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# Observations of GPS ionospheric scintillations over Wuhan during geomagnetic storms

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**Abstract.** During the two geomagnetic storms which occurred on 1 October 2002 and 22 January 2004, the strong ionospheric scintillations of the GPS L1 band were observed at Wuhan station ( $30.6^{\circ}$  N,  $114.4^{\circ}$  E,  $45.8^{\circ}$  Dip), which is situated near the northern crest of the equatorial ionosphere anomaly. We found that the intense scintillations were associated with the main phases of the storms and were co-located with the enhancement of the equatorial ionization anomaly (EIA); the co-existence of large- and small-scale irregularities at post-midnight was also found. The results may be relevant regarding the influence of the equatorial ionospheric eastward electric field during geomagnetic storms. On the other hand, GPS L1 band scintillations were not observed during the other two similar storms on 16 July 2003 and 20 November 2003. One of the reasons is probably that the sporadic E layer observed at the storms inhibited the generation of spread F by changing the Pedersen conductivity and suppressing the upward plasma drift.

**Keywords.** Ionosphere (Ionospheric irregularities; Mid-latitude ionosphere) – Magnetospheric physics (Storms and substorms)

## 1 Introduction

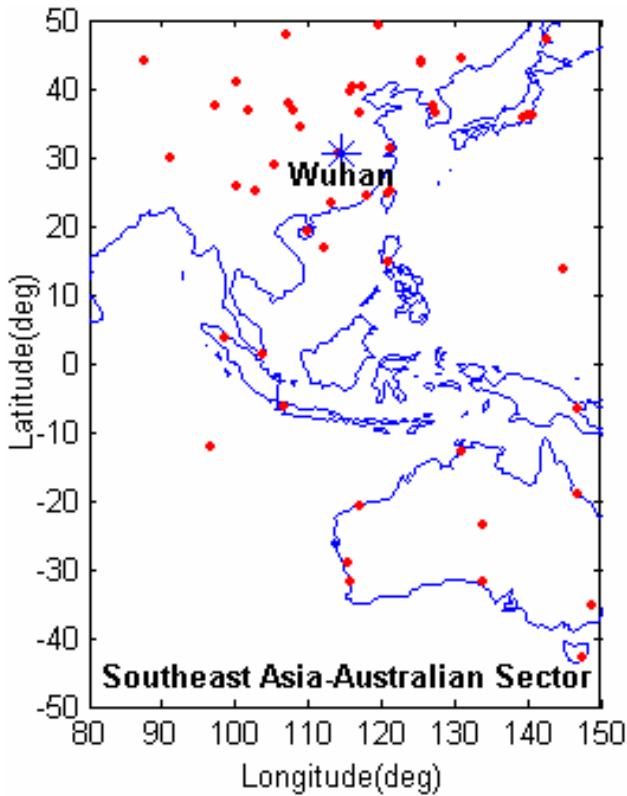
Sometimes when a radio signal acts on the disturbed ionosphere, the received signal will show rapid fluctuations in amplitude and phase which are not consistent with the source strength or modulation. This so-called scintillation is attributable to electron density irregularities primarily in the ionospheric E and F regions (Yeh and Liu, 1982). Scintillations observed at low latitudes are primarily controlled by the generation and growth of irregularities over the magnetic equator, where the effects of upward plasma  $E \times B$  drift

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and more effective ion recombination at lower altitudes result in a steep gradient of electron density on the bottom-side F region. When the altitude of the F region is high enough or the bottom-side background electron density gradients are large enough to overcome recombination effects, the Rayleigh-Taylor instability mechanism initiates a growth in plasma fluctuations. An upward moving bubble of depleted plasma is produced, which rises and eventually transforms itself into a plethora of smaller irregularities. By examining quantitatively the influence of the pre-midnight sporadic E (Es) layer on the daily and hourly development of equatorial spread F (ESF), and calculating changes in the flux tube-integrated Pedersen conductivity as it affects the growth rate of the Rayleigh-Taylor instability, Andrew et al. (2002) concluded that Es layers might change the Pedersen conductivity ratio and the growth rate of ESF. At auroral area, structured F-region density gradients are the primary instability source (Basu et al., 1993). Contrary to the high and low latitudes, the mid-latitude ionosphere is generally regarded as a less active scintillation environment, especially at L-band frequencies, and is considered to be absent of scintillations.

Recent studies indicate that the mid-latitude region is more complicated than previously thought. Foster et al. (2002, 2005) observed plumes of greatly enhanced total electron content (TEC) during the intense geomagnetic storms of 31 March 2001 and 30–31 October 2003. By observing GPS L1 amplitude scintillations in the mid-latitude region, Ledvina et al. (2002) concluded that this storm-enhanced density (SED) structure plays a major role in the formation of irregularities. Moreover, Basu et al. (2005) found two distinct classes of plasma processes that produce mid-latitude ionospheric irregularities associated with auroral plasma process and SED gradients during the geomagnetic storm of 30 October 2003.

Some mechanisms for mid-latitude irregularities were also investigated. Perkins (1973) presented a model that was used to explain the mid-latitude structures. Later, Hamza (1999)



**Fig. 1.** The dots show the geographic distribution of the GPS receivers that were used in our TEC calculation. The “\*” marks the location of Wuhan station.

revisited the Perkins instability and introduced two new elements. Recently, Tsunoda and Cosgrove (2001) pointed out that the F layer and the sporadic E layer in the nighttime mid-latitude ionosphere must be considered electro-dynamically as a coupled system (Haldoupis et al., 1996; Tsunoda, 1998; Cosgrove and Tsunoda, 2001, 2002a, 2002b, 2003, 2004; Cosgrove et al., 2004).

In this paper, we investigate both the amplitude and phase scintillations of the GPS L1 band observed at the low-middle latitude station Wuhan ( $30.6^{\circ}$  N,  $114.4^{\circ}$  E,  $45.8^{\circ}$  Dip) during a period of intense geomagnetic storms, together with GPS TEC maps and Digisonde measurements. The observational results show that intensified scintillations are seen associated with the enhancement of the equatorial ionization anomaly (EIA) during the main phase of the geomagnetic storm. In addition, no scintillation is observed to be accompanied by the appearance of an Es layer during whole process of a geomagnetic storm. It seems that the existence of the Es layer inhibits the generation of spread-F. Current theories are used to interpret the observed results.

The paper is organized as follows. In Sect. 2, observational methods and data are described. In Sect. 3, analyzed results are given. Discussion and conclusions are presented in Sect. 4.

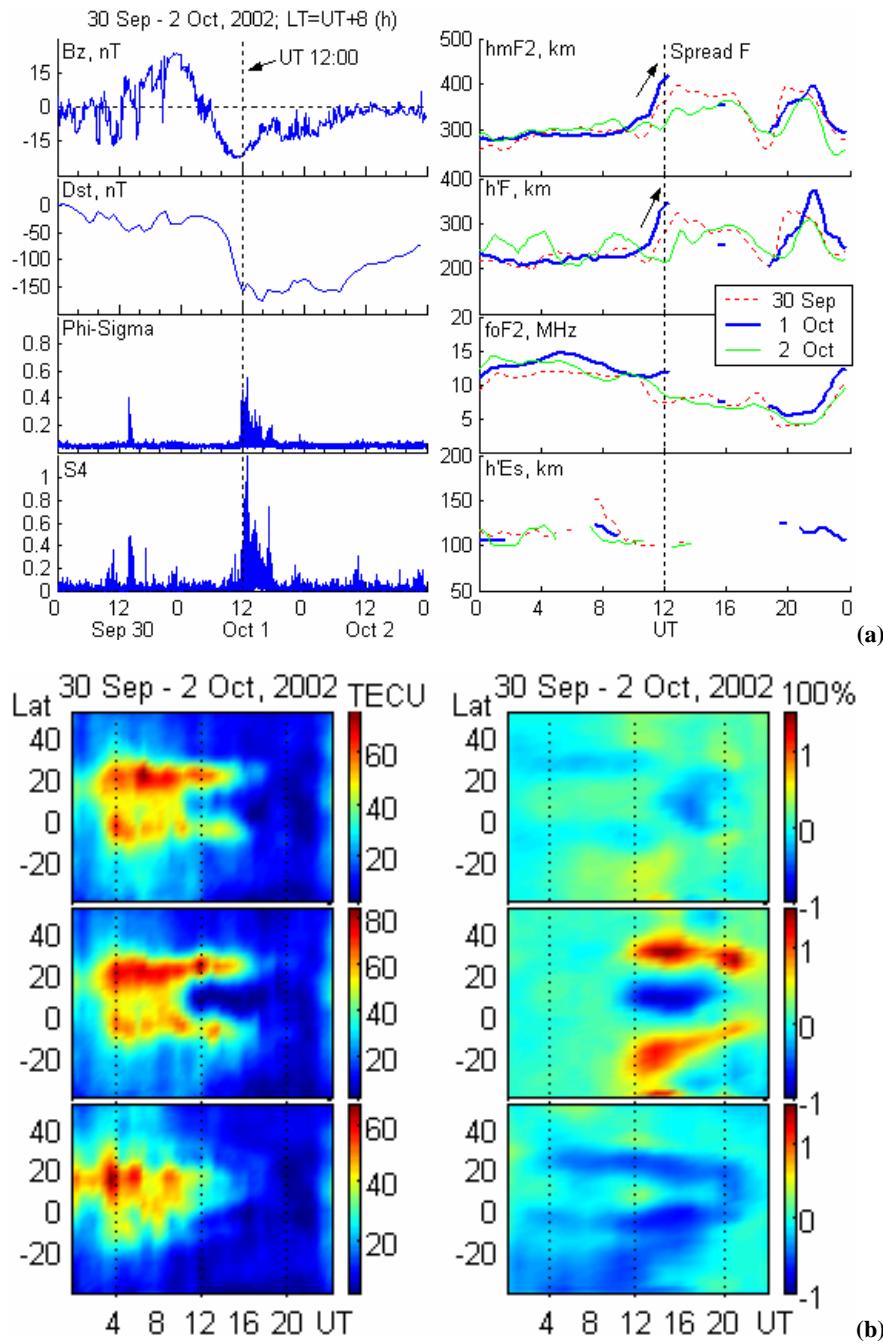
## 2 Data resources

In September 2002, a modified GPS receiver for recording the amplitude and phase of the L1 signal was set up at Wuhan. The receiver is a NovAtel Euro4 dual-frequency receiver version with OEM4 card and special firmware, which also yields the value of ionospheric TEC. The amplitude scintillation is monitored by computing the S4 index, which is the standard deviation of the received power normalized by its mean value. In our study, it is derived from the detrended received signal intensity. A fixed choice of a 0.1 Hz, 3 dB-cutoff frequency, for both phase and amplitude filtering, has been used here. A low phase noise, oven-controlled crystal oscillator performs the phase measurements by monitoring the standard deviation of the detrended carrier phase. Then phase scintillation (Phi-Sigma) is computed (Van et al., 1993). In all the cases in this paper, only the signals coming from satellites with an elevation angle greater than  $30^{\circ}$  and with a time of lock greater than 180 s were taken into account. The virtual height of the F layer,  $h'F$ , the peak height of F layer,  $hmF2$ , the critical frequency of the F layer,  $foF2$ , and virtual height of the sporadic E layer,  $h'Es$ , are measured by Digisonde-256 at the Wuhan station.

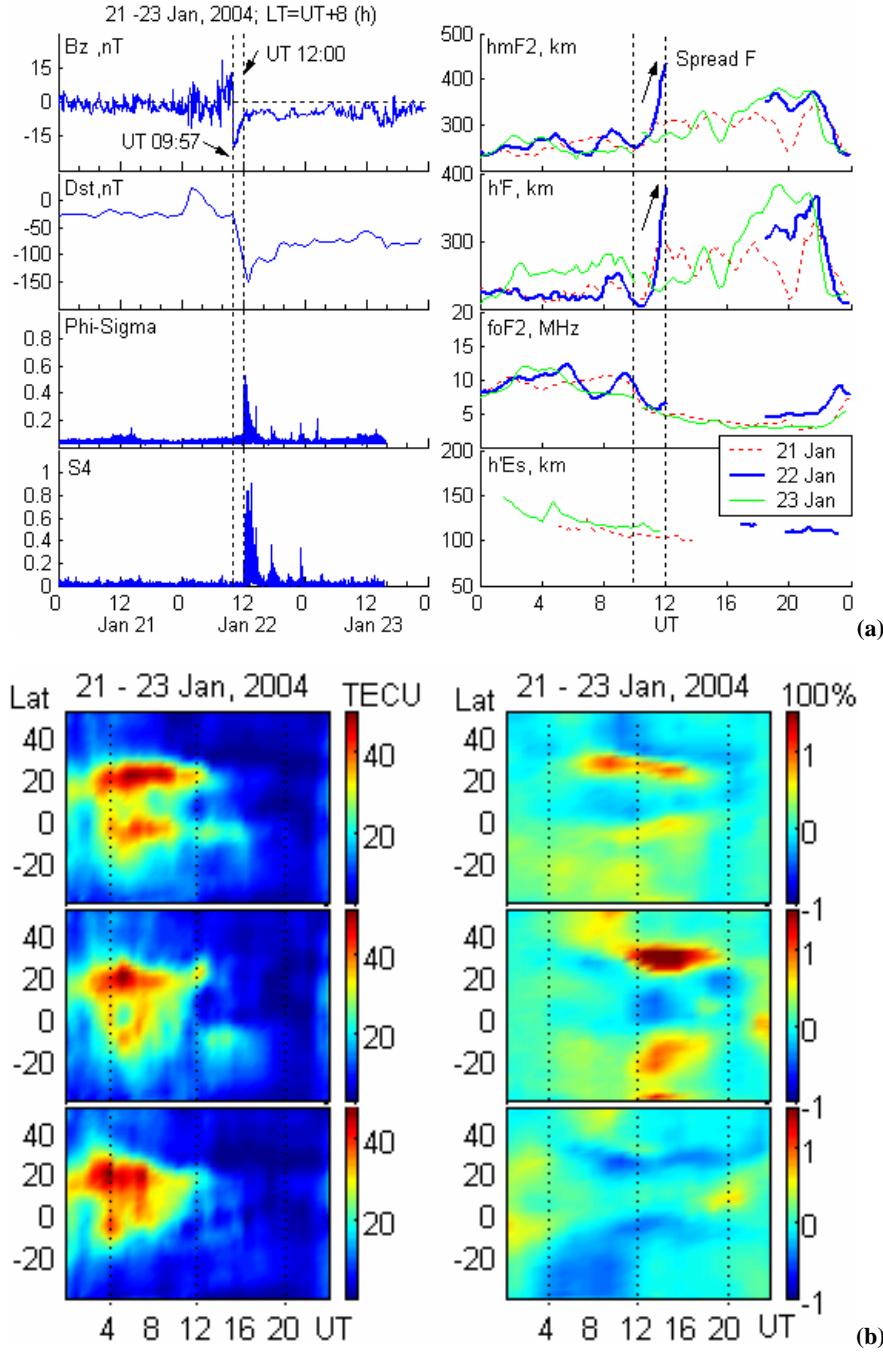
GPS data for TEC maps are obtained from a GPS network in the China, Southeast Asia and Australia regions. Due to the fact that observations at lower elevation angles have passed through a longer path through the ionosphere than those at high elevation, the slant TEC values were approximated into the equivalent vertical TEC values by assuming a thin-shell ionosphere at 400 km. Then through the GPS network, a nearest interpolation (by setting the value of an interpolated point to the value of the nearest existing data point) was employed to yield a TEC variation along geographical longitude  $110^{\circ}$  E ( $LT \approx UT + 8$ ), with a spatial resolution of about  $2.5^{\circ}$ . Thus, we obtained our TEC (Latitude versus UT) map (Zhao et al., 2005). The sampling interval of the TEC data for mapping is 30 min. The distribution of the GPS receivers is given in Fig. 1. Since vertical electron density distribution is generally not constant along the ray path, it should be pointed out that the calculation of TEC is perhaps only a rough approximation. It has been estimated that TEC can be in error by 10–20% when an elevation-dependent scaling factor is used at low elevation angles ( $5$ – $10^{\circ}$ ), and in regions of significant TEC gradients (Klobuchar et al., 1993). But for studying large-scale TEC variation, especially under conditions of intense geomagnetic disturbance, this approximation is enough for identifying common storm features (Ho et al., 1998). To show more clearly the TEC variation, we subtracted the 27-day smooth average value of TEC, and then obtained its relative deviation (RTEC):

$$RTEC(UT) = \frac{TEC(UT) - \overline{TEC(UT)}}{\overline{TEC(UT)}}.$$

Geomagnetic storms are characterized by a prolonged depression of the horizontal component (H) of the Earth's



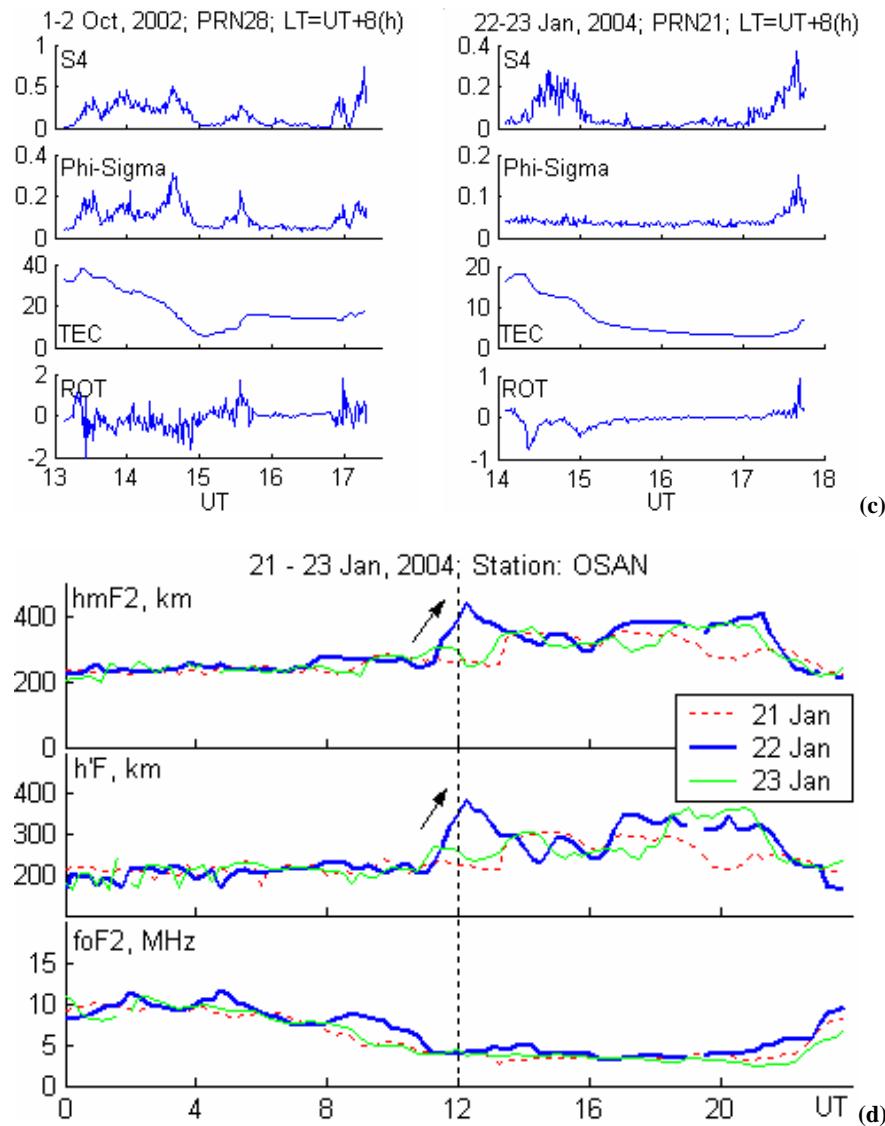
**Fig. 2.** (a)  $B_z$  component of the interplanetary magnetic field in geocentric-solar-magnetospheric (GSM) coordinates provided by the Advanced Composition Explorer (ACE) spacecraft (top panel, left) and  $D_{st}$  index (middle panel, left) obtained from CDAWeb and SPIDR; Phase scintillation index ( $\Phi$ -Sigma, middle panel, left) and Amplitude scintillation index ( $S4$ , bottom panel, left) measured from all GPS satellite signals observed at Wuhan; F layer peak height ( $hmF2$ , top panel, right), F layer virtual height ( $h'F$ , middle panel, right), F layer critical frequency ( $foF2$ , middle panel, right) and sporadic E layer virtual height ( $h'Es$ , bottom panel, right) obtained from Digisonde measurements at Wuhan for the period 30 September–2 October 2002. (b) Latitudinal variation of TEC along geographic longitude 110° E for 30 September (top panel, left), 1 October (middle pane, left), 2 October (bottom panel, left), 2002, and RTEC along geographic longitude 110° E for 30 September (top panel, right), 1 October (middle pane, right), 2 October 2002 (bottom panel, right), show the equatorial ionization anomaly (EIA) development.



**Fig. 3.** (a)  $B_z$  component of IMF,  $D_{st}$  index, Phi-Sigma and S4 index;  $hmF2$ ,  $h'F$ ,  $foF2$  and  $h'E_s$  for the period 21–23 January 2004. (b) Latitudinal variation of TEC (left) and RTEC (right) along geographic longitude 110° E for 21 January (top panel), 22 January (middle panel) and 23 January 2004 (bottom panel), show the EIA development.

magnetic field. The depression in H is characterized by the geomagnetic index  $D_{st}$  which lies between  $-100$  to  $-200$  nT for intense storms and less ( $-200$ ) for very intense storms. The southward turning of the interplanetary magnetic field (IMF)  $B_z$  ensures the transportation of solar wind energy into the Earth's magnetosphere, which is the primary cause

for the formation of geomagnetic storms (Allen et al., 1989; Tsurutani et al., 1992; Gonzalez et al., 1999). The  $D_{st}$  index is obtained from the Space Physics Interactive Data Resource and IMF  $B_z$  is obtained from Coordinated Data Analysis Web.



**Fig. 3.** Continued. (c) Plots at every 1-min interval of (1) S4 index (2) Phi-Sigma index (3) TEC (4) rate of TEC (ROT) for PRN 28 on 1–2 October, 2002 and PRN 21 on 22–23 January 2004. (d) Increases in  $hmF2$  and  $h'F$  observed at OSAN ( $37.1^\circ N, 127.0^\circ E$ ) on 22 January 2004 provide evidence of the rise in the F layer (Data obtained from SAO database).

### 3 Results

#### 3.1 Storm 1 and 2: 1 October 2002 and 22 January 2004

Here we combine the scintillations of two similar events in constructing a picture of the prompt effects associated with magnetospheric disturbances on a low-middle latitude. In Fig. 2a, the left four panels show the  $B_z$  component of IMF,  $D_{st}$ , phase and amplitude scintillation indices;  $hmF2$ ,  $h'F$ ,  $foF2$  and  $h'E$ s are presented in the right four panels for 30 September–2 October 2002. Figure 3a shows the same parameters for 20–22 January 2004. Both events are marked by major negative IMF  $B_z$  and a post-sunset steep height

increase of  $h'F$  and  $hmF2$ . The uplifting of F layer results in spread F and hence intense amplitude and phase scintillations.

The enhancement of EIA over the longitude  $110^\circ E$  is observed during each event. As shown in Figs. 2b and 3b, in comparing the 3 days, a clear EIA enhancement is seen at about 10:00 UT during the two storms, and the EIA enhancement of storm 1 and 2 extending to about 22:00 and 18:30 UT respectively relative to the quiet day. This behavior is likely associated with the intensification of the fountain effect, due to the penetration of an eastward electric field into the equatorial ionosphere, which forced a large amount of plasma to be uplifted at the magnetic equator and

subsequently diffused along the magnetic field line toward higher latitudes. For both events, Digisonde station at Wuhan registered a very significant increase of  $hmF2$  and  $h'F$ , followed by the occurrence of strong spread F (seen as diffuse F-region traces in ionograms). On 22 January 2004, Digisonde station at OSAN ( $37.1^{\circ}$ N,  $127.0^{\circ}$ E) simultaneously observed the same increase (Fig. 3d) in  $hmF2$  and  $h'F$ .

Figure 3c presents the results for PRN 28 on 1 October 2002 and PRN 21 on 22 January 2004. A close comparison of S4, TEC and ROT (rate of change of TEC per minute, Pi et al., 1997; Basu et al., 1999) for single PRN indicates that the enhanced scintillation structures correspond well with fluctuations of ROT. From Fig. 3c we can find that at post-midnight, moderate scintillations still exist with TEC fluctuations.

### 3.2 Storm 3 and 4: 16 July and 20 November 2003

In comparison to intense scintillation activity during storms 1 and 2, no scintillation was observed during the storms occurred on 16 July 2003 and 20 November 2003. Figures 4 and 5 show IMF  $B_z$ ,  $D_{st}$ , Phi-Sigma, S4,  $hmF2$ ,  $h'F$ ,  $foF2$ ,  $h'Es$  indices and TEC maps for the two events, respectively.

As shown in Fig. 4a, during storm 3, it is evident that as time advances, IMF  $B_z$  turned southward at around 02:00 UT on 16 July 2003. Except for several northward excursions, it remained negative for several hours.  $D_{st}$  reached its maximum  $-117$  nT at 12:00 UT on 16 July. Comparing the 3 days, Fig. 4b shows that the EIA began to intensify at 10:00 UT and this lasted until 18:00 UT during the storm day, similar to storms 1 and 2. But measurements of  $hmF2$  and  $h'F$  show a weak increase which started at 10:20 UT and finished at 17:00 UT, and the sporadic E layer ( $h'Es$ ) was present when the F layer rose. The weak increase of the F layer was probably caused by a storm-induced enhancement of the eastward electric field, but the increase was not apparent; this is presumably the reason why the concurrent existence of the Es layer inhibited the upward plasma drift.

For storm 4, Fig. 5a illustrates that IMF  $B_z$  turned southward at 02:20 UT on 20 November 2003, and then turned northward at 10:00 UT on 20. After approximately 50 min, IMF  $B_z$  became further southward at 10:48 UT on 20 November, and remained negative until 00:10 UT on 21 November.  $D_{st}$  became much depressed and reached its maximum at  $-472$  nT at 19:00 UT on 20 November. The EIA is also shown to be enhanced at around 11:00 UT on 20 November and this ended at about 17:00 UT on 20 November (Fig. 5b).  $hmF2$  and  $h'F$  show two large increments during 11:00–14:00 UT and 14:00–17:00 UT, respectively. It is likely that an eastward penetration electric field to the first elevation of the F layer when IMF  $B_z$  turned southward, and the second elevation which during the midnight sector could be the result of a disturbance dynamo electric field (Blanc and Richmond, 1980; Fejer and Emmert, 2003).

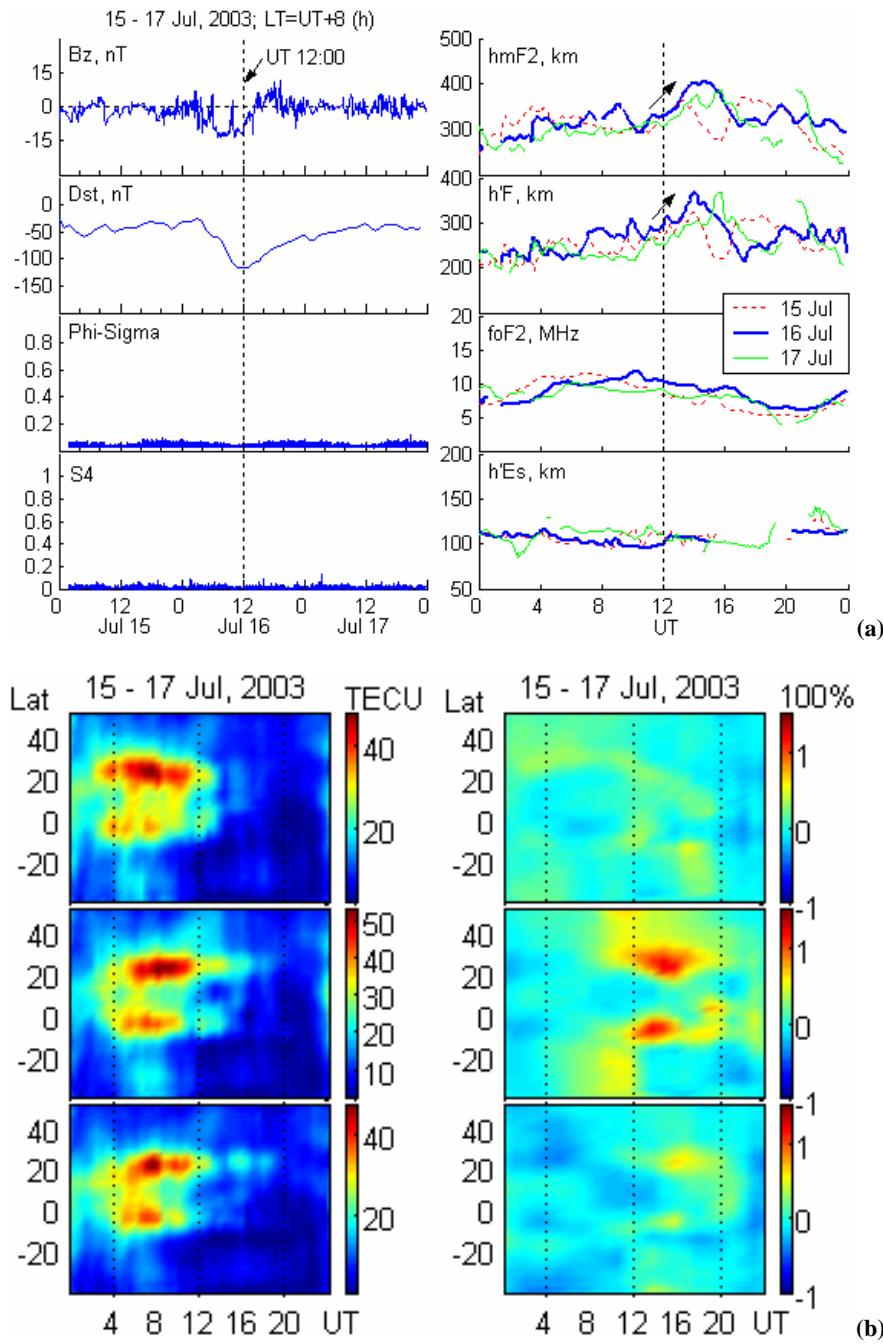
The sporadic E layer ( $h'Es$ ) also exists when the F layer rose during storm 4.

For these two events, the EIA intensification can be likely seen during the storm day as storms 1 and 2. The height variation of the F layer is not apparent relative to the adjacent 2 days for storm 3. For storm 4,  $hmF2$  and  $h'F$  show two increments. But the scintillation activity was essentially dormant. The S4 index and phase scintillation index (Phi-Sigma) rarely exceed 0.1. The reason is presumably that the existence of the Es layer in the post-sunset inhibits the upward plasma drift and leads to a significant suppression of spread F.

## 4 Discussions and conclusions

Some excellent work on ionospheric scintillations during storms has been published. Dabas et al. (1989), Aarons et al. (1991, 1997) found that  $D_{st}$  variations are the best criteria to determine the development of the equatorial irregularities. Numerous works present the experimental and theoretical relations between the solar wind parameters and the auroral ionosphere. The studies of Zmuda and Armstrong (1974), Iigima and Potemra (1978), Foster et al. (1989) and other authors concluded that the field-aligned currents (FAC) constitute a major interconnection between the magnetosphere and the polar ionosphere. In theory, Zakharov et al. (1989), Denisenko and Zamai (1992) had shown that the electric fields from the FAC could penetrate to the equatorial ionosphere and explain the equatorial electric field variations, as well as Kikuchi et al. (1996), Sizova (2002), using geomagnetic data, and Sizova and Pudovkin (2000), using ionospheric data had proved the penetration. According to the model of Sizova and Pudovkin (2000), electric fields of the field-aligned currents penetrate through the mid-latitude to the low-latitude ionosphere and create the additional equatorial electric fields. Under the additional eastward electric field, the height of the F region rises high enough where scintillation irregularities can be generated.

Although the virtual height of the F layer  $h'F$  observations does not differentiate between the effects of an eastward electric field or of an equatorward surge of the neutral wind, Tanaka (1986) once analyzed the storm-time ionospheric response and provided criteria for differentiating between the effects of electric fields and winds. The equatorial anomalies are created as an eastward electric field uplifts the equatorial ionosphere, which subsequently diffuses down the magnetic field to produce density enhancements at somewhat higher latitudes. The response to winds propagating equatorward from a high-latitude disturbance, however, is characterized by increases in  $h'F$  which travel from north to south with noticeable time dispersions and no change in  $foF2$  or TEC. From storms 1 to 4, significant changes in TEC, as well as an enhancement of the EIA were observed. These provide

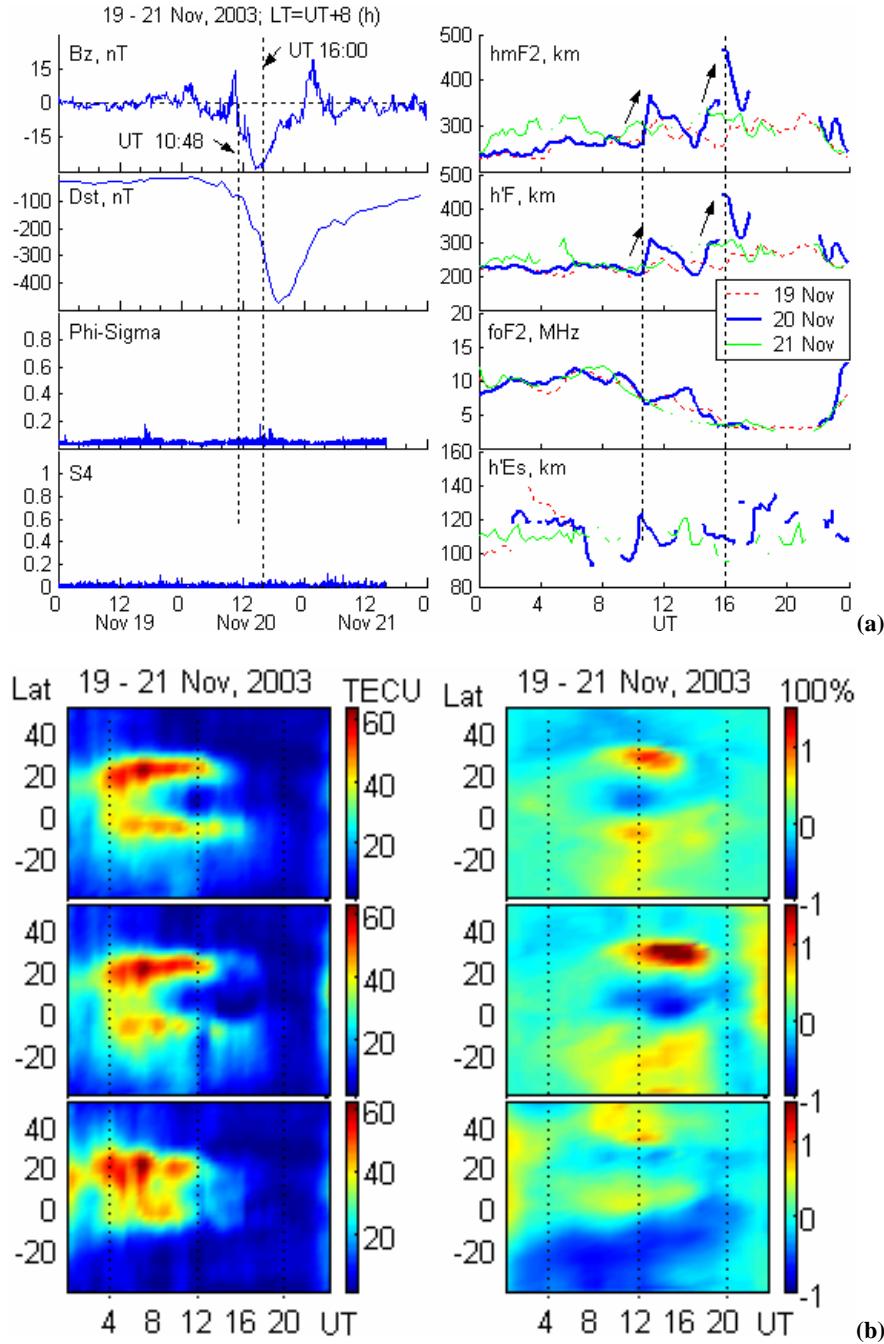


**Fig. 4.** (a)  $B_z$  component of IMF,  $D_{st}$  index, Phi-Sigma and S4 index;  $hmF2$ ,  $h'F$ ,  $foF2$  and  $h'E_s$  for the period 15–17 July, 2003. (b) Latitudinal variation of TEC (left) and RTEC (right) along geographic longitude 110° E for 15 July (top panel), 16 July (middle panel) and 17 July (bottom panel), 2003, show the EIA development.

evidences of a concentration of a storm-enhanced eastward electric field.

As we can see from  $B_z$ ,  $D_{st}$ , Phi-Sigma, S4,  $hmF2$ ,  $h'F$  and  $h'E_s$  dynamics presented in Figs. 2a–5a, the events of 16 July 2003 and 20 November 2003 did not exhibit any strong spread F or scintillation activity, while the events of

1 October 2002 and 22 January 2004 demonstrated the intense scintillation signatures associated with the presence of the spread F; it appears that the existence of the sporadic E layer has significant suppression effect on spread F and inhibits the process required for the growth of plasma irregularities. The modulation of the ESF occurrence by the Es



**Fig. 5.** (a)  $B_z$  component of IMF,  $D_{st}$  index, Phi-Sigma and S4 index;  $hmF2$ ,  $h'F$ ,  $foF2$  and  $h'Es$  for the period 19–21 November 2003. (b) Latitudinal variation of TEC (left) and RTEC (right) along geographic longitude 110° E for 19 November (top panel), 20 November (middle panel) and 21 November (bottom panel), 2003, show the EIA development.

layer was first proposed by Tsunoda (see Sultan, 1996); when a sporadic E layer develops, its contribution to the Pedersen conductivity may also become significant. Stephan et al. (2002) showed that the Es layer can change the Pedersen conductivity ratio and thus the growth rate of the gravitational Rayleigh–Taylor instability (GRT); the magnitude of

the change is dependent on the plasma densities, as well as the height and time of Es layer occurrence.

By using GPS ionospheric scintillations, TEC measurements and Digisonde data to investigate the responses of the ionosphere to the geomagnetic storms of 1 October 2002, 16 July 2003, 20 November 2003, and 22 January 2004, we

observed a prompt uplifting of the ionospheric F layer at low-middle latitude, related to the occurrence of the eastward penetration electric field or the disturbance dynamo electric field during storms. Intense scintillation occurred and even extended up to post-midnight with the co-existence of large- and small-scale irregularities during 1 October 2002 and 22 January 2004. It is proposed that an enhancement of the eastward electric field may intensify the upward drift of plasma, followed by transport to higher latitude. Then the irregularities associated with spread F can be generated, and ionospheric scintillations can be observed there. But for the storm on 16 July 2003 the uplifting of the ionospheric F layer was not apparent relative to the adjacent 2 days; for the storm on 20 November 2003,  $hmF2$  and  $h'F$  show two increments. Both of them did not generate spread F and to scintillations. The reason is probably that the sporadic E layer played a role in the generation of the ionospheric scintillation by changing the Pederson conductivity ratio and suppressing the upward plasma drift, thus reducing the growth rate of the irregularities associated with spread-F. Also, we should point out that there may be other reasons for example, the eastward electric field perturbation is not strong enough to elevate the F layer enough. The relationship between the Es layer and spread F is perhaps one of the keys to predicting ionospheric scintillation during a period of intense geomagnetic activity. This needs to be further studied and more detail investigations are required.

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