\sqrt{1 + |\nabla_x h|^2} \vec{n} = \begin{pmatrix} -\nabla_x h \\ 1 \end{pmatrix}.$$

Since the fluid pressure is defined up to a constant, we suppose that $p_{atm} = 0$. Note that equation (4) can be equivalently written as

$$\begin{aligned} \frac{p|_{z=h}}{F_0^2} &= \frac{(2\partial_z u_V - (\nabla_x h)^T D_x(u_H) (\nabla_x h))|_{z=h}}{\text{Re}(1 - |\nabla_x h|^2)}, \\ \partial_z u_H + \nabla_x u_V|_{z=h} &= \left(D_x(u_H) - \frac{2\partial_z u_V - (\nabla_x h)^T D_x(u_H) \nabla_x h}{1 - |\nabla_x h|^2} \right) |_{z=h} \nabla_x h. \end{aligned} \quad (5)$$

In this paper, we will assume that the fluid domain is relatively thin and set $h = \varepsilon \bar{h}$, $F_0^2 = \varepsilon F^2$ with $\varepsilon \ll 1$. As a result, the Froude number $F = \frac{U}{\sqrt{gH}}$ with $H = \varepsilon L$, the characteristic fluid height, is the physical Froude number as the ratio between fluid velocity and wavespeed of perturbations at the free surface. Moreover, following the method of derivation by Gerbeau and Perthame [12] to obtain viscous shallow water equations, we suppose that the friction at the bottom is small: $\gamma = \varepsilon \bar{\gamma}$. Let us define an approximate solution of Navier-Stokes system (1,2,3,4). For that purpose, we introduce (h_0, u_0) a smooth solution of the shallow water system mentioned by A. ORON, S.H. DAVIS, S.G. BANKOFF [18], see page 958-959 with $\bar{C} \rightarrow \infty$:

$$\begin{aligned} \partial_t h_0 + \text{div}_x(h_0 u_0) &= 0, \\ h_0(\partial_t u_0 + u_0 \cdot \nabla_x u_0) + \nabla_x \frac{h_0^2}{2F^2} &= \frac{1}{\text{Re}} \text{div}_x(h_0 D_x(u_0)) \\ &\quad + \frac{2}{\text{Re}} \nabla_x(h_0 \text{div}_x(u_0)) - \frac{\bar{\gamma} u_0}{\text{Re}}. \end{aligned} \quad (6)$$

For the moment, we carry out a formal computation: we will make more precise assumptions on this solution in the sequel. We present a second order approximate solutions, that means we search for an approximate solution of (1,2,3,4) in the form $h_a = \varepsilon h_0$ and

$$u_{a,H} = u_0 + u_1 z + u_2 \frac{z^2}{2}, \quad u_{a,V} = w_0 + w_1 z + w_2 \frac{z^2}{2} + w_3 \frac{z^3}{6}, \quad p_a = p_0 + p_1 z + p_2 \frac{z^2}{2}.$$

In the main Theorem we need higher order approximate solution that can be obtained in the same way, see [8] for construction of high order approximation of Navier Stokes systems with a free surface. Inserting this ansatz in the divergence free condition, one has

$$w_1 = -\operatorname{div}_x u_0, \quad w_2 = -\operatorname{div}_x u_1, \quad w_3 = -\operatorname{div}_x u_2.$$

Moreover, the boundary conditions at the bottom imposes

$$w_0 = 0, \quad u_1 = \varepsilon \bar{\gamma} u_0.$$

Next, using the second equation of (1) and boundary condition (5), one obtains the approximation of p as

$$p_a = \varepsilon h_0 - z - \frac{2\varepsilon F^2}{\operatorname{Re}} \operatorname{div}_x u_0 - \frac{2\varepsilon^3 \bar{\gamma} F^2}{\operatorname{Re}} \operatorname{div}_x u_0.$$

Next, we insert u_a and p_a into the second equation of (5): identifying order $O(\varepsilon)$ terms, one proves that

$$u_2 + \nabla_x w_1 = (D_x(u_0) + 2\operatorname{div}_x(u_0)Id) \frac{\nabla_x h_0}{h_0} - \frac{\bar{\gamma} u_0}{h_0}.$$

Remark. For the approximate solution, we can calculate an ansatz up to any order on derivatives of (h_0, u_0) and polynomial dependence with respect to z . We refer to [8] for more details. We end up at order which gives a remaining term of order (z^3) for interior equations and (ε^3) at free surface.

A Navier-Stokes type system is satisfied by the high order shallow water approximation (u_a, p_a) . The adequate ansatz satisfies

$$\begin{aligned} \partial_t u_a + u_a \cdot \nabla u_a + \frac{1}{\varepsilon F^2} \nabla p_a &= -\frac{e_3}{\varepsilon F^2} + \frac{1}{\operatorname{Re}} \operatorname{div}(D(u_a)) + \mathcal{O}(z^3), \\ \operatorname{div} u_a &= \mathcal{O}(z^3). \end{aligned} \quad (7)$$

Moreover, the boundary conditions at the bottom are *exactly* satisfied and the boundary conditions at the free surface read

$$\partial_t h_a + u_{a,H}|_{z=h_a} \nabla_H h_a - u_{a,V}|_{z=h_a} = \mathcal{O}(\varepsilon^3),$$

$$\left(\frac{D(u_a)}{\text{Re}} - \frac{p_a}{\varepsilon F^2} \right) \Big|_{z=h_a} \vec{n}_a = \mathcal{O}(\varepsilon^3).$$

For the sake of completeness, we recall the existence and uniqueness result related to the viscous shallow water system (6) that can be obtained using the same procedure than in [27]

Theorem 1 *Let $s > 0$, $u_0(0), h_0(0) - 1 \in H^{s+2}(\mathbb{R}^2)$, $\|h_0(t=0) - 1\|_{H^{s+2}(\mathbb{R}^2)} \ll 1$. Then there exists a positive time T , a unique solution (u_0, h_0) of (6) such that*

$$u_0, h_0 - 1 \in L^\infty(0, T; H^{2+s}(\mathbb{R}^2)), \quad \nabla u_0 \in L^2(0, T; H^{2+s}(\mathbb{R}^2)).$$

Furthermore, there exists a constant c such that if $\|h_0(t=0) - 1\|_{H^{s+2}(\mathbb{R}^2)} + \|u_0(t=0)\|_{H^{s+2}(\mathbb{R}^2)} \leq c$ then we can choose $T = +\infty$ with the same control on (h_0, u_0) .

Remark: the smallness assumption on $h_0(t=0) - 1$ can be replaced by $h_0(t, x) > \alpha > 0$ to prove local well posedness provided that we work in Besov spaces: see the paper of R. DANCHIN [11] for more details. However, we will need here this particular assumption to prove the well posedness of free surface Navier-Stokes equations in dimension 3 and convergence to shallow water equations. In dimension 2 for free surface Navier-Stokes equations, we just need an assumption similar to the one of [11], namely that there is no vacuum. Of course the results are only local in time: for arbitrary time interval, we need a smallness assumption on $h_0(t=0) - 1$ and $u_0(t=0)$.

For global existence of weak solutions to a related viscous shallow water system, a result has been proved in [7] assuming an extra turbulent drag term.

Remark: In this paper, we will choose $s > 4$. Indeed, a careful analysis of the remainder shows that the contained fourth order derivatives of (h_0, u_0) have to be Lipschitz.

2.2 Reduction of Navier-Stokes equations to a fixed domain

In what follows, we write the Navier-Stokes equations in a fixed domain: this is done using a Lagrangian formulation of the equations. We first search the fluid velocity and pressure in the form

$$u = u_a + \tilde{u}, \quad p = p_a + \frac{\varepsilon F^2}{\text{Re}} \tilde{p}$$

Navier-Stokes equations (1) then reads

$$\begin{aligned} \partial_t \tilde{u} + (u_a + \tilde{u}) \cdot \nabla \tilde{u} + \tilde{u} \cdot \nabla u_a + \frac{\nabla \tilde{p}}{\text{Re}} &= \frac{1}{\text{Re}} \text{div}(D(\tilde{u})) + f_a, \\ \text{div} \tilde{u} &= 0. \end{aligned} \quad (8)$$

in the fluid domain $\Omega_{\varepsilon,t} = \{(x, z) / x \in \mathbb{X}, 0 \leq z \leq \varepsilon h(x, t)\}$. The boundary conditions are then written as

$$\tilde{u}_V|_{z=0} = 0, \quad \partial_z \tilde{u}_H|_{z=0} = \varepsilon \bar{\gamma} \tilde{u}_H|_{z=0}. \quad (9)$$

At the free surface, one has

$$\begin{aligned} (D(\tilde{u}) - \tilde{p})_{z=\varepsilon h} \bar{n}_\varepsilon &= g_a, \\ \partial_t(\varepsilon h) + (u_{a,H} + \tilde{u}_H)|_{z=\varepsilon h} \cdot \nabla_x(\varepsilon h) &= u_{a,V} + \tilde{u}_V|_{z=\varepsilon h}. \end{aligned} \quad (10)$$

Next, we introduce the Lagrangian change of variable

$$\frac{dX}{dt} = u_{a,H}(t, X, Z) + \tilde{u}_H(t, X, Z), \quad \frac{dZ}{dt} = u_{a,V}(t, X, Z) + \tilde{u}_V(t, X, Z).$$

For the sake of simplicity, we assume that the initial fluid domain is $\mathbb{X} \times (0, \varepsilon)$: the rigorous justification will remain true for $|\langle h_0(t=0) \rangle - 1| \ll 1$ and $\bar{\gamma}$ sufficiently large. Denote $x_0 \in \mathbb{X}, z_0 \in (0, \varepsilon)$, the coordinates in this domain and A the Jacobian matrix of the Lagrangian change of variable

$$A = \begin{pmatrix} \frac{\partial X}{\partial x_0} & \frac{\partial X}{\partial z_0} \\ (\nabla_{x_0} Z)^T & \frac{\partial Z}{\partial z_0} \end{pmatrix}.$$

The chain rules are given by

$$\begin{pmatrix} \nabla_{x_0} \\ \partial_{z_0} \end{pmatrix} = A^T \begin{pmatrix} \nabla_x \\ \partial_z \end{pmatrix}, \quad \begin{pmatrix} \nabla_{x_0} \bar{u} & \partial_{z_0} \bar{u} \\ \nabla_{x_0} \bar{w}^T & \partial_{z_0} \bar{w} \end{pmatrix} = \begin{pmatrix} \nabla_x \bar{u} & \partial_z \bar{u} \\ \nabla_x \bar{w}^T & \partial_z \bar{w} \end{pmatrix} A.$$

The fluid height is defined implicitly as $\varepsilon h(t, X(t, x_0, \varepsilon)) = Z(t, x_0, \varepsilon)$ so that the impermeability condition is satisfied. Moreover, one has

$$\varepsilon \nabla_x h(t, X(t, x_0, \varepsilon)) = \left(\frac{\partial X}{\partial x_0} \right)^{-T} \nabla_{x_0} Z(t, x_0, \varepsilon).$$

The Lagrangian velocity is defined as $\bar{u} = \tilde{u}(t, X)$ and \bar{u} defined on the fixed thin domain $\mathbb{X} \times (0, \varepsilon)$. We further introduce the Lagrangian pressure $\bar{p} = \tilde{p}(t, X)$. In that setting, Navier-Stokes equations are written as

$$\partial_t \bar{u} + (A^{-1} \bar{u}) \cdot \nabla \bar{u}_a + \frac{A^{-T} \nabla \bar{p}}{\text{Re}} = \frac{A^{-T}}{\text{Re}} \left(\text{div}(A^T \mathcal{P}) - \nabla A : \mathcal{P} \right) + \bar{f}_a, \quad (11)$$

$$\text{div}(A^{-1} \bar{u}) = \bar{\sigma}_a, \quad (12)$$

where $\bar{f}_a = f_a(t, X)$, $\bar{u}_a = u_a(t, X)$, $(\nabla A : \mathcal{P})_i = \sum_{j,k} \partial_k A_{j,i} \mathcal{P}_{j,k}$. The matrix \mathcal{P} is the transformed deformation tensor:

$$\mathcal{P} = (\nabla \bar{u}) A^{-1} A^{-T} + A^{-T} (\nabla \bar{u})^T A^{-T}.$$

The boundary conditions at the free surface $z_0 = \varepsilon$ reads

$$\left(\mathcal{P}|_{z_0=\varepsilon} - p|_{z_0=\varepsilon} Id \right) \bar{n} = 0, \quad \bar{n} = \begin{pmatrix} -(\frac{\partial X}{\partial x_0})^{-T} \nabla_{x_0} Z(t, x_0, \varepsilon) \\ 1 \end{pmatrix}, \quad (13)$$

whereas the boundary conditions at the bottom $z = 0$ are written as

$$\bar{u}_V|_{z_0} = 0, \quad \det\left(\frac{\partial X}{\partial x_0}\right)|_{\partial_{z_0} \bar{u}_H|_{z_0=0}} = \varepsilon \bar{\gamma} \bar{u}_H|_{z_0=0}. \quad (14)$$

The main result of the paper is

Theorem 2 *There exists $C, K > 0$ so that if we assume $\|\bar{u}_0\|_2 \leq C^2 \varepsilon^{3/2}$, then there is a unique strong solution of (11)-(14) such that*

$$\sup_{(0,T)} \|\bar{u}\|_2^2 + \int_0^T \|\bar{u}\|_3^2 + \|\partial_t \bar{u}\|_1 + \|\bar{p}\|_2^2 + \|\partial_t \bar{p}\|_0^3 < K \varepsilon^{3/2}.$$

This gives the expected existence result of Navier-Stokes equations with free surface namely the system (1), (2), (3), (5) and its convergence to the solution of the shallow-water system in a Lagrangian form namely a control of the solution of (8)–(10) with respect to ε .

Corollary 1 *Let (h_0, u_0) be the solution of the viscous shallow-water system (6) such that $(h_0 - 1, u_0)$ small enough in $\mathcal{C}(0, T; H^{2+s}(\mathbb{R}^2))$ with $s > 4$. Let $(h_0^\varepsilon, u_0^\varepsilon)$ such that $\|(u_0^\varepsilon - u_a)|_{t=0}\|_2 \leq C \varepsilon^{3/2}$ and $h_0^\varepsilon = h_0$ and u_0^ε satisfying the compatibility conditions. Then System (1), (2), (3), (5) is well posed and $u^\varepsilon - u_a$ satisfy in its Lagrangian form the estimates given by Theorem 2.*

Note that there is no assumption on the time interval: the solution of Navier-Stokes system is defined on the same existence time as the solution of the viscous shallow-water system (6).

The associated linear system. Next we introduce a linear problem that will be usefull to prove the well-posedness of the *full* Navier-Stokes system. Indeed, this will be done using a fixed point argument. First, let us introduce the following Lagrangian system of coordinates

$$\frac{dX_0}{dt} = u_0(X_0), \quad \frac{dZ_0}{dt} = -Z_0 \operatorname{div}(u_0)(X_0),$$

with the initial conditions $X_0(0, \cdot) = x_0$, $Z_0(0, \cdot) = z_0$. It is easily proved that X_0, Z_0 satisfy

$$\frac{\partial X_0}{\partial z_0} = 0, \quad Z_0(t, x_0, z_0) = z_0 h_0(t, X_0(t, x_0)), \quad \det\left(\frac{\partial X_0}{\partial x_0}\right) = h_0^{-1}(t, X_0(t, x_0)).$$

In what follows, we will consider a special linear problem that is obtained if we substitute $X = X_0$ and $Z = Z_0$ in Navier-Stokes equations written in Lagrangian coordinates and drop order one differential operators. More precisely, this system is written as

$$\partial_t \bar{u} + \frac{A_0^{-T} \nabla p}{\text{Re}} = \frac{A_0^{-T}}{\text{Re}} \text{div}(A_0^T \mathcal{P}_0) + f, \quad (15)$$

with $\mathcal{P}_0 = \nabla(\bar{u}) A_0^{-1} A_0^{-T} + A_0^{-T} (\nabla \bar{u})^T A_0^{-T}$. The divergence condition reads

$$\text{div}(A_0^{-1} \bar{u}) = \sigma. \quad (16)$$

The boundary conditions at the bottom are

$$\bar{u}_V|_{z_0=0} = 0, \quad h_0^{-1} \partial_{z_0} \bar{u}_H|_{z_0=0} = \varepsilon \bar{\gamma} \bar{u}_H|_{z_0=0} + g_1, \quad (17)$$

whereas the boundary conditions at the free surface read

$$\left(A_0^T \mathcal{P}_0 - p \text{Id}\right)|_{z_0=\varepsilon} e_3 = g_2. \quad (18)$$

3 Linearized Navier-Stokes around the approximate solution of viscous shallow-water equations

This section is devoted to study the linearized Navier-Stokes equations around the approximate solution of viscous shallow-water (15,16,17,18) that was introduced in the previous section. We begin by a subsection dealing with Sobolev inequalities, Korn inequality and regularity result for Stokes problem in thin domain. In the second subsection, we give well posedness result for the full linearized system.

3.1 Anisotropic estimates in thin domains

In this part, we recall Sobolev inequalities that arises in thin domains and pay a special attention to the constants arising in such inequalities. We also consider a linear Stokes problem in a thin domain that will be useful to obtain estimates on the Sobolev norms of (\bar{u}, \bar{p}) .

3.1.1 Sobolev inequalities

In what follows, we use that mean of \tilde{u}_1 in the vertical direction is 0 to write Poincaré, Agmon and Ladyzhenskaia anisotropic inequalities in *thin* domains. Here the domain Ω is $\Omega = \mathbb{X}^n \times (0, \varepsilon)$. In this paper, we will use the following notations: $\|u\|_k = \|u\|_{H^k(\Omega)}$, for any $k \in \mathbb{N}$ with the convention $H^0(\Omega) = L^2(\Omega)$. Similarly, one denotes $|u|_k = |u|_{H^k(\mathbb{X})}$.

Proposition 1 *There exists a constant C independent of ε such that the following Ladyzhenskaia and Agmon estimates hold true:*

$$\|u\|_{L^6} \leq C\varepsilon^{-\frac{1}{3}}\|u\|_1, \quad \|u\|_{L^\infty} \leq C\varepsilon^{-\frac{1}{2}}\|u\|_2.$$

The proof of this proposition is found in [24]. From this estimate and the inequality $\|uv\|_{L^2} \leq \varepsilon^{\frac{1}{6}}\|u\|_{L^6}\|v\|_{L^6}$, one easily proves the following tame estimates for the products of functions and composition of functions.

Lemma 1 *We will use the following estimates: if $t > \frac{n}{2}$*

$$\begin{aligned} |uv|_{H^s(\mathbb{X})} &\leq C(|u|_{H^t(\mathbb{X})}|v|_{H^s(\mathbb{X})} + |u|_{H^s(\mathbb{X})}|v|_{H^t(\mathbb{X})}), \\ \|uv\|_{H^s(\Omega)} &\leq \frac{C}{\sqrt{\varepsilon}}(\|u\|_{H^t(\Omega)}\|v\|_{H^s(\Omega)} + \|u\|_{H^s(\Omega)}\|v\|_{H^t(\Omega)}), \end{aligned}$$

We have also the following property: For all $F : \mathbb{R}^p \rightarrow \mathbb{R}$ smooth enough with $F(0) = 0$ then

$$\|F(u)\|_s \leq C(\|u\|_{L^\infty})\|u\|_s.$$

We will also need the following trace theorems:

Lemma 2 *Let $u \in H^1$ is so that $u|_{z=0} = 0$ then one has $|u|_{\frac{1}{2}} \leq C\|u\|_1$, with C independent of ε . In the general case, one has $|u|_{\frac{1}{2}} \leq \frac{C}{\sqrt{\varepsilon}}\|u\|_1$.*

This result is easily proved with Fourier series, that reduces the problem to a 1d problem. The trace of a product is estimated using the following result.

Lemma 3 *Let $u \in H^{-\frac{1}{2}}$ and $v \in H^2$ then, one has*

$$|uv|_{-\frac{1}{2}} \leq \frac{C}{\sqrt{\varepsilon}}|u|_{-\frac{1}{2}}\|v\|_2.$$

The proof of this lemma is found in [10] in a fixed domain: we combine this proof and proposition 1 to obtain the estimate of lemma 3.

3.1.2 Korn inequality in thin domain

In what follows, we prove the following result.

Proposition 2 *There is a constant $C > 0$ independent of ε such that*

$$2 \|D(u)\|_{L^2(\Omega)}^2 + \varepsilon \bar{\gamma} |u_H|_{z=0}|_{L^2(\mathbb{X})}^2 \geq C \|u\|_{H^1(\Omega)}^2,$$

for all $u \in H^1(\Omega)$ so that $\operatorname{div} u = 0$ and $u_V|_{z=0} = 0$.

Proof. This is a consequence of Poincaré inequality and trace estimates that the L^2 norm of u_H and u_V are estimated as

$$\|u_V\|_{L^2(\Omega)} \leq C\varepsilon \|\partial_z u_V\|_{L^2(\Omega)}, \quad \|u_H\|_{L^2(\Omega)} \leq C(\varepsilon |u_H|_{z=0}|_{L^2(\mathbb{X})} + \varepsilon \|\partial_z u_H\|_{L^2(\Omega)}).$$

We will prove a refined inequality on the deformation tensor when $\mathbb{X} = \mathbb{T}^n$

$$2 \|D(u)\|_{L^2(\Omega)}^2 \geq C \|\nabla u\|_{L^2(\Omega)}^2.$$

For that purpose, we decompose the functions into Fourier series in the horizontal variables, $(u_H, u_V) = \sum_{k \in \mathbb{Z}^2} (u_k(z), w_k(z)) \exp(ik \cdot x)$. Note that a similar result is obtained when $\mathbb{X} = \mathbb{R}^n$, using Fourier transform. We find, using the divergence free condition

$$\begin{aligned} 2 \|D(u)\|_{L^2(\Omega)}^2 &= 2 \sum_{k \in \mathbb{Z}^2} \int_0^\varepsilon k_1^2 |u_{k,1}|^2 + k_2^2 |u_{k,2}|^2 + |k \cdot u_k|^2 \\ &+ \sum_{k \in \mathbb{Z}^2} \int_0^\varepsilon |k_2 u_{k,1} + k_1 u_{k,2}|^2 + |u'_{k,1} + \int_0^z k_1 k \cdot u_k|^2 + |u'_{k,2} + k_2 \int_0^z k \cdot u_k|^2, \\ \|\nabla u\|_{L^2(\Omega)}^2 &= \sum_{k \in \mathbb{Z}^2} \int_0^\varepsilon |k|^2 |u_k|^2 + |k \cdot u_k|^2 + \left| \int_0^z k \cdot u_k \right|^2 + |u'_k|^2. \end{aligned}$$

where $u'(z) = \frac{du}{dz}$. Setting $U_k = \int_0^z u_k$, there exists $C_k \in (0, 2)$ so that

$$\begin{aligned} &\int_0^\varepsilon 2k_1^2 |U'_1|^2 + 2k_2^2 |U'_2|^2 + 2|k \cdot U'| + |k_2 U'_1 + k_1 U'_2|^2 + |U''_1 + k_1 k \cdot U|^2 \\ &+ \int_0^\varepsilon |U''_2 + k_2 k \cdot U|^2 \geq C_k \int_0^\varepsilon |k|^2 |U|^2 + |k \cdot U'|^2 + |k \cdot U|^2 + |U''|^2, \end{aligned} \quad (19)$$

for all $U \in H^1(0, \varepsilon)^2$ so that $U|_{z=0} = 0$. We prove the inequality for each modes. We must prove that there is $C > 0$ independent of k so that $C_k \geq C$, for all $k \in \mathbb{Z}^2$. If $k = 0$, the two quantities are equal and one can choose $C_0 = 1$. In what follows, we assume that $k \neq 0$. Denote $C_k \in (0, 2)$ the

optimal value for the mode k . Using test function in $\mathcal{D}(0, \varepsilon)$, one can prove that, if $C_k \neq 1$, an optimal solution necessarily satisfies the differential system

$$D^2U_1 + k_2^2DU_1 = k_1k_2DU_2, \quad D^2U_2 + k_1^2DU_2 = k_1k_2DU_1, \quad \forall z \in (0, \varepsilon). \quad (20)$$

where D is the differential operator $D = \frac{d^2}{dz^2} - |k|^2 Id$. This space of solutions is 8-dimensional and reduced to a 6-dimensional one, using the condition $U|_{z=0} = 0$. Moreover, the eigenfunction U can be written $U(z) = \tilde{U}(|k|z)$, where $\tilde{U} \in \mathcal{S}$, the vector space defined as

$$\begin{aligned} \tilde{U}(z) = & \begin{pmatrix} -c \\ s \end{pmatrix} (a_1z + b_1 \sinh(z) + c_1(\cosh(z) - 1)) \\ & + \begin{pmatrix} s \\ c \end{pmatrix} (a_2 \sinh(z) + b_2z \sinh(z) + c_2z \cosh(z)). \end{aligned} \quad (21)$$

where we have denoted k as $k = |k|(c, s)$ and $a_i, b_i, c_i \in \mathbb{R}, i = 1, 2$. Then we have to show that there is C independent of k, ε so that $Q_2(\tilde{U}, k) \geq C Q_1(\tilde{U}, k)$, for all $\tilde{U} \in \mathcal{S}$ and Q_i the quadratic forms defined as

$$\begin{aligned} Q_2(U, k) = & \int_0^M 2c^2(U'_1)^2 + 2s^2(U'_2)^2 + 2(cU'_1 + sU'_2)^2 \\ & + \int_0^M (sU'_1 + cU'_2)^2 + (U''_1 + c(cU_1 + sU_2))^2 + (U''_2 + s(cU_1 + sU_2))^2, \\ Q_1(U, k) = & \int_0^M |U'|^2 + (cU'_1 + sU'_2)^2 + |U''|^2 + (cU_1 + sU_2)^2, \end{aligned}$$

with $M = \varepsilon|k|$. Let us denote $q_i(k), i = 1, 2$ the symmetric matrix associated to the quadratic forms $Q_i(\cdot, k), i = 1, 2$. If C_k is the minimum value for which the inequality $Q_2(\tilde{U}, k) \geq C_k Q_1(\tilde{U}, k)$ for all \tilde{U} is satisfied, then $\det(q_2(k) - C_k q_1(k)) = 0$ and one is left with an eigenvalue problem. Let us study the roots of $P(X) = \det(q_2(k) - X q_1(k))$. First, an integration by parts in the quadratic form $Q_2(\cdot, k)$ yields $Q_2(U, k) = \tilde{Q}_2(U, k) + Q_1(U, k)$

$$\tilde{Q}_2(U, k) = 2(cU_1 + sU_2)(M)(cU'_1 + sU'_2)(M).$$

The quadratic form $\tilde{Q}_2(\cdot, k)$ is a product of two linear forms $\tilde{Q}_2(U, k) = (v_1(k), a)(v_2(k), a)$ with $(v_1(k), v_2(k)) \in \mathbb{R}^6$ and $a \in \mathbb{R}^6$ represents the coordinates of \tilde{U} in the basis of \mathcal{S} . Then the matrix $\tilde{q}_2(k)$ associated to $\tilde{Q}_2(\cdot, k)$ is given by $\tilde{q}_2(k)a = (v_1(k), a)v_2(k) + (v_2(k), a)v_1(k)$. As a consequence, for any $a \in (v_1(k), v_2(k))^\perp$, one has $\tilde{q}_2(k)a = q_1(k)a$. One can prove that $v_1(k), v_2(k)$ are not colinear so that $X = 1$ is a root of P with multiplicity 4. There remains two extremal values which correspond to minimal and maximal value

of C_k . We know that $C_k \in (0, 2)$: let us show that $X = 2$ is exactly the maximum value for C_k . It is sufficient to show that there is a divergence free vector fields (u, w) so that $J(u, w) = J(u, w)^T$ and $w(0) = 0$, $J(u, w)$ being the Jacobian matrix of u, w . This is done choosing a potential flow $(u, w) = (\nabla\psi, \partial_z\psi)$ so that $\psi_z(0) = 0$ and $\Delta\psi = 0$. Split the functions into their Fourier modes, one shows that it suffices to choose $\psi_k(z) = \alpha_k \cosh(|k|z)$ with $\alpha \in l^2(\mathbb{Z}^2)$ chosen so that the H^1 regularity of u, w is satisfied. As a result, the polynom P is factorized in the form

$$P(X) = \Gamma_k(X - 1)^4(X - 2)(X - \Lambda(k))$$

and there is a unique minimal value $\Lambda(k)$ which depends continuously of $M = \varepsilon|k|$ and $\sigma = c + is \in \mathbb{S}^1$. We analyse $\Lambda(k)$ when $M \sim \varepsilon$ and $M \rightarrow \infty$. We have calculated with a formal computation software the expansion of P with respect to M at $M = 0$:

$$P(X) = M^4 \frac{8}{9} (Z - 1)^5 (Z - 2) + \mathcal{O}(M^6),$$

so that $\Lambda(k) \sim 1$ as $\varepsilon \rightarrow 0$ and k lies in a fixed compact set. Next we deal with the limit $M \rightarrow \infty$: Assume that there is U_M so that $MQ_1(U_M, k) = 1$ and $MQ_2(U_M, k) \rightarrow 0$. This is equivalent to the existence of \tilde{U}_M so that

$$\int_0^1 M^2 (c\tilde{U}_{1,M} + s\tilde{U}_{2,M})^2 + \frac{|\tilde{U}_M''|^2}{M^2} + |\tilde{U}_M'|^2 + (c\tilde{U}_{1,M}' + s\tilde{U}_{2,M}')^2 = 1,$$

$$\lim_{M \rightarrow \infty} (c\tilde{U}_{1,M}(1) + s\tilde{U}_{2,M}(1)) \frac{c\tilde{U}_{1,M}' + s\tilde{U}_{2,M}'}{M} = -1.$$

Denote $f_M = c\tilde{U}_{1,M} + s\tilde{U}_{2,M}$: it is easily proved that $f_M(1)$ and $f_M'(1)$ are estimated as

$$f_M(1)^2 = 2 \int_0^1 f_M(z) f_M'(z) dz \leq \frac{2}{M},$$

$$\frac{1}{M^2} f_M'(1)^2 \leq \frac{1}{M^2} \int_0^1 (f_M'(z))^2 dz + \frac{2}{M^2} \int_0^1 \int_z^1 |f_M'(s) f_M''(s)| ds dz \leq \frac{3}{M}.$$

This yields the contradiction with $\lim_{M \rightarrow \infty} M^{-1} f_M(1) f_M'(1) = -1$ and concludes the proof of the Korn's inequality. \square

3.2 Well posedness of linearized Navier-Stokes equations

3.2.1 A special divergence free linear problem

In this section, we consider the special linear problem obtained from (15-18), assuming $h_0 = 1$, $u_0 = 0$: the more general case $|h_0 - 1| \ll 1$ and $|u_0| \ll 1$

will be treated in the next section through a perturbation argument. In the case considered here, $A_0 = Id$ and the linear problem is

$$\partial_t \bar{u} + \frac{\nabla \bar{p}}{\text{Re}} = \frac{1}{\text{Re}} \text{div} D(\bar{u}) + f, \quad \text{div} u = 0, \quad (22)$$

with the boundary conditions at the bottom

$$\bar{u}_V|_{z_0=0} = 0, \quad \partial_{z_0} \bar{u}_H|_{z_0=0} = \varepsilon \bar{\gamma} \bar{u}_H|_{z_0=0} + g_1, \quad (23)$$

and at the free surface

$$(\partial_{z_0} \bar{u}_H + \nabla_{x_0} \bar{u}_V)|_{z_0=\varepsilon} = g_{2,1}, \quad (2\partial_{z_0} \bar{u}_V - \bar{p})|_{z_0=\varepsilon} = g_{2,2}. \quad (24)$$

We prove the existence of a weak solution. The weak formulation of (22,23,24) is written as

$$\begin{aligned} \partial_t \int_{\Omega} (\bar{u}, \phi) + \frac{1}{2 \text{Re}} \int_{\Omega} D(\bar{u}) : D(\phi) + \frac{\varepsilon \bar{\gamma}}{\text{Re}} \int_{\mathbb{X}} (\bar{u}_H|_{z_0=0}, \phi_H|_{z_0=0}) = \\ \int_{\Omega} (f, \phi) + \frac{1}{\text{Re}} \int_{\mathbb{X}} (g_2, \phi|_{z_0=\varepsilon}) - (g_1, \phi_H|_{z_0=0}), \end{aligned} \quad (25)$$

for all $\phi \in (H^1(\Omega))^3$ so that $\text{div} \phi = 0$ and $\phi_V|_{z_0=0} = 0$. In what follows, we prove the existence (and uniqueness) of a solution of (25).

Proposition 3 *Let $f \in L^2((0, T), (H^1(\Omega))')$ and $g \in L^2((0, T), H^{-\frac{1}{2}}(\mathbb{T}))$. For any $u_0 \in L^2_{div}(\Omega)$, there is a unique weak solution $u \in L^2((0, T), H^1(\Omega))$ of (25) and $\partial_t u \in L^2((0, T), (H^1(\Omega))')$. Moreover, it satisfies the energy estimates*

$$\begin{aligned} \max_{(0,t)} \|u\|_0^2 + C \int_0^t \|u\|_1^2 \leq \|u_0\|_0^2 + \\ C \int_0^t \|f\|_{(H^1(\Omega))'}^2 + |g_{2,2}|_{-\frac{1}{2}}^2 + \frac{1}{\varepsilon} (|g_{2,1}|_{-\frac{1}{2}}^2 + |g_1|_{-\frac{1}{2}}^2). \end{aligned} \quad (26)$$

Proof. The proof follows from a classical Galerkin approximation process: we only prove here the energy estimate. Substituting $\phi = u$ in the weak formulation (25), one obtains

$$\begin{aligned} \frac{\partial \|u\|_0^2}{\partial t} + \frac{1}{2} \int_{\Omega} D(u) : D(u) + \varepsilon \bar{\gamma} \int_{\mathbb{T}} |u_H|_{z=0}|_0^2 = \\ (f, u)_{(H^1)', H^1} + (g_2, u|_{z=\varepsilon})_{H^{-\frac{1}{2}}, H^{\frac{1}{2}}} - (g_1, u_H|_{z=0})_{H^{-\frac{1}{2}}, H^{\frac{1}{2}}}. \end{aligned} \quad (27)$$

Using the Korn's inequality on the functional space $(H_{div}^1(\Omega))^3$ so that $u_V|_{z=0} = 0$, one proves that there exists a constant $C > 0$ independent of ε so that

$$\|u\|_1^2 \leq C \left(\int_{\Omega} D(u) : D(u) + \varepsilon \bar{\gamma} \int_{\mathbb{T}} |u_H|_{z=0}|^2 \right).$$

Moreover, using lemma 2, one has the trace estimates

$$|u_V|_{z=\varepsilon}|_{\frac{1}{2}} \leq C \|u_V\|_1, \quad \sqrt{\varepsilon} |u_H|_{\frac{1}{2}} \leq C \|u_H\|_1.$$

Using a Young's inequality, one can prove that for any $\eta > 0$, there exists $c(\eta) > 0$ so that

$$\begin{aligned} |(g_{2,1}, u_H|_{z=\varepsilon}) - (g_1, u_H|_{z=0})_{H^{-\frac{1}{2}}, H^{\frac{1}{2}}}| &\leq \frac{c(\eta)}{\varepsilon} (|g_{2,1}|_{-\frac{1}{2}}^2 + |g_1|_{-\frac{1}{2}}^2) + \eta \|u_H\|_1^2, \\ |(g_{2,2}, u_V)|_{H^{-\frac{1}{2}}, H^{\frac{1}{2}}}| &\leq c(\eta) |g_{2,2}|_{-\frac{1}{2}}^2 + \eta \|u_V\|_1^2 \\ |(f, u)|_{(H^1)', H^1} &\leq c(\eta) \|f\|_{(H^1)'}^2 + \eta \|u\|_1^2. \end{aligned}$$

Inserting this inequality into (27) gives the estimate

$$\frac{\partial}{\partial t} \|u\|_0^2 + C \|u\|_1^2 \leq C \left(\|f\|_{(H^1)'}^2 + |g_{2,1}|_{-\frac{1}{2}}^2 + \frac{1}{\varepsilon} (|g_{2,1}|^2 + |g_1|^2) \right).$$

Then integrating this equation on the interval $(0, t)$ yields (26). From the weak formulation, one easily proves

$$\int_0^t \|\partial_t u\|_{(H^1)'}^2 \leq \|u_0\|_2^2 + \int_0^t \|f\|_{(H^1)'}^2 + |g_{2,1}|_{-\frac{1}{2}}^2 + \frac{1}{\varepsilon} (|g_{2,1}|_{-\frac{1}{2}}^2 + |g_1|_{-\frac{1}{2}}^2)$$

Using an interpolation argument, one has $u \in C((0, t), L_{div}^2)$ and the uniqueness follows, letting $f = g_i = 0$ and $u_0 = 0$ into (26). This completes the proof of the proposition. \square

This is a standard argument that there exists a pressure $p \in L^2((0, t), L^2)$ so that $\int_{\Omega} p(\cdot, t) = 0$ and

$$\begin{aligned} &\int_{\Omega} (u_t, \phi) + D(u) : D(\phi) - \int_{\Omega} p \operatorname{div} \phi + \varepsilon \bar{\gamma} \int_{\mathbb{X}} u_H|_{z=0} \phi_H|_{z=0} \\ &= \int_{\Omega} (f, \phi) + \int_{\mathbb{X}} (g_2, \phi|_{z=\varepsilon}) - (g_1, \phi_H|_{z=0}) \quad \forall \phi \in H^1, \phi_V|_{z=0} = 0. \end{aligned} \quad (28)$$

Moreover, the pressure term p satisfies

$$\int_0^t \|p\|_0^2 \leq C \left(\|u_0\|_0 + \int_0^t \|f\|_{(H^1)'}^2 + |g_{2,2}|_{-\frac{1}{2}}^2 + \frac{1}{\varepsilon} (|g_{2,1}|_{-\frac{1}{2}}^2 + |g_1|_{-\frac{1}{2}}^2) \right),$$

where C is a constant independent of ε . Next, we prove a result on the regularity of the weak solution.

Proposition 4 Let $f \in L^2((0, T), H^1)$ such that $\partial_t f \in L^2((0, T), (H^1)')$ and $g_i \in L^2((0, T), H^{\frac{3}{2}})$, $\partial_t g_i \in L^2((0, T), H^{-\frac{1}{2}})$ which satisfies the compatibility conditions

$$\begin{aligned} (\partial_{z_0} u_H + \nabla_{x_0} u_V)_{z_0=\varepsilon}(t=0, \cdot) &= g_{2,1}(0, \cdot), \\ (\partial_{z_0} u_H|_{z_0=0} - \varepsilon \bar{\gamma} u_H|_{z_0=0})(t=0, \cdot) &= g_1(0, \cdot). \end{aligned}$$

and $\operatorname{div} u_0 = 0$, $u_0 \in H^2$. Then the weak solution u verifies

$$\begin{aligned} \max_{(0,t)} (\|u\|_2^2 + \|u_t\|_0^2) + \int_0^t \|u(s, \cdot)\|_3^2 + \|u_t(s, \cdot)\|_1^2 ds \leq \\ \|u_0\|_2^2 + \|f(0, \cdot)\|_0^2 + \|g_i(0, \cdot)\|_{\frac{1}{2}}^2 \\ + \int_0^t \|f\|_1^2 + \|f_t\|_{(H^1)'}^2 + |g_{2,2}|_{\frac{3}{2}}^2 + |g_{2,2,t}|_{-\frac{1}{2}}^2 + \frac{1}{\varepsilon} (|g_{i,1}|_{\frac{3}{2}}^2 + |g_{i,1,t}|_{-\frac{1}{2}}^2) \end{aligned}$$

Proof. Let us first prove that there exists an “initial” pressure $p_0 \in H^1(\Omega)$, satisfying the mixed Dirichlet/Neuman problem

$$\begin{aligned} -\Delta p_0 &= -\operatorname{div}(\Delta u_0 + f(0)), \\ p_0|_{z_0=\varepsilon} &= 2\partial_{z_0} u_{0,V}|_{z_0=\varepsilon} - g_{2,2}(0), \\ \partial_{z_0} p_0|_{z_0=0} &= \varepsilon \bar{\gamma} \operatorname{div}_{x_0} u_{0,H}|_{z=0} + \operatorname{div}_{x_0} g_1(0) + f_V(0). \end{aligned}$$

This problem is well posed since $f(0) \in L^2$ and $g_i(0) \in H^{\frac{1}{2}}$ with a compactness argument respectively on $f \in L^2(0, t; H^1) \cap H^1((0, t); (H^1)')$ and $g \in L^2((0, t); H^{\frac{3}{2}}) \cap H^1(0, t; H^{-\frac{1}{2}})$. We split this problem into three parts: first, it is easily proved using a Lax Milgram argument that there exists a unique solution $p_{0,1} \in H^1(\Omega)$ so that

$$-\Delta p_{0,1} = -\operatorname{div}(\Delta u_0 + f(0)), \quad p_{0,1}|_{z_0=\varepsilon} = \partial_z p_{0,1}|_{z_0=0} = 0.$$

This is a straightforward computation (with Fourier series) to prove that there exists C , independent of ε , such that

$$\|p_{0,1}\|_1 \leq C(\|f(0)\|_0 + \|u_0\|_2).$$

Next we consider the problem

$$\Delta p_{0,2} = 0, \quad p_{0,2}|_{z=\varepsilon} = h, \quad \partial_z p_{0,2}|_{z=0} = 0,$$

with $h = 2\partial_z u_{0,V}|_{z_0=\varepsilon} - g_{2,2}(0) \in H^{\frac{1}{2}}$. Using Fourier series, this is equivalent to find $p_{0,2}^k$ so that

$$-\frac{d^2}{dz^2} p_{0,2}^k + |k|^2 p_{0,2}^k = 0, \quad p_{0,2}^k|_{z=\varepsilon} = h_k, \quad \frac{d}{dz} p_{0,2}^k|_{z=0} = 0.$$

The solution is easily calculated and satisfies

$$p_k(z) = \frac{\cosh|k|z}{\cosh(|k|\varepsilon)} h_k, \quad \left\| \frac{d}{dz} p_{0,2}^k \right\|^2 + |k|^2 \|p_{0,2}^k\| \leq C(|k|\varepsilon) |k| h_k^2,$$

with $C(y) = \frac{1}{\cosh^2 y} \int_0^y \cosh(2x) dx \leq C_1, \forall y \geq 0$ so that $\|p_{0,2}\|_1 \leq C_1 |h|_{\frac{1}{2}}$.

Similarly, the solution $p_{0,3}$ of

$$\Delta p_{0,3} = 0, p_{0,3}|_{z_0=\varepsilon} = 0, \partial_z p_{0,3}|_{z=0} = \varepsilon \bar{\gamma} \operatorname{div}_{x_0} u_{0,H}|_{z_0=0} + \operatorname{div}_{x_0} g_1(0) + f_V(0)|_{z_0=0},$$

satisfies $\|p_{0,3}\|_1 \leq C_1 (\|u_0\|_1 + |g_1(0)|_{\frac{1}{2}} + \|f_2(0)\|_0)$. Let us define an initial condition $\tilde{u}_0 \in L^2(\Omega)$ (that plays the role of $\partial_t u(0)$) as

$$\tilde{u}_0 = -\nabla p_0 + \Delta u_0 + f(0).$$

This satisfies the following estimate

$$\|\tilde{u}_0\|_0 + \|p_0\|_1 \leq C_1 \left(\|u_0\|_2 + \|f(0)\|_0 + |g_1(0), g_2(0)|_{\frac{1}{2}} \right),$$

with C_1 independent of ε . We apply the previous proposition with this initial condition that lies in the space of initial conditions of weak solutions $\{u \in L^2 / \operatorname{div} u = 0, u_V|_{z_0=0} = 0\}$ and f, g are replaced by (f_t, g_t) . There exists a unique weak solution $\tilde{u} \in L^2(0, t; H^1(\Omega))$ associated to this problem and that satisfies the energy estimate

$$\begin{aligned} \max_{(0,t)} \|\tilde{u}\|_0 + \int_0^t \|\tilde{u}\|_1 \leq C_1 \left((\|u_0\|_2^2 + \|f(0)\|_0^2 + |(g_1, g_2)(0)|_{\frac{1}{2}}^2 \right. \\ \left. + \int_0^t \|f_t\|_{(H^1)'}^2 + |g_{2,2,t}|_{-\frac{1}{2}}^2 + \frac{1}{\varepsilon} (|g_{1,t}|_{-\frac{1}{2}} + |g_{2,1,t}|_{-\frac{1}{2}}^2) \right). \quad (29) \end{aligned}$$

Associated to this weak solution is the pressure pressure term $\tilde{p} \in L^2$ (which plays the role of p_t) that satisfies the estimate $L^2(0, t; L^2)$

$$\begin{aligned} \int_0^t \|\tilde{p}\|_{L^2}^2 \leq C_1 \left((\|u_0\|_2^2 + \|f(0)\|_0^2 + |g_1(0), g_2(0)|_{\frac{1}{2}}^2 \right. \\ \left. + \int_0^t \|f_t\|_{(H^1)'}^2 + |g_{2,2,t}|_{-\frac{1}{2}}^2 + \frac{1}{\varepsilon} (|g_{1,t}|_{-\frac{1}{2}} + |g_{2,1,t}|_{-\frac{1}{2}}^2) \right). \quad (30) \end{aligned}$$

Then, let us consider $u_1 = u_0 + \int_0^t \tilde{u}$, $p_1 = p_0 + \int_0^t \tilde{p}$. Integrating on the time intervals $(0, t)$ the weak formulation on \tilde{u} and using the definition of p_0 , one can prove that u_1 is a weak solution of the initial problem with

$u_1(t=0) = u_0$. Then one finds that $u = u_1$, $p = p_1$ and $u_t = \tilde{u}$, $p_t = \tilde{p}$ satisfies the estimate (29). This concludes the proof of the proposition. \square

Finally, we check the space regularity of the weak solution. One can prove the proposition

Proposition 5 *Let $f \in L^2(0, T; H^1) \cap H^1(0, T; (H^1)')$,*

$$g_i \in L^2(0, T; H^{\frac{3}{2}}) \cap H^1(0, T; H^{-\frac{1}{2}})$$

and $u_0 \in H^2$ satisfying the compatibility conditions

$$\partial_{z_0} u_H + \nabla_{x_0} u_V|_{z_0=\varepsilon}(t=0) = g_{2,1}(0), \quad (\partial_{z_0} u_H - \varepsilon \bar{\gamma} u_H)|_{z_0=0}(t=0) = g_1(0).$$

Then there exists a constant C independent of ε such that the weak solution $u \in L^2(0, T; H^3(\Omega)) \cap H^1(0, T; H^1(\Omega))$, $p \in L^2(0, T; H^2(\Omega)) \cap H^1(0, T; L^2(\Omega))$ satisfies

$$\begin{aligned} & \max_{(0,T)} \|u\|_2^2 + \int_0^T \|u\|_3^2 + \|p\|_2^2 + \|u_t\|_1^2 + \|p_t\|_0^2 \\ & \leq C_1 \left(\|u_0\|_2^2 + \|f(0)\|_0^2 + |(g_1(0), g_2(0))|_{\frac{1}{2}}^2 \right. \\ & \quad \left. + \int_0^T \|f\|_1^2 + \|f_t\|_{(H^1)'}^2 + |g_{2,2}|_{\frac{3}{2}}^2 + |g_{2,2,t}|_{-\frac{1}{2}}^2 \right. \\ & \quad \left. + \frac{1}{\varepsilon} \int_0^T (|g_{1,1}|_{\frac{3}{2}}^2 + |g_{1,t}|_{-\frac{1}{2}}^2) \right). \end{aligned} \quad (31)$$

3.2.2 The full linear problem

In this part, we consider the full problem with a divergence term and $h_0 \neq 1$. Let us first deal with the case $h_0 = 1$, $u_0 = 0$ (steady state) and consider the linear problem (22,23,24) with the divergence free condition replaced by

$$\operatorname{div} u = h, \quad (32)$$

with $h \in L^2(0, T; H^2) \cap H^1(0, T; L^2)$. In order to prove the well posedness of (22,23,24) with the divergence condition (32), we will apply proposition 4: for that purpose, we introduce $u_h = \nabla \phi$ so that

$$\Delta \phi = h, \quad \frac{\partial \phi}{\partial z}|_{z=0} = \frac{\partial \phi}{\partial z}|_{z=\varepsilon} = 0.$$

One can prove, with Fourier series, that there exists a constant $C > 0$ independent of ε such that

$$\|\phi\|_4 + \|\phi_t\|_2 \leq C(\|h\|_2 + \|h_t\|_0).$$

Next, we consider \tilde{f}, \tilde{g}_i such that $\tilde{f} = f + \Delta u_h - u_{h,t}$ and

$$\tilde{g}_1 = g_1 + \varepsilon \bar{\gamma} u_h|_{z=0}, \quad \tilde{g}_{1,2} = g_{1,2} - 2\partial_{xz}\phi|_{z=\varepsilon}, \quad \tilde{g}_{2,2} = g_{2,2} + 2\partial_{zz}\phi.$$

This is a straightforward computation to prove that there exists a constant C independent of ε such that

$$\|\tilde{f}\|_1 \leq C(\|f\|_1 + \|h\|_2 + \|h_t\|_0),$$

$$\begin{aligned} |\tilde{g}_{1,2}|_{\frac{3}{2}} + |\tilde{g}_{1,t}|_{-\frac{1}{2}} &\leq C(|g_{1,2}|_{\frac{3}{2}} + |g_{1,t}|_{-\frac{1}{2}} + \|h\|_2 + \|h_t\|_0), \\ |\tilde{g}_{2,2}|_{\frac{3}{2}} + |\tilde{g}_{2,t}|_{-\frac{1}{2}} &\leq C\left(|g_{2,2}|_{\frac{3}{2}} + |g_{2,t}|_{-\frac{1}{2}} + \frac{1}{\sqrt{\varepsilon}}(\|h\|_2 + \|h_t\|_0)\right). \end{aligned}$$

The estimate on $\|\tilde{f}_t\|_{(H^1)'}$ is obtained using the linear Navier-Stokes equations: more precisely, one has $\tilde{f}_t = f_t + \Delta u_{h,t} - u_{h,tt}$ and

$$\operatorname{div} u_{h,tt} = h_{tt} = \operatorname{div} u_{tt} = \operatorname{div}\left(f_t + \frac{1}{Re} \operatorname{div}(D(u_t)) - \frac{1}{Re} \nabla p_t\right).$$

Then, one obtains the estimate of $u_{h,tt}$ (and thus \tilde{f}_t) in $L^2((0, T), (H^1)')$ norm, using elliptic regularity. Next, we apply proposition 4 on the divergence free problem for the fluid velocity $u - u_h$ and pressure p . Finally, we have proved the well-posedness of the linear problem (22,23,24) with (32) and the solution (u, p) satisfies the energy estimates

$$\begin{aligned} \max_{(0, T)} (\|u\|_2^2 + \|u_t\|_0^2) + \int_0^T (\|u\|_3^2 + \|u_t\|_1^2 + \|p\|_2^2 + \|p_t\|_0^2) \\ \leq C\left(\|u_0\|_2 + \|f(0)\|_0^2 + |(g_1, g_2)(0)|_{\frac{1}{2}}^2\right. \\ \left.+ \int_0^T (\|f\|_1^2 + \|f_t\|_{(H^1)'}^2 + |g_{2,2}|_{\frac{3}{2}}^2 + |g_{2,2,t}|_{-\frac{1}{2}}^2\right. \\ \left.+ \frac{1}{\varepsilon} \int_0^T (\|h\|_2^2 + \|h_t\|_0^2 + |g_{1,2}|_{\frac{3}{2}}^2 + |g_{1,t}|_{-\frac{1}{2}}^2)\right). \quad (33) \end{aligned}$$

Let us now deal with the full problem (15,16,17,18). For the sake of simplicity, we assume $\sigma = 0$, the inhomogeneous case $\sigma \neq 0$ is treated as previously. In order to deal with the more classical divergence free condition, we introduce

the Jacobian decomposition of the fluid velocity $\bar{u} = A_0 \tilde{u}$. The linear problem then reads (we also drop a linear term in the form $L_0(u, \nabla u, x, t)$)

$$A_0^T \partial_t (A_0 \tilde{u}) + \frac{1}{\text{Re}} \nabla p = \frac{1}{\text{Re}} \text{div}(\tilde{\mathcal{P}}_0) + f, \quad (34)$$

with $\tilde{\mathcal{P}}_0 = A_0^T A_0 \nabla(\tilde{u}) A_0^{-1} A_0^{-T} + (\nabla \tilde{u})^T$. The divergence condition reads

$$\text{div} \tilde{u} = 0. \quad (35)$$

The boundary conditions at the bottom are written as

$$\tilde{u}_V|_{z_0=0} = 0, \quad h_0^{-1} \partial_{z_0} \tilde{u}_H|_{z_0=0} = \varepsilon \bar{\gamma} \tilde{u}_H|_{z_0=0} + g_1, \quad (36)$$

whereas the boundary conditions at the free surface read

$$\left(\tilde{\mathcal{P}}_0 - p \text{Id} \right)|_{z_0=\varepsilon} e_3 = g_2. \quad (37)$$

Similarly to the special linear problem treated in the previous section, one can obtain a weak formulation of (34-37):

$$\begin{aligned} & \partial_t \int_{\Omega} (A_0 \tilde{u}, A_0 \phi) + \frac{1}{\text{Re}} \int_{\Omega} \nabla \phi : \tilde{\mathcal{P}}_0 + \frac{\varepsilon \bar{\gamma}}{\text{Re}} \int_{\mathbb{X}} h_0^{-1} \left(\frac{\partial X_0}{\partial x_0} \tilde{u}_H|_{z_0=0}, \frac{\partial X_0}{\partial x_0} \phi_H|_{z_0=0} \right) \\ &= \int_{\Omega} (A_{0,t} \phi, A_0 \tilde{u}) + (f, \phi) + \frac{1}{\text{Re}} \int_{\mathbb{X}} (g_2, \phi_H|_{z_0=\varepsilon}) - h_0^{-1} \left(\frac{\partial X_0}{\partial x_0} g_1, \frac{\partial X_0}{\partial x_0} \phi_H|_{z=0} \right). \end{aligned}$$

for all $\phi \in H^1(\Omega)$ so that $\text{div} \phi = 0$ and $\phi_V|_{z_0=0}$. One prove the following theorem

Theorem 3 *There is $\eta_1 > 0, \eta_2 > 0$ such that for any $(h_0, u_0) \in W^{2,\infty}$ be a smooth solution of (6) which first satisfies $\|h_0(t=0) - 1\|_{\infty} < \eta_1$ and either*

- *(small data, large time) $\forall t \in (0, T), \|h_0(t, \cdot) - 1, u_0(t, \cdot)\|_{\infty} \leq \eta_2$ or,*
- *(large data, small time) $T \sup_{t \in (0, T)} \|h_0(t, \cdot) - 1, u_0(t, \cdot)\|_{W^{2,\infty}} \leq \eta_2$*

Then for all source terms $f, g_i, i = 1, 2$ and h so that

$$\begin{aligned} & f \in L^2((0, T); H^1(\Omega)) \cap H^1((0, T); (H^1)'(\Omega)), \\ & g_i \in L^2((0, T); H^{\frac{3}{2}}(\mathbb{X})) \cap H^1((0, T); H^{-\frac{1}{2}}(\mathbb{X})), \\ & h \in L^2((0, T); H^2(\Omega)) \cap H^1((0, T); L^2(\Omega)). \end{aligned}$$

There exists a unique weak solution (u, p) of (22,23,24,32) and this solution satisfies energy estimates (33) .

Proof. The proof of this theorem follows the one of proposition 4: one first constructs an approximate solution using a Galerkin method and then prove uniform estimates on this solution. The only difference here is that we have to prove the Korn's inequality when $(h_0, u_0) \approx (1, 0)$. Let us first assume that $h_0(t = 0) = 1$. In the case of small data, one can write that

$$\int_{\Omega} \nabla \tilde{u} : \tilde{\mathcal{P}}_0 = \frac{1}{2} \int_{\Omega} D(\tilde{u}) : D(\tilde{u}) + \mathcal{O}(\eta_2) \|\nabla \tilde{u}\|^2, \\ \int_{\mathbb{X}} h_0^{-1} \left| \frac{\partial X_0}{\partial x_0} \tilde{u}_H \Big|_{z_0=0} \right|^2 = (1 + \mathcal{O}(\eta_2)) \int_{\mathbb{W}} |\tilde{u}_H|_{z_0=0}|^2$$

Then, one can use the Korn's inequality proved in the case $h_0 = 1$ to show that for η_2 sufficiently small, there is $C > 0$ independent of ε so that

$$\int_{\Omega} \nabla \tilde{u} : \tilde{\mathcal{P}}_0 + \varepsilon \bar{\gamma} \int_{\mathbb{X}} h_0^{-1} \left| \frac{\partial X_0}{\partial x_0} \tilde{u}_H \Big|_{z_0=0} \right|^2 \geq C \|\nabla \tilde{u}\|^2.$$

From the Korn inequality, we deduce energy estimates on the approximate solutions: this proves the existence of a weak solution. The proof of regularity is similar to the case of $h_0 = 1, u_0 = 0$: indeed the main task was to prove a Korn's inequality in this setting, which is done mainly using the fact that the Lagrangian change of variable is close to identity uniformly with respect to ε . The case of large data and small time is simply treated by considering that the Lagrangian change of variable is close to identity for T small enough and one can prove a similar Korn's inequality in this case. The more general case $h_0(t = 0) \neq 1$ is treated by perturbation, simply through the change of variable $(x, z) \mapsto (x, zh_0(t = 0, x))$. \square

4 Well posedness and convergence of the full Navier-Stokes equations in Lagrangian form.

We write an iterative scheme to prove the well posedness of free surface Navier Stokes equations with Navier slip condition at the bottom together with the convergence of this set of equations to a viscous shallow water model. We will use Tychonov fixed point theorem, as in [10], with now the functional space

$$\mathbb{X} = \left\{ (u, p) \in L^2((0, T); H^3 \times H^2) / \|(u, p)\|_{\mathbb{X}} \leq C(\|u_0\|_2 + \varepsilon^{3/2}) \right\}, \\ \|(u, p)\|_{\mathbb{X}}^2 = \sup_{(0, T)} \|u\|_2^2 + \int_0^T \|u\|_3^2 + \|\partial_t u\|_1^2 + \|p\|_2^2 + \|\partial_t p\|_0^2.$$

The iterative scheme is written as

$$\begin{aligned} \partial_t u^{k+1} + A_0^{-1} u^{k+1} \cdot \nabla u_a(X_0) + \frac{A_0^{-T} \nabla p^{k+1}}{R_e} \\ = \frac{A_0^{-T}}{R_e} \left(\operatorname{div}(A_0^T \mathcal{P}_0^{k+1}) - \nabla A_0 : \mathcal{P}_0^{k+1} \right) + F^k, \\ \operatorname{div}(A_0^{-1} u^{k+1}) = H_k, \end{aligned}$$

with

$$H_k = \sigma_a(X_0, Z_0) + (\sigma_a(X^k, Z^k) - \sigma_a(X_0, Z_0)) + \operatorname{div}((A_0^{-1} - A_k^{-1})u^k)$$

and

$$\begin{aligned} F^k &= (A_0^{-1} - A_k^{-1})u^k \cdot \nabla u_a(X^k, Z^k) + A_0^{-1}u^k \cdot \nabla (u_a(X_0, Z_0) - u_a(X^k, Z^k)) \\ &+ \frac{1}{R_e} \left((A_0^{-T} - A_k^{-T}) \nabla p^k + (A_k^{-T} - A_0^{-T}) (\operatorname{div}(A_k^T \mathcal{P}^k) - \nabla A_k : \mathcal{P}^k) \right) \\ &+ \frac{1}{R_e} A_0^{-T} \left(\nabla A_0 : \mathcal{P}_0^k - \nabla A_k : \mathcal{P}^k + \operatorname{div}(A_k^T \mathcal{P}^k - A_0^T \mathcal{P}_0^k) \right) \\ &+ (f_a(X^k, Z^k) - f_a(X_0, Z_0)) + f_a(X_0, Z_0). \end{aligned}$$

This set of equation is completed with boundary conditions at the bottom:

$$u_V^{k+1}|_{z=0} = 0, \quad \frac{1}{h_0} \partial_z u_H^{k+1}|_{z=0} = \varepsilon \bar{\gamma} u_H^{k+1}|_{z=0} + G_k^1, \quad (38)$$

with $G_k^1 = \varepsilon \bar{\gamma} \left(\left(\frac{\det(\partial_x X_0)}{\det(\partial_x X^k)} - 1 \right) u_H^k \right)|_{z=0}$ and boundary conditions at the free surface

$$(A_0^T \mathcal{P}_0^{k+1} - p^{k+1} Id)|_{z=\varepsilon} e_3 = G_k^2, \quad (39)$$

where

$$G_k^2 = A_0^T \left((g_a(X^k, Z^k) - g_a(X_0, Z_0)) + g_a(X_0, Z_0) + (\mathcal{P}_0^k - \mathcal{P}^k) n + (\mathcal{P}_0^k - p^k Id)(n_0 - n) \right)|_{z=\varepsilon}$$

with

$$\begin{aligned} \mathcal{P}^k &= (\nabla u^k) A_k^{-1} A_k^{-T} + A_k^{-T} (\nabla u^k)^T A_k^{-T}, \\ \mathcal{P}_0^k &= (\nabla u^k) A_0^{-1} A_0^{-T} + A_0^{-T} (\nabla u^k)^T A_0^{-T}. \end{aligned}$$

The function u^{k+1} has the initial condition $u^{k+1}|_{t=0} = u_0$. Using the well posedness results obtained in the previous section, this set of equation has a unique solution and the weak continuity of the map is done. We first prove that (u^k, p^k) is bounded, uniformly on the time interval $(0, T)$ and uniformly with respect to ε , in a suitable norm provided that the initial condition is sufficiently small. The proof is done in 5 steps: estimates of the Lagrangian

coordinates (more precisely the deviation from the “shallow water” coordinates), the boundary terms, the divergence term and nonlinear terms in momentum equations.

Step 1: Estimates of Lagrangian coordinates. Let us recall that Lagrangian coordinates satisfy the set of equations

$$\begin{aligned}\frac{d}{dt}(X^k - X_0) &= u_{a,H}(X^k, Z^k) - u_{a,H}(X_0, Z_0) + u_H^k, \\ \frac{d}{dt}(Z^k - Z_0) &= u_{a,V}(X^k, Z^k) - u_{a,V}(X_0, Z_0) + u_V^k.\end{aligned}$$

Let us first compute the basic L^2 estimate

$$\begin{aligned}\frac{d}{dt}(\|X^k - X_0\|_0^2 + \|Z^k - Z_0\|_0^2) &\leq C(\|u_a(t, \cdot)\|_{W^{1,\infty}})(\|X^k - X_0\|_0^2 + \|Z^k - Z_0\|_0^2) \\ &\quad + \|u^k\|_0(\|X^k - X_0\|_0 + \|Z^k - Z_0\|_0).\end{aligned}$$

This is a straightforward consequence of Gronwall inequality that

$$\sup_{t \in (0, T)} (\|X^k - X_0\|_0^2 + \|Z^k - Z_0\|_0^2) \leq T e^{C(T+1)} \int_0^T \|u^k(t, \cdot)\|_0^2 dt,$$

with $C = C(\sup_{t \in (0, T)} (\|u_a\|_{W^{1,\infty}}))$. There is no difficulty to extend this estimate to higher order space derivatives:

$$\sup_{t \in (0, T)} (\|X^k - X_0\|_s^2 + \|Z^k - Z_0\|_s^2) \leq T e^{C(T+1)} \int_0^T \|u^k(t, \cdot)\|_s^2 dt,$$

with $C = C(\sup_{t \in (0, T)} (\|u_a\|_{W^{1,\infty}}))$. Next, we compute boundary terms at the bottom:

Step 2: Estimates of nonlinear terms at the bottom. In order to estimate the boundary term at the bottom, we will use lemma 2 and 3. The nonlinear term at the bottom reads

$$G_1^k = \varepsilon \bar{\gamma} u_H^k|_{z=0} \left(\frac{\det(\partial_x X_0)}{\det(\partial_x X^k)} - 1 \right) |_{z=0} = \varepsilon \bar{\gamma} u_H^k|_{z=0} \mathcal{F}(\partial_x X^k - \partial_x X_0).$$

We first compute an estimate for $|g_1^k|_{\frac{3}{2}}$. Let us fix $t_0 > 1$ so as $H^{t_0}(\mathbb{T}^2) \subset$

$L^\infty(\mathbb{T}^2)$. We then have

$$\begin{aligned}
|G_1^k|_{\frac{3}{2}} &\leq \varepsilon \bar{\gamma} \left(|u_H^k|_{z=0|_{t_0}} |\mathcal{F}(\partial_x X^k - \partial_x X_0)|_{\frac{3}{2}} + |u_H^k|_{z=0|_{\frac{3}{2}}} |\mathcal{F}(\partial_x X^k - \partial_x X_0)|_{t_0} \right) \\
&\leq \varepsilon \bar{\gamma} C (|\partial_x X^k - \partial_x X_0|_\infty) (|u_H^k|_{z=0|_{t_0}} |\partial_x X^k - \partial_x X_0|_{t_0} \\
&\quad + |u_H^k|_{z=0|_{\frac{3}{2}}} |\partial_x X^k - \partial_x X_0|_{t_0}) \\
&\leq \varepsilon \bar{\gamma} C (|X^k - X_0|_{t_0+1}) \left(\frac{\|u_H^k\|_{t_0+\frac{1}{2}} \|X^k - X_0\|_3}{\sqrt{\varepsilon}} + \frac{\|u_H^k\|_2 \|X^k - X_0\|_{t_0+\frac{3}{2}}}{\sqrt{\varepsilon}} \right) \\
&\leq \bar{\gamma} C \left(\frac{\|X^k - X_0\|_{t_0+\frac{3}{2}}}{\sqrt{\varepsilon}} \right) (\|u^k\|_{t_0+\frac{1}{2}} \|X^k - X_0\|_3 \\
&\quad + \|u^k\|_2 \|X^k - X_0\|_{t_0+\frac{3}{2}}). \tag{40}
\end{aligned}$$

Then it is easily seen, choosing $t_0 = \frac{3}{2}$, that

$$\int_0^T |G_1^k(t, \cdot)|_{\frac{3}{2}}^2 dt \leq \bar{\gamma} C \left(\sqrt{\frac{T}{\varepsilon}} \|u^k\|_{L^2((0,T),H^3)} \right) T \sup_{t \in (0,T)} (\|u^k(t, \cdot)\|_2^2) \int_0^T \|u^k(t, \cdot)\|_3^2 dt.$$

Using lemma 3, we next compute an estimate of $|\partial_t g_1^k|_{-\frac{1}{2}}$. One has

$$\partial_t G_1^k = \varepsilon \bar{\gamma} \left(\partial_t u_H^k|_{z=0} \mathcal{F}(\partial_x X^k - \partial_x X_0) + u_H|_{z=0} \partial_1 \mathcal{F}(\partial_x X^k - \partial_x X_0) (\partial_{xt}^2 (X^k - X_0)) \right)$$

Then, the following inequality holds:

$$\begin{aligned}
|\partial_t G_1^k|_{-\frac{1}{2}} &\leq \varepsilon \left(\frac{|\partial_t u^k|_{-\frac{1}{2}}}{\sqrt{\varepsilon}} \|\mathcal{F}(\partial_x X^k - \partial_x X_0)\|_2 \right. \\
&\quad \left. + |u_H|_{z=0} \mathcal{G}(\partial_x X^k - \partial_x X_0) (\partial_x u_H^k + \mathcal{F}_1(\partial_x X^k - \partial_x X_0))|_{-\frac{1}{2}} \right) \\
&\leq C \left(\sqrt{\frac{T}{\varepsilon}} \|u^k\|_{L^2((0,T),H^3)} \right) (\|u^k\|_1 + \|\partial_t u^k\|_1) \|X^k - X_0\|_3 \\
&\quad + \sqrt{\varepsilon} |\partial_x u^k|_{-\frac{1}{2}} \|\mathcal{F}(\partial_x X^k - \partial_x X_0) u_H\|_2 \tag{41} \\
&\leq C \left(\sqrt{\frac{T}{\varepsilon}} \|u^k\|_{L^2((0,T),H^3)} \right) \left((\|u^k\|_1 + \|\partial_t u^k\|_1) \|X^k - X_0\|_3 + \|u^k\|_2 \|X^k - X_0\|_3 \right).
\end{aligned}$$

So that we obtain the estimate

$$\int_0^T |\partial_t G_1^k|_{-\frac{1}{2}}^2 \leq C \left(\sqrt{\frac{T}{\varepsilon}} \|u^k\|_{L^2((0,T),H^3)} \right) T \left(\int_0^T \|\partial_t u^k\|_1^2 + \|u^k\|_3^2 \right)^2.$$

Next we compute interior estimates of “divergence” terms: this is done using

proposition 1 and tame estimates of lemma 1.

Step 3: Estimate of divergence terms. One part of the divergence term $H_1^k = \text{Tr}((A_k^{-1} - A_0^{-1})\nabla u^k)$ satisfies

$$\begin{aligned}\|H_1^k\|_2 &\leq \frac{1}{\sqrt{\varepsilon}}C(\|\partial_x X^k - \partial_x X_0\|_\infty)\|X^k - X_0\|_3\|\nabla u^k\|_2, \\ \|H_1^k\|_2 &\leq \frac{1}{\sqrt{\varepsilon}}C\left(\sqrt{\frac{T}{\varepsilon}}\|u^k\|_{L^2((0,T),H^3)}\right)\|X^k - X_0\|_3\|u^k\|_3.\end{aligned}$$

As a consequence, one has

$$\int_0^T \|H_1^k(t, \cdot)\|_2^2 dt \leq \frac{C\left(\sqrt{\frac{T}{\varepsilon}}\|u^k\|_{L^2((0,T),H^3)}\right)}{\varepsilon} \left(\int_0^T \|u^k\|_3^2\right)^2.$$

Next, we compute an estimate of $\|\partial_t H_1^k\|_0$. On the one hand, one has

$$\|\text{Tr}((A_k^{-1} - A_0^{-1})\nabla \partial_t u^k)\|_0 \leq \frac{1}{\sqrt{\varepsilon}}C(\|\partial_x X^k - \partial_x X_0\|_\infty)\|X^k - X_0\|_3\|\partial_t u^k\|_1,$$

and

$$\begin{aligned}\|\text{Tr}((\partial_t A_k - \partial_t A_0)\nabla u^k)\|_0 &\leq \|\nabla u^k\|_\infty\|C(\partial_{x,t} X^k - \partial_{x,t} X_0)\|_0, \\ \|\text{Tr}((\partial_t A_k - \partial_t A_0)\nabla u^k)\|_0 &\leq \frac{1}{\sqrt{\varepsilon}}\frac{\|u^k\|_3}{\sqrt{\varepsilon}}C(\|\partial_x X^k - \partial_x X_0\|_\infty \\ &\quad + \|\partial_x u^k\|_\infty)(\|X^k - X_0\|_1 + \|u^k\|_1).\end{aligned}$$

So that we obtain

$$\int_0^T \|\partial_t H_1^k\|_0^2 \leq \frac{T}{\varepsilon}C\left(\sqrt{\frac{T}{\varepsilon}}\|u^k\|_{L^2((0,T),H^3)}\right) \int_0^T \|\partial_t u^k\|_1^2 dt \int_0^T \|u^k(t, \cdot)\|_3^2 dt.$$

This concludes the estimate of nonlinear divergence terms.

Step 4: Estimates of "interior" terms. The interior terms are estimated with proposition 1 and lemma 1. We first compute an estimate of $\|(A_0^{-T} - A_k^{-T})\nabla p^k\|_1$. One has

$$\begin{aligned}\|(A_0^{-T} - A_k^{-T})\nabla p^k\|_1 &\leq \|A_0^{-T} - A_k^{-T}\|_\infty\|p^k\|_2 + \|\nabla(A_0^{-T} - A_k^{-T})\nabla p^k\|_0 \\ &\leq \|A_0^{-T} - A_k^{-T}\|_\infty\|p^k\|_2 + \varepsilon^{\frac{1}{6}}\frac{\|\nabla(A_0^{-T} - A_k^{-T})\|_1}{\sqrt{\varepsilon}}\frac{\|p^k\|_2}{\sqrt{\varepsilon}} \\ &\leq \frac{C(\|\partial_x X^k - \partial_x X_0\|_\infty)}{\sqrt{\varepsilon}}\|X^k - X_0\|_3\|p^k\|_2.\end{aligned}\tag{42}$$

As a consequence, we obtain the estimate

$$\int_0^T \|(A_0^{-T} - A_k^{-T})\nabla p^k\|_1^2 \leq \frac{T}{\varepsilon} C \left(\sqrt{\frac{T}{\varepsilon}} \|u^k\|_{L^2((0,T),H^3)} \right) \int_0^T \|u^k(t, \cdot)\|_3^2 dt \int_0^T \|p^k\|_2^2.$$

Similarly, the time derivative of this term satisfies

$$\begin{aligned} \int_0^T \|\partial_t((A_0^{-T} - A_k^{-T})\nabla p^k)\|_{(H^1)'}^2 & \\ & \leq \frac{T}{\varepsilon} C \left(\sqrt{\frac{T}{\varepsilon}} \|u^k\|_{L^2((0,T),H^3)} \right) \int_0^T \|u^k(t, \cdot)\|_3^2 dt \int_0^T \|\partial_t p^k(t, \cdot)\|_0^2 dt \\ & \quad + \frac{T}{\varepsilon} C \left(\sqrt{\frac{T}{\varepsilon}} \|u^k\|_{L^2((0,T),H^3)} \right) \int_0^T \|\partial_t u^k(t, \cdot)\|_1^2 dt \int_0^T \|p^k(t, \cdot)\|_2^2 dt. \end{aligned}$$

The other nonlinear term with high derivatives that we consider is

$$f^k = A_k^{-T} \operatorname{div}(A_k^T \mathcal{P}^k) - A_0^{-T} \operatorname{div}(A_0^T \mathcal{P}_0^k).$$

This is a lengthy but straightforward computation to show that

$$\begin{aligned} \int_0^T \|f^k\|_1^2 + \|\partial_t f^k\|_{(H^1)'}^2 & \\ & \leq \frac{T}{\varepsilon} C \left(\sqrt{\frac{T}{\varepsilon}} \|u^k\|_{L^2((0,T),H^3)} \right) \int_0^T \|u^k\|_3^2 \left(\int_0^T \|u^k\|_3^2 + \|\partial_t u^k\|_1^2 \right). \end{aligned}$$

Step 5: estimates of boundary term at the free surface. We have to deal with the nonlinear term

$$G_2^k = A_0^T \left((\mathcal{P}^k - \mathcal{P}_0^k) \mathbf{n} + (\mathcal{P}^k - p^k)(\mathbf{n}_0 - \mathbf{n}) \right) \Big|_{z=\varepsilon}.$$

On the one hand, one has

$$|(\mathcal{P}^k - \mathcal{P}_0^k) \mathbf{n}|_{z=\varepsilon} \Big|_{\frac{3}{2}} \leq C \left(\sqrt{\frac{T}{\varepsilon}} \|u^k\|_{L^2((0,T),H^3)} \right) |\mathcal{P}^k - \mathcal{P}_0^k|_{\frac{3}{2}}.$$

Moreover, this is a straightforward computation to show that

$$|\mathcal{P}^k - \mathcal{P}_0^k|_{\frac{3}{2}} \leq C \left(\frac{\|X^k - X_0\|_3}{\sqrt{\varepsilon}} \right) \frac{\|X^k - X_0\|_3 \|u^k\|_3}{\varepsilon}.$$

As a consequence, one obtains

$$\int_0^T |(\mathcal{P}^k - \mathcal{P}_0^k)n|_{z=\varepsilon}|_{\frac{3}{2}}^2 \leq \frac{1}{\varepsilon^2} C \left(\sqrt{\frac{T}{\varepsilon}} \|u^k\|_{L^2((0,T),H^3)} \right) \left(\int_0^T \|u^k\|_3^2 \right)^2.$$

For the term $(\mathcal{P}^k - p^k)(n - n_0)|_{z=\varepsilon}$, we first remark that $(n - n_0)|_{z=0} = 0$. As a consequence, we obtain the estimate

$$\begin{aligned} & \int_0^T |(\mathcal{P}^k - p^k)(n - n_0)|_{z=\varepsilon}|_{\frac{3}{2}}^2 \\ & \leq \frac{1}{\varepsilon} C \left(\sqrt{\frac{T}{\varepsilon}} \|u^k\|_{L^2((0,T),H^3)} \right) \int_0^T \|u^k\|_3^2 \left(\int_0^T \|u^k\|_3^2 + \|p^k\|_2^2 \right). \end{aligned}$$

Hence, using the estimates on the linear problem and using that remaining terms in the Navier-Stokes type system (7) satisfied by the high order shallow water approximation (u_a, p_a) satisfy

$$\begin{aligned} & \int_0^T \left(\|f_a(X_0, Z_0)\|_1^2 + \|\partial_t f_a(X_0, Z_0)\|_{(H^1)'}^2 + |g_{2,2}^a(X_0, Z_0)|_{\frac{3}{2}}^2 + |\partial_t g_{2,2}^a(X_0, Z_0)|_{-1/2}^2 \right) \\ & + \frac{1}{\varepsilon} \int_0^T \left(\|H_a(X_0, Z_0)\|_2^2 + \|\partial_t H_a(X_0, Z_0)\|_0^2 + |g_1^a(X_0, Z_0), g_{2,1}^a(X_0, Z_0)|_{\frac{3}{2}}^2 \right. \\ & \quad \left. + |g_{1,t}^a(X_0, Z_0), \partial_t g_{2,1}^a(X_0, Z_0)|_{-1/2}^2 \right) \leq C\varepsilon^4. \end{aligned}$$

We have then proved the following proposition.

Proposition 6 *Assume that $\|u_0\|_2 \leq \frac{C-1}{C^2} \varepsilon^{\frac{3}{2}}$, then for ε small enough, one has, for any $k \geq 1$*

$$\|(u^k, p^k)\|_X^2 \leq C(\|u_0\|_2^2 + \varepsilon^3).$$

At this step, there are two strategies: Either we use a Banach fixed point argument and in this case we need estimates for higher Sobolev norms that are obtained in a way similar to our previous calculation. In this case, we can prove that the sequence (u^k, p^k) is a Cauchy sequence respectively in the space $\mathcal{C}(0, T; H^{2+\eta}(\Omega))$ and $L^2((0, T), H^{2+\eta}(\Omega))$ with $0 < \eta < 1/2$, see [3] for further references. In this case, the result is not sharp. To obtain the well posedness in critical spaces, one has to use Tychonoff fixed point procedure as in [10]. There is no difficulty here to follow both strategies to obtain the well posedness result in our situation since the estimate in high norms is proved.

5 Conclusion

In this paper, we have proved the well-posedness of free surface Navier-Stokes equations with a slip condition at the bottom and convergence to a viscous shallow water model in the shallow water scaling $\varepsilon \rightarrow 0$, ε being the aspect ratio. In contrast to our mathematical justification of a shallow water model for a fluid flowing down an inclined plane, we dropped the tension surface effect (which is important in order to deal with realistic situations). This is due to the fact that we worked in the more suitable Lagrangian coordinates. There is no difficulty to extend this result for fluids flowing down an inclined plane when the slope θ is asymptotically small, $\theta \sim \varepsilon$ and when the uniform flow is *stable* (see [8] for more details).

When the uniform flow is unstable, it is known that instabilities so called roll-waves appear: these are periodic travelling waves. *Small amplitude* roll-waves are proved to exist both in viscous shallow water equations [15] and free surface Navier-Stokes equations [17]. However, the result of NISHIDA et al. is proved in the presence of *surface tension* whereas roll-waves are proved to exist in viscous shallow water equations without surface tension and are of large amplitude when the viscosity is small. In order to obtain a larger range of validity of shallow water equations, this would be of interest to prove the existence of roll-waves in Navier-Stokes equations without surface tension that are close to “shallow water roll-waves” in the shallow water scaling.

References

- [1] G. ALLAIN. Un problème de Navier-Stokes avec surface libre et tension superficielle. *Ann. Fac. Sci. Toulouse*, 7 (1985), 29–56.
- [2] G. ALLAIN. Small time existence for the Navier-Stokes equations with a free surface. *Appl. Math. Optim.*, 16 (1987), 37–50.
- [3] J.T. BEALE. The Initial Value Problem for the Navier-Stokes Equations with a Free Surface. *Comm. Pure and Appl. Maths*, 34 (1981), 359–392.
- [4] J.T. BEALE. Large Time Regularity of Viscous Surface Waves. *Arch. Rational Mech Anal*, 84 (1984), 307–352.
- [5] M. BOUTONET, L. CHUPIN, P. NOBLE, J.-P. VILA. Shallow water flows for arbitrary topography. *Communication in Mathematical Sciences* 6, 1, (2008) 531–536.

- [6] D. BRESCH. Shallow-Water Equations and Related Topics. Handbook of Differential Equations: Evolutionary Equations. Eds. C. Dafermos and M. Pokorný (2008).
- [7] D. BRESCH, B. DESJARDINS. Existence of global weak solutions for a 2D viscous shallow water equations and convergence to the quasi-geostrophic model. *Commun. Math. Phys.*, 238, 1-2, (2003), p. 211–223.
- [8] D. BRESCH, P. NOBLE. Mathematical justification of a shallow water model, *Methods and Appl. Anal.*, Volume 14, Number 2 (2007), 87–118.
- [9] I.D. CHUESHOV, G. RAUGEL, A.M. REKALO. Interface boundary value problem for the Navier-Stokes equations in thin two-layer domains. *J. Differential Equations*, 208 (2005), no. 2, 449–493.
- [10] D. COUTAND, S. SHKOLLER. Unique solvability of the free-boundary Navier-Stokes equations with surface tension. arXiv:math/0212116.
- [11] R. DANCHIN Well-Posedness in Critical Spaces for Barotropic Viscous Fluids with Truly Not Constant Density. *Comm. Partial Differential Equations* , 32 (2007), no. 7-9, 1373-1397.
- [12] J.-F. GERBEAU, B. PERTHAME, Derivation of shallow water equations and numerical validation *DCDS*. Ser. B, Vol. 1, Num. 1, (2001), 89–102.
- [13] D. HOFF. Uniqueness of weak solutions of the Navier-Stokes equations of multidimensional, compressible flow. *SIAM J. Math. Anal.* 37 (2006), no. 6, 1742–1760.
- [14] D. IFTIMIE, G. RAUGEL, G.R. SELL. Navier-Stokes equations in thin 3D domains with Navier boundary conditions. *Indiana Univ. Math. J.* 56 (2007), no. 3, 1083–1156.
- [15] J.H. MERKIN, D.J. NEEDHAM. An infinite period bifurcation arising in roll waves down an open inclined channel. *Proc. Roy. Soc. London Ser. A* 405 (1986), no. 1828, 103–116.
- [16] T. NISHIDA, Y. TERAMOTO, H.A. WIN. Navier-Stokes flows down an inclined plane: downward periodic motion. *J. Math. Kyoto Univ.* 33-3 (1993), 787–801.
- [17] T. NISHIDA, Y. TERAMOTO, H. YOSHIKI. Hopf bifurcation in viscous incompressible flow down an inclined plane. *J. Math. Fluid Mech.* 7 (2005), no. 1, 29–71.

- [18] A. ORON, S.H. DAVIS, S.G. BANKOFF. Long-scale evolution of thin liquid films. *Rev. Mod. Phys.*, 69 (1997), 931–980.
- [19] V.A. SOLONNIKOV. Solvability of a problem on the motion of a viscous incompressible fluid bounded by a free surface, *Izv. Akad. Nauk. SSSR Ser. Mat.* 41 (1977), 1388–1424 (In russian) ; English transl.: *Math. USSR-Izv* 11, 11 (1977), 1323–1358.
- [20] V.A. SOLONNIKOV. On an unsteady flow of a finite mass of a liquid bounded by a free surface. *Zap. Nauchn. Sem. LOMI*, 152 (1986), 137–157 (in Russian) ; English transl.: *J. Soviet Math.* 40 (1988), 672–686.
- [21] L. SUNDBYE. Existence for the Cauchy Problem for the Viscous Shallow Water Equations. *Rocky Mountain Journal of Mathematics*, 1998, 28 (3), 1135–1152.
- [22] A. TANI. Small time existence for the three dimensional incompressible Navier-Stokes equations with free surface. *Arch. Rational Mech. Anal.* (1996), 299–331.
- [23] A. TANI, N. TANAKA. Large time existence of surface waves in incompressible viscous fluids with or without surface tension. *Arch. Rational Mech. Anal.*, 130 (1995), 303–314.
- [24] R. TEMAM, M. ZIANE. Navier-Stokes equations in three-dimensional thin domains with various boundary conditions. *Adv. Differential Equations*, 1 (1996), no. 4, 499–546.
- [25] Y. TERAMOTO. The initial value problem for a viscous incompressible flow down an inclined plane. *Hiroshima Math. J.*, 15 (1985), 619–643.
- [26] Y. TERAMOTO. On the Navier-Stokes flow down an inclined plane. *J. Math. Kyoto Univ.*, 32 (1992), 593–619.
- [27] W. WANG, C.-J. XU. The Cauchy problem for viscous shallow water equations. *Rev. Mat. Iberoamericana*, 21, no. 1 (2005), 1–24.